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Physical Chemistry 2

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Lecture 8. The Rigid Rotor (2)

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Topics in Lecture 8

Rigid rotor model

Spherical harmonics

IR spectroscopy

In Atkins' Physical Chemistry (11th ed.),

7F Rotational motion

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$$\hat{H} = -\frac{\hbar^2}{2\mu} \left[\frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] = -\frac{\hbar^2}{2I} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

Therefore, we can derive \hat{L}^2 operator from the relation

$$\hat{H} = \frac{\hat{L}^2}{2I}, \quad \hat{L}^2 = -\hbar^2 \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

The equation is fully angular. It depends to only angular variables.

$$\hat{H}Y(\theta,\phi) = -\frac{\hbar^2}{2I} \left[\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y(\theta,\phi)}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2 Y(\theta,\phi)}{\partial\phi^2} \right]$$

The eigenfunction of this equation is very well-known special function, called spherical harmonics.



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Assume that

$$\hat{H}Y(\theta,\phi) = -\frac{\hbar^2}{2I} \left[\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y(\theta,\phi)}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2 Y(\theta,\phi)}{\partial\phi^2} \right] = -\frac{\hbar^2 l(l+1)}{2I} Y(\theta,\phi)$$

Therefore

$$\frac{1}{Y} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y(\theta, \phi)}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y(\theta, \phi)}{\partial \phi^2} \right] = -l(l+1)$$

$$\sin\theta \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y}{\partial\theta} \right) + \frac{\partial^2 Y}{\partial\phi^2} = -l(l+1)Y\sin^2\theta$$

Separation of variables: $Y(\theta, \phi) = \Theta(\theta)\Phi(\phi)$



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Plug in and divide by $\Theta\Phi$.

$$\left\{ \frac{1}{\Theta} \left[\sin \theta \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) \right] + l(l+1)\sin^2 \theta \right\} + \frac{1}{\Phi} \frac{\partial^2 \Phi}{\partial \phi^2} = 0$$

Set separation constant

$$\frac{1}{\Theta} \left[\sin \theta \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) \right] + l(l+1)\sin^2 \theta = m^2, \qquad \frac{1}{\Phi} \frac{\partial^2 \Phi}{\partial \phi^2} = -m^2$$

The ϕ equation is easy: we have solved this already.

$$\Phi(\phi) = e^{im\phi}, \quad (m \in \mathbb{Z})$$

The θ equation is:

$$\sin\theta \frac{d}{d\theta} \left(\sin\theta \frac{d\Theta}{d\theta} \right) + (l(l+1)\sin^2\theta - m^2)\Theta = 0$$



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The solution of this equation is

$$\sin\theta \frac{d}{d\theta} \left(\sin\theta \frac{d\Theta}{d\theta} \right) + (l(l+1)\sin^2\theta - m^2)\Theta = 0 \implies \Theta(\theta) = AP_l^m(\cos\theta)$$

where P_l^m is the associated Legendre function.

$$P_{l}^{m}(x) = (-1)^{m} (1 - x^{2})^{m/2} \left(\frac{d}{dx}\right)^{m} P_{l}(x)$$

$$P_{l}(x) = \frac{1}{2^{l} l!} \left(\frac{d}{dx}\right)^{l} (x^{2} - 1)^{l}$$

They are defined from the Legendre polynomial P_l , where Legendre polynomials are defined by Rodrigues formula.

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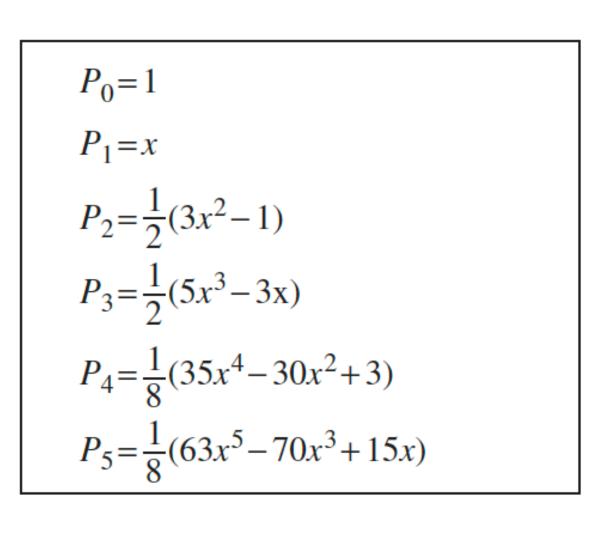
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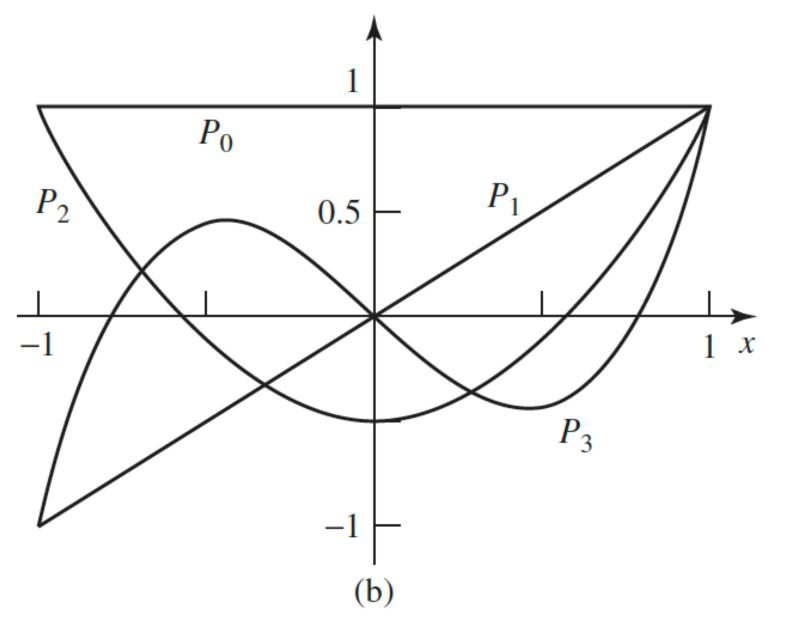
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The Legendre polynomials are well-known special functions.

Table 4.1: The first few Legendre polynomials, $P_{\ell}(x)$: (a) functional form, (b) graph.





Spherical harmonics

Table 4.2: Some associated Legendre functions, $P_{\ell}^{m}(\cos\theta)$: (a) functional form, (b) graphs of $r = |P_{\ell}^{m}(\cos\theta)|$ (in these plots r tells you the magnitude of the function in the direction θ ; each figure should be rotated about the z axis).

Table 4.3: The first few spherical harmonics, $Y_{\ell}^{m}(\theta, \phi)$.

$$Y_0^0 = \left(\frac{1}{4\pi}\right)^{1/2} \qquad Y_2^{\pm 2} = \left(\frac{15}{32\pi}\right)^{1/2} \sin^2\theta e^{\pm 2i\phi}$$

$$Y_1^0 = \left(\frac{3}{4\pi}\right)^{1/2} \cos\theta \qquad Y_3^0 = \left(\frac{7}{16\pi}\right)^{1/2} (5\cos^3\theta - 3\cos\theta)$$

$$Y_1^{\pm 1} = \mp \left(\frac{3}{8\pi}\right)^{1/2} \sin\theta e^{\pm i\phi} \qquad Y_3^{\pm 1} = \mp \left(\frac{21}{64\pi}\right)^{1/2} \sin\theta (5\cos^2\theta - 1) e^{\pm i\phi}$$

$$Y_2^0 = \left(\frac{5}{16\pi}\right)^{1/2} (3\cos^2\theta - 1) \qquad Y_3^{\pm 2} = \left(\frac{105}{32\pi}\right)^{1/2} \sin^2\theta \cos\theta e^{\pm 2i\phi}$$

$$Y_2^{\pm 1} = \mp \left(\frac{15}{8\pi}\right)^{1/2} \sin\theta \cos\theta e^{\pm i\phi} \qquad Y_3^{\pm 3} = \mp \left(\frac{35}{64\pi}\right)^{1/2} \sin^3\theta e^{\pm 3i\phi}$$

By multiplying associated Legendre function and $\Phi(\phi)$, we yield spherical harmonics.

$$Y_l^m(\theta,\phi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{im\phi} P_l^m(\cos\theta)$$



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The energy level of rigid rotor is

$$\hat{H}Y(\theta,\phi) = -\frac{\hbar^2}{2I} \left[\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y(\theta,\phi)}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2 Y(\theta,\phi)}{\partial\phi^2} \right] = -\frac{\hbar^2 l(l+1)}{2I} Y(\theta,\phi)$$

Therefore,

$$E_J = \frac{\hbar^2}{2I}J(J+1)$$
 $(J=0,1,2,...)$

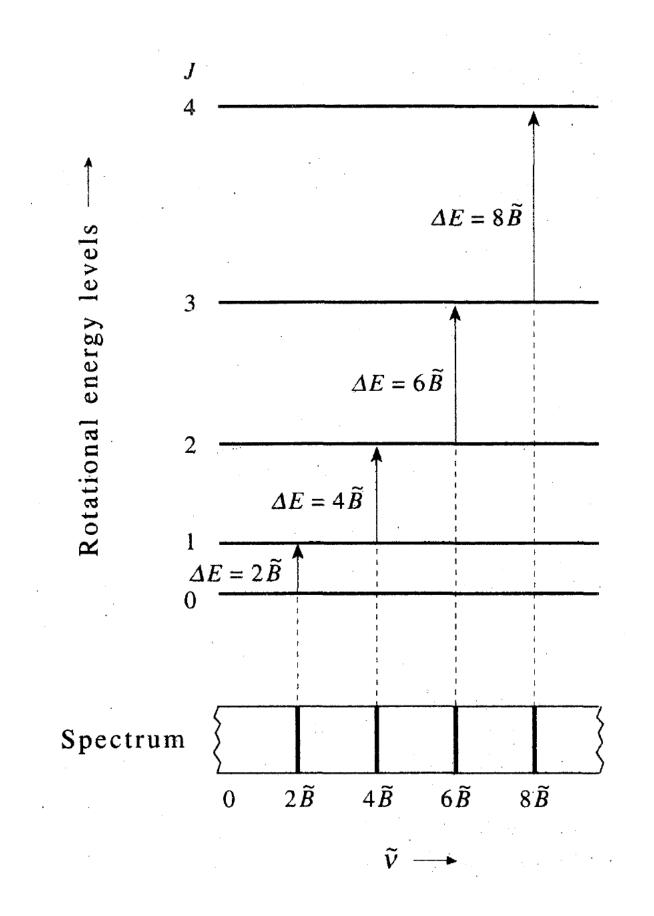
Selection rules are

$$\Delta J = \pm 1$$

$$\Delta E = E_{J+1} - E_J = \frac{h^2}{4\pi^2 I} (J+1)$$

$$\nu = \frac{h}{4\pi^2 I} (J+1) = 2B(J+1), B = \frac{h}{8\pi^2 I}$$

$$\tilde{\nu} = 2\tilde{B}(J+1), \quad \tilde{B} = \frac{h}{8\pi^2 cI}$$



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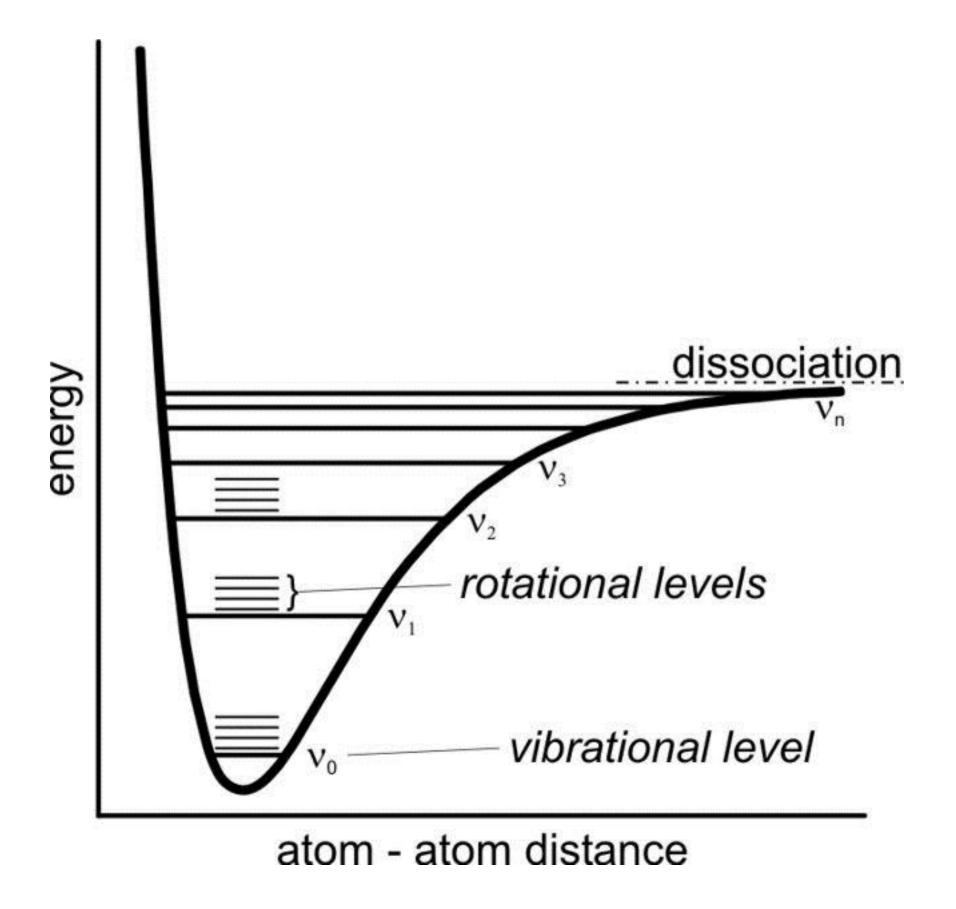
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IR spectroscopy

Vibrational-rotational energy level



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