

1.1) Show that the Dirac matrices must be even dimensional. Hint: show that the eigenvalues of β are all ± 1 , and that $\text{Tr } \beta = 0$. To show that $\text{Tr } \beta = 0$, consider, e.g., $\text{Tr } \alpha_1^2 \beta$. Similarly, show that $\text{Tr } \alpha_i = 0$.

$$\alpha_i \alpha_j + \alpha_j \alpha_i = 2\delta_{ij} \Rightarrow \alpha_i^2 = 1 \quad \forall i$$

$$\alpha_i \beta + \beta \alpha_i = 0$$

$$\beta \alpha_i \beta + \beta^2 \alpha_i = 0, \quad \beta^2 = 1$$

$$\underbrace{\text{tr } \beta \alpha_i \beta + \text{tr } \alpha_i}_{\text{tr } \alpha_i} = 0 \Rightarrow \text{tr } \alpha_i = 0 \quad \forall i$$

$$M^2 = 1$$

$$= (U^\dagger M_x U)^2$$

$$= U^\dagger M_x^2 U$$

$$\therefore U U^\dagger = 1$$

$$= M_x^2$$

$$= \text{diag}(\lambda_1^2, \lambda_2^2, \dots) \Rightarrow \lambda_i^2 = 1 \quad \forall i$$

$$\therefore \alpha_{ix} = \underbrace{\text{diag}(1, \dots, 1)}_{N_+}, \underbrace{\text{diag}(-1, \dots, -1)}_{N_-}$$

$$\text{tr } \alpha_i = 0 \Rightarrow N_+ = N_- \quad \therefore \mathcal{D}(\alpha) = 2N_+ : \text{even-dimensional}$$

$$\underbrace{\alpha_i \alpha_i \beta + \alpha_i \beta \alpha_i}_1 = 0$$

$$\text{tr } \beta + \underbrace{\text{tr } \alpha_i \beta \alpha_i}_{\text{tr } \beta} = 0 \Rightarrow \text{tr } \beta = 0$$

$$\left. \begin{array}{l} \text{tr } \beta = 0 \\ \beta^2 = 1 \end{array} \right\} \mathcal{D}(\beta) = 2N_+ : \text{even-dimensional similarly to } \alpha_i$$

Since we have already eliminated the case of 2×2 , the minimum dimension should be 4×4 .

1.2) With the hamiltonian of eq. (1.32), show that the state defined in eq. (1.33) obeys the abstract Schrödinger equation, eq. (1.1), if and only if the wave function obeys eq. (1.30). Your demonstration should apply both to the case of bosons, where the particle creation and annihilation operators obey the commutation relations of eq. (1.31), and to fermions, where the particle creation and annihilation operators obey the anticommutation relations of eq. (1.38).

$$H = \int d^3x a^\dagger(\mathbf{x}) \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{x}) \right) a(\mathbf{x}) + \frac{1}{2} \int d^3x d^3y V(\mathbf{x} - \mathbf{y}) a^\dagger(\mathbf{x}) a^\dagger(\mathbf{y}) a(\mathbf{y}) a(\mathbf{x}). \quad (1.32)$$

$$|\psi, t\rangle = \int d^3x_1 \dots d^3x_n \psi(\mathbf{x}_1, \dots, \mathbf{x}_n; t) a^\dagger(\mathbf{x}_1) \dots a^\dagger(\mathbf{x}_n) |0\rangle, \quad (1.33)$$

$$i\hbar \frac{\partial}{\partial t} |\psi, t\rangle = H |\psi, t\rangle, \quad (1.1)$$

$$i\hbar \frac{\partial}{\partial t} \psi = \left[\sum_{j=1}^n \left(-\frac{\hbar^2}{2m} \nabla_j^2 + U(\mathbf{x}_j) \right) + \sum_{j=1}^n \sum_{k=1}^{j-1} V(\mathbf{x}_j - \mathbf{x}_k) \right] \psi, \quad (1.30)$$

$$\begin{aligned} [a(\mathbf{x}), a(\mathbf{x}')] &= 0, \\ [a^\dagger(\mathbf{x}), a^\dagger(\mathbf{x}')] &= 0, \\ [a(\mathbf{x}), a^\dagger(\mathbf{x}')] &= \delta^3(\mathbf{x} - \mathbf{x}'), \end{aligned} \quad (1.31)$$

$$\begin{aligned} \{a(\mathbf{x}), a(\mathbf{x}')\} &= 0, \\ \{a^\dagger(\mathbf{x}), a^\dagger(\mathbf{x}')\} &= 0, \\ \{a(\mathbf{x}), a^\dagger(\mathbf{x}')\} &= \delta^3(\mathbf{x} - \mathbf{x}'), \end{aligned} \quad (1.38)$$

See code_1 and code_3.

1.3) Show explicitly that $[N, H] = 0$, where H is given by eq. (1.32) and N by eq. (1.35).

$$H = \int d^3x a^\dagger(\mathbf{x}) \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{x}) \right) a(\mathbf{x}) + \frac{1}{2} \int d^3x d^3y V(\mathbf{x} - \mathbf{y}) a^\dagger(\mathbf{x}) a^\dagger(\mathbf{y}) a(\mathbf{y}) a(\mathbf{x}). \quad (1.32)$$

$$N = \int d^3x a^\dagger(\mathbf{x}) a(\mathbf{x}) \quad (1.35)$$

See code_3.