

Ch 1: Problems

1.1) Show that the Dirac matrices must be even-dimensional. Hint: Show that the eigenvalues of β are all ± 1 and that $\text{tr } \beta = 0$. To show that $\text{tr } \beta = 0$, consider e.g. $\text{tr } \alpha_i^2 \beta = 0$. Similarly, show that $\text{tr } \alpha_i = 0$.

Solution

$$\alpha_i \alpha_j + \alpha_j \alpha_i = 2\delta_{ij} \Rightarrow \alpha_i^2 = 1 \quad \forall i$$

$$\alpha_i \beta + \beta \alpha_i = 0$$

$$\beta \alpha_i \beta + \beta^2 \alpha_i = 0, \quad \beta^2 = 1$$

$$\underbrace{\text{tr } \beta \alpha_i \beta} + \text{tr } \alpha_i = 0 \Rightarrow \text{tr } \alpha_i = 0 \quad \forall i$$

$$\text{tr } \alpha_i$$

$$M^2 = 1 = (U^\dagger M_* U)^2 = U^\dagger M_*^2 U$$

$$\therefore U U^\dagger = 1 = M_*^2 = \text{diag}(\lambda_1^2, \lambda_2^2, \dots) \Rightarrow \lambda_i^2 = 1 \quad \forall i$$

$$\therefore \alpha_{i*} = \text{diag}(\underbrace{1, \dots, 1}_{N_+}, \underbrace{-1, \dots, -1}_{N_-})$$

$$\text{tr } \alpha_i = 0 \Rightarrow N_+ = N_-$$

$$\therefore D(\alpha) = 2N_+ : \text{even-dimensional}$$

$$\underbrace{\alpha_i \alpha_i}_1 \beta + \alpha_i \beta \alpha_i = 0$$

$$\text{tr } \beta + \underbrace{\text{tr } \alpha_i \beta \alpha_i}_{\text{tr } \beta} = 0 \Rightarrow \text{tr } \beta = 0$$

$$\left. \begin{array}{l} \text{tr } \beta = 0 \\ \beta^2 = 1 \end{array} \right\} D(\beta) = 2N_+ : \text{even-dimensional} \\ \text{similarly to } \alpha_i$$

Since we have already eliminated the case of 2×2 , the minimum dimension should be 4×4 .

1.2) w/ the hamiltonian

$$H = \int d^3x \, a^\dagger(\vec{x}) \left[-\frac{\hbar^2}{2m} \vec{\nabla}^2 + U(\vec{x}) \right] a(\vec{x})$$

$$+ \frac{1}{2} \int d^3x \, d^3y \, V(\vec{x} - \vec{y}) \, a^\dagger(\vec{x}) \, a^\dagger(\vec{y}) \, a(\vec{y}) \, a(\vec{x})$$

show that the state defined by

$$|\Psi(t)\rangle = \int d^3x_1 \dots d^3x_n \Psi(\vec{x}_1, \dots, \vec{x}_n, t) a^\dagger(\vec{x}_1) \dots a^\dagger(\vec{x}_n) |0\rangle$$

obey the Schrödinger equation

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = H |\Psi(t)\rangle$$

iff the wavefunction obeys

$$i\hbar \frac{\partial}{\partial t} \Psi = \left\{ \sum_{j=1}^n \left[-\frac{\hbar^2}{2m} \nabla_j^2 + U(\vec{x}_j) \right] + \sum_{j=1}^n \sum_{k=1}^{j-1} V(\vec{x}_j - \vec{x}_k) \right\} \Psi$$

Your demonstration should apply both to the case of bosons, where the particle creation and annihilation operators obey the commutation relations

$$[a(\vec{x}), a(\vec{x}')] = 0$$

$$[a^\dagger(\vec{x}), a^\dagger(\vec{x}')] = 0$$

$$[a(\vec{x}), a^\dagger(\vec{x}')] = \delta^3(\vec{x} - \vec{x}')$$

and to the fermions, where the particle creation and annihilation operators obey

the anticommutation relations

$$\{a(\vec{k}), a(\vec{k}')\} = 0$$

$$\{a^\dagger(\vec{k}), a^\dagger(\vec{k}')\} = 0$$

$$\{a(\vec{k}), a^\dagger(\vec{k}')\} = \delta^3(\vec{k} - \vec{k}')$$

Solution See code-1 and code-3

1.3) Show explicitly that $[N, H] = 0$, where H is given in the previous problem and

$$N = \int d^3x \, a^\dagger(\vec{x}) a(\vec{x})$$

Solution See code-2.

