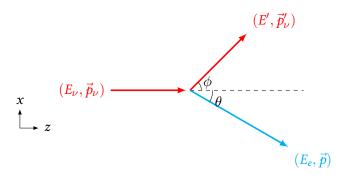
$\nu_{\mu}e$ kinematics & dynamics

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A neutrino collides with an electron at rest. Write down the 4-momenta of the neutrino and the electron in the lab frame. Suppose the neutrino is massless. Before collision we have $(\hbar=c=1)$

$$\mathbb{p}_{\nu} = (E_{\nu}, 0, 0, E_{\nu}) \tag{1}$$

$$\mathbb{p}_e = (m,0,0,0) \tag{2}$$

, where m is the electron rest mass, and E_{ν} is the total energy of the incident neutrino. After collision, we have

$$p'_{\nu} = (E', E' \sin \phi, 0, E' \cos \phi) \tag{3}$$

$$p'_{e} = (E_{e}, -p\sin\theta, 0, p\cos\theta)$$
 (4)

The total 4-momentum before collision is

$$\mathbb{P} = (E_{\nu} + m, 0, 0, E_{\nu}) \tag{5}$$

The total 4-momentum after collision is

$$\mathbb{P}' = (E' + E_e, -p\sin\theta + E'\sin\phi, 0, p\cos\theta + E'\cos\phi)$$
 (6)

Since total 4-momentum is conserved before and after collision, we have equations

$$E' + E_e = E_{\nu} + m \tag{7}$$

$$-p\sin\theta + E'\sin\phi = 0 \tag{8}$$

$$p\cos\theta + E'\cos\phi = E_{\nu} \tag{9}$$

Rearranging equations (8) and (9), squaring, and adding them, we eliminate ϕ and obtain

$$E'^2 = E_{\nu}^2 - 2E_{\nu}p\cos\theta + p^2 \tag{10}$$

If we eliminate E' with Eq. (7), and employ the energy-momentum relation $E_e^2 = p^2 + m^2$ once, we obtain

$$p\cos\theta = (E_e - m)\left(1 + \frac{m}{E_\nu}\right) \tag{11}$$

By squaring Eq. (11), and employing the energy-momentum relation again, we arrive at the answer

$$\frac{E_e}{m} = \frac{\left(1 + \frac{m}{E_\nu}\right)^2 + \cos^2 \theta}{\left(1 + \frac{m}{E_\nu}\right)^2 - \cos^2 \theta}$$
(12)

If we choose to eliminate θ and E_e , we get

$$E' = \frac{E_{\nu} m}{E_{\nu} (1 - \cos \phi) + m} \tag{13}$$

Rearranging and dividing Eq. (8) by Eq. (9), we have

$$\cot \theta = \frac{\frac{E_{\nu}}{E'} - \cos \phi}{\sin \phi} = \left(1 + \frac{E_{\nu}}{m}\right) \tan \left(\frac{\phi}{2}\right) \tag{14}$$

, where we have used Eq. (13) for E_{ν}/E' .

For a neutrino scattering angle $0 \le \phi \le \pi$, the range of the electron scattering angle is

$$0 \le \theta \le \frac{\pi}{2} \tag{15}$$

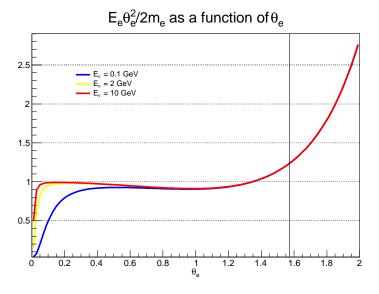
kinematical cut I

For small θ , Eq. (12) becomes $E_e\theta^2/2m\approx 1$. However, since we have the *exact* formula for E_e , we can actually plot the function

$$f(\theta) = \frac{E_e \theta^2}{2m} = \frac{\left(1 + \frac{m}{E_\nu}\right)^2 + \cos^2 \theta}{\left(1 + \frac{m}{E_\nu}\right)^2 - \cos^2 \theta} \cdot \frac{\theta^2}{2}$$
(16)

and see how close it is to 1.

kinematical cut II



$\nu_{\mu}e$ dynamics I

The differential cross section of $\nu_{\mu}e$ scattering is ¹

$$\frac{d\sigma}{d\cos\theta} = \sigma_0 \frac{4E_{\nu}^2 (m_e + E_{\nu})^2 \cos\theta}{\left[(m_e + E_{\nu})^2 - E_{\nu}^2 \cos^2\theta \right]^2} \times \left[g_1^2 + g_2^2 \left(1 - \frac{2m_e E_{\nu} \cos^2\theta}{(m_e + E_{\nu})^2 - E_{\nu}^2 \cos^2\theta} \right)^2 - g_1 g_2 \frac{2m_e^2 \cos^2\theta}{(m_e + E_{\nu})^2 - E_{\nu}^2 \cos^2\theta} \right]$$
(17)

, where

$$\sigma_0 = \frac{2G_F^2 m_e^2}{\pi} \approx 88.06 \times 10^{-46} \text{ cm}^2$$

$$g_1 = -\frac{1}{2} + \sin^2 \theta_W \approx -0.27$$

$$g_2 = \sin^2 \theta_W \approx 0.23$$

¹Giunti C. and Kim C. W., Fundamentals of Neutrino Physics and Astrophysics, 2007

$u_{\mu}e$ dynamics II

