

Lions Running Ramps:

Q. In the original derivation of the Fokker-Planck equation, we have

$$\Delta \psi = \frac{\partial \psi}{\partial t} + \frac{\partial \psi}{\partial x} v + \frac{1}{2} \frac{\partial^2 \psi}{\partial x^2} v^2$$

But now, we have $\Delta \psi = 0$

$$\Delta \psi(x, x_2) = \frac{\partial \psi}{\partial t} + \left\{ \frac{\partial \psi}{\partial x_1} v_1 + \frac{1}{2} \frac{\partial^2 \psi}{\partial x_1^2} v_1^2 \right\} + \left\{ \frac{\partial \psi}{\partial x_2} v_2 + \frac{1}{2} \frac{\partial^2 \psi}{\partial x_2^2} v_2^2 \right\}$$

$$\frac{\partial \psi}{\partial t}(x, x_2, t) = 0 \left(\frac{\partial v_1^2}{\partial x_1^2} + \frac{\partial v_2^2}{\partial x_2^2} \right) \psi(x, x_2, t)$$

Q. $\begin{cases} x_1 = p + q \\ x_2 = p - q \end{cases} \quad \begin{cases} p = \frac{x_1 + x_2}{2} \\ q = \frac{x_1 - x_2}{2} \end{cases}$

$$\frac{\partial \psi}{\partial x_1} = \frac{\partial \psi}{\partial p} \frac{\partial p}{\partial x_1} + \frac{\partial \psi}{\partial q} \frac{\partial q}{\partial x_1}$$

$$= \frac{1}{2} \frac{\partial \psi}{\partial p} + \frac{\partial \psi}{\partial q}$$

$$\frac{\partial \psi}{\partial x_2} = \frac{\partial \psi}{\partial p} \frac{\partial p}{\partial x_2} + \frac{\partial \psi}{\partial q} \frac{\partial q}{\partial x_2}$$

$$= \frac{1}{2} \frac{\partial \psi}{\partial p} - \frac{\partial \psi}{\partial q}$$

$$\left(\frac{\partial^2 \psi}{\partial x_1^2} + \frac{\partial^2 \psi}{\partial x_2^2} \right) = \frac{1}{4} \frac{\partial^2 \psi}{\partial p^2} + \frac{\partial^2 \psi}{\partial q^2} + \frac{1}{4} \frac{\partial^2 \psi}{\partial p^2} - \frac{\partial^2 \psi}{\partial q^2} = \frac{1}{2} \frac{\partial^2 \psi}{\partial p^2}$$

So, the Fokker-Planck equation becomes:

$$\frac{\partial \psi(p, q, t)}{\partial t} = \frac{1}{2} \frac{\partial^2 \psi(p, q, t)}{\partial p^2} + 20 \frac{\partial \psi(p, q, t)}{\partial q}$$

If we now assume that the probability distribution decomposes as:

$$P(R, r, t) = P_{\text{com}}(R, t) \cdot P_{\text{RR}}(r, t).$$

And substituting in the Fokker-Planck equation

$$\frac{\partial P_{\text{com}}}{\partial t} \cdot P_{\text{RR}} + P_{\text{com}} \cdot \frac{\partial P_{\text{RR}}}{\partial t} =$$

$$= \frac{1}{2} D \cdot P_{\text{RR}} \cdot \frac{\partial^2}{\partial R^2} P_{\text{com}} + 2D P_{\text{com}} \frac{\partial^2}{\partial r^2} P_{\text{RR}} = 0$$

$$\Rightarrow \left\{ \frac{\partial P_{\text{com}}(R)}{\partial t} \frac{1}{P_{\text{com}}(R)} - \frac{1}{2} \frac{1}{P_{\text{com}}(R)} \frac{\partial^2}{\partial R^2} P_{\text{com}}(R) \right\}$$

$$= \left\{ - \frac{\partial P_{\text{RR}}(r)}{\partial t} \frac{1}{P_{\text{RR}}(r)} + 2D \frac{1}{P_{\text{RR}}(r)} \frac{\partial^2}{\partial r^2} P_{\text{RR}}(r) \right\}$$

$$P(R) = R(r) = C - \text{const.}$$

$$\Rightarrow \left\{ \begin{aligned} \frac{\partial P_{\text{com}}(R)}{\partial t} - \frac{1}{2} D \frac{\partial^2}{\partial R^2} P_{\text{com}}(R) &= C P_{\text{com}}(R) \\ \frac{\partial P_{\text{RR}}(r)}{\partial t} - 2D \frac{\partial^2}{\partial r^2} P_{\text{RR}}(r) &= -C P_{\text{RR}}(r) \end{aligned} \right.$$

+ RR equations are verified if both $P_{\text{com}}(R)$ and $P_{\text{RR}}(r)$ obey separate diffusion equations, with diffusion constants $\frac{1}{2}$ and $2D$, respectively.

$$\left\{ \begin{aligned} \frac{\partial P_{\text{com}}^{(i)}}{\partial t} &= \frac{1}{2} D \frac{\partial^2}{\partial R^2} P_{\text{com}}^{(i)} \\ \frac{\partial P_{\text{RR}}^{(i)}}{\partial t} &= 2D \frac{\partial^2}{\partial r^2} P_{\text{RR}}^{(i)} \end{aligned} \right.$$

which is verified if $C=0$.

Let's now verify that $P(R, r, t)$ does indeed decompose in P and r .

1. Solution to diffusion equation with initial condition $P(x, t=0) = \delta(x)$.

$$P(x, t) = \frac{1}{\sqrt{4\pi Dt}} \cdot e^{-\frac{x^2}{4Dt}}$$

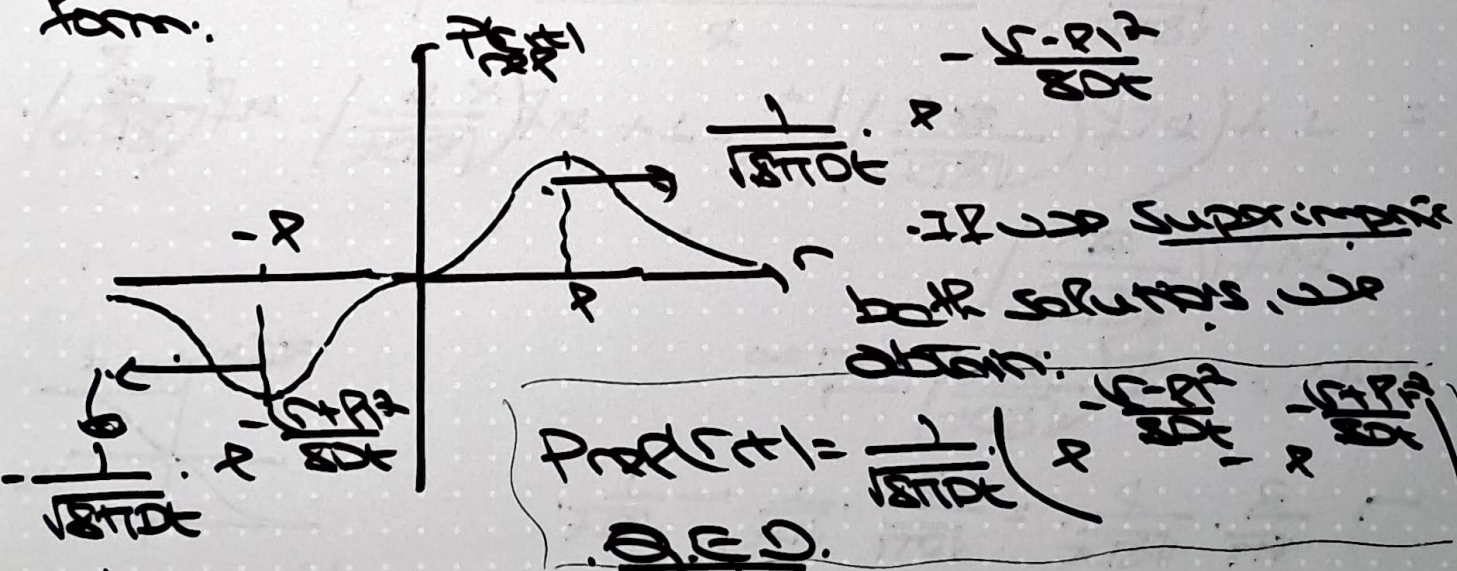
$$= \frac{1}{\sqrt{4\pi Dt}} \cdot e^{-\frac{x^2}{4Dt}}$$

However this ~~does not~~ satisfy the boundary condition: $P(x=0, t) = 0$.

In order to enforce this, we use the image method.

Introduce negative probability for $x < 0$, to force the total function to be odd \rightarrow Antisymmetric around $x=0$, with $P(x=0, t) = 0$.

Total solution with zero flux boundary: $\frac{\partial P}{\partial x} = 0$ at $x=0$.



$$P(x, t) = \frac{1}{\sqrt{4\pi Dt}} \left(e^{-\frac{x^2}{4Dt}} - e^{-\frac{x^2}{4Dt}} \right)$$

Q.E.D.

2. $J \rightarrow$ probability current density: $\frac{\partial P}{\partial t}$ (time).

In 1D, this is just 2 points. So we have units of Prob. - probability w/2 per unit of time. At time t .

Page 2

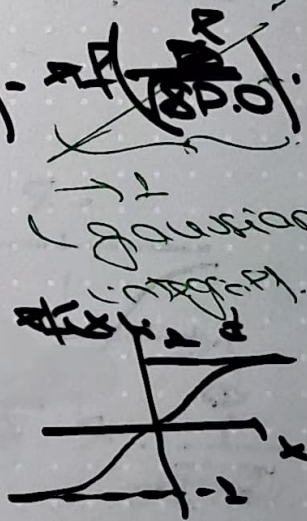
$$\begin{aligned}
 & \frac{d}{dx} \left(\frac{1}{\sqrt{8\pi\sigma^2}} \right) = \frac{1}{\sqrt{8\pi\sigma^2}} \cdot \left(-\frac{1}{2} \cdot \frac{2}{\sigma^2} \right) = -\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \\
 & = \frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \cdot \left(-\frac{1}{2} \cdot \frac{2}{\sigma^2} \right) = -\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \\
 & = \frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \cdot \left(-\frac{1}{2} \cdot \frac{2}{\sigma^2} \right) = -\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \\
 & = \frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \cdot \left(-\frac{1}{2} \cdot \frac{2}{\sigma^2} \right) = -\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \\
 & = \frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} \cdot \left(-\frac{1}{2} \cdot \frac{2}{\sigma^2} \right) = -\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}}
 \end{aligned}$$

Survival probability:

$$\begin{aligned}
 & S(t) = 1 - \int_0^t f(x) dx \\
 & = 1 - \frac{1}{\sqrt{8\pi\sigma^2}} \int_0^t \frac{1}{\sigma^2} dx \\
 & = 1 - \frac{1}{\sqrt{8\pi\sigma^2}} \cdot \left(\frac{1}{\sigma^2} \cdot x \right) \Big|_0^t \\
 & = 1 - \frac{1}{\sqrt{8\pi\sigma^2}} \cdot \left(\frac{1}{\sigma^2} \cdot t \right) \\
 & = 1 - \frac{t}{\sigma^2 \sqrt{8\pi\sigma^2}} \\
 & = 1 - \frac{t}{\sigma^2 \sqrt{8\pi\sigma^2}}
 \end{aligned}$$

$$\frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}} = \frac{1}{\sigma^2 \sqrt{8\pi\sigma^2}}$$

Maclaurin Series
of "exp" function
→
We know it keep first
two first terms



Maclaurin Series
of "exp" function

exp(x)