

# Euler Class Topological Insulators

Theory and Numerical Results for Multi-Band Models

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Foundations Theoretical Framework

Results I:  $k \cdot p$  Analysis of 3-Band Models

Results II: Bulk-Boundary Correspondence in a 4-Band Model

Outlook

# Foundations   Theoretical Framework

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# Historical Perspective

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- 2005–2007: The field exploded with the discovery of  $\mathbb{Z}_2$  Topological Insulators, protected by time-reversal symmetry.
- 2017–Present: The frontier moved to classifying more subtle phases, leading to the concepts of higher-order and fragile topology.

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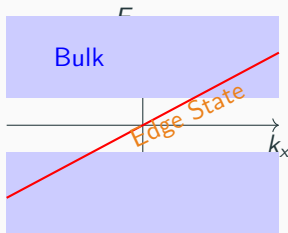
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- The gap must close at the edge to support these states.



# The Role of Symmetry in Band Theory

Symmetry constrains the form of the Hamiltonian and the wavefunctions.

- A tight-binding model can be Fourier transformed into the Bloch picture:

$$\hat{H} = \sum_{\mathbf{k} \in BZ} \sum_{\alpha, \beta} |\phi_{\alpha, \mathbf{k}}\rangle H_{\alpha\beta}(\mathbf{k}) \langle \phi_{\beta, \mathbf{k}}|$$

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- $C_2T$  **symmetry** is an anti-unitary symmetry  $\mathcal{A} = UK$  that is crucial for real topological phases. It satisfies:

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- **Key Consequence:** Within a plane left invariant by the symmetry (the  $C_2T$  plane), the Hamiltonian is forced to be **real and symmetric**:

$$\mathbf{1} \cdot \tilde{H}^*(\mathbf{k}) \cdot \mathbf{1} = \tilde{H}(\mathbf{k})$$

This reality condition makes the Chern number vanish and the **Euler class** the relevant invariant.

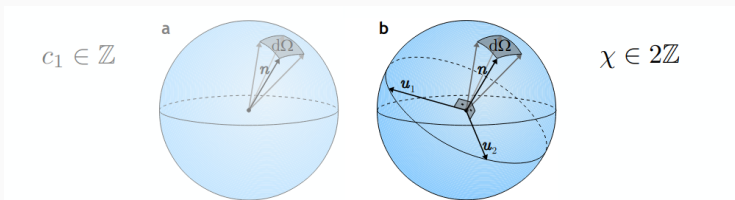
# The Euler Class: Topology of Real Bundles

## Chern Class (Complex Case)

- Describes complex vector bundles.
- Hamiltonian:  $H(\mathbf{k}) = \mathbf{h}(\mathbf{k}) \cdot \boldsymbol{\sigma}$
- Curvature:  
$$F_{ij} = \frac{1}{2} \mathbf{n} \cdot (\partial_i \mathbf{n} \times \partial_j \mathbf{n})$$
- Invariant:  $c_1 \in \mathbb{Z}$

## Euler Class (Real Case, $C_2 T^2 = +1$ )

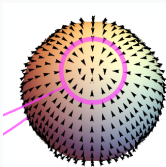
- Describes real oriented vector bundles.
- Hamiltonian:  $H(\mathbf{k}) = 2\mathbf{n}\mathbf{n}^T - 1$
- Euler Form:  
$$Eu_{ij} = \mathbf{n} \cdot (\partial_i \mathbf{n} \times \partial_j \mathbf{n})$$
- Invariant:  $\chi \in 2\mathbb{Z}$



# Euler Class from Nodal Points Patches

## Nodal Points Vortices:

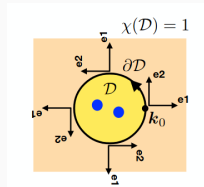
- Per the Poincaré-Hopf theorem, the Euler class is the sum of the vorticities (winding numbers) of the vector field  $\mathbf{n}(\mathbf{k})$  around its zeros.
- These zeros are the **nodal points** between bands.
- The number of stable nodal points is directly related to the Euler class:  $\#NP = 2|\chi|$ .



## Patch Euler Number:

- The total Euler class can be computed by summing the local contributions from patches  $\mathcal{D}_n$  around each nodal point.

$$\begin{aligned}\chi(\mathcal{D}) &= \frac{1}{\pi} \sum_n \left[ \int_{\mathcal{D}_n} Eu - \oint_{\partial \mathcal{D}_n} a \right] \\ &= \sum_n W_n \in \mathbb{Z}\end{aligned}$$



## Results I: k.p Analysis of 3-Band Models

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# k.p Expansion of the Square Lattice Model

## Method:

- Start with a 3-band tight-binding Hamiltonian on the square lattice.
- Perform a  $k \cdot p$  expansion around the  $\Gamma$  point to get a low-energy effective Hamiltonian.

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## Model Hamiltonian:

$$H = \begin{pmatrix} -\frac{9}{4} + s_1^2 & s_1 s_2 & (1 - c_1 - c_2)s_1 \\ s_1 s_2 & \frac{7}{4} + s_2^2 & (1 - c_1 - c_2)s_2 \\ (1 - c_1 - c_2)s_1 & (1 - c_1 - c_2)s_2 & -2 + (c_1 + c_2 - 1)^2 \end{pmatrix}$$

where  $c_i = \cos(\pi k_i)$ ,  $s_i = \sin(\pi k_i)$

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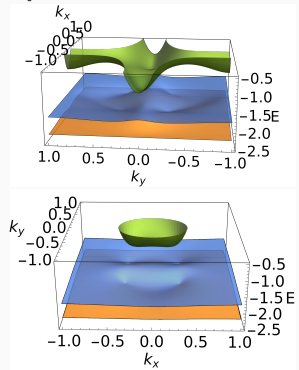
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where  $c_i = \cos(\pi k_i)$ ,  $s_i = \sin(\pi k_i)$

## k.p Approximated Hamiltonian:

$$H_{kp} = \begin{pmatrix} -\frac{9}{4} + k_1^2 \pi^2 & k_1 k_2 \pi^2 & \frac{k_1 \pi}{2} (-2 + \pi^2 (k_1^2 + k_2^2)) \\ k_1 k_2 \pi^2 & \frac{7}{4} + k_2^2 \pi^2 & \frac{k_2 \pi}{2} (-2 + \pi^2 (k_1^2 + k_2^2)) \\ \frac{k_1 \pi}{2} (\dots) & \frac{k_2 \pi}{2} (\dots) & -2 + \frac{1}{4} (-2 + \pi^2 (k_1^2 + k_2^2))^2 \end{pmatrix}$$

## Original vs. Perturbed Spectrum:

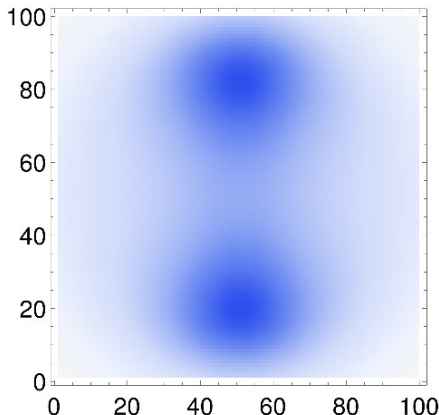


# Euler Form of the Square Lattice Model

## The Euler Form:

- The Euler form is the integrand of the Euler class ( $\chi = \frac{1}{2\pi} \int \mathbf{e} d^2k$ ).
- It can be viewed as the density of "topological charge" in the Brillouin zone.
- Its peaks indicate the momentum-space regions that contribute most to the non-trivial topology.

Plotting the Euler form inside the rectangle...



*Figure: Euler form for the gapped square lattice model, showing strong peaks that integrate to  $\chi = 2$ .*

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- Apply  $k \cdot p$  expansion to a 3-band Kagome lattice model.

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## Model Hamiltonian:

$$H = \begin{pmatrix} \frac{4}{5}c_1 & c_{1-2} + c_{1+2} & c_2 + c_{21+2} \\ c_{1-2} + c_{1+2} & \frac{4}{5}c_2 & c_1 + c_{1+22} \\ c_2 + c_{21+2} & c_1 + c_{1+22} & \frac{4}{5}c_{1+2} \end{pmatrix}$$

where  $c_i = \cos(\pi k_i)$ ,  $c_{1\pm 2} = \cos(\frac{\pi}{2}(k_1 \pm k_2))$ ,  
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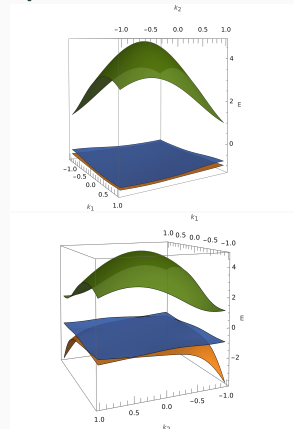
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## k.p Approximated Hamiltonian:

$$H_{kp} = \begin{pmatrix} \frac{4}{5}(1 - \frac{\pi^2 k_1^2}{2}) & 2 - \frac{\pi^2(k_1^2 + k_2^2)}{4} & \dots \\ 2 - \frac{\pi^2(k_1^2 + k_2^2)}{4} & \frac{4}{5}(1 - \frac{\pi^2 k_2^2}{2}) & \dots \\ \dots & \dots & \frac{4}{5}(1 - \frac{\pi^2(k_1 + k_2)^2}{2}) \end{pmatrix}$$

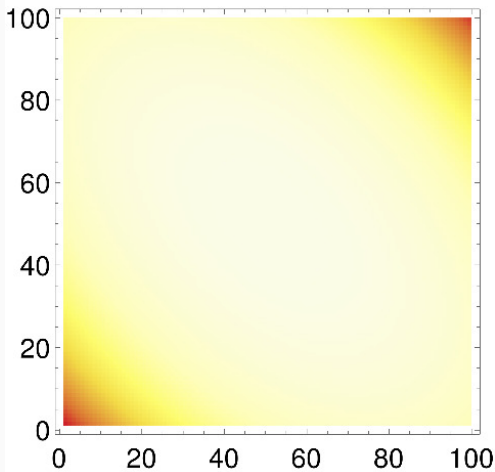
## Original vs. Perturbed Spectrum:



# Euler Form of the Kagome Lattice Model

## Topological Charge:

- The Euler form for the Kagome model shows a different distribution of topological charge.
- The integral confirms the Euler class is  $\chi = 1$ .
- This confirms  $k \cdot p$  theory is a valid tool for analyzing fragile topology in different geometries.



*Figure: Euler form for the gapped Kagome model. The charge is concentrated near the BZ corners.*



## **Results II: Bulk-Boundary Correspondence in a 4-Band Model**

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# The 4-Band Model: Bulk Properties

**Model Hamiltonian  $H(k_1, k_2)$ :**

$$\begin{pmatrix} s_2 & s_1 & -1/4 & 1 - c_1 - c_2 \\ s_1 & -s_2 & -1 + c_1 + c_2 & 1/4 \\ -1/4 & -1 + c_1 + c_2 & s_2 & s_1 \\ 1 - c_1 - c_2 & 1/4 & s_1 & -s_2 \end{pmatrix}$$

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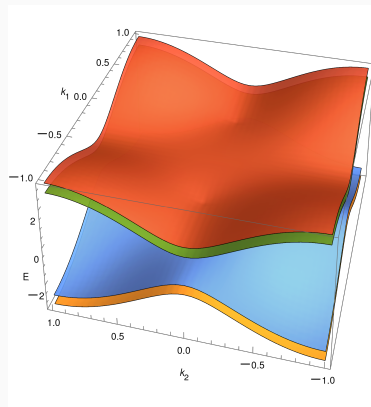
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where  $c_i = \cos(k_i\pi)$ ,  $s_i = \sin(k_i\pi)$

**Bulk Analysis:**

- A Wilson loop calculation confirms the model has a **non-trivial bulk invariant**.
- The 3D bulk band structure shows the system is fully gapped.



*Figure: 3D Bulk Spectrum.*

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# Ribbon Spectrum 1D Edge States

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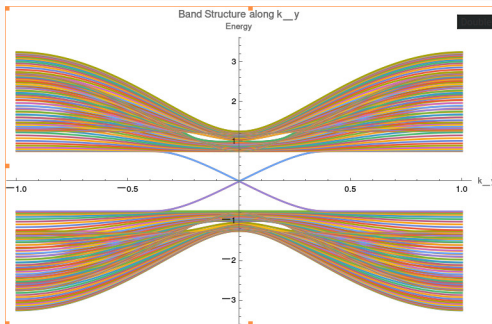
- Simulation on a ribbon geometry (open boundaries in one direction).

## Result:

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## Interpretation:

- These are the **1D edge states** predicted by the non-trivial bulk. They are the first signature of the bulk-boundary correspondence.



*Figure: Ribbon spectrum showing protected 1D edge states crossing the bulk gap.*

## Method:

- We simulate a 2D ribbon of the 4-band model, finite in  $x$  direction, infinite and periodic in  $y$  direction

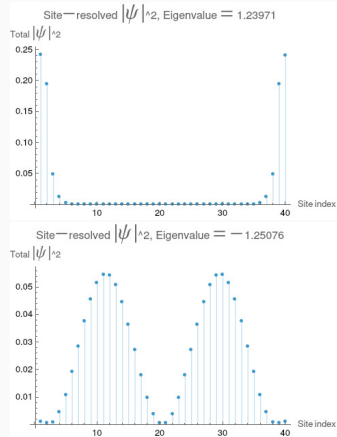
# Finite Flake Simulation Corner States

## Method:

- We simulate a 2D ribbon of the 4-band model, finite in  $x$  direction, infinite and periodic in  $y$  direction

## Result:

- The low-energy states show a remarkable concentration of their probability density  $|\psi|^2$  at the **edges**.



*Figure: The resolved site  $|\psi|^2$  shows a clear localization at the edges of the system (with 40 sites)*



# Outlook

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- **4-Band Model:** Provides a complete demonstration of the higher-order bulk-boundary correspondence, from the gapped bulk to 1D edge states and finally to localized 0D corner modes.
- **Outlook:** This work validates the theoretical framework of fragile topology and provides tools to identify and engineer these phases in real materials and metamaterials.



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