

# Case Study # 1: 1D Transient Heat Diffusion

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## 1 Problem Description

The problem of 1D unsteady heat diffusion in a slab of unit length with a zero initial temperature and both ends maintained at a unit temperature can be described by:

$$\frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} \text{ subject to } \begin{cases} T(x, 0^-) = 0 & \text{for } 0 \leq x \leq 1 \\ T(0, t) = T(1, t) = 1 & \text{for } t > 0 \end{cases} \quad (1)$$

and has the well known analytical solution:

$$T^*(x, t) = 1 - \sum_{k=1}^{\infty} \frac{4}{(2k-1)\pi} \sin[(2k-1)\pi x] \exp[-(2k-1)^2 \pi^2 t]. \quad (2)$$

In addition to the analytical solution, several numerical methods can be employed to solve the diffusion equation. Two of these methods, both an explicit and an implicit scheme, are derived in the following section. A Python script was then used to obtain results for a 21 point mesh (N=21) along the slab, and the Root Mean Square error,

$$\text{RMS} = \frac{1}{N^{\frac{1}{2}}} \sqrt{\sum_{i=1}^N [T_i^n - T^*(x_i, t_n)]^2} \quad (3)$$

was obtained for  $s(= \Delta t / \Delta x^2) = 1/6, 0.25, 0.5$ , and  $0.75$ , at  $t = 0.03, 0.06$ , and  $0.09$  using both the explicit and the implicit methods.

## 2 Solution Algorithms

The Taylor-series (TS) method can be used on this equation to derive a finite difference approximation to the PDE. Applying the definition of the derivative,

$$f'(x) \approx \frac{f(x + \epsilon) - f(x)}{\epsilon} \quad (4)$$

to Eqn. (1) yields

$$\begin{aligned} \frac{\partial T}{\partial t} &= \frac{T(x, t + \Delta t) - T(x, t)}{\Delta t} \\ &= \frac{T_i^{k+1} - T_i^k}{\Delta t}. \end{aligned} \quad (5)$$

From the definition of the Taylor series,

$$f(x + \epsilon) = f(x) + \epsilon f'(x) + \frac{\epsilon^2}{2} f''(x) + \dots \quad (6)$$

which, when applied to  $T_i^{k+1}$  and  $T_i^k$  gives

$$T_{i+1} = T_i + \Delta x \frac{\partial T_i}{\partial x} + \frac{\Delta x^2}{2} \frac{\partial^2 T_i}{\partial x^2} + O(\Delta x^3) \quad (7)$$

and

$$T_{i-1} = T_i - \Delta x \frac{\partial T_i}{\partial x} + \frac{\Delta x^2}{2} \frac{\partial^2 T_i}{\partial x^2} - O(\Delta x^3). \quad (8)$$

Adding Eqn. (7) and Eqn. (8) yields

$$T_{i+1} + T_{i-1} = 2T_i + \Delta x^2 \frac{\partial^2 T_i}{\partial x^2} + O(\Delta x^4) \quad (9)$$

which can be rearranged as the approximation for the second order term from Eqn. (1),

$$\frac{\partial^2 T}{\partial x^2} = \frac{T_{i+1} + 2T_i + T_{i-1}}{\Delta x^2} + O(\Delta x^4), \quad (10)$$

and can also be combined with the the above equations to form

$$\frac{T_i^{k+1} - T_i^k}{\Delta t} \approx \frac{T_{i+1}^k - 2T_i^k + T_{i-1}^k}{\Delta x^2}. \quad (11)$$

This result can be arranged to form Forward-Time, Centered-Space (FTCS) explicit and implicit schemes.

## 2.1 Explicit

Eqn. (11) can be rearranged to form the explicit scheme, which is

$$T_i^{k+1} = sT_{i+1}^k + (1 - 2s)T_i^k + sT_{i-1}^k, \quad (12)$$

where

$$s = \frac{\alpha \Delta t}{\Delta x^2}, \quad (13)$$

and  $\alpha$  is the thermal diffusivity of the material.

This scheme can be implemented to solve the problem computationally. As pseudocode, this looks like:

```

while  $t \leq t_{end}$  do
   $i \leftarrow 1$ 
  for  $i$  in  $N - 1$  do
     $T_{k+1}[i] = sT_k[i + 1] + (1 - 2s)T_k[i] + sT_k[i - 1]$ 
     $i \leftarrow i + 1$ 
  end for
   $T_k = T_{k+1}$ 
   $t \leftarrow t + dt$ 
end while

```

where  $N$  is the number of elements in your mesh (and  $i = 0$  is the first element in the mesh). Each element in the interior is looped over (the boundary conditions remain constant), and the time marches forward until the designated end time has been reached.

## 2.2 Implicit

Eqn. (11) can also be rearranged to form the implicit scheme, which is

$$T_i^k = -sT_{i+1}^{k+1} + (1 + 2s)T_i^{k+1} - sT_{i-1}^{k+1}. \quad (14)$$

Where again,

$$s = \frac{\alpha \Delta t}{\Delta x^2}, \quad (15)$$

and  $\alpha$  is the thermal diffusivity of the material.

Eqn. (14) can be rewritten as  $[A]T^{k+1} = T^k$ , where matrix  $A$  is tridiagonal. This tridiagonal system of  $N$  unknowns may be written as  $a_i x_{i-1} + b_i x_i + c_i x_{i+1} = d_i$ , where  $a_1 = 0$  and  $c_N = 0$ . In matrix form, this appears as:

$$\begin{bmatrix} b_1 & c_1 & & 0 \\ a_2 & b_2 & c_2 & \\ & a_3 & b_3 & \ddots \\ & & \ddots & \ddots & c_{N-1} \\ 0 & & & a_N & b_N \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ \vdots \\ x_N \end{bmatrix} = \begin{bmatrix} d_1 \\ d_2 \\ d_3 \\ \vdots \\ d_N \end{bmatrix} \quad (16)$$

This equation can be quickly solved using the tridiagonal matrix algorithm (also known as the Thomas algorithm), which consists of a forward sweep followed by back substitution. The forward sweep consists of modifying the coefficients as follows, denoting the new modified coefficients with primes:

$$c'_i = \begin{cases} \frac{c_i}{b_i} & ; i = 1 \\ \frac{c_i}{b_i - a_i c'_{i-1}} & ; i = 2, 3, \dots, N-1 \end{cases} \quad (17)$$

and

$$d'_i = \begin{cases} \frac{d_i}{b_i} & ; i = 1 \\ \frac{d_i - a_i d'_{i-1}}{b_i - a_i c'_{i-1}} & ; i = 2, 3, \dots, N. \end{cases} \quad (18)$$

The solution is then obtained by back substitution:

$$\begin{aligned} x_N &= d'_N \\ x_i &= d'_i - c'_i x_{i+1} \quad ; i = N-1, N-2, \dots, 1. \end{aligned} \quad (19)$$

In pseudocode, looks like:

```

form  $[A]$ 
while  $t \leq t_{end}$  do
   $T_{k+1} = \text{TDMA}(T_k, [A])$ 
   $T_k = T_{k+1}$ 
   $t \leftarrow t + dt$ 
end while

```

## 3 Results

A Python script was used to obtain results for a 21 point mesh ( $N=21$ ), and the Root Mean Square error was obtained for  $s (= \Delta t / \Delta x^2) = 1/6, 0.25, 0.5$ , and  $0.75$ , at  $t = 0.03, 0.06$ , and  $0.09$  using both the explicit and the implicit methods. The RMS between the implicit and analytic solutions and the explicit and analytic solutions are shown in Tables 1 through 4. The temperature along the slab for each  $s$  at each time are available as Plots 1 through 4. These plots show that the numerical solutions are, except for the case of large timesteps, very accurate when compared to the analytical solution.

t	Explicit RMS	Implicit RMS
0.03	4.17E-03	3.33E-03
0.06	1.00E-03	3.20E-04
0.09	2.22E-03	9.09E-04

Table 1. RMS results from the numerical simulations compared to the analytic solution for  $s = 1/6$

t	Explicit RMS	Implicit RMS
0.03	1.77E-03	2.15E-03
0.06	1.30E-03	5.83E-04
0.09	1.07E-03	8.99E-04

Table 2. RMS results from the numerical simulations compared to the analytic solution for  $s = 0.25$

t	Explicit RMS	Implicit RMS
0.03	5.25E-03	3.63E-03
0.06	3.72E-03	1.47E-03
0.09	3.04E-03	1.88E-03

Table 3. RMS results from the numerical simulations compared to the analytic solution for  $s = 0.5$

t	Explicit RMS	Implicit RMS
0.03	4.15E+02	5.18E-03
0.06	1.79E+07	2.37E-03
0.09	9.82E+11	2.85E-03

Table 4. RMS results from the numerical simulations compared to the analytic solution for  $s = 0.75$

The RMS for the explicit solution tended to be about an order of magnitude higher than that of the implicit solution, E-04 compared to E-03, for values of  $s = 1/6$  and 0.25, while the RMS tended to be about the same, around E-03, for  $s =$  and 0.5. At  $s = 0.75$ , however, the explicit solution becomes unstable. While the RMS for the implicit solution was still around E-03, the RMS for the explicit solution grew rapidly, to E+11.

#### 4 Discussions

The solutions show a symmetric parabolic curve which approaches the boundary condition temperature,  $T = 1$ , as  $t \rightarrow 1$ . Smaller timesteps tended to lead to more accurate results. More specifically, the overall error was proportional to the step size. As  $s$  grows too large, however, the explicit so-

lution grows unstable and produces a meaningless noncontinuous, nonphysical solution. This numerical instability is common when using unsuitably large timesteps with Euler methods.

While the explicit solution works well for solving the 1D unsteady heat diffusion for low timesteps, care must be taken when using this method for larger timesteps. If larger timesteps are required, it is inappropriate to use the explicit method and the implicit method should instead be employed.

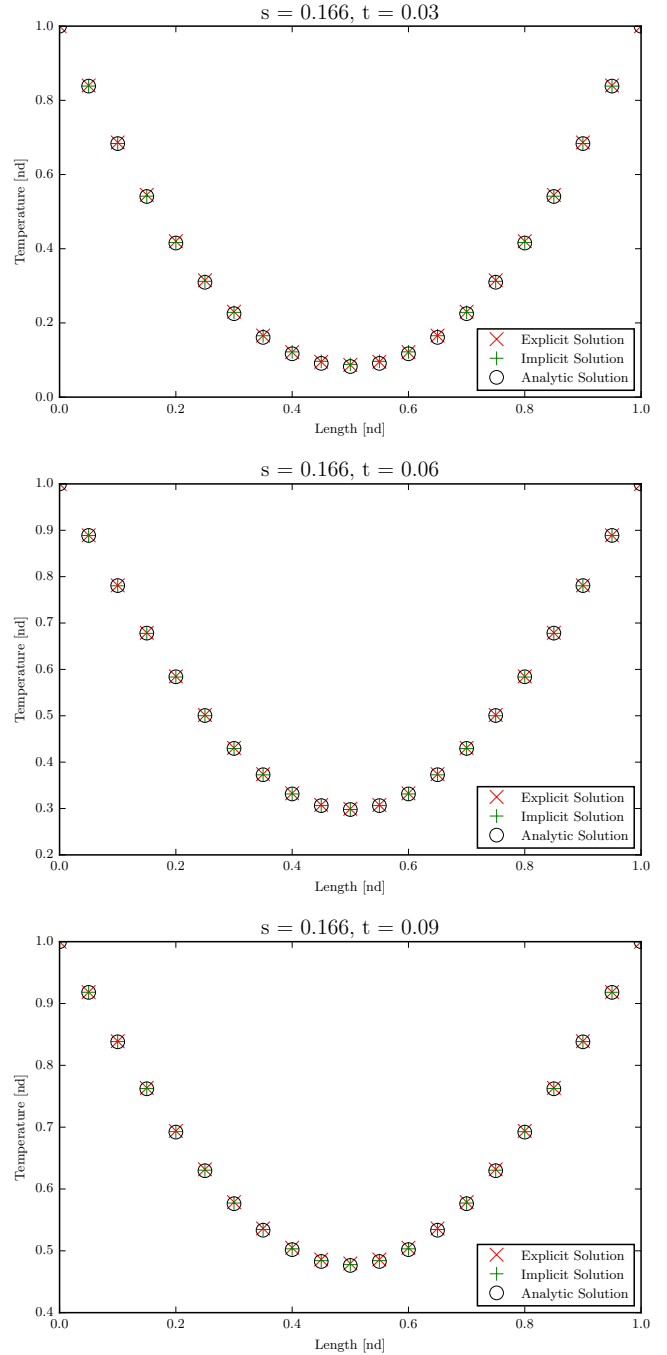


Fig. 1. Results for  $s = 1/6$

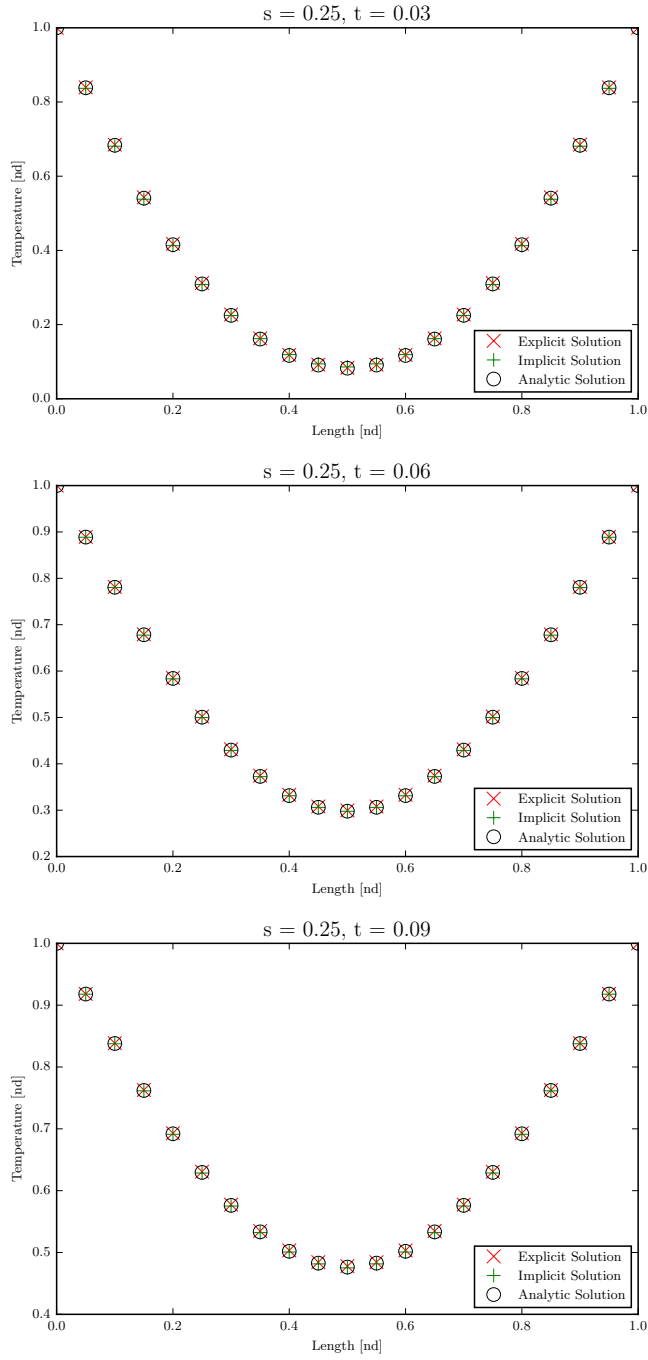


Fig. 2. Results for  $s = 0.25$

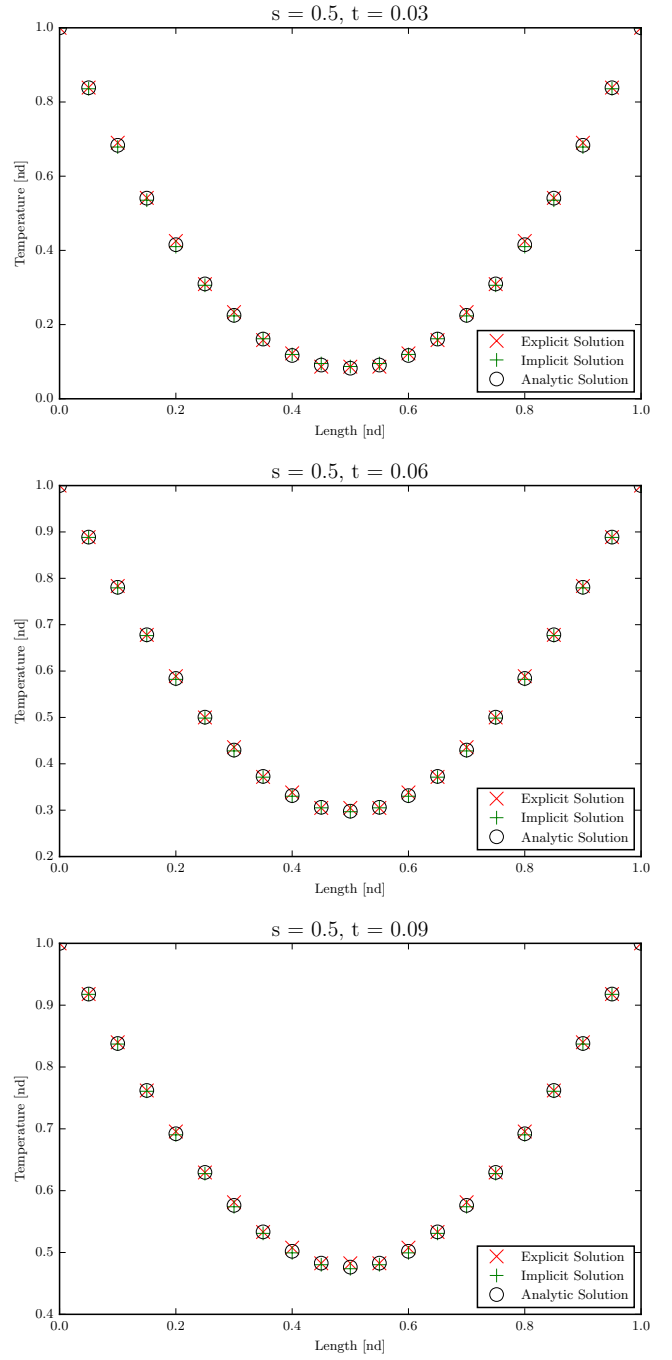


Fig. 3. Results for  $s = 0.5$

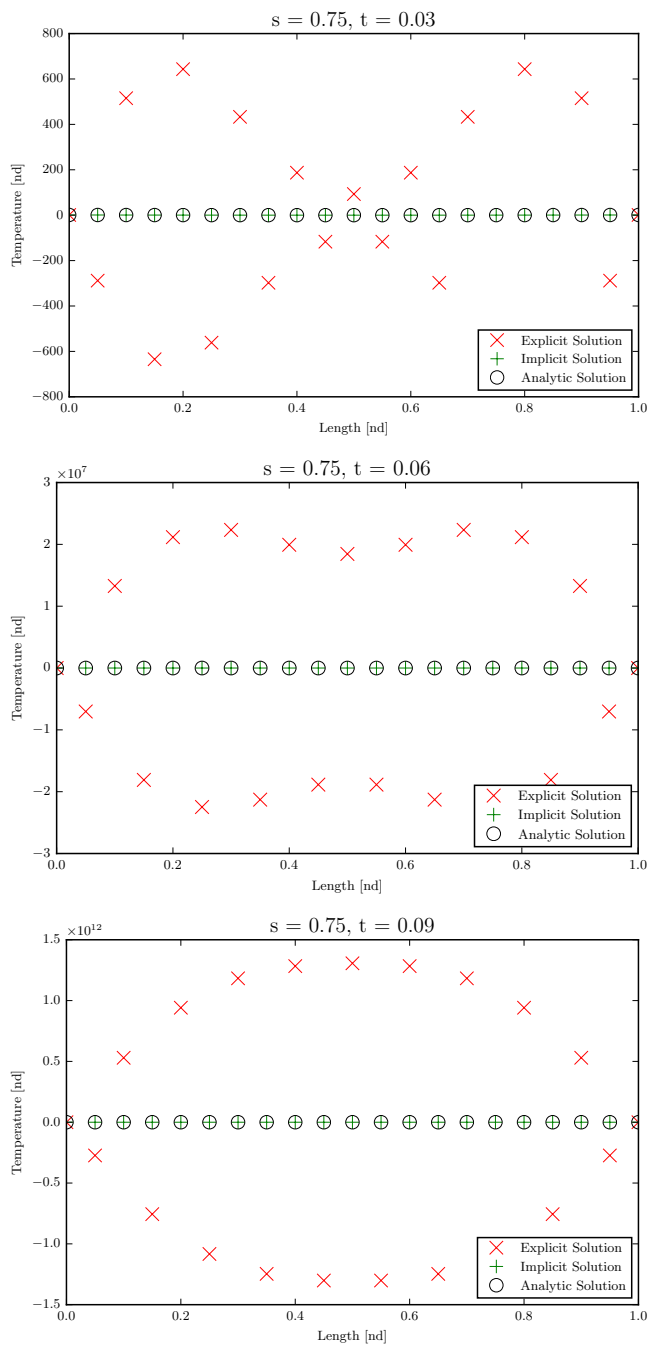


Fig. 4. Results for  $s = 0.75$

## Appendix A: Python Code

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 import os
4
5 # Configure figures for production
6 WIDTH = 495.0 # the number latex spits out
7 FACTOR = 1.0 # the fraction of the width the figure should occupy
8 fig_width_pt = WIDTH * FACTOR
9
10 inches_per_pt = 1.0 / 72.27
11 golden_ratio = (np.sqrt(5) - 1.0) / 2.0 # because it looks good
12 fig_width_in = fig_width_pt * inches_per_pt # figure width in inches
13 fig_height_in = fig_width_in * golden_ratio # figure height in inches
14 fig_dims = [fig_width_in, fig_height_in] # fig dims as a list
15
16
17 def Solver(s, t_end, show_plot=False):
18     # Problem Parameters
19     L = 1. # Domain length [n.d.]
20     T0 = 0. # Initial temperature [n.d.]
21     T1 = 1. # Boundary temperature [n.d.]
22     N = 21
23
24     # Set-up Mesh
25     x = np.linspace(0, L, N)
26     dx = x[1] - x[0]
27
28     # Calculate time-step
29     dt = s * dx ** 2.0
30
31     # Initial Condition with boundary conditions
32     T_initial = [T0] * N
33     T_initial[0] = T1
34     T_initial[N - 1] = T1
35
36     # Explicit Numerical Solution
37     T_explicit = Explicit(list(T_initial), t_end, dt, s)
38
39     # Implicit Numerical Solution
40     T_implicit = Implicit(list(T_initial), t_end, dt, s)
41
42     # Analytical Solution
43     T_analytic = list(T_initial)
44     for i in range(0, N):
45         T_analytic[i] = Analytic(x[i], t_end)
46
47     # Find the RMS
48     RMS = RootMeanSquare(T_implicit, T_analytic)
49     ExplicitRMS = RootMeanSquare(T_explicit, T_analytic)
50
51     # Format our plots
52     plt.figure(figsize=fig_dims)
53     # plt.axis([0, L, T0, T1])
54     plt.xlabel('Length [nd]')
55     plt.ylabel('Temperature [nd]')
56     plt.title('s = ' + str(s)[:5] + ', t = ' + str(t_end)[:4])
57
58     # ...and finally plot
59     plt.plot(x, T_explicit, 'xr', markersize=9, label='Explicit Solution')
60     plt.plot(x, T_implicit, '+g', markersize=9, label='Implicit Solution')
61     plt.plot(x, T_analytic, 'ob', markersize=9, mfc='none', label='Analytic Solution')
62     plt.legend(loc='lower right')
63
64     # Save plots
65     save_name = 'proj_1_s_' + str(s)[:5] + '_t_' + str(t_end) + '.pdf'
66     try:
```

```

67         os.mkdir('figures')
68     except Exception:
69         pass
70
71     plt.savefig('figures/' + save_name, bbox_inches='tight')
72     if show_plot:
73         plt.show()
74     plt.clf()
75
76     return RMS, ExplicitRMS
77
78
79 def Explicit(Told, t_end, dt, s):
80     """
81     This function computes the Forward-Time, Centered-Space (FTCS) explicit
82     scheme for the 1D unsteady heat diffusion problem.
83     """
84     N = len(Told)
85     time = 0.
86     Tnew = list(Told)
87
88     while time <= t_end:
89         for i in range(1, N - 1):
90             Tnew[i] = s * Told[i + 1] + (1 - 2.0 * s) * Told[i] + s * Told[i - 1]
91
92         Told = list(Tnew)
93         time += dt
94
95     return Told
96
97
98 def Implicit(Told, t_end, dt, s):
99     """
100     This function computes the Forward-Time, Centered-Space (FTCS) implicit
101     scheme for the 1D unsteady heat diffusion problem.
102     """
103     N = len(Told)
104     time = 0.
105
106     # Build our 'A' matrix
107     a = [-s] * N
108     a[0], a[-1] = 0, 0
109     b = [1 + 2 * s] * N
110     b[0], b[-1] = 1, 1      # hold boundary
111     c = a
112
113     while time <= t_end:
114         Tnew = TDMA solver(a, b, c, Told)
115
116         Told = list(Tnew)
117         time += dt
118
119     return Told
120
121
122 def RootMeanSquare(a, b):
123     """
124     This function will return the RMS between two lists (but does no checking
125     to confirm that the lists are the same length).
126     """
127     N = len(a)
128
129     RMS = 0.
130     for i in range(0, N):
131         RMS += (a[i] - b[i]) ** 2.
132
133     RMS = RMS ** (1. / 2.)

```

```

134     RMS /= N**(1./2.)
135
136     return RMS
137
138
139 def TDMA_solver(a, b, c, d):
140     """
141     Tridiagonal Matrix Algorithm (a.k.a Thomas algorithm).
142     """
143     N = len(a)
144     Tnew = list(d)
145
146     # Initialize arrays
147     gamma = np.zeros(N)
148     xi = np.zeros(N)
149
150     # Step 1
151     gamma[0] = c[0] / b[0]
152     xi[0] = d[0] / b[0]
153
154     for i in range(1, N):
155         gamma[i] = c[i] / (b[i] - a[i] * gamma[i - 1])
156         xi[i] = (d[i] - a[i] * xi[i - 1]) / (b[i] - a[i] * gamma[i - 1])
157
158     # Step 2
159     Tnew[N - 1] = xi[N - 1]
160
161     for i in range(N - 2, -1, -1):
162         Tnew[i] = xi[i] - gamma[i] * Tnew[i + 1]
163
164     return Tnew
165
166
167 def Analytic(x, t):
168     """
169     The analytic answer is 1 - Sum(terms). Though there are an infinite
170     number of terms, only the first few matter when we compute the answer.
171     """
172     result = 1
173     large_number = 1E6
174
175     for k in range(1, int(large_number) + 1):
176         term = ((4. / ((2. * k - 1.) * np.pi)) *
177                 np.sin((2. * k - 1.) * np.pi * x) *
178                 np.exp(-(2. * k - 1.) ** 2. * np.pi ** 2. * t))
179
180         # If subtracting the term from the result doesn't change the result
181         # then we've hit the computational limit, else we continue.
182         # print '{0} {1}, {2:.15f}'.format(k, term, result)
183         if result - term == result:
184             return result
185         else:
186             result -= term
187
188
189 def main():
190     """
191     Main function to call solver over assigned values and create some plots to
192     look at the trends in RMS compared to s and t.
193     """
194     # Loop over requested values for s and t
195     s = [1. / 6., .25, .5, .75]
196     t = [0.03, 0.06, 0.09]
197
198     RMS = []
199     with open('results.dat', 'w+') as f:
200         for i, s_ in enumerate(s):

```



```

201     sRMS = [0] * len(t)
202     for j, t_ in enumerate(t):
203         sRMS[j], ExplicitRMS = Solver(s_, t_, False)
204         f.write('{0:.3f} {1:.2f} {2:.2e} {3:.2e} \n'.format(s_, t_, sRMS[j], ExplicitRMS))
205         # print i, j, sRMS[j]
206     RMS.append(sRMS)
207
208     # Convert to np array to make this easier...
209     RMS = np.array(RMS)
210
211     # Check for trends in RMS vs t
212     plt.figure(figsize=fig_dims)
213     plt.plot(t, RMS[0], '.r', label='s = 1/6')
214     plt.plot(t, RMS[1], '.g', label='s = .25')
215     plt.plot(t, RMS[2], '.b', label='s = .50')
216     plt.plot(t, RMS[3], '.k', label='s = .75')
217     plt.xlabel('t')
218     plt.ylabel('RMS')
219     plt.title('RMS vs t')
220     plt.legend(loc='best')
221
222     save_name = 'proj_1_rms_vs_t.pdf'
223     plt.savefig('figures/' + save_name, bbox_inches='tight')
224     plt.clf()
225
226     # Check for trends in RMS vs s
227     plt.figure(figsize=fig_dims)
228     plt.plot(s, RMS[:, 0], '.r', label='t = 0.03')
229     plt.plot(s, RMS[:, 1], '.g', label='t = 0.06')
230     plt.plot(s, RMS[:, 2], '.b', label='t = 0.09')
231     plt.xlabel('s')
232     plt.ylabel('RMS')
233     plt.title('RMS vs s')
234     plt.legend(loc='best')
235
236     save_name = 'proj_1_rms_vs_s.pdf'
237     plt.savefig('figures/' + save_name, bbox_inches='tight')
238     plt.clf()
239
240 if __name__ == "__main__":
241     main()

```