



Technical Mechanics

Trigonometric Identities

$\sin(-\theta)$	$-\sin(\theta)$
$\cos(-\theta)$	$\cos(\theta)$
$\tan(-\theta)$	$-\tan(\theta)$
$\sin(\frac{\pi}{2} - \theta)$	$\cos(\theta)$
$\cos(\frac{\pi}{2} - \theta)$	$\sin(\theta)$
$\tan(\frac{\pi}{2} - \theta)$	$\cot(\theta)$
$\sin(\pi - \theta)$	$\sin(\theta)$
$\cos(\pi - \theta)$	$-\cos(\theta)$
$\tan(\pi - \theta)$	$-\tan(\theta)$
$\sin(\theta \pm \frac{\pi}{2})$	$\pm \cos(\theta)$
$\cos(\theta \pm \frac{\pi}{2})$	$\mp \sin(\theta)$
$\sin(\theta + \pi)$	$-\sin(\theta)$
$\cos(\theta + \pi)$	$-\cos(\theta)$
$\sin(\alpha \pm \beta)$	$\sin \alpha \cos \beta \pm \cos \alpha \sin \beta$
$\cos(\alpha \pm \beta)$	$\cos \alpha \cos \beta \mp \sin \alpha \sin \beta$

deg / rad	$30^\circ / (\frac{\pi}{6})$	$45^\circ / (\frac{\pi}{4})$	$60^\circ / (\frac{\pi}{3})$	$90^\circ / (\frac{\pi}{2})$
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sin	$\frac{1}{2}$	$\frac{1}{\sqrt{2}}$	$\frac{\sqrt{3}}{2}$	1
cos	$\frac{\sqrt{3}}{2}$	$\frac{1}{\sqrt{2}}$	$\frac{1}{2}$	0
tan	$\frac{1}{\sqrt{3}}$	1	$\sqrt{3}$	—

Vector Identities

$\vec{a} \times \vec{b}$	$\ a\ \ b\ \sin(\theta) \vec{n}$	Area of parallelogram
$\vec{a} \times (b + c)$	$\vec{a} \times \vec{b} + \vec{a} \times \vec{c}$	
$a \cdot (b \times c)$	$b \cdot (c \times a)$	$c \cdot (a \times b)$

Calculus Rules

$\frac{dy}{dx}$	$\frac{dy}{du} \frac{du}{dx}$
$\frac{d(uv)}{dx}$	$u \frac{dv}{dx} + v \frac{du}{dx}$
\vec{v}	$\frac{d(\vec{r})}{dt}$
ω	$\frac{d\Theta}{dt}$

TODO: Add acceleration formulae!

Cartesian Coordinates

\vec{r}	$x\vec{e}_x + y\vec{e}_y + z\vec{e}_z$
\vec{v}	$\dot{x}\vec{e}_x + \dot{y}\vec{e}_y + \dot{z}\vec{e}_z$

Cylindrical Coordinates

\vec{r}	$\rho\vec{e}_\rho + z\vec{e}_z$
\vec{v}	$\dot{\rho}\vec{e}_\rho + \rho\dot{\varphi}\vec{e}_\varphi + \dot{z}\vec{e}_z$
\vec{e}_ρ	$\cos(\varphi)\vec{e}_x + \sin(\varphi)\vec{e}_y$
\vec{e}_φ	$\frac{d(\vec{e}_\rho)}{d\varphi} = -\sin(\varphi)\vec{e}_x + \cos(\varphi)\vec{e}_y$
$\frac{d\vec{e}_\varphi}{d\varphi}$	$-\cos(\varphi)\vec{e}_x - \sin(\varphi)\vec{e}_y = -\vec{e}_\rho$
$\frac{d\vec{e}_\rho}{d\varphi}$	\vec{e}_φ
$\dot{\vec{e}}_\varphi$	$-\dot{\varphi}\vec{e}_\rho$
$\dot{\vec{e}}_\rho$	$\dot{\varphi}\vec{e}_\varphi$

Rigid Bodies

$$\forall P, Q \in \kappa$$

$$|r_{PQ}| = |r_Q - r_P| = \text{Constant}$$

SdpG

$$e = \frac{r_Q - r_P}{|r_Q - r_P|}$$

$$\vec{v}_Q \cdot \vec{e}_{PQ} = \vec{v}_P \cdot \vec{e}_{PQ}$$

$$v'_Q = v'_P$$

Movement across a plane

- All velocities are parallel to the plane
- All points along a normal to the plane have the same velocity
- Either a translation or a rotation at any point in time

Types of Movement (2D)

Translation	$\forall P, Q \in \kappa^2 \mid v_P = v_Q$
Rotation	$\exists P, Q \in \kappa^2 \mid v_P \neq v_Q, \omega_P = \omega_Q$

Center of Rotation - intersection of all lines perpendicular to the velocities of points.

Rollen ohne Gleiten - Center of rotation is at the point of contact with the ground

$$\vec{v}_P = (\omega \vec{e}_z) \times \vec{r}_{CP}$$

$$\text{Scalar : } \vec{v}_P = \omega d$$

$\vec{\omega}$ - Unit vector in the direction of axis of rotation, so that positive rotation is anti-clockwise.

Movement in 3D space

- Simultaneous translation & rotation possible as “schraubung”.

Rotation Axis - $\forall P, Q \in \overrightarrow{e_{\text{axis}}}\lambda, v_P = v_Q$

TODO: Parametric equation of points along the rotational axis

Starrkörperformel

$$\vec{v}_P = \vec{v}_B + \vec{\omega} \times \vec{r}_{BP}$$

Every point in a rigid body rotates around every other point in the body with the same angular velocity.

Invariants of motion in space

$I_1 = \vec{\omega} \forall P \in \kappa$	Current angular velocity same regardless of the reference point
$I_2 = \vec{\omega} \cdot v_P \forall P \in \kappa$	Component of velocity in the direction of the rotation axis is the same for all points in the body. Called "Translation velocity"

Kinematic: $\{\vec{v}_B, \vec{\omega}\}$

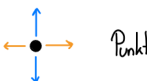
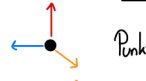

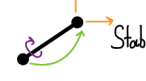


Types of Movement (3D)

Translation	$\vec{\omega} = 0 \forall P \in \kappa$
Kreiselung	$\vec{\omega} \neq 0 \wedge I_2 = 0 \therefore \exists P \in \kappa \mid v_P = 0$
Schraubung	$I_2 \neq 0$







Degrees of freedom

Minimum number of coordinates to clearly determine the state of a system.

$$f = n - b$$

2D (Ebene)	3D
 Punkt $f = 2$	 Punkt $f = 3$
 Stab $f = 2 \cdot 2 - 1 = 3$	 Stab $f = 2 \cdot 3 - 1 = 5$
 Starrer Körper $f = 3$	 Starrer Körper $f = 6$

IMPORTANT: The joint at a roller / pivot must be accounted for too! (using the n-Gelenk formula)

b von	2D	b von	2D
 Auflager (beidseitig)	1	 Gelenk (2 SKs verbunden)	2
 Gelenk (festlager)	2	 Rollen ohne gleiten	2
 Einspannung	3	 Gelenk (n SK verbunden)	$(n-1) \cdot 2$

b-value for slider depends on how it's connected, usually 1

Forces

$$\sum \text{Inner Forces} = 0$$

Static Equivalence

$$\mathcal{P}_{\text{tot}}(G_1) = \mathcal{P}_{\text{tot}}(G_2)$$

Therefore in a rigid body:

$$R_1 = R_2$$

$$(M_B)_1 = (M_B)_2$$

Furthermore, two forces are equivalent if they have the same magnitude and line of action.

Forces with lines of action going through the same point have only a resultant force - no moment:



$$M_P = 0, R \neq 0$$

Moments

A moment is a concept for describing the capacity of a force to rotate an object around an arbitrary center of rotation with units Nm .

The moment of a force around the center of rotation O in vector form is:

$$M_O = \vec{r}_{OP} \times \vec{F}$$

The resulting moment lies along the axis of rotation and describes the angular direction of the rotation caused by the moment.

Alternatively, it can be expressed as a scalar with the perpendicular distance from the line of action of F to O: d:

$$M_O = dF$$

Transformation of moments

The moment of a force can be transformed with respect to a different point using the following formula:

$$M_A = M_B + r_{AB} \times R$$

Torque

Perpendicular distance between lines of action d:

$$R = 0$$

$$M_P = dF$$

Whenever the resultant force is 0, the moment around all points in the body is the same.

Dynamic

The dynamic of a force group with respect to a point O describes the entire set of forces on the body:

$$\{R, M_O\}$$

Where R is the resultant force and M_O is the resultant moment around O: $\sum \vec{r}_{OP_i} \times \vec{F}_i$

The following invariants apply to the dynamic:

- $I_1 = R \forall P \in \kappa$
- $I_1 = R \cdot M_O \forall P \in \kappa$ - the component of the resultant in the direction of the moment with respect to the same point is the same for all points

Power

The rate of transfer of energy.

Due to the work done by a force $\int_C \vec{F} d\vec{s}$, the power exerted by a force at a point in time can be expressed as:

$$\mathcal{P} = \vec{F} \cdot \vec{v}$$

- *Accelerating force* ($\frac{\pi}{2} < \alpha \leq \pi$) - A force with a positive component in the direction of the velocity is contributing kinetic energy to the object and increasing the power
- *Braking Force* ($0 < \alpha < \frac{\pi}{2}$) - Reduces the kinetic energy of the object
- A force perpendicular to the velocity of an object does not contribute to its power until the object begins moving with a component in the direction of the perpendicular force.

Total power of a rigid body

The total power of an object is the sum of powers for each force acting on the body:

$$\mathcal{P}_{\text{tot}} = \sum_{i=1}^n \vec{F}_i \cdot \vec{v}_i$$

When the kinematic $\{v_B, \omega\}$ and dynamic $\{R, M_B\}$ with respect to a point B are known, we can calculate the total power thanks to the rigid body formula and the “pacman” identity:

$$\mathcal{P}_{\text{tot}} = R \cdot v_B + M_B \cdot \omega$$

Parallel Forces

When all forces acting on a body point in the same direction, they can be written as:

$$\vec{F}_i = F_i \vec{e}$$

Where \vec{e} is the unit vector of their common direction.

The dynamic with reference to a point O is:

$$R = \sum \vec{F}_i = \vec{e} \sum F_i$$

$$M = \sum \vec{r}_i \times \vec{F}_i = \sum \vec{r}_i F_i \times \vec{e}$$

The sum $\sum \vec{r}_i F_i$ is called the dipole moment (sometimes written as N) of a set of parallel forces. It is independent of the point O, as long as it's consistent for each force included.

Center of Forces

This is the point on which a pivot can be placed and no resultant moment would act on the body. In other words, an equal and opposite resultant force can act on this point resulting in a net 0 dynamic.

It is unique to a dipole moment, the direction of \vec{e} is irrelevant (although gravity always acts in the same direction, so in practice this fact isn't too important).

The position vector of the center of forces from the point O can be calculated using a dipole moment with **the same point O** as its origin:

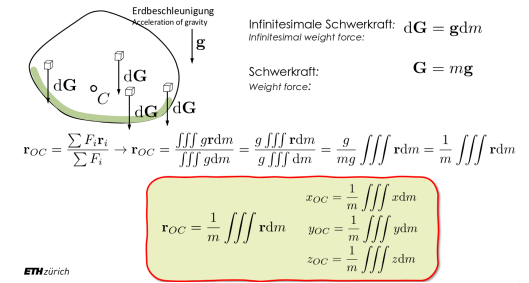
$$r_{OC} = \frac{\sum \vec{r}_i F_i}{\sum F_i}$$

For a dipole moment with the specific direction of the forces \vec{e} , an entire line can serve as a center of forces:

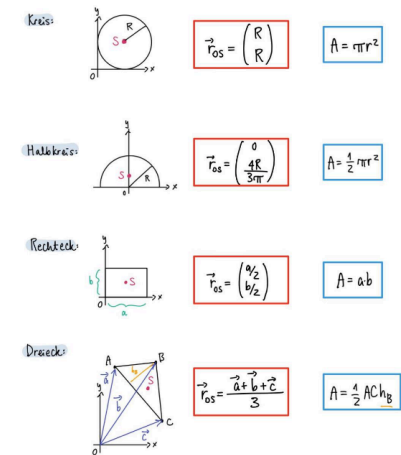
$$\vec{r}_{OC}(\lambda) = \vec{r}_{OC} + \lambda \vec{e}$$

Center of Mass

This is the average location of all the weight of an object. It can be calculated using a volume integral over the density of the body:

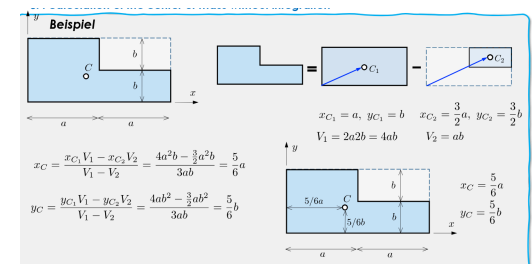


Where dm is effectively the density at the point of each differential.



Integration is a linear transformation - a center of mass can be calculated as the sum of separate integrals (which of course can also be negative):

$$\vec{r}_{OC} = \frac{\sum \iiint \vec{r} dm}{\sum m_i}$$



LTD: Check general formula

$$\vec{r}_C = \frac{\sum m_i \vec{r}_{Ci}}{\sum m_i}$$

Rest

A system is at rest when all of its velocities are 0.

- Instantaneous rest - $\vec{v}_p = 0 \forall p \in \kappa \mid t = t_0$
- State of rest - $\vec{v}_p = 0 \forall p \in \kappa, \forall t$

Fundamental Theorem of Statics

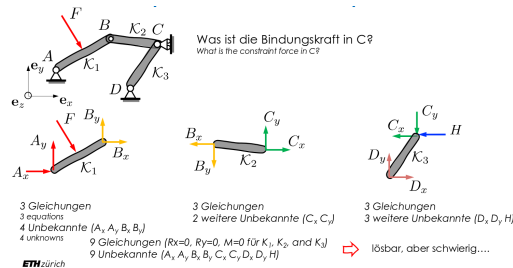
This also means that the resultant force and moment around any point (the moment is always the same if $R = 0$) is 0:

$$\vec{R} = 0$$

$$\vec{M} = 0$$

Constraints exert equal and opposite forces to prevent an object moving through them.

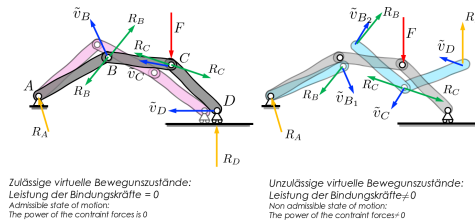
Forces in a system at rest with multiplied bodies can be solved by including forces at constraints:



These systems of equations can be solved through Gaussian elimination.

Virtual Motion

A system can be modelled as in virtual motion when virtual velocities are modelled at each point, such that the total power of constraint forces is 0:



- For example, velocity at the slider is perpendicular to its constraint force, and equal and opposite constraint forces at joints cancel each other's accelerating and braking power out.
- There can be no virtual velocity at a pivot, as it exerts constraint forces in both the x and y component - they are orthogonal and the power of this force would never be 0.
- There can however exist a rotational velocity around a pivot.

These virtual velocities, denoted as \tilde{v} can be found using techniques in the kinematics of rigid bodies.

Theorem of Virtual Power

A system is at rest when the virtual total power is 0 for every virtual state of motion:

$$\mathcal{P}_{\text{tot}} = \sum_{i=1}^n \vec{F}_i \cdot \tilde{v}_i = 0$$

This is useful to calculate a few external / constraint forces, by strategically allowing virtual motion which involves that force. Other forces with and equal and opposite reaction can be ignored as the resultant at that point is 0.

If many forces in a system are needed, then a full analysis using the theorem of statics is more appropriate.

Framework

The constraints in a framework can be calculated by removing one of the rods, whose compression (or tension, this becomes apparent if a negative value is calculated) is acting as a constraint force.

LTD: Force cut

Power of Torque

It is useful to solve statics problems with pulleys using the virtual power of a torque (the cable around a pulley has the same tension throughout):

$$\mathcal{P} = \vec{\omega} \cdot \vec{M}$$

As usual in a torque, the origin of the moment is irrelevant.

LTD: Do not assume all tensions are equal in pulleys exam!