

## PHYSICS-I

Department of Physics, IIT Guwahati.

Course No: PH 101

Mid-Semester Examination

Date: 16 Sept., 2018

Time: 2-4 pm

Total Marks: 40

### General Instructions

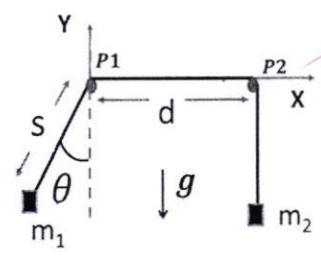
- Make sure that there are **six sheets** (including this) in this **Question-cum-Answer Booklet**.
- Write your **Name** and **Roll Numbers** on **every sheet** in the space provided.
- You must write the answers **ONLY IN THE SPACE PROVIDED** for the given **question**. Answers written **elsewhere WILL NOT be evaluated!**
- NO** extra **answer-sheets** will be provided!
- Supplementary sheets provided are **ONLY** for rough work.
- It is advised that you first solve the problems on the supplementary sheet, and then copy the key steps in the space provided for that problem in **this Question-cum-Answer booklet**.
- Be **legible!** Also, make sure that your answers are systematic, logically as well as mathematically connected.
- Question 1 is on the rear side of **this sheet**. The answers to **each** of the remaining questions (Q2-Q6) should be contained in a single sheet.
- Marks for each question is indicated below the question number.

Student's Name :	Rahul.D
Roll No:	180102054
Signature:	Rahul

Signature of the Invigilator:	Jyoti 16/9/18
-------------------------------	---------------

<b>Q1. Report the degree of freedom (DOF) of the following systems in the space provided.</b> <b>No justification required!</b>			
a)	Consider a <b>thin</b> uniform rigid rod of mass $m$ and length $l$ , under the following situations.		
[3×1]	The rod is sliding down such a way that one end of it always maintains contact with a vertical wall while the other end slips on a horizontal floor. <b>The whole motion is on the <math>xy</math> –plane.</b>	DOF=	1 ✓
	The rod is free to move any fashion but constrained <b>only to the <math>xy</math> –plane.</b>	DOF=	3 ✓
	The rod is constrained to move <b>inside</b> a larger spherical shell of radius $R$ ( $2R > l$ ) such a way that both ends of the rod always maintain contact with the inner surface of the shell.	DOF=	3 ✓
b)	Water molecule ( $H_2O$ ) is composed of two hydrogen atoms bonded to an oxygen atom, with its average O–H bonds measuring $1 \text{ \AA}$ , and with an equilibrium H–O–H angle of $104.5^\circ$ . Obtain the DOF of a water molecule that is free to move in <b>three dimensions</b> under the following models.		
[4×1]	<b>Model#1:</b> The O–H bonds, or bond lengths, are flexible. But the H–O–H bond angle is fixed at $104.5^\circ$	DOF=	8 ✓
	<b>Model#2:</b> The O–H bonds are rigid at $1 \text{ \AA}$ , and the H–O–H bond angle is fixed at $104.5^\circ$ .	DOF=	6 ✓
	<b>Model#3:</b> The O–H bonds are rigid at $1 \text{ \AA}$ , but the H–O–H bond angle is free to change.	DOF=	7 ✓
	<b>Model#4:</b> The O–H bonds as well as the H–O–H bond angle is free to change.	DOF=	9 ✓

7

Name:	Rahul. P	Roll No:	180102054
Q2	Two point masses $m_1$ and $m_2$ are connected by a light inextensible string (of length $l$ ), which runs over two massless pulleys of <b>negligible radius</b> , as shown in the figure. The mass $m_1$ can oscillate on the $xy$ -plane, while $m_2$ can move only in the vertical direction. The first pulley $P_1$ is at the origin, and the 2 <sup>nd</sup> pulley $P_2$ is on the $x$ -axis, at distance $d$ from the first. Gravity acts downwards as shown. Obtain the Lagrangian and the Euler-Lagrange equations for the system, employing $S$ and $\theta$ (marked in the figure) as the generalized coordinates.		
[7]			

It is given that  $S$  and  $\theta$  should be used as generalized coordinates

Let cartesian coordinates of  $m_1$  &  $m_2$  be  $(x_1, y_1)$  &  $(x_2, y_2)$

$$x_1 = -s \sin \theta$$

$$x_2 = d$$

$$\dot{x}_1 = -[\dot{s} \sin \theta + s \cos \theta \dot{\theta}]$$

$$\dot{x}_2 = 0$$

$$y_1 = -s \cos \theta$$

$$y_2 = -(L-d-s)$$

$$\dot{y}_1 = -[\dot{s} \cos \theta - s \sin \theta \dot{\theta}] \quad \dot{y}_2 = -\dot{s}$$

$$\text{Kinetic energy for system } (T) = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2)$$

$$T = \frac{1}{2} m_1 ([-\dot{s} \sin \theta - s \cos \theta \dot{\theta}]^2 + [3\dot{\theta} \sin \theta - \dot{s} \cos \theta]^2) + \frac{1}{2} m_2 (0^2 + \dot{s}^2)$$

$$T = \frac{1}{2} m_1 (\dot{s}^2 + s^2 \dot{\theta}^2) + \frac{1}{2} m_2 \dot{s}^2$$

$$V = [-m_1 g s \cos \theta] + [-m_2 g (L-d-s)] \quad \left[ \text{considering gravitational potential energy zero at origin} \right]$$

$$L = T - V$$

$$L = \frac{1}{2} m_1 (\dot{s}^2 + s^2 \dot{\theta}^2) + \frac{1}{2} m_2 \dot{s}^2 + m_1 g s \cos \theta + m_2 g (L-d-s)$$

$$\frac{\partial L}{\partial s} = m \ddot{\theta}^2 s + m_1 g \cos \theta - m_2 g \quad \frac{\partial L}{\partial \dot{s}} = m_1 \dot{s} + m_2 \dot{s}$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{s}} \right) - \frac{\partial L}{\partial s} = 0 \quad [1^{st} \text{ E-L eq}]$$

$$\Rightarrow m_1 \ddot{s} + m_2 \ddot{s} - m \ddot{\theta}^2 s - m_1 g \cos \theta + m_2 g = 0 \rightarrow \textcircled{1}$$

$$\frac{\partial L}{\partial \theta} = m_1 g s (-\sin \theta) \quad \frac{\partial L}{\partial \dot{\theta}} = m_1 s^2 \dot{\theta}$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0 \quad [2^{nd} \text{ E-L eq}]$$

$$\Rightarrow m_1 [2s \dot{s} \dot{\theta} + s^2 \ddot{\theta}] + m_1 g s \sin \theta = 0 \rightarrow \textcircled{2}$$



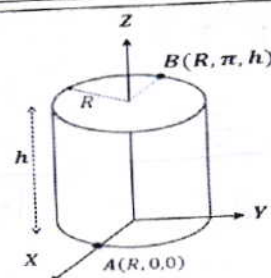


Name:	Rahul.D	Roll No:	180102054
-------	---------	----------	-----------

Q3

[6]

Find the geodesic on the surface of a cylinder of radius  $R$  and height  $h$ , employing cylindrical coordinate system. Find the total length of this path between points  $A(R, 0, 0)$  and  $B(R, \pi, h)$ . Coordinate of  $A$  and  $B$  are given in cylindrical coordinate system.



6

On considering cylindrical coordinate system, the ~~path~~ any point on the measured path can be shown as  $(R, \theta, z)$  where  $R$  is constant.

$$dr = R d\theta \hat{\theta} + dz \hat{z}$$

$$\text{thus } ds = \sqrt{(R d\theta)^2 + (dz)^2}$$

geodesic means the path length should be minimum  
so we are considering a ~~path~~ variable which denotes path length ( $S$ ).

$$S = \int_A^B \sqrt{(R d\theta)^2 + (dz)^2}$$

$$S = \int_A^B \sqrt{R^2 + \left(\frac{dz}{d\theta}\right)^2} d\theta$$

$$F(\dot{z}, \theta, z) = \sqrt{R^2 + \dot{z}^2} \quad \text{where } \dot{z} = \frac{dz}{d\theta}$$

For the path to be geodesic the function should

obey Euler-Lagrange equation

$$\text{i.e. } \frac{d}{dt} \left( \frac{\partial F}{\partial \dot{z}} \right) - \frac{\partial F}{\partial z} = 0$$

$$\frac{d}{dt} \left( \frac{\dot{z}}{\sqrt{R^2 + \dot{z}^2}} \right) = 0 \Rightarrow 0 = 0$$

$$\Rightarrow \frac{\dot{z}}{\sqrt{R^2 + \dot{z}^2}} = k$$

$$\frac{\dot{z}^2}{R^2 + \dot{z}^2} = k^2$$

$$\frac{\dot{z}^2}{R^2} = \frac{k^2}{1-k^2}$$

$$\dot{z}^2 = \frac{R^2 k^2}{(1-k^2)}$$

$$\dot{z} = \frac{Rk}{\sqrt{1-k^2}}$$

$$\frac{dz}{d\theta} = \frac{Rk}{\sqrt{1-k^2}}$$

$$\int_0^{\pi} dz = \frac{Rk}{\sqrt{1-k^2}} \int_0^{\pi} d\theta$$

$$z = \frac{Rk}{\sqrt{1-k^2}} \theta$$

on substituting ~~h~~ & coordinate of B

$$h = \frac{Rk}{\sqrt{1-k^2}} \pi$$

$$\text{Thus } \frac{Rk}{\sqrt{1-k^2}} = \frac{h}{\pi}$$

on substituting the value of constant.

$$\text{then } z = \frac{h}{\pi} \theta$$

$$\dot{z} = \frac{h}{\pi}$$

$$S = \int_0^{\pi} \sqrt{R^2 + \left(\frac{h}{\pi}\right)^2} d\theta$$

$$S = \sqrt{R^2 + \left(\frac{h}{\pi}\right)^2} [\theta]_0^{\pi}$$

$$S = \sqrt{R^2 + \left(\frac{h}{\pi}\right)^2} (\pi)$$

$$S = \sqrt{R^2 \pi^2 + h^2}$$

→ path length

Name:	Rahul. D	Roll No:	180102054
-------	----------	----------	-----------

6

Q4	A point particle of mass $m$ is constrained to move on the surface of a rigid-fixed sphere of radius $R$ . Obtain the Hamiltonian and Hamilton's equations of motion of the mass $m$ in <b>spherical polar</b> coordinate system. The particle is acted upon <b>only</b> by the forces of constraint.
----	---

The particle is constrained to move on rigid fixed sphere of radius  $R$ , so on considering the spherical polar coordinate system, the coordinate of particle is  $(R, \theta, \phi)$

$R \rightarrow$  constant  $\theta$  &  $\phi$  variables

Velocity of a particle in spherical polar coordinate system

$$\vec{v} = \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} + r \sin \theta \dot{\phi} \hat{\phi}$$

In this case  $r = R$  so  $\dot{r} = 0$ .

$$\text{thus } \vec{v} = R \dot{\theta} \hat{\theta} + R \sin \theta \dot{\phi} \hat{\phi}$$

$$\text{Kinetic energy (T)} = \frac{1}{2} m v^2 = \frac{m}{2} [(R \dot{\theta})^2 + (R \sin \theta \dot{\phi})^2]$$

Since the particle is acted upon only by forces of constraint, the potential at all points on the surface of sphere remains constant and let us assume it as  $C$

$$\text{thus } V = C$$

$$L = T - V$$

$$L = \frac{m}{2} R^2 \dot{\theta}^2 + \frac{m}{2} R^2 \sin^2 \theta \dot{\phi}^2 - C$$

$$P_{\theta} = \frac{\partial L}{\partial \dot{\theta}} = m R^2 \dot{\theta} \quad \Rightarrow \quad \dot{\theta} = \frac{P_{\theta}}{m R^2}$$

$$P_{\phi} = \frac{\partial L}{\partial \dot{\phi}} = m R^2 \sin^2 \theta \dot{\phi} \quad \Rightarrow \quad \dot{\phi} = \frac{P_{\phi}}{m R^2 \sin^2 \theta}$$

$$\cancel{H = \sum \dot{q}_i \frac{\partial L}{\partial \dot{q}_i}} \quad H = \sum P_j \dot{q}_j - L$$

$$H = m R^2 \dot{\theta}^2 + m R^2 \sin^2 \theta \dot{\phi}^2 - \frac{m R^2 \dot{\theta}^2}{2} - \frac{m}{2} R^2 \sin^2 \theta \dot{\phi}^2 + C$$

$$H = \frac{m R^2 \dot{\theta}^2}{2} + \frac{m R^2 \sin^2 \theta \dot{\phi}^2}{2} + C$$

$$H = \frac{m R^2}{2} \left( \frac{P_{\theta}}{m R^2} \right)^2 + \frac{m R^2 \sin^2 \theta}{2} \left( \frac{P_{\phi}}{m R^2 \sin^2 \theta} \right)^2 + C$$

$$H = \frac{P_{\theta}^2}{2 m R^2} + \frac{P_{\phi}^2}{2 m R^2 \sin^2 \theta} + C$$

4

The required equations are  $\dot{p}_j = -\frac{\partial H}{\partial q_j}$   $\dot{q}_j = \frac{\partial H}{\partial p_j}$

$$\dot{p}_\phi = 0 \quad \dot{\phi} = \frac{2 p_\phi}{2 m R^2 \omega^2 \theta} = \frac{p_\phi}{m R^2 \omega^2 \theta}$$

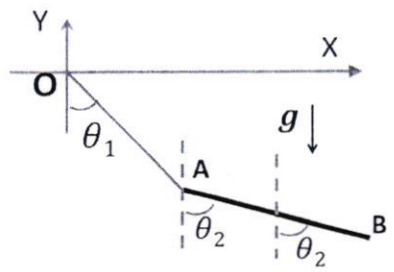
$$\dot{p}_\theta = -\frac{\partial H}{\partial \theta} = \frac{p_\phi^2 \omega^2 \theta \omega^2 c^2 \theta}{m R^2} \quad \dot{\theta} = \frac{\partial H}{\partial p_\theta} = \frac{p_\theta}{m R^2}$$

1+

1+



Name:	Rahul.D	Roll No:	180102054
-------	---------	----------	-----------

Q5	A thin uniform rod, AB, of mass $m$ and length $l$ is suspended to the ceiling (at O), by one end, through a light inextensible string of length $d$ (see figure). Assume that the rod can only move on the $xy$ -plane, and without any slackening of the string. Obtain the Lagrangian of the system using $\theta_1$ and $\theta_2$ (marked in the figure) as the generalized coordinates. The system is under gravity. Euler-Lagrange equations are <b>not required</b> !	
[6]		

It is given that the generalized coordinates to be used is  $\theta_1$  and  $\theta_2$ .

Since the ~~the~~ string is light, ~~its~~ kinetic and potential energy is neglected.

Kinetic energy (rigid body) =  $\frac{1}{2} m v_c^2$  + Kinetic energy in COM frame.

where COM  $\rightarrow$  centre of mass

$v_c \rightarrow$  Velocity of COM

~~Since~~ Rod is uniform so COM will be at  $\frac{1}{2}$  dist from end.

$$\vec{r}_{\text{com}} = (d \sin \theta_1 + \frac{1}{2} l \sin \theta_2) \hat{i} - (d \cos \theta_1 + \frac{1}{2} l \cos \theta_2) \hat{j}$$

$$\vec{v}_{\text{com}} = (d \cos \theta_1 \dot{\theta}_1 + \frac{1}{2} l \cos \theta_2 \dot{\theta}_2) \hat{i} - (d(-\sin \theta_1) \dot{\theta}_1 + \frac{1}{2} (-\sin \theta_2) \dot{\theta}_2) \hat{j}$$

$$\vec{v}_{\text{com}} = (d \cos \theta_1 \dot{\theta}_1 + \frac{1}{2} l \cos \theta_2 \dot{\theta}_2) \hat{i} + (d \sin \theta_1 \dot{\theta}_1 + \frac{1}{2} l \sin \theta_2 \dot{\theta}_2) \hat{j}$$

~~T~~ Kinetic energy (T) of system =  $\frac{1}{2} m ((d \dot{\theta}_1)^2 + (\frac{1}{2} l \dot{\theta}_2)^2)$

$$T = \frac{m}{2} \left[ (d \dot{\theta}_1 \dot{\theta}_1 + \frac{1}{2} l \dot{\theta}_2 \dot{\theta}_2)^2 + (d \dot{\theta}_1 \dot{\theta}_1 + \frac{1}{2} l \dot{\theta}_2 \dot{\theta}_2)^2 \right] + \frac{1}{2} \frac{m l^2}{12} \dot{\theta}_2^2$$

$$T = \frac{m}{2} \left[ (d \dot{\theta}_1)^2 + (\frac{1}{2} \dot{\theta}_2)^2 + d l \cos(\theta_1 - \theta_2) \right] + \frac{m l^2 \dot{\theta}_2^2}{24}$$

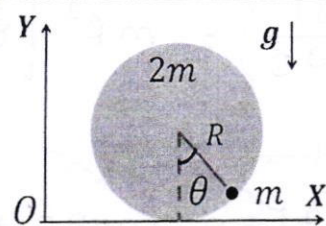
$$V = -m g \left[ d \cos \theta_1 + \frac{1}{2} l \cos \theta_2 \right]$$

$$L = \frac{m}{2} \left[ (d \dot{\theta}_1)^2 + (\frac{1}{2} \dot{\theta}_2)^2 + d l \cos(\theta_1 - \theta_2) \right] + \frac{m l^2 \dot{\theta}_2^2}{24} + m g \left[ d \cos \theta_1 + \frac{1}{2} l \cos \theta_2 \right]$$



Name:	Rahul.D	Roll No:	180102054.
-------	---------	----------	------------

- Q6 A uniform disk of mass  $2m$  and radius  $R$  has a point mass  $m$  fixed close to its edge. The disk can roll without sliding along the  $x$ -axis, under gravity. Obtain the Euler-Lagrange (E-L) equations for the system. Using the E-L equation find the frequency of small angle oscillations in  $\theta$  (as marked in the figure). Hint: For small angle oscillations,  $\theta^2 \approx 0$  and  $\dot{\theta}^2 \approx 0$ .
- [8]



As the point mass is rigidly fixed to disk, it can be considered as a rigid body.

DOF = 6 - 5 = 1 [  $y=0$ , rotationally allowed in one axis  
condition of no slipping, move on  $x$  axis. ]

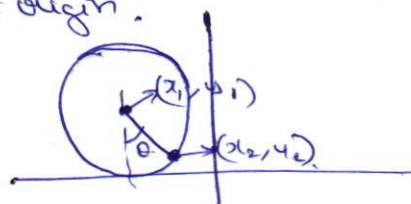
Consider that the disk was initially at origin.  
Now on ~~move~~ rotating it by  $\theta$

$$x_1 = -R\theta \quad y_1 = R$$

$$\dot{x}_1 = -R\dot{\theta} \quad \dot{y}_1 = 0$$

$$x_2 = -R\theta + R\sin\theta \quad y_2 = R - R\cos\theta$$

$$\dot{x}_2 = -R\dot{\theta} + R\cos\theta\dot{\theta} \quad \dot{y}_2 = R\sin\theta\dot{\theta}$$



KE of rigid body = KE of COM + KE (in COM frame).

Thus ~~T~~ for the sys

$$KE_{\text{system}} (T) = \frac{1}{2} 2m (R\dot{\theta})^2 + \frac{1}{2} \frac{2m R^2 \dot{\theta}^2}{2} + \frac{1}{2} m [(-R\dot{\theta} + R\cos\theta\dot{\theta})^2 + (R\sin\theta\dot{\theta})^2]$$

$$T = mR^2\dot{\theta}^2 + \frac{mR^2\dot{\theta}^2}{2} + \frac{m}{2} [2R^2\dot{\theta}^2 - 2R^2\dot{\theta}^2\cos\theta]$$

$$= \frac{3}{2} mR^2\dot{\theta}^2 + mR^2\dot{\theta}^2(1 - \cos\theta)$$

$$T = \frac{5}{2} mR^2\dot{\theta}^2 - mR^2\dot{\theta}^2\cos\theta$$

$$V = 2mgR + mg(R - R\cos\theta)$$

$$L = T - V$$

$$L = \frac{5}{2} mR^2\dot{\theta}^2 - mR^2\dot{\theta}^2\cos\theta - 3mgR + mgR\cos\theta$$

$$\frac{\partial L}{\partial \theta} = mgR(-\sin\theta) + mR^2\dot{\theta}^2\sin\theta$$

$$\frac{\partial L}{\partial \theta} = 5mR^2\dot{\theta} - 2mR^2\dot{\theta}\cos\theta$$

$$\frac{\partial L}{\partial \dot{\theta}} = mR\sin\theta [R\dot{\theta}^2 - g]$$

$$\frac{\partial L}{\partial \dot{\theta}} = mR^2\dot{\theta} [5 - 2\cos\theta]$$

E-L equation  $\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) - \frac{\partial L}{\partial q_j} = 0$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) = mR^2 \left[ \ddot{\theta} [5 - 2\cos\theta] + 2\sin\theta \dot{\theta}^2 \right]$$

Thus ~~E-L equation~~

$$mR^2 [5\ddot{\theta} - 2\cos\theta \ddot{\theta} + 2\sin\theta \dot{\theta}^2] + mR \sin\theta [g - R\dot{\theta}^2] = 0$$

~~It is given~~ It is given that for small ~~oscillation~~ oscillation  $\dot{\theta}^2 \approx 0$ .

so.  $mR^2 \ddot{\theta} [5 - 2\cos\theta] + mR \sin\theta g = 0$ .

$$\ddot{\theta} = - \frac{mR \sin\theta g}{mR^2 (5 - 2\cos\theta)}$$

$$= - \frac{g \sin\theta}{R(5 - 2\cos\theta)}$$

~~as  $\theta^2 \rightarrow$~~

as  $\theta^2 \approx 0$   $\theta \rightarrow 0 \Rightarrow \sin\theta \rightarrow \theta$  &  $\cos\theta \rightarrow 1$

$$\ddot{\theta} \approx - \frac{g \theta}{5R}$$

for SHM.

$$\ddot{\theta} = -\omega^2 \theta$$

on comparing  $\omega^2 = \frac{g}{5R}$

$$\omega = \sqrt{\frac{g}{5R}}$$

$$\text{frequency} = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{g}{5R}}$$