Cavity QED Systems: Spectral Properties

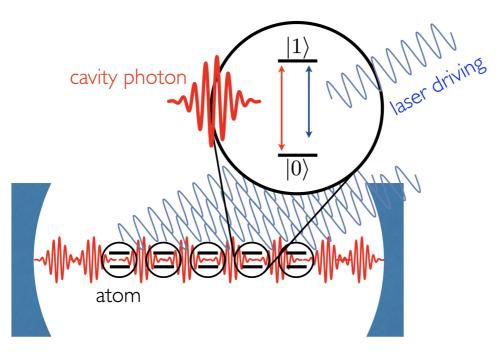


Figure 1: Originally, a cavity-QED system was a system of atoms interacting with quantized light in the cavity.

Two-level systems (atoms) are described by Qubits, denoted by symbol S.

```
In[*]:= Let[Qubit, S]

In[*]:= Dimension[S]

Out[*]:= 2

Bosonic modes (cavity photons) are described by Bosons, denoted by symbol c.

In[*]:= Let[Boson, c]

In[*]:= {Bottom[c], Top[c], Dimension[c]}

Out[*]:= {0, 5, 6}
```

Jaynes-Cummings Hamiltonian

The qubit has transition frequence Ω .

$$In[\cdot]:=$$
 Let[Real, Ω]
Hq = (Ω / 2) S[3]
 $Out[\cdot]=$
 $\frac{\Omega S^{z}}{2}$

The photon has frequency ω .

$$In[\circ]:=$$
 Let[Real, ω]
$$Hc = \omega * Dagger[c] **c$$

$$Out[\circ]=$$

$$\omega c^{\dagger}c$$

The qubit-photon coupling in the rotating-wave approximation (RWA).

Construct the total Hamiltonian.

In[
$$\circ$$
]:= HH = Hq + Hc + Hg
Out[\circ] =
$$\omega \ c^{\dagger}c + g \ \left(c \ S^{+} + c^{\dagger} \ S^{-}\right) + \frac{\Omega \ S^{z}}{2}$$

Basic Method

Examine the matrix representation of the Hamiltonian.

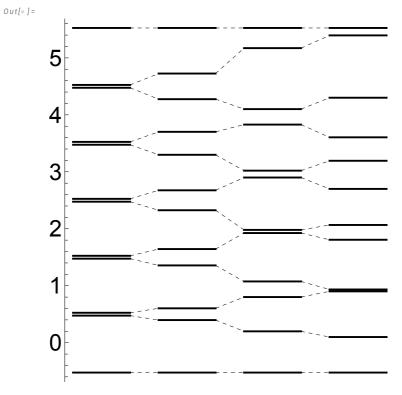
Out[•] // MatrixForm =

Write a small utility function to use to investigate the spectrum of the Hamiltonian.

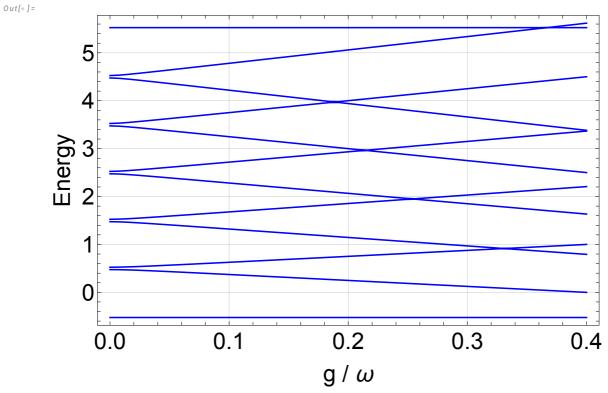
```
In[@]:= spectrum[gv_] := Block[
        \{\omega = 1,
         \Omega = 1.05,
          g},
        g = gv;
        val = Sort@Eigenvalues[Normal@mat]
```

Now, examine the level structure of the Hamiltonian and its evolution as coupling g varies.

```
In[0]:= LevelsPlot[{spectrum[0.], spectrum[0.1], spectrum[0.3], spectrum[0.4]},
      ImageSize → Medium]
```



Note that there seems to be crossing of energy levels. Let us examine the spectrum more closely by varying g continuously.



Parity Conservation

The numerical results above suggest that energy levels cross each other. This is unusual unless there is a symmetry.

Question: Is there a symmetry? What is it if any?

Indeed, the parity is conserved.

Then, let us take a look at the basis again.

$$\begin{aligned} & & & \text{In[=]:= bs = Basis[\{c,S\}]} \\ & & \text{Out[=]=} \\ & & & & \left\{ \left| 0_c 0_S \right\rangle, \, \left| 0_c 1_S \right\rangle, \, \left| 1_c 0_S \right\rangle, \, \left| 1_c 1_S \right\rangle, \, \left| 2_c 0_S \right\rangle, \\ & & & & \left| 2_c 1_S \right\rangle, \, \left| 3_c 0_S \right\rangle, \, \left| 3_c 1_S \right\rangle, \, \left| 4_c 0_S \right\rangle, \, \left| 4_c 1_S \right\rangle, \, \left| 5_c 0_S \right\rangle, \, \left| 5_c 1_S \right\rangle \right\} \end{aligned}$$

Note that each basis state is an eigenstate of the parity operator, each with different parities.

Let us group the basis elements according to their parities.

$$\label{eq:constraints} \begin{split} & \text{out} [\circ \] := \ ps = GroupBy[bs, ParityValue[\{c, S\}]] \\ & \text{out} [\circ \] := \\ & \quad \langle \left| \ 1 \rightarrow \left\{ \left| \ 0_c 0_S \right\rangle, \ \left| \ 1_c 1_S \right\rangle, \ \left| \ 2_c 0_S \right\rangle, \ \left| \ 3_c 1_S \right\rangle, \ \left| \ 4_c 0_S \right\rangle, \ \left| \ 5_c 1_S \right\rangle \right\}, \\ & \quad -1 \rightarrow \left\{ \left| \ 0_c 1_S \right\rangle, \ \left| \ 1_c 0_S \right\rangle, \ \left| \ 2_c 1_S \right\rangle, \ \left| \ 3_c 0_S \right\rangle, \ \left| \ 4_c 1_S \right\rangle, \ \left| \ 5_c 0_S \right\rangle \right\} \left| \right\rangle \end{split}$$

Rearrange the basis so that basis states of the same parity come together, and then examine again the matrix.

0

new // MatrixForm Out[]//MatrixForm= 0 0 0 $\frac{1}{2}$ (6 ω + Ω) 0 0 0 0 0 0 0 2 g

0

The resulting matrix is indeed block diagonal!

Therefore, it is more efficient to handle each block separately.

5 ω +

Out[0]=

Summary

Functions

- LevelsPlot
- Sort
- Parity, ParityValue
- MatrixIn
- Qubit, Dimension
- Boson, Bottom, Top, Dimension
- Matrix, ExpressionFor
- ProperValues, ProperStates, ProperSystem

Related Links

- Mahn-Soo Choi, Advanced Quantum Technologies 12, 2000085 (2020), "Exotic Quantum States of Circuit Quantum Electrodynamics in the Ultra-Strong Coupling Regime."
- C. Dongni, S. Luo, Y.-D. Wang, S. Chesi, and Mahn-Soo Choi, Physical Review A 105, 022627 (2022), "Geometric manipulation of a decoherence-free subspace in atomic ensembles."