

Notes on Kinematic Fit

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Assume that we have measured a momentum of two particles, which decay from a mother particle with an exact momentum P_0 . In the real world, every measurement inherently carries some errors. Consequently, the measured momentum may not satisfy momentum conservation:

$$P_0 = P_1 + P_2; \quad P_{1,meas} + P_{2,meas} \neq P_0$$
 (1)

However, since these measurements did not incorporate our prior knowledge from physics, we can make a more informed estimate of the measured parameters. If the momentum resolution of particles 1 and 2 is well-known, then we can express χ^2 as

$$\chi^2 = \frac{(P_1 - P_{1,meas})^2}{\sigma_1^2} + \frac{(P_2 - P_{2,meas})^2}{\sigma_2^2}$$
 (2)



By incorporating the *Kinematic Constraints*, specifically *momentum conservation*, into our example, we introduce additional terms known as *Lagrange Multiplier* to Equation (3):

$$\chi^2 = \frac{(P_1 - P_{1,meas})^2}{\sigma_1^2} + \frac{(P_2 - P_{2,meas})^2}{\sigma_2^2} + 2\lambda(\mathbf{P_{1,meas}} + \mathbf{P_{2,meas}} - \mathbf{P_0})$$
(3)

We then proceed to evaluate the conditions for local minima, i.e. setting the partial derivatives equal to zero:

$$\frac{1}{2}\frac{\partial\chi^2}{\partial P_1} = \frac{(P_1 - P_{1,meas})}{\sigma_1^2} + \lambda = 0 \tag{4}$$

$$\frac{1}{2}\frac{\partial \chi^2}{\partial P_2} = \frac{(P_2 - P_{2,meas})}{\sigma_2^2} + \lambda = 0$$
 (5)

$$\frac{1}{2}\frac{\partial \chi^2}{\partial \lambda} = (P_1 + P_2 - P_0) = 0 \tag{6}$$



By solving the equations 4,5,6, we obtain the following expressions:

$$P_{1} = \frac{\sigma_{2}^{2} P_{1,meas} - \sigma_{1}^{2} P_{2,meas} + \sigma_{1}^{2} P_{0}}{\sigma_{1}^{2} + \sigma_{2}^{2}}$$
(7)

$$P_{2} = \frac{\sigma_{1}^{2} P_{2,meas} - \sigma_{2}^{2} P_{1,meas} + \sigma_{2}^{2} P_{0}}{\sigma_{1}^{2} + \sigma_{2}^{2}}$$
(8)

$$\lambda = \frac{P_{1,meas} + P_{2,meas} - P_0}{\sigma_1^2 + \sigma_2^2}.$$
 (9)

Now, we have obtained the 'corrected' measurements with minimized χ^2 , which incorporates momentum conservation. Let us delve into the interpretation of these equations.



In a straightforward interpretation, λ can be viewed as a kind of 'normalized variance' of the kinematic constraint. It quantifies the error in momentum conservation $(P_{1,meas} + P_{2,meas} - P_0)$ relative to the overall resolution $(\sigma_1^2 + \sigma_2^2)$. Equation (4) implies that

$$P_1 = P_{1,meas} + \sigma_1^2 \lambda, \tag{10}$$

suggesting that the corrected momentum(P_1) is essentially the measured momentum($P_{1,meas}$) augmented by a term proportional to the detector resolution and the normalized error of the kinematic constraint. Thus, we can assert that we have applied a statistically fair correction to the momentum, taking into account both the detector resolution and kinematic constraints



Fitting in General

Assume that you have a set of measurements, $\mathbf{m} = \{m_1, m_2 \dots m_N\}$, and some unmeasured data, $\mathbf{u} = \{u_1, u_2 \dots u_J\}$ to be estimated. Kinematic constraints can be represented by sets of equations $\mathbf{f} = \{f_1(m_1, m_2, \dots m_N, u_1, u_2, \dots u_N), f_2, \dots f_K\}$. We will iteratively solve the problem by guessing the best parameter for each step and checking χ^2 . Let \mathbf{m}^0 denote our initial measured data, and \mathbf{m} represent the 'guess' of the data in each iterative step.

$$\chi^{2}(\mathbf{m}) = (\mathbf{m}^{0} - \mathbf{m})^{T} V^{-1}(\mathbf{m}^{0} - \mathbf{m}) + 2\lambda^{T} \mathbf{f}(\mathbf{m}, \mathbf{u}).$$
 (11)

Here, the Lagrange multiplier $\lambda = \{\lambda_1, \lambda_2, \dots \lambda_K\}$ is not just a number but a column vector with k elements, corresponding to each kinematic constraint in \mathbf{f} . Our task is to minimize χ^2 to obtain the best guesses in this is also fixed by the property of the p

χ^2 Minimization

By (partially)differentiating with respect to all variables involved, we obtain the gradients of χ^2 . Setting all of them to zero indicates that we have reached a minimum point of χ^2 . We have 3 sets of gradient equations:

$$\nabla_{\mathbf{m}} = -2V^{-1}(\mathbf{m}^{0})(\mathbf{m}^{0} - \mathbf{m}) + 2\mathbf{F}_{\mathbf{m}}^{T}(\mathbf{m}, \mathbf{u})\lambda = 0$$
 (12)

$$\nabla_{\mathbf{u}} = 2\mathbf{F}_{\mathbf{u}}^{T}(\mathbf{m}, \mathbf{u})\lambda = 0 \tag{13}$$

$$\nabla_{\lambda} = \mathbf{f}(\mathbf{m}, \mathbf{u}). \tag{14}$$

Here, the subscripts denote partial derivatives. i.e. $((\mathbf{F}_m)_{ki} \equiv \frac{\partial f_k}{\partial m_i})$.



Processing Iterative Steps.

We can express the following equations based on the ones provided above:

$$V^{-1}(\mathbf{m}^{0})(\mathbf{m}^{\nu+1} - \mathbf{m}^{0}) + (\mathbf{F}_{m}^{T})^{\nu} \lambda^{\nu+1} = 0$$
 (15)

$$(\mathbf{F}_{u}^{T})^{\nu}\lambda^{\nu+1} = 0 \tag{16}$$

$$\mathbf{f}^{\nu} + \mathbf{F}_{m}^{\nu}(\mathbf{m}^{\nu+1} - \mathbf{m}^{\nu}) + \mathbf{F}_{u}^{\nu}(\mathbf{u}^{\nu+1} - \mathbf{u}^{\nu}) = 0.$$
 (17)

Equation (17) is not a direct consequence of Equation (13) but rather a linear approximation to proceed with our iteration step. Expanding the ∇_{λ} term with a Taylor series leads to this equation. Note that, as our parameters \mathbf{m} and \mathbf{u} are updated during the step, our constraint matrix \mathbf{f} should also be updated during the iteration. Here, λ should be indexed as $\nu+1$ since it is a parameter to be guessed in the next step.



Solving the Equation(1)

Multiplying **V** to Equation (15) leads to:

$$\mathbf{m}^{\nu+1} - \mathbf{m}^0 = -V(\mathbf{m}^0)(\mathbf{F}_m^T)^{\nu} \lambda^{\nu+1}. \tag{18}$$

Substituting Equation (18) into Equation (17), we get:

$$\mathbf{F}_{u}^{\mathbf{v}}(\mathbf{u}^{\mathbf{v}+1} - \mathbf{u}^{\mathbf{v}}) = -\mathbf{f}^{\mathbf{v}} - \mathbf{F}_{m}^{\mathbf{v}}(-V(\mathbf{m}^{0})(\mathbf{F}_{m}^{T})^{\mathbf{v}}\lambda^{\mathbf{v}+1} + \mathbf{m}^{0} - \mathbf{m}^{\mathbf{v}})$$
$$= S\lambda^{\mathbf{v}+1} - R \tag{19}$$

where $S \equiv \mathbf{F}_{\mathbf{m}}^{\mathbf{v}} V(\mathbf{m}^0) (\mathbf{F}_{\mathbf{m}}^T)^{\mathbf{v}}$ and $R \equiv \mathbf{f}^{\mathbf{v}} + \mathbf{F}_{\mathbf{m}}^{\mathbf{v}} (\mathbf{m}^0 - \mathbf{m}^{\mathbf{v}})$. Multiplying $(\mathbf{F}_u^T)^{\mathbf{v}} S^{-1}$ and substituting Equation (16), we get:

$$(\mathbf{F}_{u}^{T})^{\nu} S^{-1} \mathbf{F}_{u}^{\nu} (\mathbf{u}^{\nu+1} - \mathbf{u}^{\nu}) = (\mathbf{F}_{u}^{T})^{\nu} \lambda^{\nu+1} - (\mathbf{F}_{u}^{T})^{\nu} S^{-1} R.$$
(20)



Solving the Equation(2)

Then we naturally obtain:

$$\mathbf{u}^{\nu+1} = \mathbf{u}^{\nu} - ((\mathbf{F}_{u}^{T})^{\nu} S^{-1} \mathbf{F}_{u}^{\nu})^{-1} (\mathbf{F}_{u}^{T})^{\nu} S^{-1} R.$$
 (21)

and from Equation (19)

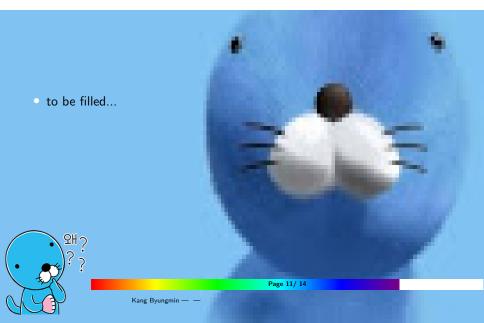
$$\lambda^{\nu+1} = S^{-1}(\mathbf{F}_{u}^{\nu}(\mathbf{u}^{\nu+1} - \mathbf{u}^{\nu}) + R) \tag{22}$$

. For a summary, we have obtained all equations to proceed to the next step. All other matrices in the equation can be calculated from parameters of the current step, and χ^2 can be evaluated from (11) .

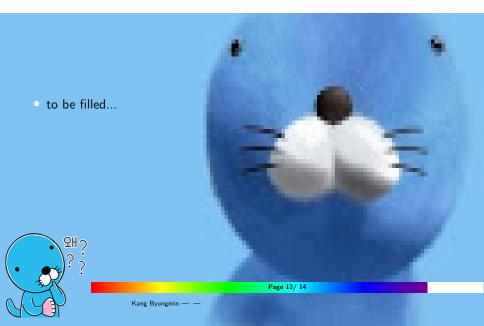
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\left\{ \begin{array}{l}
\mathbf{u}^{\vee+1} = \mathbf{u}^{\vee} - ((\mathbf{F}_{u}^{T})^{\vee} S^{-1} \mathbf{F}_{u}^{\vee})^{-1} (\mathbf{F}_{u}^{T})^{\vee} S^{-1} R & (21) \\
\lambda^{\vee+1} = S^{-1} (\mathbf{F}_{u}^{\vee} (\mathbf{u}^{\vee+1} - \mathbf{u}^{\vee}) + R) & (22) \\
\mathbf{m}^{\vee+1} = \mathbf{m}^{0} - V(\mathbf{m}^{0}) (\mathbf{F}_{m}^{T})^{\vee} \lambda^{\vee+1} & (18)
\end{array} \right.$$



Evolution of Variance Matrix



Pull distribution



Example: 3-C Fit

Assuming a decay of $\Lambda \to p\pi^-$, let's represent the momentum of those particles in spherical coordinates:

$$\vec{P}_{\Lambda} = (P_{\Lambda}, \theta_{\Lambda}, \phi_{\lambda}), \cdots \tag{23}$$

Then the kinematic constraints can be expressed as:

$$\begin{pmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \end{pmatrix} = \begin{pmatrix} -P_{\Lambda} \sin \theta_{\Lambda} \cos \phi_{\Lambda} + P_{p} \sin \theta_{p} \cos \phi_{p} + P_{\pi} \sin \theta_{\pi} \cos \phi_{\pi} \\ -P_{\Lambda} \sin \theta_{\Lambda} \sin \phi_{\Lambda} + P_{p} \sin \theta_{p} \sin \phi_{p} + P_{\pi} \sin \theta_{\pi} \sin \phi_{\pi} \\ -P_{\Lambda} \cos \theta_{\Lambda} + P_{p} \cos \theta_{p} + p_{\pi} \cos \theta_{\pi} \\ -\sqrt{P_{\Lambda}^{2} + m_{\Lambda}^{2}} + \sqrt{P_{p}^{2} + m_{p}^{2}} + \sqrt{P_{\pi}^{2} + m_{\pi}^{2}} \end{pmatrix}.$$

$$(24)$$



Example: 3-C Fit

We have unmeasured and measured variables as:

$$\mathbf{u} = \{P_{\Lambda}\}; \quad \mathbf{m} = \{\theta_{\Lambda}, \phi_{\Lambda}, P_{\rho}, \theta_{\rho}, \phi_{\rho}, P_{\pi}, \theta_{\pi}, \phi_{\pi}\}. \tag{25}$$

Since there is 1 unmeasured variable with 4 kinematical constraints, this is a 4-1 = 3-Constrained fit. Let us substitute Eq (24) and (25) into Eq (11)s and its resulting equations. We get $\mathbf{F_u}$ and $\mathbf{F_m}$ as

$$\mathbf{F_{u}} = \begin{pmatrix} \frac{\partial f_{1}}{\partial P_{\Lambda}} \\ \frac{\partial f_{2}}{\partial P_{\Lambda}} \\ \frac{\partial f_{3}}{\partial P_{\Lambda}} \\ \frac{\partial f_{4}}{\partial P_{\Lambda}} \end{pmatrix}; \quad \mathbf{F_{m}} = \begin{pmatrix} \frac{\partial f_{1}}{\partial \theta_{\Lambda}} & \cdots & \frac{\partial f_{1}}{\partial \phi_{\pi}} \\ \frac{\partial f_{2}}{\partial \theta_{\Lambda}} & \cdots & \frac{\partial f_{2}}{\partial \phi_{\pi}} \\ \frac{\partial f_{3}}{\partial \theta_{\Lambda}} & \cdots & \frac{\partial f_{3}}{\partial \phi_{\pi}} \\ \frac{\partial f_{4}}{\partial \theta_{\Lambda}} & \cdots & \frac{\partial f_{4}}{\partial \phi_{\pi}} \end{pmatrix}$$
(26)

We have all the matrices to calculate in each step. By applying an appropriate variance matrix and employing χ^2 selection criteria, we can kinematically fit the particles.