

# Transport- and Measure-Theoretic Approaches for Modeling, Identifying, and Forecasting Dynamical Systems

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**Yunan Yang**, Cornell University

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Kantorovich Initiative Seminar Series. Online.

List of works:

- Optimal transport for parameter identification of chaotic dynamics via invariant measures. 2023. *SIADS*.
- Learning dynamics on invariant measures using PDE-constrained optimization. 2023. *Chaos*.
- Measure-Theoretic Time-Delay Embedding. arXiv:2409.08768.
- Invariant Measures in Time-Delay Coordinates for Unique Dynamical System Identification. arXiv:2412.00589.

# Collaborators



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(f) Maria Oprea  
(Cornell)



(g) Romit Malik (PSU)

# **Data-Driven Modeling of Dynamical Systems**

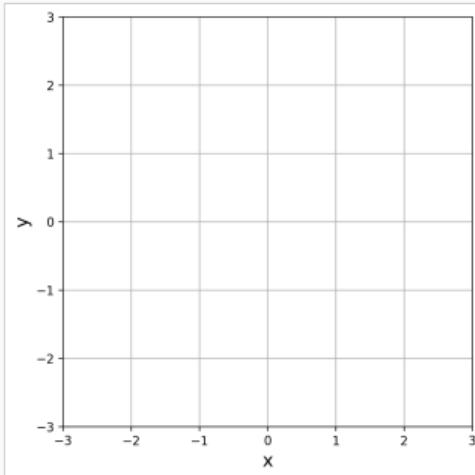
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# Data-Driven Modeling for Dynamical System

$X$



State space



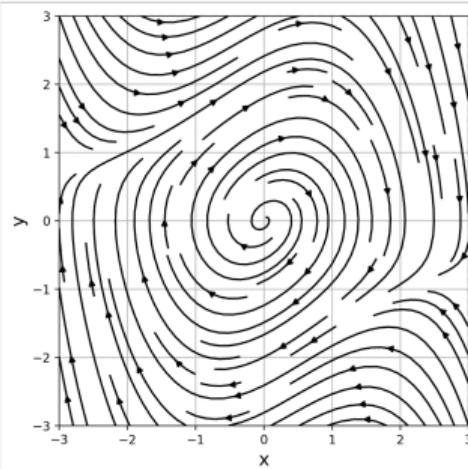
$\dot{x} = v(x)$

data-driven  
modeling

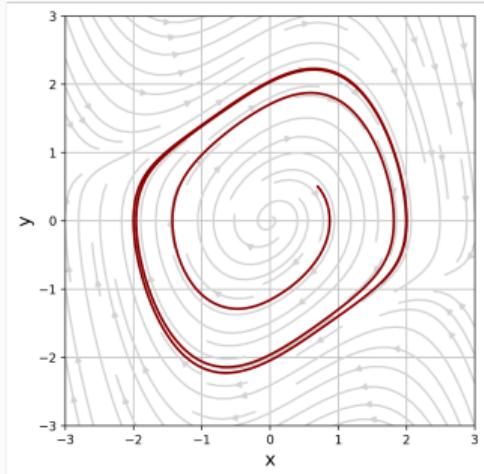
$\{x(t_k)\}_{k=0}^{N-1}$



Evolution rule



Trajectory samples



## Parameter Identification

A general parameterized dynamical system may take the form

$$\dot{\mathbf{x}} = \begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{z} \end{pmatrix} = v(x, y, z; \underbrace{\sigma, \rho, \beta}_{\theta}) \approx v(\mathbf{x}, \theta)$$

where the mathematical approximation  $v \approx v(\cdot, \theta)$  is given by

- polynomials, e.g., SINDy [Brunton et al., 2016], [Schaeffer-Tran-Ward, 2018]
- other basis functions, e.g., piecewise polynomials, RBFs, Fourier, etc.
- neural networks [many references], and so on,

where  $\theta$  corresponds to **expansion coefficients, neural network weights**, etc.

# Unique Challenges for Chaotic Systems: Chaos

**Challenge One:** The initial condition of the system is unknown.

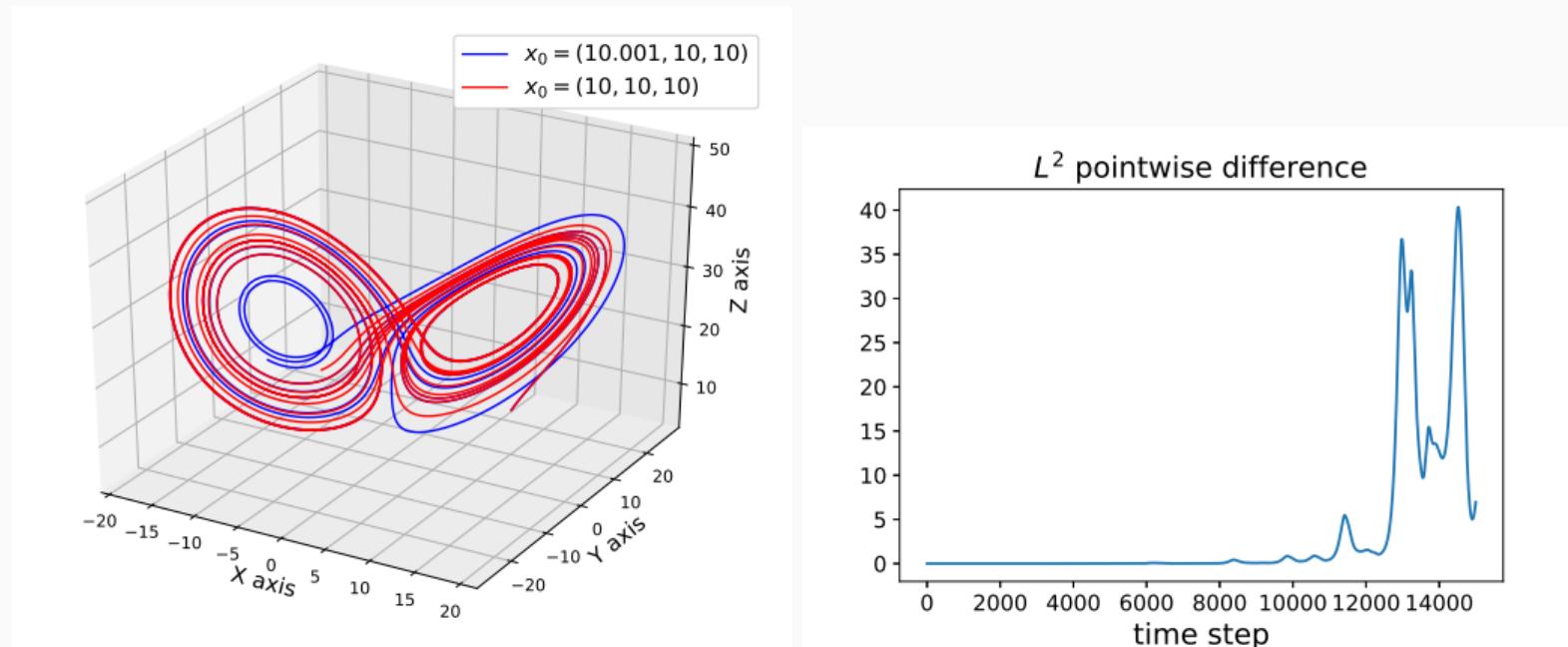


Figure: The comparison between  $\mathbf{x}_0 = [10.001, 10, 10]$  and  $\mathbf{x}_0 = [10, 10, 10]$ .

# Unique Challenges for Chaotic Systems: Noises

**Challenge Two:** The time trajectories contain noise.

No noise

$$\dot{\mathbf{x}} = f(\mathbf{x}).$$

Extrinsic noise

$$\mathbf{x}_\gamma = \mathbf{x} + \gamma, \dot{\mathbf{x}} = f(\mathbf{x}).$$

Intrinsic noise

$$\dot{\mathbf{x}} = f(\mathbf{x}) + \omega.$$

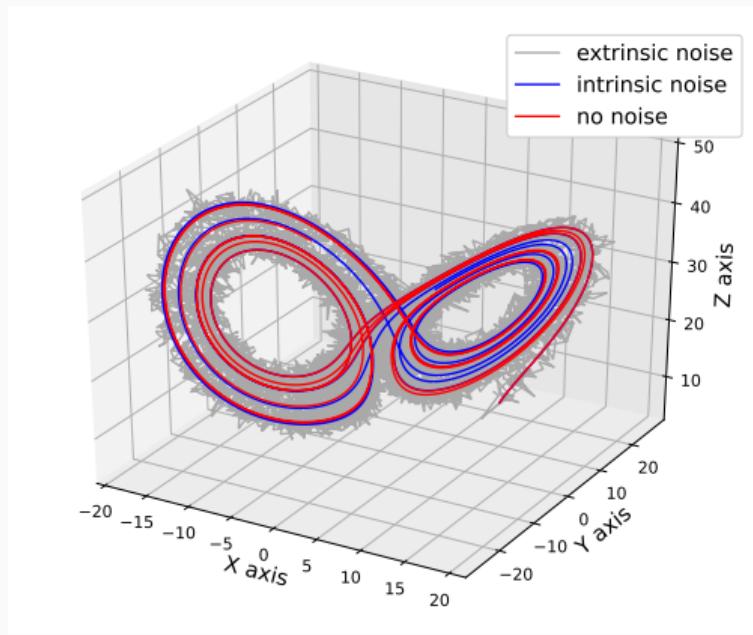


Figure: The comparison among the three cases.

# Unique Challenges for Chaotic Systems: Poor Data Quality

**Challenge Three:** Cannot measure the Lagrangian particle velocity flow

Measurements  $\{\mathbf{x}_i\}$  are not good enough to estimate the particle velocity  $\dot{\mathbf{x}}$  evaluated at  $\{\mathbf{x}_i\}$

$$\hat{v} \approx \frac{\mathbf{x}_{i+1} - \mathbf{x}_i}{t_{i+1} - t_i}$$

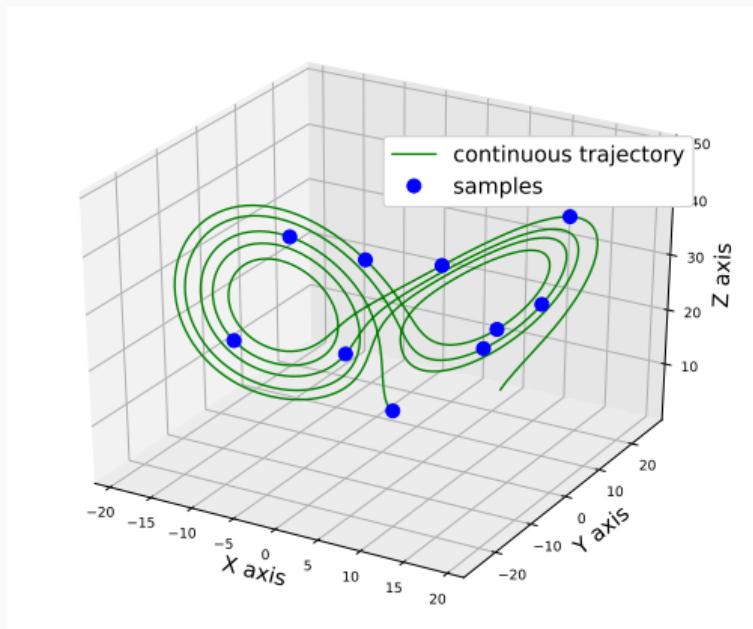


Figure: The continuous trajectory vs the samples

## The Eulerian Approach

Often, chaotic systems admit well-defined **statistical properties**:

$$\mu_{x,T}(B) = \frac{1}{T} \int_0^T \mathbb{1}_B(\mathbf{x}(s)) ds = \frac{\int_0^T \mathbb{1}_B(\mathbf{x}(s)) ds}{\int_0^T \mathbb{1}_{\mathbb{R}^d}(\mathbf{x}(s)) ds},$$

where  $\mathbf{x}(t)$  is a trajectory starting with  $\mathbf{x}(0) = x$ , and  $\mu_{x,T}$  is called the *occupation measure*. We call  $\mu^*$  a **physical measure** if  $\lim_{T \rightarrow \infty} \mu_{x,T} = \mu^*$  for  $x \in U$ ,  $\text{Leb}(U) > 0$ .

**Data Change:** take  $\mu^*$  as **observation data** instead of the **trajectory**  $\mathbf{x}(t)$ .

**Model Change:**  $\mu^*$  is the **steady-state** solution to the continuity equation:

$$\frac{\partial \rho(\mathbf{x}, t)}{\partial t} + \nabla \cdot (\mathbf{v}(\mathbf{x}, \theta) \rho(\mathbf{x}, t)) = 0.$$

# Road map: from Lagrangian to Eulerian

ODE model  $\dot{\mathbf{x}} = \mathbf{v}(\mathbf{x})$ , observe  $\{\mathbf{x}(t_i)\}_i$



Occupation measure

$$\mu_{\mathbf{x}, T}(B) = \frac{1}{T} \int_0^T \mathbb{1}_B(\mathbf{x}(s)) ds$$

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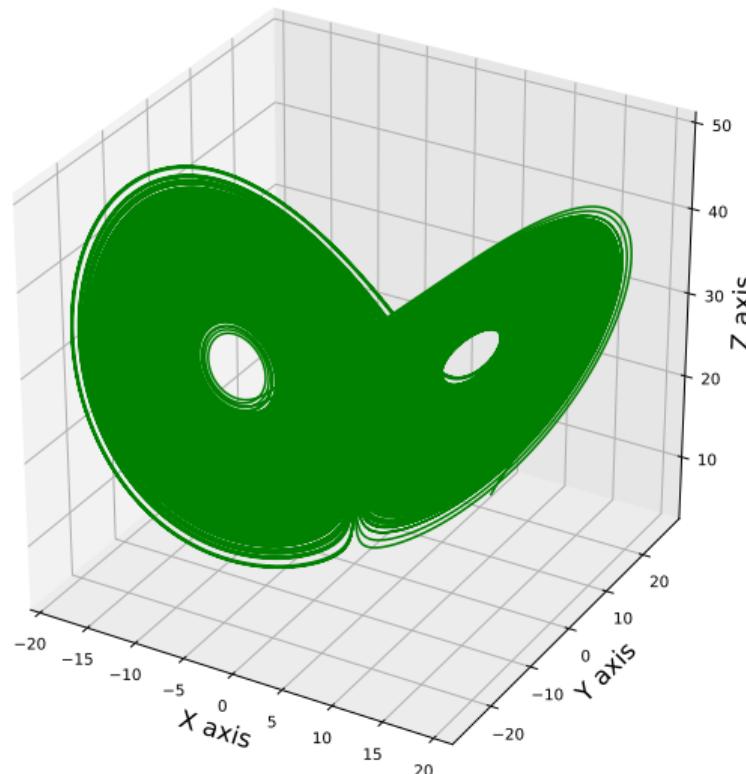
physical measure  $\mu^*$



**Stationary** distributional solutions of

$$\frac{\partial \rho(\mathbf{x}, t)}{\partial t} + \nabla \cdot (\mathbf{v}(\mathbf{x}, \theta) \rho(\mathbf{x}, t)) = 0.$$

Lorenz system (without noise)



## The New Method – A PDE-Constrained Optimization Problem

We treat the parameter identification problem for the dynamical system as a PDE-constrained optimization problem:

$$\theta = \operatorname{argmin}_{\theta} d(\rho^*, \rho(\theta)),$$

$$\text{s.t. } \frac{\partial \rho}{\partial t} = -\nabla \cdot (v(\mathbf{x}, \theta) \rho(\mathbf{x}, t)) \left[ + \frac{1}{2} \frac{\partial^2 D_{ij} \rho}{\partial x_i \partial x_j} \right] = 0.$$

$\rho^*$  : the observed occupation measure converted from time trajectories

$\rho(\theta)$  : the distributional steady-state solution of the PDE

$d$  : an appropriate metric that captures the essential differences, e.g.,  $W_2$  metric

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The gain is to work with a much **More Stable** inverse problem!

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Next: an objective function comparing distributions

# Optimal Transport

- Monge (1781)
- Kantorovich (1975)
- Brenier, Caffarelli, Gangbo, McCann, Benamou, Otto, Villani, Figalli, etc. (1990s - present)

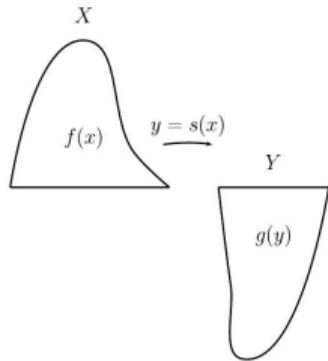


Figure: Proposed by Monge in 1781

# Optimal Transport

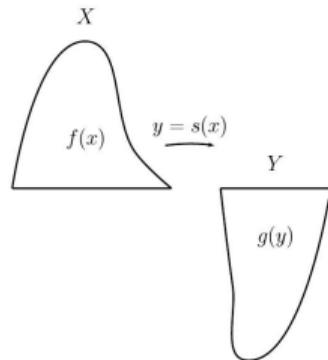


Figure: Proposed by Monge in 1781

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- **Data Assimilation** (Reich, Vidard, Bocquet...)
- **Hyperbolic Model Reduction** (Mula, Peherstorfer,Ravela)
- **Image Processing** (T. Chan, Peyré, C. B. Schönlieb...)
- **Inverse Problems** (Bao, Marzouk, Engquist, Singer, Y,...)
- **Machine Learning** (Cuturi,Peyré, Solomon, ...)
- **Sampling** (Marzouk, Rigollet, Chewi, ...)
- **And more**

# The Wasserstein Distance

## Definition of the Wasserstein Distance

For  $f, g \in \mathcal{P}(\Omega)$  ( $f, g \geq 0$  and  $\int f = \int g = 1$ ), the Wasserstein distance is formulated as

$$W_p(f, g) = \left( \inf_{T \in \mathcal{M}} \int |x - T(x)|^p f(x) dx \right)^{\frac{1}{p}} \quad (1)$$

$\mathcal{M}$ : the set of all maps that rearrange the distribution  $f$  into  $g$ .

The commonly used cases include  $p = 1$  and  $p = 2$ .

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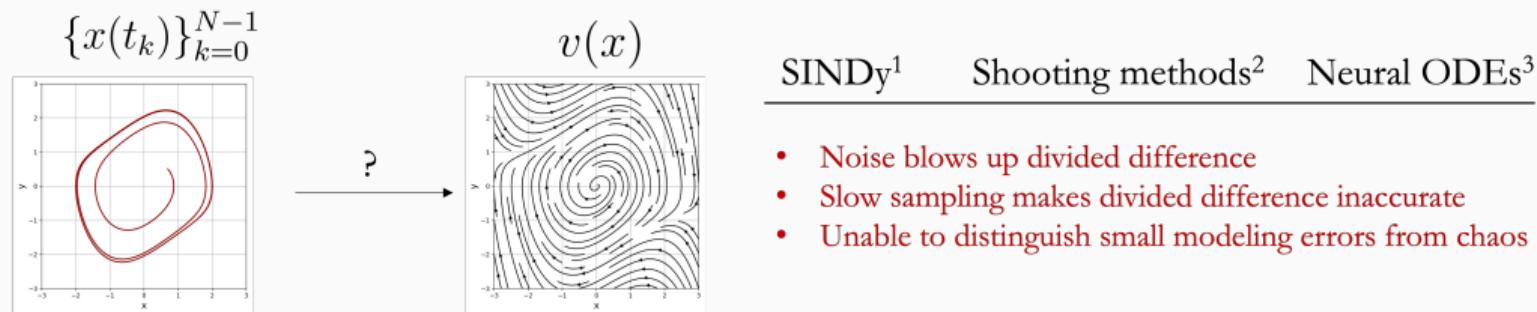
## Properties of $W_2$ as the loss function

(1) Provide better optimization landscape for Nonlinear Inverse Problems:

$$\theta^* = \operatorname{argmin}_{\theta} W_2^2(\rho(\theta), \rho^*)$$

(2) Robust in Inversion with Noisy Data (equivalent to  $H^{-1}$  norm)

# Recap: our approach from Lagrangian view to Eulerian perspective



- Noise blows up divided difference
- Slow sampling makes divided difference inaccurate
- Unable to distinguish small modeling errors from chaos

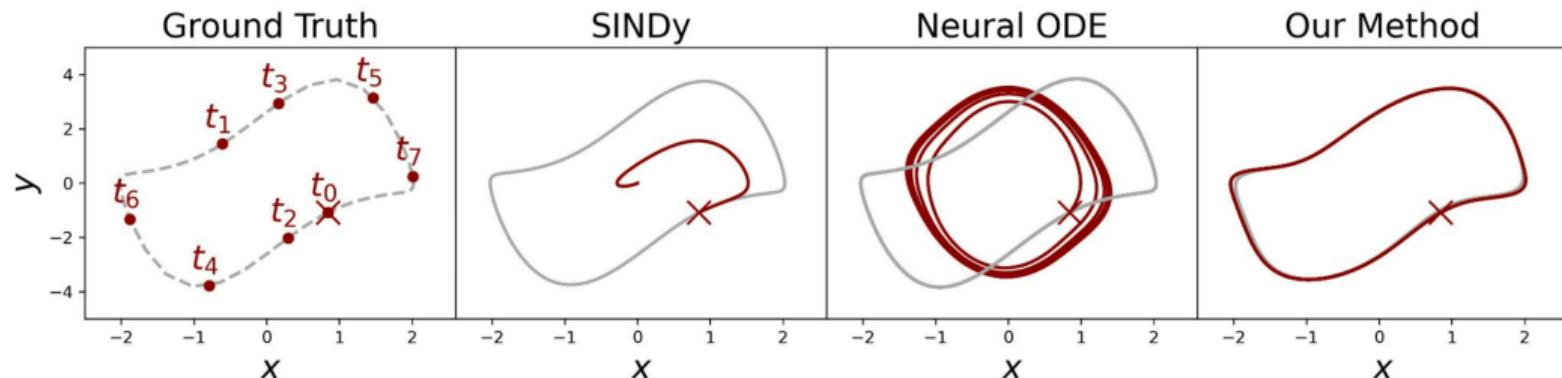
$Data$ $\rho^* := \frac{1}{N} \sum_{k=0}^{N-1} \delta_{x(t_k)}$	$Forward\ Model$ $\theta \mapsto \rho(\theta)$	$Objective\ Function$ $\min_{\theta \in \Theta} \mathcal{J}(\rho(\theta), \rho^*)$
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<sup>1</sup>Brunton, S. L., Proctor, J. L., & Kutz, J. N. (2016). Discovering governing equations from data by sparse identification of nonlinear dynamical systems. *Proceedings of the national academy of sciences*, 113(15), 3932-3937.

<sup>2</sup>Michalik, C., Hannemann, R., & Marquardt, W. (2009). Incremental single shooting—a robust method for the estimation of parameters in dynamical systems. *Computers & Chemical Engineering*, 33(7), 1298-1305.

<sup>3</sup>Chen, R. T., Rubanova, Y., Bettencourt, J., & Duvenaud, D. K. (2018). Neural ordinary differential equations. *Advances in neural information processing systems*, 31.

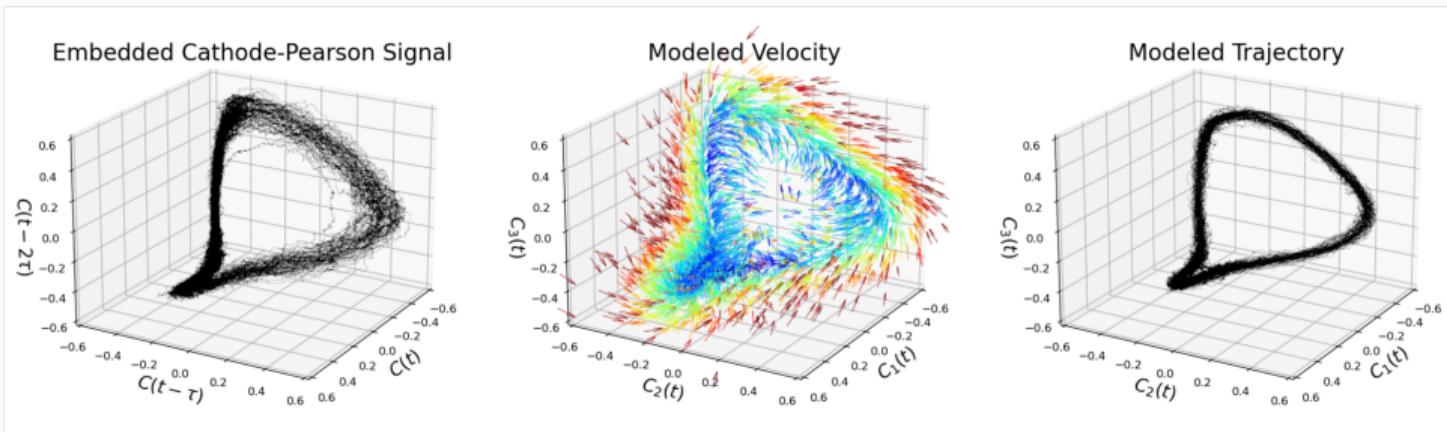
# Comparison with Other Methods



Method	Sampling Freq.	Wall-Clock Time (s)	Error
SINDy	10.00	$2 \cdot 10^{-2}$	$5.6 \cdot 10^{-3}$
Neural ODE	10.00	$5 \cdot 10^2$	$5.32 \cdot 10^{-3}$
Ours	10.00	$5 \cdot 10^2$	$1.14 \cdot 10^{-1}$

Method	Sampling Freq.	Wall-Clock Time (s)	Error
SINDy	0.25	$10^{-2}$	3.52
Neural ODE	0.25	$5 \cdot 10^2$	1.81
Ours	0.25	$5 \cdot 10^2$	$6.79 \cdot 10^{-2}$

# Application to Real-World Data: Hall-Effect Thruster

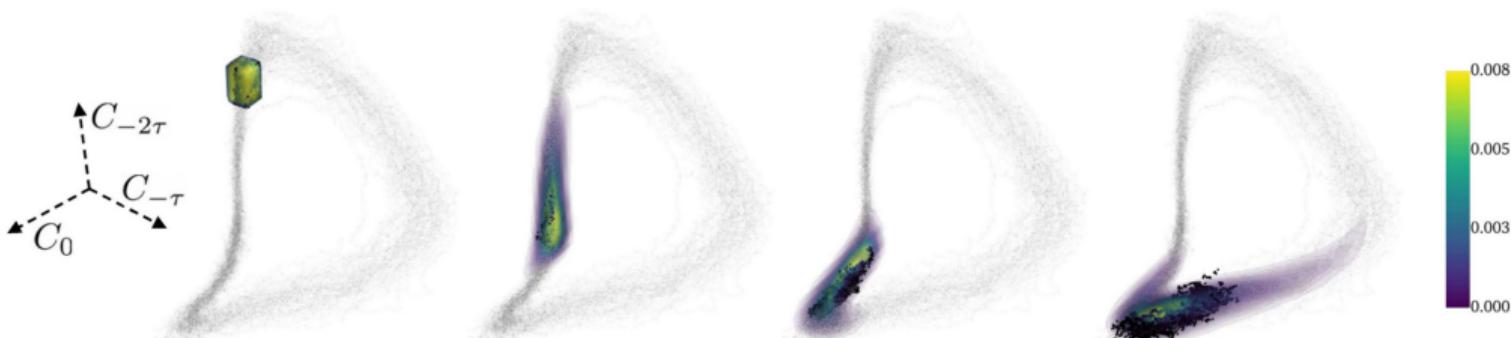


time = 0.00

time = 0.17

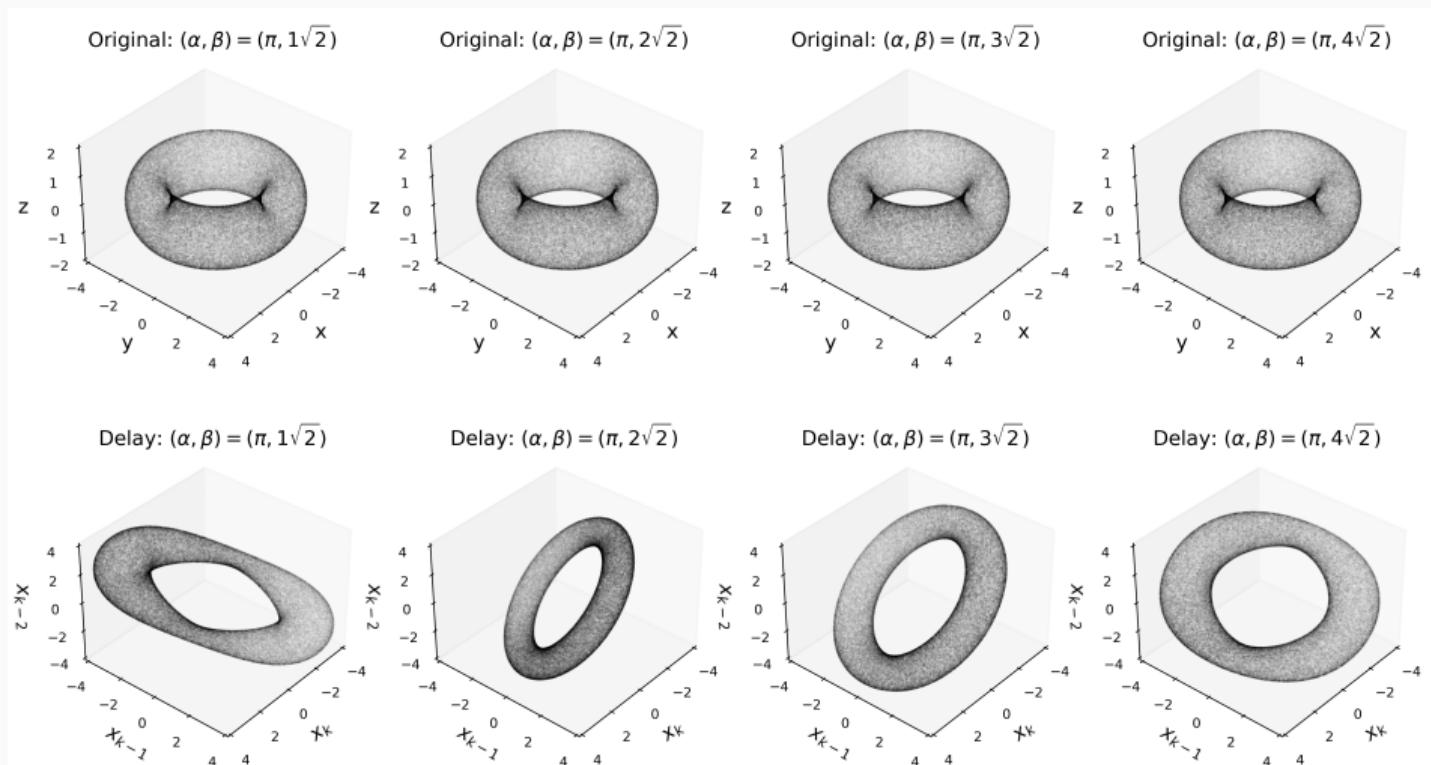
time = 0.33

time = 0.50



# Limitation: Nonuniqueness

$$T_{\#}\mu = \mu \text{ & } S_{\#}\mu = \mu \not\Rightarrow T = S$$



# Invariant Measures in Time-Delay Coordinates for *Unique* Dynamical System Identification

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# Takens' Embedding Theorem

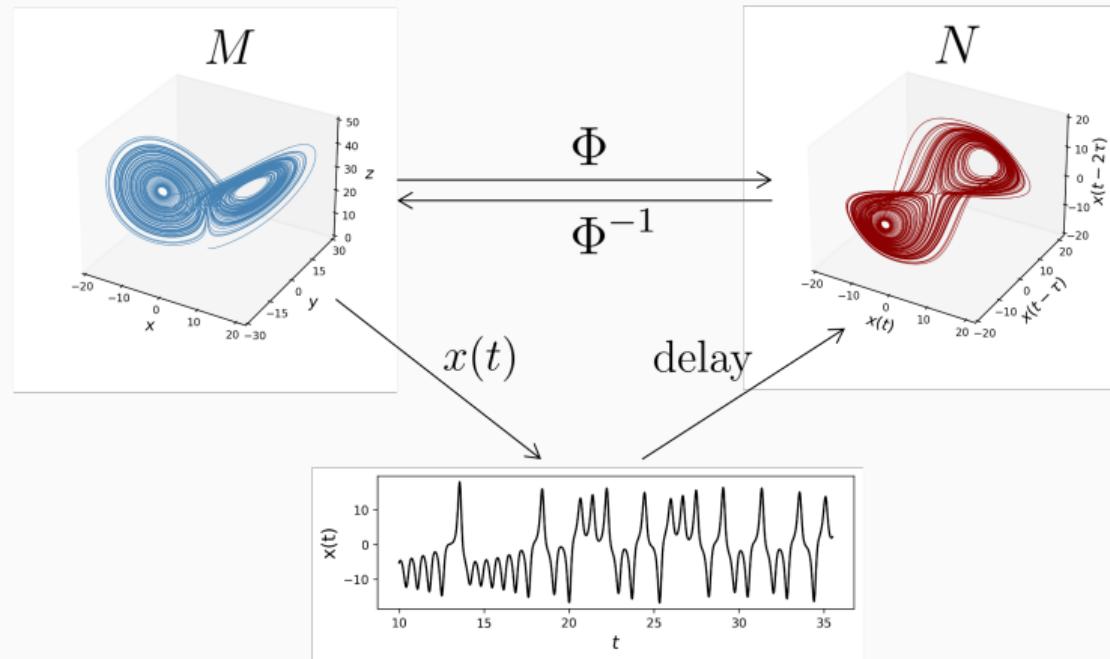
## Theorem (Takens, 1981)

Let  $M$  be a compact manifold of dimension  $m$ . For pairs  $(y, T)$ , where  $T \in C^2(M, M)$  and  $y \in C^2(M, \mathbb{R})$ , it is a generic property that the mapping  $\Phi_{(y,T)} : M \rightarrow N \subseteq \mathbb{R}^{2d+1}$  given by  $\Phi_{(y,T)}(\mathbf{x}) := (y(\mathbf{x}), y(T(\mathbf{x})), \dots, y(T^{2m}(\mathbf{x})))$  is an embedding of  $M$  in  $\mathbb{R}^{2d+1}$ .

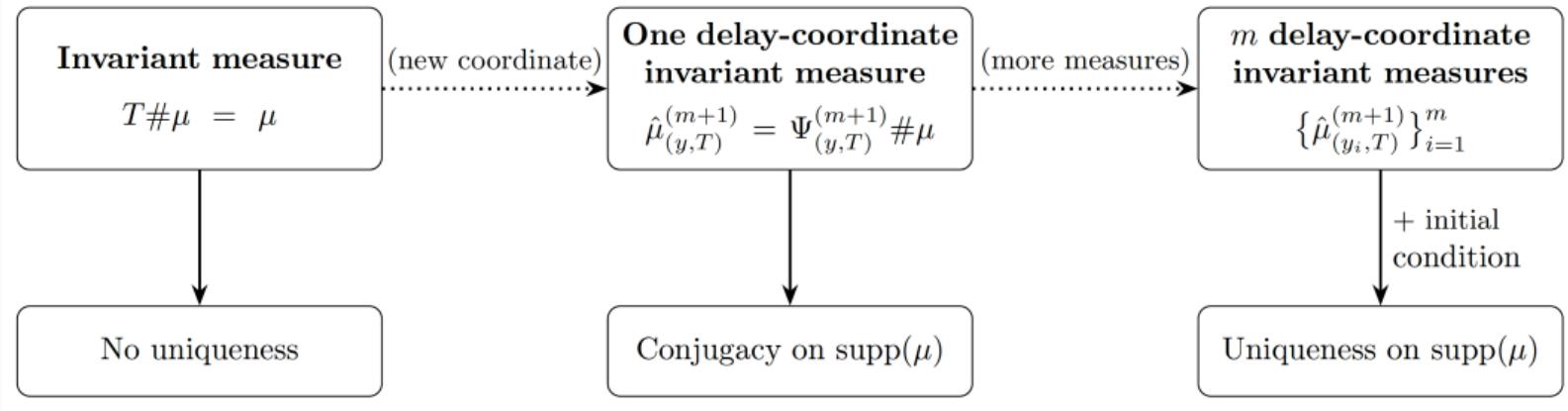
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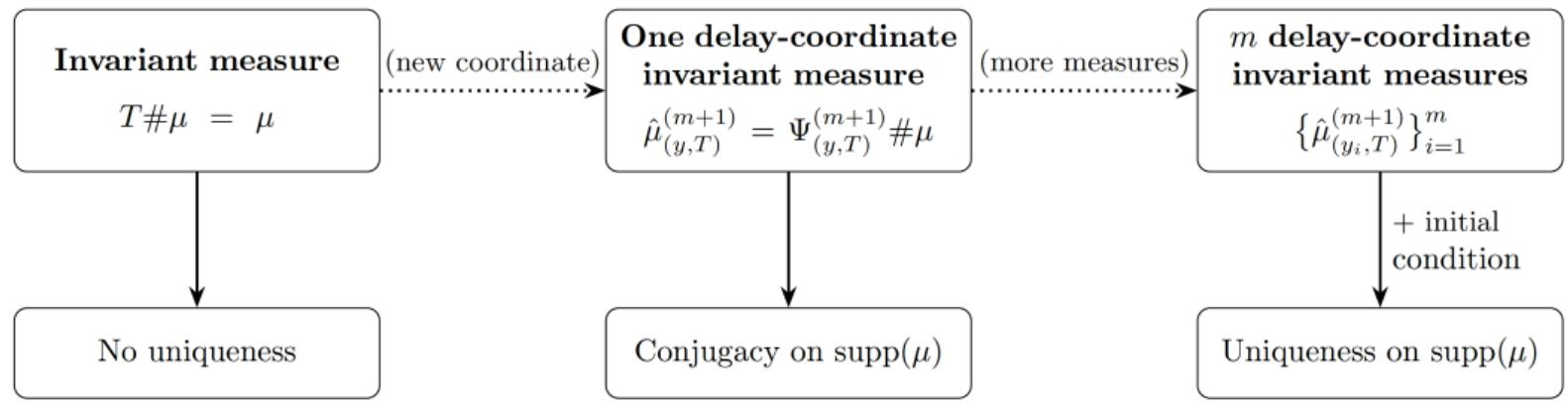
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# Invariant Measures in Time-Delay Coordinates for Uniqueness

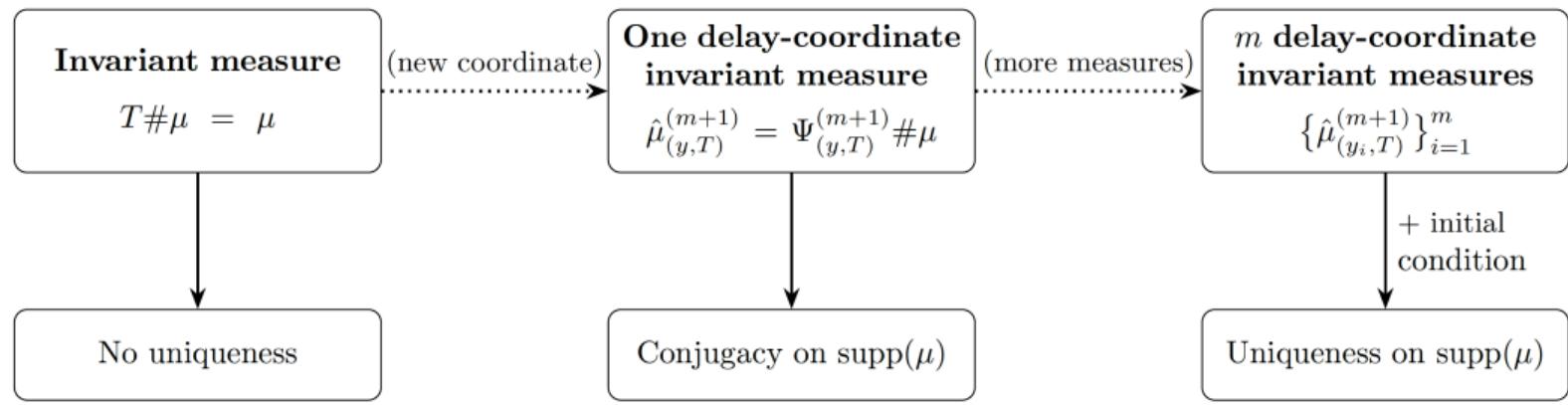


# Invariant Measures in Time-Delay Coordinates for Uniqueness



**Theorem 1.** The equality  $\hat{\mu}_{(y,T)}^{(m+1)} = \hat{\nu}_{(y,S)}^{(m+1)}$  implies  $T|_{\text{supp}(\mu)}$  and  $S|_{\text{supp}(\nu)}$  are topologically conjugate, for almost every  $y \in C^1(U, \mathbb{R})$ .

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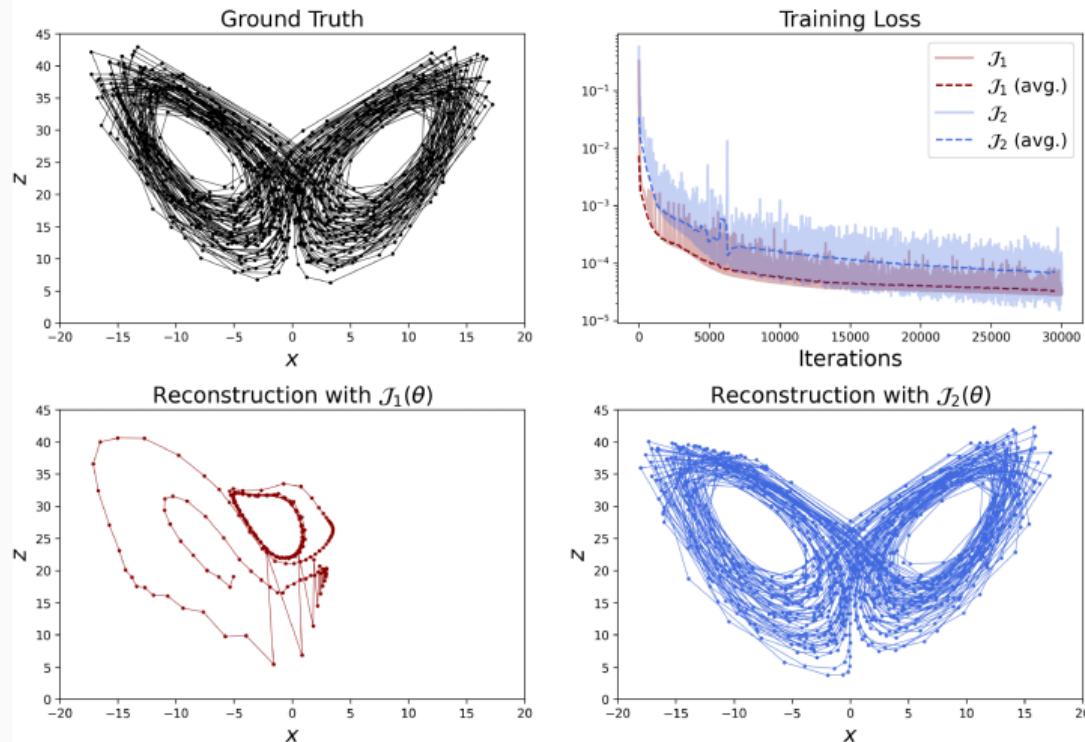
**Theorem 2.** The conditions below imply that  $T = S$  on  $\text{supp}(\mu)$ , for a.e.  $Y \in C^1(U, \mathbb{R}^m)$ :

1. there exists  $x^* \in B_{\mu,T} \cap \text{supp}(\mu)$ , such that  $T^k(x^*) = S^k(x^*)$  for  $1 \leq k \leq m-1$ , and
2.  $\hat{\mu}_{(y_j,T)}^{(m+1)} = \hat{\mu}_{(y_j,S)}^{(m+1)}$  for  $1 \leq j \leq m$ , where  $Y := (y_1, \dots, y_m)$  is a vector-valued observable.

# Numerical Example

$$\mathcal{J}_1(\theta) := \mathcal{D}(T_\theta \# \mu^*, T^* \# \mu^*), \quad \mathcal{J}_2(\theta) := \mathcal{D}(T_\theta \# \mu^*, T^* \# \mu^*) + \mathcal{D}(\Psi_\theta \# \mu^*, \Psi^* \# \mu^*).$$

$\Psi_\theta$  is the delay map based on  $T_\theta$ , and  $\Psi^*$  is the true delay map.



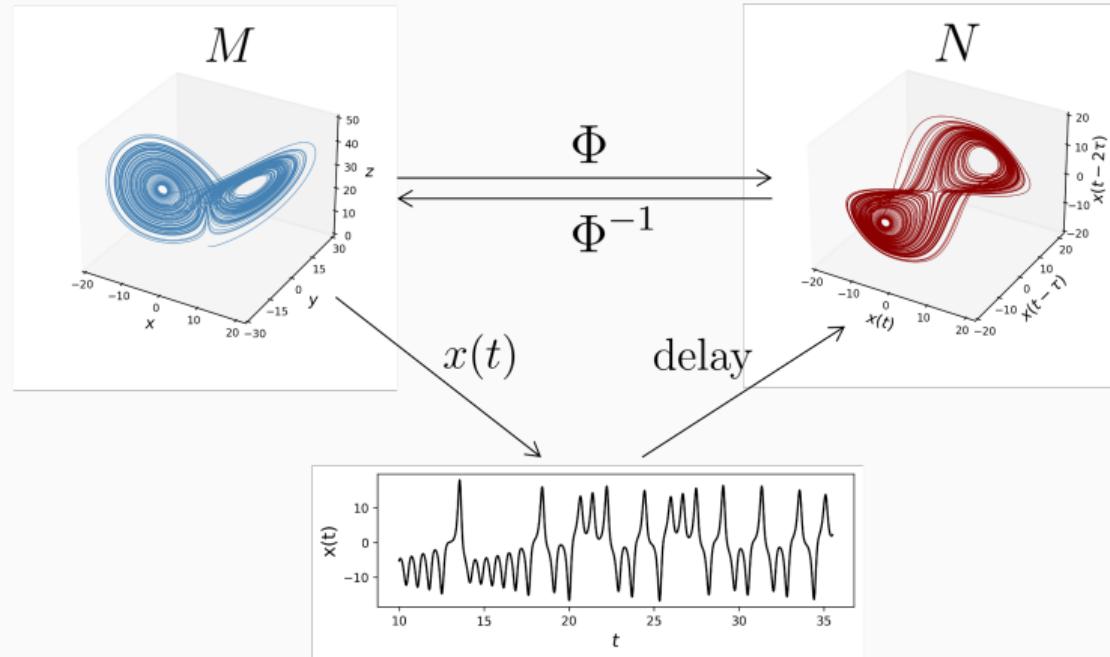
## **Embedding Over the Probability Space $\mathcal{P}_2(M)$**

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# Takens' Embedding Theorem (Again)

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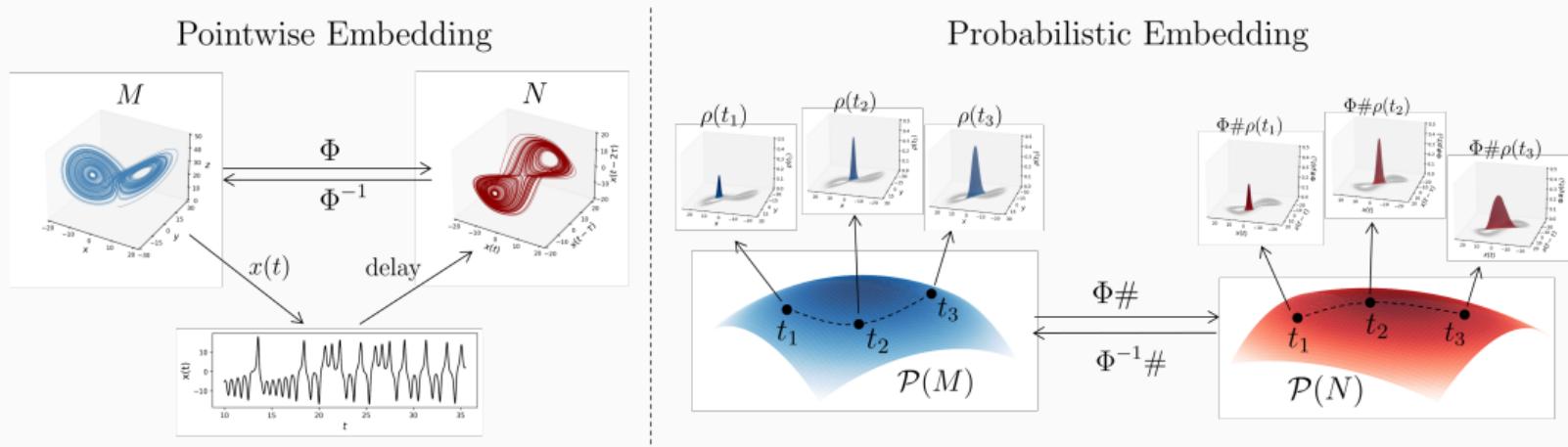
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# Measure-Theoretic Embedding

## Pointwise embedding ( $\Phi$ )

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### Definition (Differentiability of operator $\mathcal{P}_2(M) \rightarrow \mathcal{P}_2(N)$ )

A map  $\Psi : \mathcal{P}_2(M) \rightarrow \mathcal{P}_2(N)$  is differentiable if for all  $\mu \in \mathcal{P}_2(M)$  there is a bounded linear operator  $d\Psi_\mu : T_\mu \mathcal{P}_2(M) \rightarrow T_{\Psi(\mu)} \mathcal{P}_2(N)$  s.t. for any differentiable curve  $t \mapsto \mu_t$  through  $\mu$ , the curve  $t \mapsto \Psi(\mu_t)$  is differentiable with velocity  $v_t$  and  $d\Psi_{\mu_t}(v_t) = \frac{d}{dt} \Psi(\mu_t)$ .

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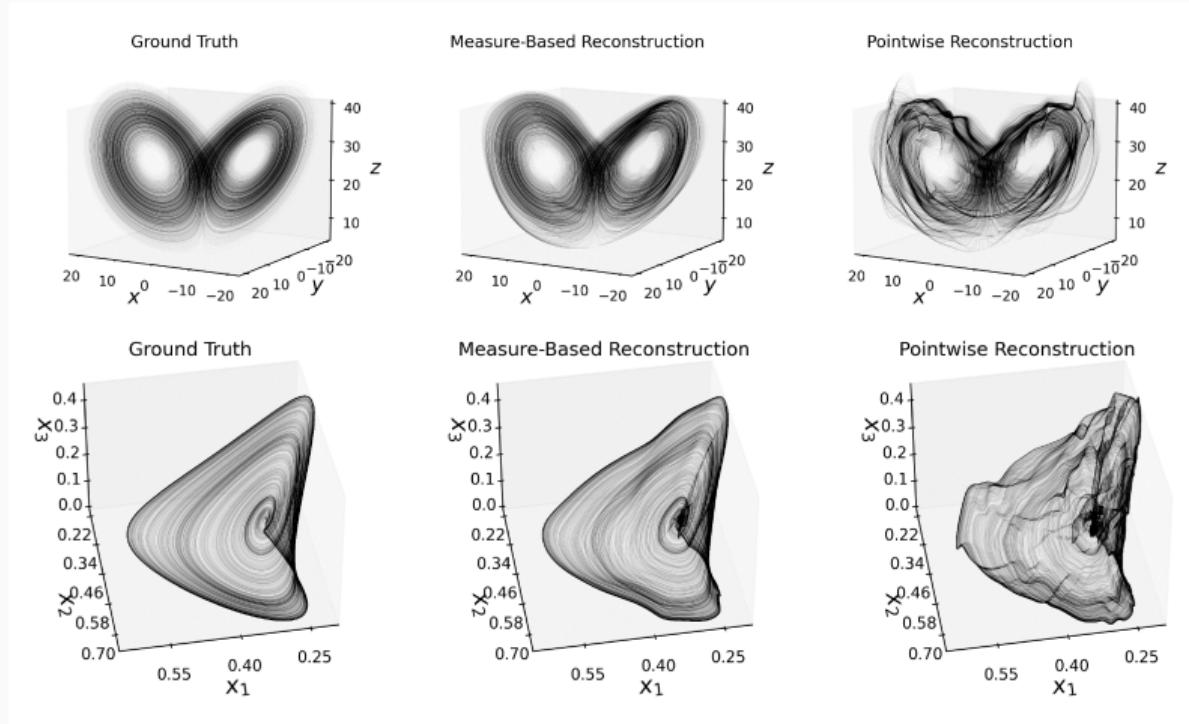
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### Theorem (Our Main Result)

If  $\Phi : M \rightarrow N$  is an embedding between differentiable manifolds, then the map  $\Phi\# : \mathcal{P}_2(M) \rightarrow \mathcal{P}_2(N)$  is also an embedding.

# Numerical Example

$$\underbrace{\mathcal{L}_p(\theta)}_{\text{pointwise loss}} = \frac{1}{N} \sum_{i=1}^N \|x_i - \mathcal{R}_\theta(\Phi(x_i))\|_2^2, \quad \underbrace{\mathcal{L}_m(\theta)}_{\text{measure-theoretic loss}} = \frac{1}{K} \sum_{i=1}^K \mathcal{D}(\mu_i, \mathcal{R}_\theta \# (\Phi \# \mu_i)).$$



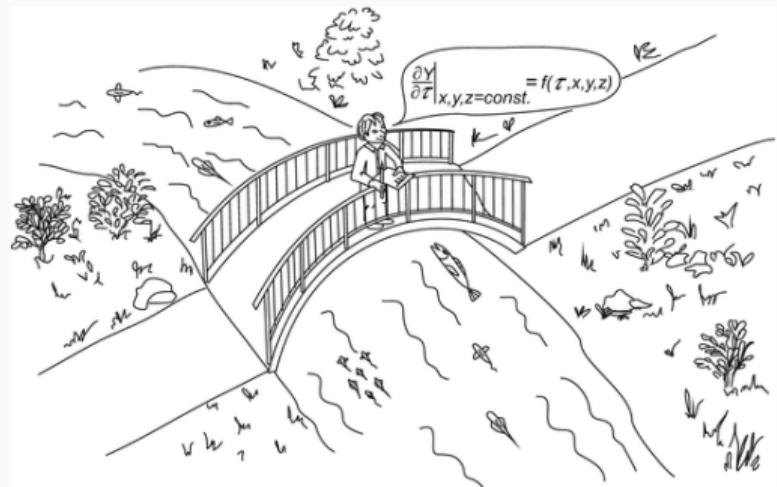
## **Conclusion**

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# Conclusions



(a) Lagrangian view



(b) Eulerian view

[Bird-Stewart-Lightfoot, Transport Phenomena, 2002]

# Conclusions

## Summaries

- From **Lagrangian** to **Eulerian** to tackle chaos (ODE  $\Rightarrow$  PDE problem)
- Using **optimal transport** to study dynamical system
  1. Invariant measure matching
  2. Invariant measure in time-delay coordinate matching
  3. Generalize **pointwise** embedding to **measure-theoretic** embedding

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- Using **optimal transport** to study dynamical system
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  3. Generalize **pointwise** embedding to **measure-theoretic** embedding

## Outlook

There is great potential for using optimal transport in data-driven modeling of dynamical systems.

# Acknowledgments

Research support from



Thank you for the attention!