## DEPARTMENT OF PHYSICS INDIAN INSTITUTE OF TECHNOLOGY, MADRAS

## PH1020 Physics II

## Tutorial 1 (22.1.2018)

1. Consider an infinite uniformly-charged plate occuping the xoy plane, carrying a surface charge density  $\sigma$ , with a circular hole of radius a centered at the origin. Find the force on a charge Q lying on the z-axis. If the charge is negative, (i.e., Q < 0), discuss the force on it when released at a distance x close to the origin (i.e., z << a).

Ans

The Bield Pt:  $\vec{Y} = z\hat{e}_z$ 

The source pts: (v', 4) = v'cos 4'êx + v'sin 4'êy

Element of charge da! = or'dr'd4

The Jonce on  $Q = \frac{Q}{4\pi\epsilon_0} \int_{0}^{\infty} \frac{d^{2}\pi}{(z\hat{e}_z - v'\cos\varphi'\hat{e}_x - v'\sin\varphi'\hat{e}_y)} v'dv'd\varphi'$ 

Integrations over wsp' and sing' vanish

The force on 
$$Q$$

$$= \overline{FQ} = \frac{Q}{4\pi\epsilon_0} \frac{2\pi}{a} \int_{a}^{\infty} \frac{\partial z}{(z'+v''a')^3/a} v' dv'$$

$$=\frac{Q}{4\pi E_0} 2\pi - z \frac{1}{2} \int_{a}^{\infty} \frac{d(y'^2 + z^2)}{(z^2 + y'^2)^3 \omega}$$

$$= -\frac{Q}{2\varepsilon_0} \frac{-Z}{\sqrt{Z^2 + Y^2 Z^2}} \frac{dz}{dz} = \frac{Q-Z}{2\varepsilon_0} \frac{dz}{\sqrt{Z^2 + \alpha^2}}$$

(Hote: At large distances,  $\overrightarrow{F_Q} = \frac{Q - e_Z}{2E_0}$ , which is exactly same as corresponding to the Silvation in which the hole is absent! So, at large distances, Q does not see the hole!)

For zzza and Qzo

$$\vec{F}_{Q} = \frac{Q - z}{2\epsilon_{0}a} \frac{\hat{e}_{z}}{(1 + \frac{z^{2}}{a^{2}})^{1/2}} = \frac{Q - z}{2\epsilon_{0}a} \left(1 + \frac{z^{2}}{a^{2}}\right)^{-1/2} \hat{e}_{z}$$

$$\simeq \frac{Q - z}{2\epsilon_{0}a} \hat{e}_{z}$$

$$\simeq \frac{Q - z}{2\epsilon_{0}a} \hat{e}_{z}$$

The motion of charge is simple harmonic.

2. Determine the electric field at the center of a sphere of radius R that carries a charge on its surface with charge density  $\sigma = \vec{k} \cdot \vec{r}$ , where  $\vec{k}$  is a constant vector.

Ans

$$\vec{dE} = \vec{k} \cdot \vec{Y} = kR \cos \theta'$$

$$\vec{dE} = \vec{l} \quad (kR \cos \theta)$$

$$4.71 E_0 R^2$$

$$\left(R^2 \sin \theta d\theta d\phi'\right) (-\hat{e}_Y)$$

(-cosóêz+sinóêxoy)

Where 
$$\hat{e}_{xoy}$$
 is the appropriate unit vector in the XOY plane  $\times$  and  $y$  components of the pleibilic field vanish.

$$\vec{dE} = 2 \times \frac{kR \cos \theta}{4\pi \epsilon_0 R^2} \left( \frac{R^2 \sin \theta}{4\pi \epsilon_0 R^2} \right) \left( -\cos \theta \right) \hat{e}_z$$

$$\vec{E} = \int d\vec{e} |_{Het} = \int \frac{e^{2\pi i}}{\int \frac{E}{2\pi \epsilon_0 R^2}} \frac{R^2 \sin \theta}{4\pi \epsilon_0 R^2} \left( \frac{R^2 \sin \theta}{4\pi \epsilon_0 R^2} \right) \left( \frac{R^2 \sin \theta}{4\pi \epsilon_0 R^2} \right) \hat{e}_z$$

$$\vec{\varphi} = 0 \ \theta = 0$$

$$= -\frac{kR}{\varepsilon_0} \hat{e}_z \int \cos^2 \theta' dG \cos \theta' = \frac{kR}{\varepsilon_0} \hat{e}_z \left(\frac{\cos^3 \theta}{3}\right)^{\frac{\pi}{2}} \hat{e}_z$$

$$= -\frac{kR}{\varepsilon_0} \hat{e}_z \int \cos^2 \theta' dG \cos \theta' = \frac{kR}{\varepsilon_0} \hat{e}_z \left(\frac{\cos^3 \theta}{3}\right)^{\frac{\pi}{2}} \hat{e}_z$$

$$= -\frac{kR}{3\varepsilon_0} \hat{e}_z$$

3. A fixed charge +q at a point O is surrounded by a continuous distribution of charge whose density  $\rho(<0)$  is a function only of the distance r from O. The total negative charge  $exceeds\ q$  in magnitude. A point charge +q', with a mass m, is free to move on a line passing through O. (a) Obtain the condition for +q' to be in equilibrium at a distance  $r_0$  from O. (b) If q' is released at a point very close to  $r_0$ , find the force acting on it.

Ans
(a) 
$$4\pi r^{2} E_{r} = \begin{cases} 2 + 4\pi \int_{0}^{r} s^{2} p(s) ds \end{cases}^{2} \frac{1}{E_{0}}$$

For  $2'$  to be in equilibrium at  $Y = Y_{0}$ ,  $E_{r0} = 0$ 

$$\Rightarrow 2 + 4\pi \int_{0}^{r} s^{2} p(s) ds = 0 \text{ Which is the required condition.}$$

(b) How let us apply the Grouss Theorem to a sphere of readous  $Y = Y_0 + S$   $4\pi (Y_0 + S)^2 E_T = \begin{cases} 2 + 4\pi \int_0^{Y_0 + S} s^2 p(s) \, ds \end{cases} \frac{1}{E_0}$ 

$$= \begin{cases} 2 + 4\pi \int_{0}^{\infty} s^{2} \rho(s) ds + 4\pi \int_{0}^{\infty} s^{2} \rho(s) ds \end{cases} = \begin{cases} 2 + 4\pi \int_{0}^{\infty} s^{2} \rho(s) ds + 4\pi \int_{0}^{\infty} s^{2} \rho(s) ds \end{cases} = \begin{cases} 4\pi r_{0}^{2} \delta \\ \rho(r_{0}) \end{cases} = \begin{cases} 4\pi r_{0}^{2} \delta \\ \rho(r$$

i.e 
$$E_r = \frac{r_0^2}{(r_0 + s)^2} \frac{SP(r_0)}{\varepsilon_0} \simeq \frac{SP(r_0)}{\varepsilon_0}$$

This means the equation of motion for the charge is  $m\ddot{s} = 2/8P(r\dot{o})/E_{D}$ 

Since, 920 The motion of the charge 2' is simple harmonic. Such oscillations are known as gradual oscillations.