

1. A cylinder of length $2L$ and radius a is centred at the origin, with the z -axis as its symmetry axis. The cylinder is uniformly polarized with polarization $\mathbf{P} = P_0 \hat{e}_z$ where P_0 is a constant. (i) Find the bound charge densities ρ_b and σ_b . (ii) Find the electric field at all points on the positive z -axis, and verify that it satisfied the appropriate boundary condition at $z = L$ (iii) Find the electric field at the origin, and sketch its magnitude as a function of ratio a/L .
2. Consider a uniform spherical free charge distribution of radius a and charge density ρ_0 . This region is filled with a medium of dielectric constant K_1 , and surrounded by a medium of dielectric constant K_2 . Find (i) the bound volume charge density everywhere in space, and (ii) the bound surface charge density on the surface of the sphere.
3. A capacitor is formed of two concentric conducting spheres of radii a and b ($a < b$), and the space between is filled with a substance. The dielectric constant of the substance at a distance r from the centre is $\frac{c+r}{r}$, where c is constant. The outer sphere is earthed and the inner sphere is charged. Calculate the capacitance of the system.
4. An insulated spherical conductor in air carries a charge q . The conductor is now surrounded by a concentric spherical shell of dielectric of radii b and c , ($c > b$), whose dielectric constant is a function $k(r)$ of the radial distance r from the centre. Calculate the electrostatic energy.
5. A dielectric sphere of radius R contains a uniform distribution of free charge with charge density ρ_f . Find the potential at the centre of the sphere.
6. A sphere of radius R and dielectric constant k , centered at the origin of coordinates, is placed in a constant field E_0 directed along the z -axis. The corresponding electrostatic potential is given by $\phi(r, \theta, \varphi) = (-E_0 r + b_1 r^{-2}) \cos \theta$ outside the sphere, and $\phi(r, \theta, \varphi) = (b_2 r) \cos \theta$ inside the sphere. Find (i) the constants b_1 and b_2 , in terms of k , E_0 and R , (ii) the electric field at all points in space, (iii) the polarization \mathbf{P} of the sphere, and the dipole moment of the sphere about the origin, (iv) the volume and surface densities of the bound charge in the sphere.

TUTORIAL SHEET - 4

T4-1

SOLUTIONS

1. (i) Since $\rho_f = 0$, $\rho_b = 0$



$$\sigma_b = \vec{P} \cdot \hat{n} = \begin{array}{l} P_0 \text{ through the top surface} \\ -P_0 \text{ through the bottom surface} \\ 0 \text{ through the curved surface} \end{array}$$

(ii) The electric field is due to the bound surface charges in the top and bottom surfaces alone. \vec{E} at a height z above the center of a circular disc of charge density σ_b in the xy plane with its center at the origin is,

$$\vec{E} = \frac{\sigma_b}{2\epsilon_0} \left[1 - \frac{z}{(a^2 + z^2)^{1/2}} \right] \hat{e}_z$$

Here $z > L$

$$\begin{aligned} \vec{E} &= \frac{P_0}{2\epsilon_0} \left[1 - \frac{(z-L)}{[a^2 + (z-L)^2]^{1/2}} \right]^{1/2} - \left[1 + \frac{(z+L)}{[a^2 + (z+L)^2]^{1/2}} \right] \hat{e}_z \\ &= \frac{P_0}{2\epsilon_0} \left[\frac{z+L}{[a^2 + (z+L)^2]^{1/2}} - \frac{(z-L)}{[a^2 + (z-L)^2]^{1/2}} \right] \hat{e}_z \end{aligned}$$

|||, for points inside the cylinder, $z < L$

$$\vec{E} (z < L) = -\frac{P_0}{2\epsilon_0} \left[2 - \frac{(L-z)}{[a^2 + (L-z)^2]^{1/2}} - \frac{(L+z)}{[a^2 + (L+z)^2]^{1/2}} \right] \hat{e}_z$$

At just outside the top surface, $z = L$ and the field is

$$\vec{E}_{out} = \frac{P_0}{\epsilon_0} \left[\frac{L}{(a^2 + 4L^2)^{1/2}} \right] \hat{e}_z$$

The field just inside the cylinder is obtained by substituting L for z in the expression for $\vec{E} (z < L)$

$$E_{in} = -\frac{P_0}{\epsilon_0} \left[1 - \frac{L}{(a^2 + 4L^2)^{1/2}} \right] \hat{e}_z \quad \text{which is consistent with boundary condition}$$

(iii) At the origin the \vec{E} is obtained by putting $z = 0$

$$\vec{E} = -\frac{P_0}{\epsilon_0} \left[1 - \frac{L}{(a^2 + L^2)^{1/2}} \right] \hat{e}_z = -\frac{P_0}{\epsilon_0} \left[1 - \frac{1}{[1 + (a/L)^2]^{1/2}} \right] \hat{e}_z$$



T4-2

② As charge density is uniform, \vec{E} , \vec{D} and \vec{P} have only radial components.

a) From the Gauss' law, in the region $0 < r < a$

$$4\pi r^2 D_r = \frac{4}{3} \pi r^3 \rho_0 \Rightarrow D_{r < a} = \frac{\rho_0 r}{3} \Rightarrow E_{r < a} = \frac{\rho_0 r}{3\epsilon_0 k_1}$$

$$\therefore P_{r < a} = D_{r < a} - \epsilon_0 E_{r < a} = \frac{\rho_0 r}{3} \left[1 - \frac{1}{k_1} \right]$$

$$\Rightarrow \rho_b = -\vec{\nabla} \cdot \vec{P} = -\rho_0 \left[1 - \frac{1}{k_1} \right]$$

For $r > a$

$$4\pi r^2 D_r = \frac{4}{3} \pi a^3 \rho_0 \Rightarrow D_{r > a} = \frac{\rho_0 a^3}{3 r^2}$$

$$\Rightarrow E_{r > a} = \frac{\rho_0 a^3}{3\epsilon_0 r^2 k_2}$$

$$\therefore P_{r > a} = \frac{\rho_0 a^3}{3 r^2} \left[1 - \frac{1}{k_2} \right] \text{ and } \rho_b = 0$$

$$(b) \quad \sigma_b = |P_{r > a} - P_{r < a}|_{r=a} = \frac{\rho_0 a}{3} \left[\frac{1}{k_1} - \frac{1}{k_2} \right]$$

As the free charge density is zero, D_r is continuous at $r=a$ but, P_r is not, leading to the bound charge density.

T4-3

③ Apply Gauss theorem considering a sphere of radius r :

$$4\pi r^2 D_r = Q \quad a \leq r \leq b$$

where Q is the charge on the inner sphere and D_r is the radial component of \vec{D} , and the other components of \vec{D} vanish. Therefore.

$$\epsilon E_r = \frac{Q}{4\pi r^2} ; E_\theta, E_\phi = 0$$

$$\Rightarrow E_r = \frac{Q}{4\pi\epsilon_0 r(c+r)} = -\frac{dV(r)}{dr} = \frac{-Q}{4\pi\epsilon_0 c} \left[\frac{1}{c+r} - \frac{1}{r} \right]$$

$$\therefore, V(r) = \frac{Q}{4\pi\epsilon_0 c} \ln\left(\frac{c+r}{r}\right) + A$$

Given: $V = V_1$ at $r = a$ and $V = 0$ at $r = b$

$$\therefore V_1 = \frac{Q}{4\pi\epsilon_0 c} \ln\left[\frac{b(a+c)}{a(b+c)}\right]$$

\therefore The capacitance C is given by the relation

$$\frac{V_1}{Q} = \frac{1}{C} = \frac{1}{4\pi\epsilon_0 c} \ln\left[\frac{b(a+c)}{a(b+c)}\right]$$

$$\text{or } C = \frac{4\pi\epsilon_0 c}{\ln\left[\frac{b(a+c)}{a(b+c)}\right]}$$

T4-4

④ Let 'a' be the radius of the insulated spherical conductor.

$$V(\vec{r}) = \frac{a}{4\pi\epsilon_0 r} \text{ and the electrostatic}$$

$$\text{energy is } W_1 = \frac{a^2}{8\pi\epsilon_0} \int_a^\infty \frac{dr}{r^2}$$

Now, let us place the dielectric shell around spherical conductor

$$E_r = \begin{cases} \frac{a}{4\pi\epsilon_0 r^2} & , a \leq r \leq b \\ \frac{a}{4\pi\epsilon_0 k(r) r^2} & b \leq r \leq c \\ \frac{a}{4\pi\epsilon_0 r^2} & r \geq c \end{cases}$$

$$\text{and } D_r = \frac{a}{4\pi r^2}$$

Now the electrostatic energy is given by

$$W_2 = \frac{1}{2} \int \vec{D} \cdot \vec{E} d\tau$$

$$= \frac{a^2}{32\pi^2} \int_a^b \frac{4\pi r^2 dr}{\epsilon_0 r^4} + \frac{a^2}{32\pi^2} \int_b^c \frac{4\pi r^2 dr}{\epsilon_0 k(r) r^4} + \frac{a^2}{32\pi^2} \int_c^\infty \frac{4\pi r^2 dr}{\epsilon_0 r^4}$$

$$= \frac{a^2}{8\pi\epsilon_0} \left[\int_a^b \frac{dr}{r^2} + \int_b^c \frac{dr}{k(r) r^2} + \int_c^\infty \frac{dr}{r^2} \right]$$

$$\text{Then } W_1 - W_2 = \frac{a^2}{8\pi\epsilon_0} \int_b^c \left[1 - \frac{1}{k(r)} \right] \frac{dr}{r^2}$$

Thus there is a loss of energy. when dielectric intervenes an electrostatic system, the electrostatic energy of the system decreases!!

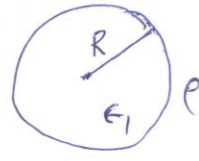
(5)

$$V(0) - V(R) = - \int_R^0 \vec{E}(\vec{r}) \cdot d\vec{r}$$

$$V(0) - \frac{Q}{4\pi\epsilon_0 R} = \frac{1}{4\pi\epsilon_1} \int_0^R \frac{Qr}{R^3} dr$$

$$= \frac{Q}{4\pi\epsilon_1 R}$$

$$\therefore V(0) = \frac{Q}{4\pi R} \left(\frac{1}{\epsilon_0} + \frac{1}{2\epsilon_1} \right)$$



~~Electric field is zero~~

(6)

The electrostatic potential

$$\Phi(r, \theta, \phi) = \begin{cases} (-E_0 r + b_1 r^{-2}) \cos \theta, & r > R \\ b_2 r \cos \theta, & r < R \end{cases}$$

The corresponding electric field is $\vec{E} = -\nabla \Phi$

$$\vec{E} = \begin{cases} E_0 \hat{e}_z + \frac{b_1}{r^3} (2 \cos \theta \hat{e}_r + \sin \theta \hat{e}_\theta) & r > R \\ -b_2 \hat{e}_z & r < R \end{cases} \rightarrow \text{①}$$

The displacement \vec{D} can be computed using $\vec{D} = \epsilon_0 k \vec{E}$

for $r < R$ and $\vec{D} = \epsilon_0 \vec{E}$ for $r > R$.

$$\therefore \frac{\vec{D}}{\epsilon_0} = \begin{cases} E_0 \hat{e}_z + \frac{b_1}{r^3} (2 \cos \theta \hat{e}_r + \sin \theta \hat{e}_\theta) & \text{for } r > R \\ -b_2 k \hat{e}_z & \text{for } r < R \end{cases}$$

(a) To calculate b_1, b_2 constants we use the condition at $r=R$ and arbitrariness (θ, ϕ) . The normal component of \vec{D} is $\vec{D} \cdot \hat{e}_r$ to be continuous as there is no free charge on the interface. The tangential component of \vec{E}

As $\vec{D} \times \hat{e}_r$ should be continuous. (use $\hat{e}_r \times \hat{e}_z = -\sin\theta \hat{e}_\phi$)

$$\left(E_0 + 2 \frac{b_1}{R^3}\right) \cos\theta = -b_2 k \cos\theta$$

$$\left(-E_0 + \frac{b_1}{R^3}\right) \sin\theta \hat{e}_\theta = b_2 \sin\theta \hat{e}_\phi$$

we obtain

$$b_1 = \left(\frac{k-1}{k+2}\right) E_0 R^3, \quad b_2 = -\left(\frac{3}{k+2}\right) E_0$$

(b) Inserting the b_1 & b_2 in eqn ① we get

$$\vec{E} = \begin{cases} E_0 \hat{e}_z + E_0 \left(\frac{k-1}{k+2} \frac{R^3}{r^3}\right) (2\cos\theta \hat{e}_r + \sin\theta \hat{e}_\theta) & \text{for } r > R \\ \frac{3}{k+2} E_0 \hat{e}_z & \text{for } r < R \end{cases}$$

(c) The polarization is given by $\vec{P} = \vec{D} - \epsilon_0 \vec{E}$ and for $r < R$

$$\vec{P} = \frac{3(k-1)}{(k+2)} \epsilon_0 E_0 \hat{e}_z$$

\therefore Dipole moment of the sphere is

$$\vec{p} = \frac{4}{3} \pi R^3 \vec{P}$$

(d) The volume and bound charge densities are

$$\rho_b(\theta, \phi) = -\nabla \cdot \vec{P} = 0, \quad \sigma_b(\theta, \phi) = \vec{P} \cdot \hat{n} = \vec{P} \cdot \hat{e}_r$$

$$\boxed{\sigma_b(\theta, \phi) = \frac{3(k-1)}{k+2} \epsilon_0 E_0 \cos\theta}$$