

Now the amplitude for a process is given by

$$A = \langle f | H_s | i \rangle$$

so for process (a) $A_a = \langle \frac{3}{2}, +\frac{3}{2} | H_s | \frac{3}{2}, +\frac{3}{2} \rangle$

for process (b) things are a bit more complicated

$$A_b = \left(\sqrt{\frac{1}{3}} \langle \frac{3}{2}, \frac{1}{2} | - \sqrt{\frac{2}{3}} \langle \frac{1}{2}, -\frac{1}{2} | \right) H_s \left(\sqrt{\frac{1}{3}} | \frac{3}{2}, \frac{1}{2} \rangle - \sqrt{\frac{2}{3}} | \frac{1}{2}, \frac{1}{2} \rangle \right)$$

$$= \frac{1}{3} \langle \frac{3}{2}, -\frac{1}{2} | H_s | \frac{3}{2}, -\frac{1}{2} \rangle - \frac{\sqrt{2}}{3} \langle \frac{3}{2}, -\frac{1}{2} | H_s | \frac{1}{2}, \frac{1}{2} \rangle \\ - \frac{\sqrt{2}}{3} \langle \frac{1}{2}, -\frac{1}{2} | H_s | \frac{3}{2}, -\frac{1}{2} \rangle + \frac{2}{3} \langle \frac{1}{2}, -\frac{1}{2} | H_s | \frac{1}{2}, -\frac{1}{2} \rangle$$

But strong interaction conserves strong isospin
thus only terms with the same initial and final
isospins survive!

$$\therefore A_b = \frac{1}{3} \langle \frac{3}{2}, -\frac{1}{2} | H_s | \frac{3}{2}, -\frac{1}{2} \rangle + \frac{2}{3} \langle \frac{1}{2}, -\frac{1}{2} | H_s | \frac{1}{2}, -\frac{1}{2} \rangle$$

Since the energy states of the strong Hamiltonian are
only dependent on I_s and not I_{s3} we can
write A_b as:

$$A_b = \frac{1}{3} A_a + \frac{2}{3} \langle \frac{1}{2}, -\frac{1}{2} | H_s | \frac{1}{2}, -\frac{1}{2} \rangle$$

similarly for process (c)

$$A_c = \left(\sqrt{\frac{1}{3}} \langle \frac{3}{2}, -\frac{1}{2} | - \sqrt{\frac{2}{3}} \langle \frac{1}{2}, -\frac{1}{2} | \right) H_s \left(\sqrt{\frac{2}{3}} | \frac{3}{2}, -\frac{1}{2} \rangle + \sqrt{\frac{1}{3}} | \frac{1}{2}, -\frac{1}{2} \rangle \right) \\ = \frac{\sqrt{2}}{3} \langle \frac{3}{2}, -\frac{1}{2} | H_s | \frac{3}{2}, -\frac{1}{2} \rangle - \frac{\sqrt{2}}{3} \langle \frac{1}{2}, -\frac{1}{2} | H_s | \frac{1}{2}, -\frac{1}{2} \rangle$$

where again since strong interactions conserve isospin only
the same initial and final isospins survive