## Statistical Physics (3rd tierce exam)

## Name:

## ID:

	11).				
	1. (a)	(b)	(c)	(d)	
U	2. (a)	(b)	(c)	(d)	
	3. (a)	(b)	(c)	(d)	
	4. (a)	(b)	(c)	(d)	
	5. (a)	(b)	(c)	(d)	
	6. (a)	(b)	(c)	(d)	
	7. (a)	(b)	(c)	(d)	
	8. (a)	(b)	(c)	(d)	
	9. (a)	(b)	(c)	(d)	
	10. (a)	(b)	(c)	(d)	
	11. (a)	(b)	(c)	(d)	
	12. (a)	(b)	(c)	(d)	
	13. (a)	(b)	(c)	(d)	
	14. (a)	(b)	(c)	(d)	
	15. (a)	(b)	(c)	(d)	

- 1. Ideal gas equation PV = NKT is true when the system is in classical domain (high temperature). To understand this macroscopic phenomena, we have to consider microscopic statistical description: gas constituent particles with
  - (a) Maxwell-Boltzmann (MB) distribution function
  - (b) Bose-Einstein (BE) distribution function
  - (c) Fermi-Dirac (FD) distribution function
  - (d) none of the above
- 2. According to ideal gas equations PV = NKT,  $U = \frac{3}{2}NKT$ , at zero temperature (T = 0), pressure P and internal energy U will be zero but it is not true because low temperature is quantum domain, where gas constituent particles follow
  - (a) classical distribution MB
  - (b) quantum distribution FD/BE
  - (c) equipartition of energy law
  - (d) none of the above
- 3. Considering compact star like white dwarf as electron gas only (ignoring its other components), then for statistical mechanical description of white dwarf, which distribution is required?
  - (a) MB
  - (b) **B**E
  - (c) FD
  - (d) none of the above
- 4. Fermion with energy  $\epsilon$  in a gas with temperature T and chemical potential/Fermi energy  $\mu$  will follow FD distribution, whose mathematical form will be  $f_{FD} =$ 
  - (a)  $1/\{\exp(\frac{\epsilon-\mu}{KT})+1\}$
  - (b)  $1/\{\exp(\frac{\epsilon-\mu}{KT}) 1\}$ (c)  $\exp(-\frac{\epsilon-\mu}{KT})$

  - (d) none of the above
- 5. At Fermi energy, i.e.  $\epsilon = \mu$ 
  - (a)  $f_{FD}(\epsilon = \mu) = 0$
  - (b)  $f_{FD}(\epsilon = \mu) = 1/2$ (c)  $f_{FD}(\epsilon = \mu) = 1$

  - (d) none of the above
- 6. At  $\epsilon \ll \mu$  or  $\epsilon \to 0$ 
  - (a)  $f_{FD} = 0$
  - (b)  $f_{FD} = 1/2$
  - (c)  $f_{FD} = 1$
  - (d) none of the above
- 7. At  $\epsilon \gg \mu$  or  $\epsilon \to \infty$ 
  - (a)  $f_{FD} = 0$
  - (b)  $f_{FD} = 1/2$
  - (c)  $f_{FD} = 1$
  - (d) none of the above
- 8. At T = 0
  - (a)  $f_{FD} = 0$  for  $\epsilon < \mu$  and  $f_{FD} = 1$  for  $\epsilon > \mu$

(b) 
$$f_{FD}=1$$
 for  $\epsilon<\mu$  and  $f_{FD}=0$  for  $\epsilon>\mu$  (c)  $f_{FD}=1$  (d) none of the above

- 9. No of electrons N in metal can be expressed as

$$N = 2 \int_0^\infty \frac{d^3 x d^3 p}{h^3} \frac{1}{e^{\beta(\epsilon - \mu)} + 1}$$
 (1)

with electron's energy  $\epsilon = \frac{p^2}{2m}$ . Here  $\beta = 1/(KT)$ ,  $\mu = \frac{p_F^2}{2m}$ . At T=0 we get number density as (a)  $\frac{N}{V} = \frac{8\pi}{5h^3}(2m\mu)^{5/2}$  (b)  $\frac{N}{V} = \frac{8\pi}{4h^3}(2m\mu)^{4/2}$  (c)  $\frac{N}{V} = \frac{8\pi}{3h^3}(2m\mu)^{3/2}$  (d) none of the above

- 10. Total energy (internal energy) of electron gas can be expressed as

$$U = 2 \int_0^\infty \frac{d^3 x d^3 p}{h^3} \frac{6}{e^{\beta(\epsilon - \mu)} + 1} . \tag{2}$$

In the same condition, described in earlier question, energy density of electron gas will be

- (a)  $\frac{U}{V} = \frac{8\pi}{5h^3} \frac{(2m\mu)^{5/2}}{2m}$
- (a)  $\frac{\bar{v}}{V} = \frac{8\pi}{5h^3} \frac{(2m\mu)^4}{2m}$ (b)  $\frac{U}{V} = \frac{8\pi}{4h^3} \frac{(2m\mu)^{4/2}}{2m}$ (c)  $\frac{U}{V} = \frac{8\pi}{3h^3} \frac{(2m\mu)^{3/2}}{2m}$ (d) none of the above

- 11. Average energy of non-relativistic fermion (or boson or any particle) in classical domain is  $\epsilon_{av} = \frac{U}{N} = \frac{3}{2}KT$ , according to which  $\epsilon_{av} = 0$  at T = 0, but according quantum relations (discussed in earlier 2 questions), at T=0, average energy of non-relativistic fermion is

  - (a)  $\epsilon_{av} = \frac{3}{3}\mu$ (b)  $\epsilon_{av} = \frac{3}{5}\mu$ (c)  $\epsilon_{av} = \mu$

  - (d) none of the above
- 12. Grand canonical potential  $\Phi$  or Pressure times volume of electron gas can be expressed as

$$-\Phi = PV = 2 \int_0^\infty \frac{d^3x d^3p}{h^3} \frac{pv/3}{e^{\beta(\epsilon - \mu)} + 1} , \qquad (3)$$

In the same condition, described in earlier questions, pressure of electron gas will be

- (a)  $P = \frac{8\pi}{15h^3} \frac{(2m\mu)^{5/2}}{m}$ (b)  $P = \frac{8\pi}{12h^3} \frac{(2m\mu)^{4/2}}{m}$ (c)  $P = \frac{8\pi}{9h^3} \frac{(2m\mu)^{3/2}}{m}$ (d) none of the above

- 13. For non-relativistic fermion (or boson or any particle) in classical domain, we get relation between internal energy density and pressure as  $\frac{U}{V} = \frac{3}{2}P$  by fusing two equations  $U = \frac{3}{2}NKT$  and PV = NKT. Similar kind of relation between internal energy density and pressure for degenerate (i.e. at T=0) and non-relativistic electron gas will be

- (a)  $\frac{U}{V} = \frac{1}{3}P$ (b)  $\frac{U}{V} = \frac{3}{2}P$ (c)  $\frac{U}{V} = \frac{2}{3}P$ (d) none of the above
- 14. Example of degenerate and relativistic electron gas is
  - (a) White dwarf
  - (b) Neutron star
  - (c) black hole
  - (d) none of the above
- 15. Example of degenerate and relativistic neutron gas is
  - (a) White dwarf
  - (b) Neutron star
    - (c) black hole
    - (d) none of the above