Problem 1.

Orbital Transfer Review: remind yourself: (discuss/justify decisions)

All three of the transfers below can be accomplished with a Hohmann transfer. The amount of ΔV required for a Hohmann transfer is

$$\Delta v_1 = \sqrt{\frac{\mu}{r_1}} \left(\sqrt{\frac{2r_2}{r_1 + r_2}} - 1 \right),$$

$$\Delta v_2 = \sqrt{\frac{\mu}{r_2}} \left(1 - \sqrt{\frac{2r_1}{r_1 + r_2}} \right),$$

$$\Delta v_{total} = \Delta v_1 + \Delta v_2.$$

Additionally, part (b) requires an inclination change. The ΔV budget for a inclination change for a circular orbit is

$$\Delta v_i = 2v \sin\left(\frac{\Delta i}{2}\right).$$

The assumption of a circular orbit should be fine here, as both HST and ISS orbits have very low eccentricity.

(a) Determine the ΔV required to move from a 200km coplanar parking orbit to the HST orbit

$$r_1 = 200km + 6,371km$$

$$r_2 = 569km + 6,371km$$

$$\Delta v_{total} = \Delta v_{Hohmann}$$

$$= 0.210km/s.$$

(b) Determine the ΔV required to move from the ISS orbit to the HST orbit

$$r_{ISS} = 414.1km + 6,371km$$

 $r_{HST} = 569km + 6,371km$
 $i_{ISS} = 0.9014rad$
 $i_{HST} = 0.4969rad$
 $v_{HST} = 7.59km/s$

$$\Delta v_{total} = \Delta v_{Hohmann} + \Delta v_{PlaneChange}$$
$$= 0.086 + 3.049$$
$$= 3.135 km/s,$$

where we've again assumed a circular orbit. The plane change should take place after the Hohmann transfer, as the orbital velocity is lower at HST orbit, leading to a smaller required ΔV .

(c) Determine the ΔV required to deorbit from the HST orbit (must choose your de-orbit orbital params)

$$r_{deorbit} = 100km + 6,371km$$

$$r_{HST} = 569km + 6,371km$$

$$\Delta v = \sqrt{\frac{\mu}{r_1}} \left(\sqrt{\frac{2r_2}{r_1 + r_2}} - 1 \right),$$

$$= 0.136km/s,$$

where a height of 100km should be sufficient to cause the vehicle to deorbit rapidly.

Problem 2.

Eclipse durations (text 5.3.2): an important design aspect of your solar array system is the relative durations of eclipse and insolation. Using the algorithm given in section 5.3.2,

(a) Compute the eclipse period for ISS, for HST, and for a typical GPS satellite (choose one)

I ran the algorithm in section 5.3.2 over one complete orbit at 0.1 degree increments, and for sun positions of one complete year at day increments. The mean and standard error of the mean was calculated for eclipse period, as was the total time in eclipse and the total number of eclipses over the one year period. These results are presented in Table 1.

Satellite	Mean (min)	SEM (min)	Total (hrs)	Count
HST ISS	33.97 34.99	0.10 0.09	414.45 426.98	732 732
GPS	46.09	2.55	90.64	118

Table 1: Eclipse statistics for the HST, ISS, and a GPS satellite (Navstar 43)

Problem 3.

Lets say we lose control of your spacecraft after it has undocked from HST, but before it has de-orbited.

(a) Estimate the orbital lifetime of your spacecraft (text 5.3.4) following loss of communications: assume HST circular orbit, average solar activity.

The lifetime of an uncontrolled space vehicle is approximately

$$\tau \sim \frac{e_0^2}{2B} \left(1 - \frac{11}{6} e_0 + \frac{29}{16} e_0^2 + \frac{7}{8} \frac{H}{a_0} \right)$$

where e_0 and a_0 are the initial values of eccentricity and semi-major axis once control has ceased, H is the scale height of the atmosphere near perigee and B is given by

$$B \sim \sqrt{\frac{\mu}{a_o^3}} \frac{AC_D}{M} \rho_{p_0} a_0 e_0 I_1 \left(\frac{a_0 e_0}{H}\right) exp\left(-e_0 \left(1 + \frac{a_0}{H}\right)\right)$$

Assuming a small $e \approx 0.0002468$, $C_D = 2.2$, $\rho \approx 6.5$ E-13 kg/m³, $a \approx 6920$ E+3 m, $H \approx 72500$ m, and A = 16m², we find a time, $\tau = 718.3$ days, or approximately 2 years.

(b) Would it make any difference to the decay timescale whether your spacecraft was tumbling or not?

Tumbling has no affect on any of the parameters in τ , except for the area, A. A tumbling craft would have a continuously changing A, which could increase or decrease the lifetime of the vehicle.

Problem 4.

Geostationary orbits (text 5.6):

(a) Using the linearized solution to Keplers equation given in eqns 5.27-5.29, plot ground-track fluctuations as longitude vs latitude. Describe the results.

The longitude and latitude can be calculated by

$$\lambda = \Omega + \omega - \frac{3}{2} \frac{\delta a}{A} t \sqrt{\frac{\mu}{A^3}} + 2e \sin\left(t \sqrt{\frac{mu}{A^3}}\right)$$
$$\theta = i \sin\left(\omega + t \sqrt{\frac{mu}{A^3}}\right).$$

Figure 1 shows the Longitude and Latitude for Westar 1, America's first domestic and commercially launched geostationary communications satellite. Here $i=14.15^o, \omega=282^o, \Omega=336.8^o, e=7.333\text{E-4}, A=42164.5\text{km}$, and $\delta a=42164.5-42269=-104.5\text{km}$. The groundtrack is shown here for 24 hours, which is roughly equivalent to one orbit. The latitudinal variation is a simple oscillation, while the longitudinal variations include an oscillation and a small drift rate from the small δa value.

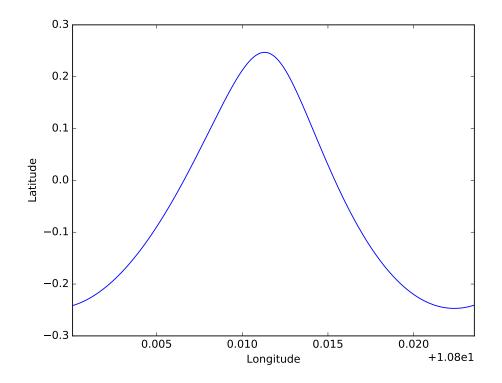


Figure 1: Ground Track Fluctuations as Longitude vs Latitude

(b) Define deadband and control limit-cycle in the context of GEO station keeping. The dead band of a GEO satellite is the maximum allowable error in λ , and the control limit-cycle describes how often a burn is required in order to return this error to the

opposite side of the dead band.

(c) Consider a GEO satellite with nominal longitude of -100deg, and an onboard propellant system capable of providing a total ΔV of 200m/s. For a maximum longitudinal error magnitude of 0.22deg, for how long can the satellite station-keep?

The required ΔV required to keep the satellite within λ_{max} degrees of error with a limit cycle of T is

$$\Delta V = 4\sqrt{\frac{rf\lambda_{max}}{3}} = 0.083 \text{ m/s}$$

$$T = 4\sqrt{\frac{r\lambda_{max}}{3f}} = 0.326 \text{ years},$$

where f = +8.10 E-9 m/s², $\lambda_{max} = 0.22$ deg, and r = 42E6 m. The length of time a satellite can station keep with fuel ΔV_{total} is then

$$T_{stationkeep} = \Delta V_{total} / \left(\frac{\Delta V}{T}\right) = 782.957 \text{ years},$$

which simply means that the lognitudinal station keeping is not the limiting factor of the satellite's lifespan.

Problem 5.

John Karasinski

Two spacecraft in elliptical Earth orbit with the orbital parameters as follows. Compute the relative position and velocity vectors.

(a) h= 52,059 km²/s, e=0.0257240, i=60deg,
$$\Omega$$
=40deg, ω =30deg, θ =40deg

(b) h= 52,362 km²/s, e=0.0072696, i=50deg,
$$\Omega$$
=40deg, ω =120deg, θ =40deg

We can calculate the Earth-centric position and velocity vectors with the sv_from_coe() function developed in Homework 4. They are

(a)
$$[\mathbf{r}_1 \ \mathbf{v}_1] = [-266.7, 3865.7, 5426.1, -6.4, -3.6, 2.4]$$

(b)
$$[\mathbf{r}_2 \ \mathbf{v}_2] = [-5890.7, -2979.7, 1792.2, 0.9, -5.2, -5.5],$$

Where \mathbf{r} is in km, and \mathbf{v} is in km/s. To calculate \mathbf{r} we can simply take the difference

$$\mathbf{r}_{rel} = \mathbf{r}_2 - \mathbf{r}_1$$

= $[-5623.9, -6845.5, -3633.9] \text{ km}$

The relative velocity vector can be caluclated by

$$\mathbf{v}_2 = \mathbf{v}_1 + \Omega \times \mathbf{r}_{rel} + \mathbf{v}_{rel}$$

where

$$\Omega = \frac{\mathbf{r} \times \mathbf{v}}{\mathbf{r}^2}
= \begin{vmatrix} i & j & k \\ -266.7 & 3865.7 & 5426.1 \\ -6.4 & -3.6 & 2.4 \end{vmatrix} \cdot \frac{1}{(6667.7)^2}
= \frac{28979.7i - 34536.6j + 26029.5k}{(6667.7)^2}
= [0.00065, -0.00077, 0.00058] \text{ rad/s}$$

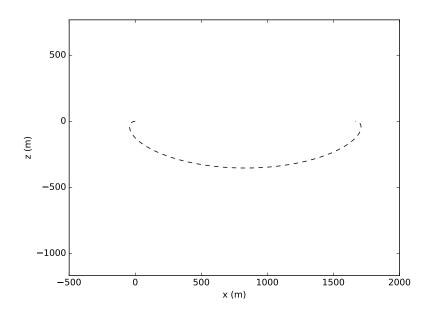
Rearranging, we have,

$$\begin{aligned} \mathbf{v}_{rel} &= \mathbf{v}_2 - \mathbf{v}_1 - \Omega \times \mathbf{r}_{rel} \\ &= [0.9, -5.2, -5.5] - [-6.4, -3.6, 2.4] - \begin{vmatrix} i & j & k \\ 0.00065 & -0.00077 & 0.00058 \\ -5623.9 & -6845.5 & -3633.9 \end{vmatrix} \\ \mathbf{v}_{rel} &= [0.58855, -0.69663, 0.91435] \text{ km/s}. \end{aligned}$$

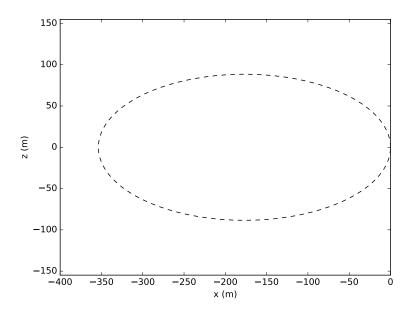
Problem 6.

Fly-around relative trajectories: for the lost EVA toolbox example considered in lecture, generate the relative motion plot for 1 orbital period, given initial conditions of:

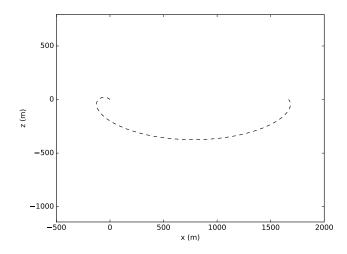
(a) Release relative velocity = (-0.1, 0, 0) m/s (prolate cycloid)



(b) Release relative velocity = (0, 0, 0.1) m/s (ellipse)

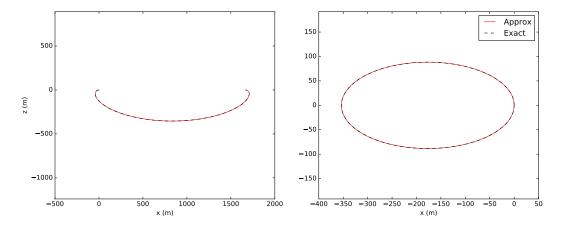


(c) Release relative velocity = (-0.1, 0, 0.1) m/s (initially 45deg backwards and up; describe subsequent motion)



This motion looks very similar to the prolate cycloid, except that altitude is gained during the very beginning of the orbit.

(d) For a and b, plot the trajectory with and without the nt≪1 assumption. Discuss.



The cos and sin parts of the solution were expanded to their first two Taylor series components. Compared to the exact solutions, these approximate solutions performed very well. There is a minimal difference between the two solutions for one orbit.

(e) How about a release relative velocity = (0, 0.1, 0) m/s? Would you see the toolbox again or not?

Motion along the y axis simply acts as a harmonic oscillator. If the orbit was sufficiently high enough, so that drag and other orbit perturbations were minimal, you may indeed see the toolbox again. For ISS orbit, however, drag would dominate, and you would not see the toolbox again.

Problem 7.

For your HST re-boost spacecraft, assume:

- Launch: drop-off circular orbit at 200km, in-plane with HST, 65deg phase angle behind HST
- Phasing: 4-orbit phasing to point S1, 30km behind and 10km below HST
- Homing: Hohmann S1 to co-orbit waiting point S2, 1km behind HST
- Closing: Cycloid close waiting point S3, 200m behind HST
- (a) compute the required ΔV and elapsed time for each phase, and for the total rendezvous to S3

To get to approximately 30km and 10km below HST, we need to phase our orbits. The amount of change in phase after each orbit is

$$\Delta\theta = -3\pi \frac{\Delta a}{a}$$

$$= -3\pi \frac{550 \text{ km} - 200 \text{ km}}{550 \text{ km} + 6371 \text{ km}}$$

$$= -27.3 \text{ degree},$$

so we'll begin our homing phase after two orbits. For the homing burn, we will burn from $\Delta a = -10 \mathrm{km}$ to a waiting point $S1 = 1 \mathrm{km}$ behind the HST. We have $\theta_f = (180^o/\pi) \cdot (10/6921) = 0.0827^o$ and $v_T = 7.591 \mathrm{km s^{-1}}$. With this we get $\theta_i = \theta_f + .195^o = .278^o$. $2\pi(\frac{0.278^o}{360^o}) \cdot 6921 = 33.58 \mathrm{km}$, which is approximately 30km. This means we have to perform a burn with

$$\Delta v = -\frac{1}{2} \frac{\Delta a}{a} \sqrt{\frac{\mu}{a}}$$
$$= 5.482 \text{ m/s},$$

which will take a time, t_H , of

$$t_{H} = \pi \sqrt{\frac{a_{H}^{3}}{\mu}}$$

$$= \pi \sqrt{\frac{\left(a\left[1 + \frac{\Delta a}{2a}\right]\right)^{3}}{\mu}}$$

$$= 2862.0 \text{ s.}$$

We use a cycloidal approach to close, with a total Δv requirement of

$$\Delta v = 2 \times \frac{\Delta x}{6\pi \cdot k} \sqrt{\frac{\mu}{a^3}}$$
$$= 0.093 \text{ m/s},$$

where $\Delta x = 800$ m. Here k = 1 has been chosen to minimize the required time to approach. This will take one full orbital period, for a total time of 5728.8 seconds.

This brings the total Δv requirements to 5.575 m/s, and the total time to 20048.4 seconds, or 3.5 orbits.

Phase	ΔV	Time (s)
Phasing	_	11457.6
Homing	5.482	2862.0
Closing	0.093	5728.8
Total	5.575	20048.4

Table 2: ΔV and time requirements for each phase of rendezvous

(b) compute the view-angle to HST, measured from the orbit-tangent (for sensor acquisition)

The view-angle is simply calculated by

$$\theta = \arctan\left(\frac{10 \text{ km}}{30 \text{ km}}\right)$$
$$= 18.43^{\circ}$$

(c) plot the total quantitative relative motion (like Walter Fig. 8.26)