

Physics 224

The Interstellar Medium

Lecture #19: B-fields, Star Formation, Feedback, Cosmic Rays

Outline

- Part I: Magnetic Fields
- Part II: Star Formation
- Part III: Feedback
- Part IV: Cosmic Rays

Magnetic Fields in the ISM

Observational Tracers:

- Synchrotron emission - from charged particles interacting with the magnetic field.
- Faraday Rotation - different phase velocities of right & left circularly polarized light in the presence of B-field leads to rotation of polarization angle
- Polarization - of starlight due to dust grains aligned along B-field or of dust emission from aligned grains
- Zeeman splitting - splitting of fine structure levels in atoms/ molecules due to interaction of electron magnetic moment and B-field

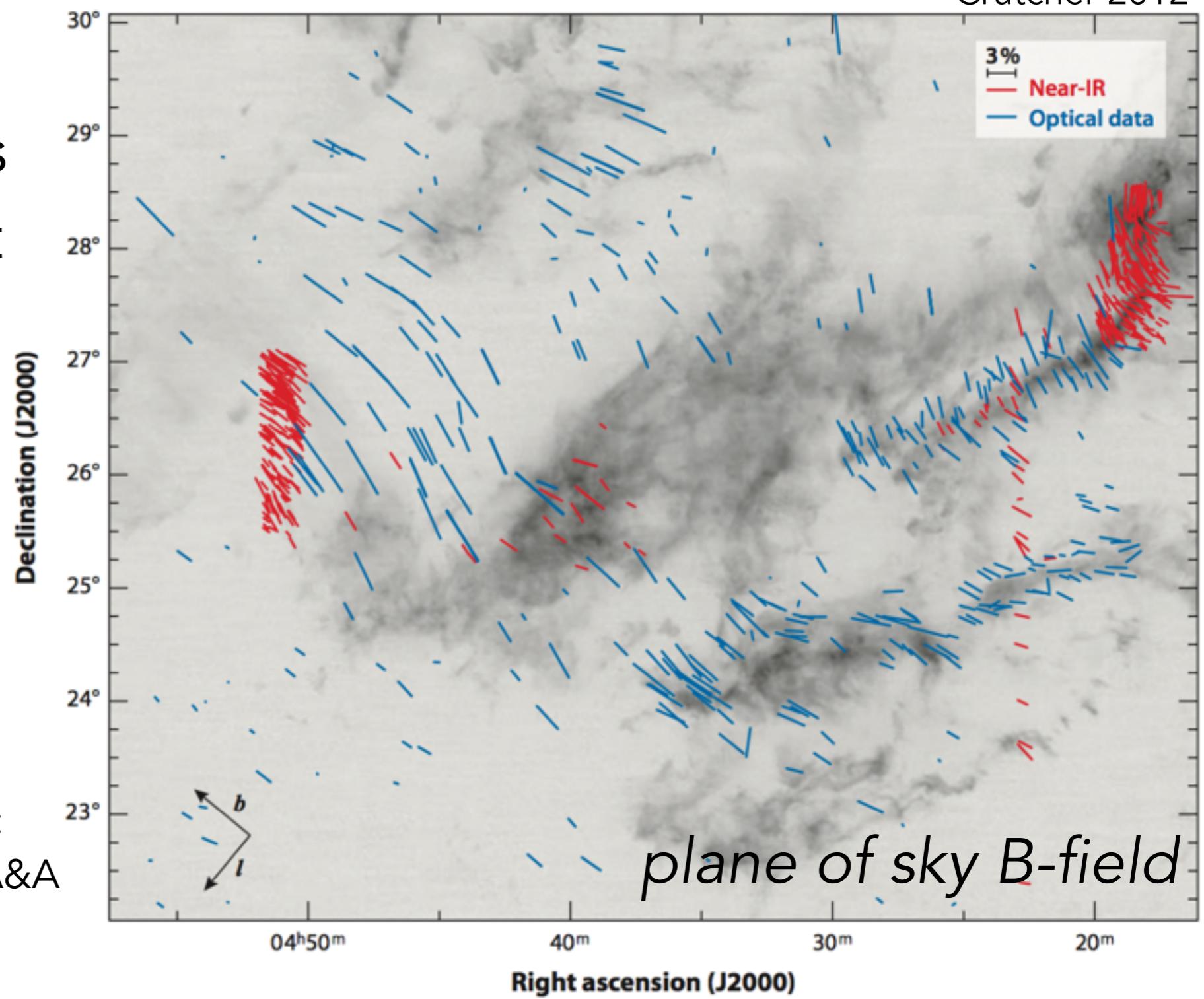
Magnetic Fields in the ISM

Crutcher 2012

Polarization of
starlight in Taurus
due to alignment
of dust grains
with the B-field.

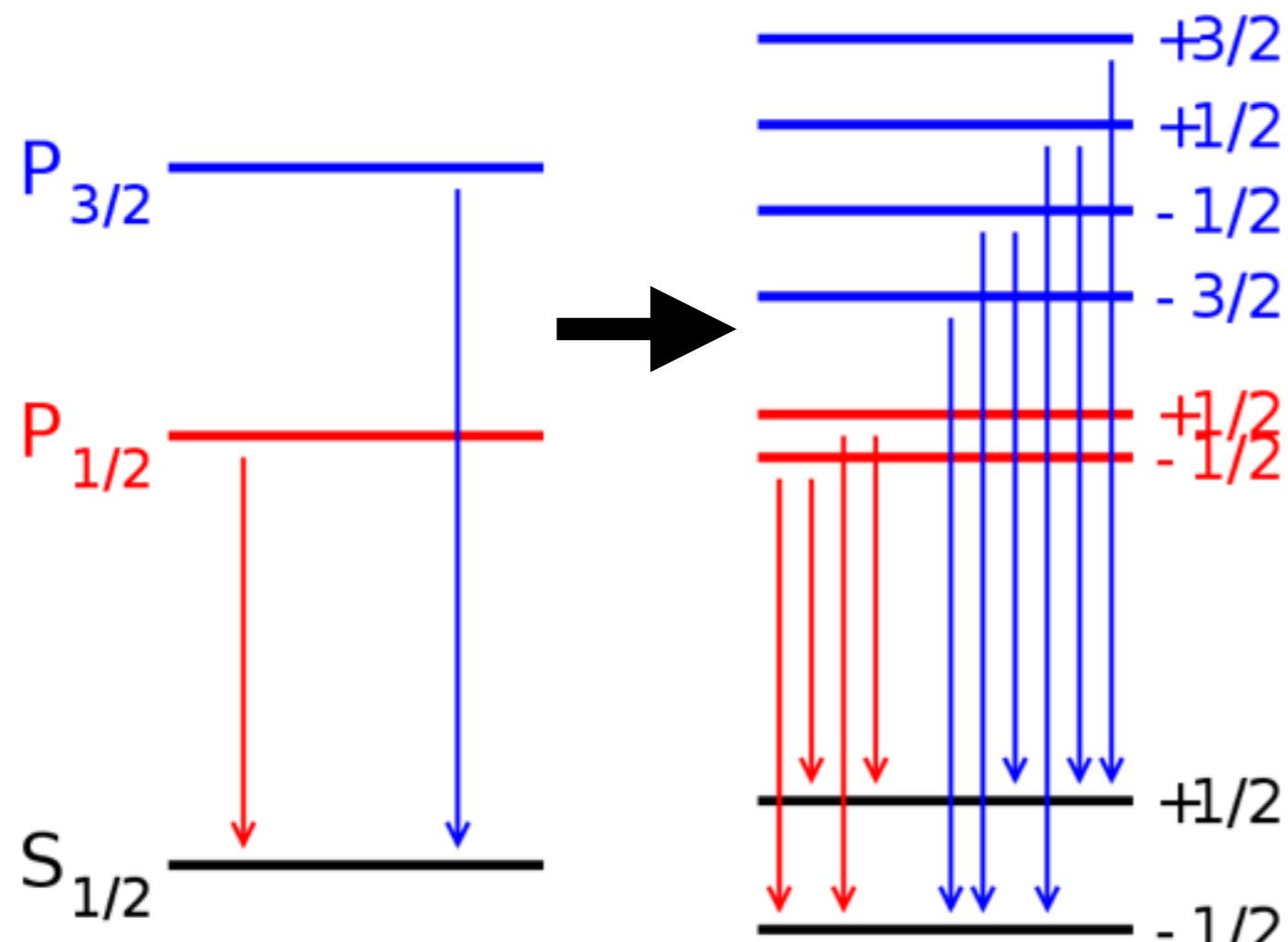
Magnetic fields review:
Crutcher 2012 ARA&A

Grain alignment review:
Andersson et al. 2015 ARA&A



Magnetic Fields in the ISM

Zeeman Effect



Zeeman splitting is largest when there is an unpaired electron in outer shell:
e.g. HI, OH, CN, CH,
CCS, SO, and O₂

Even then, energy shift is small.

But, shifted levels produce different circular polarizations.

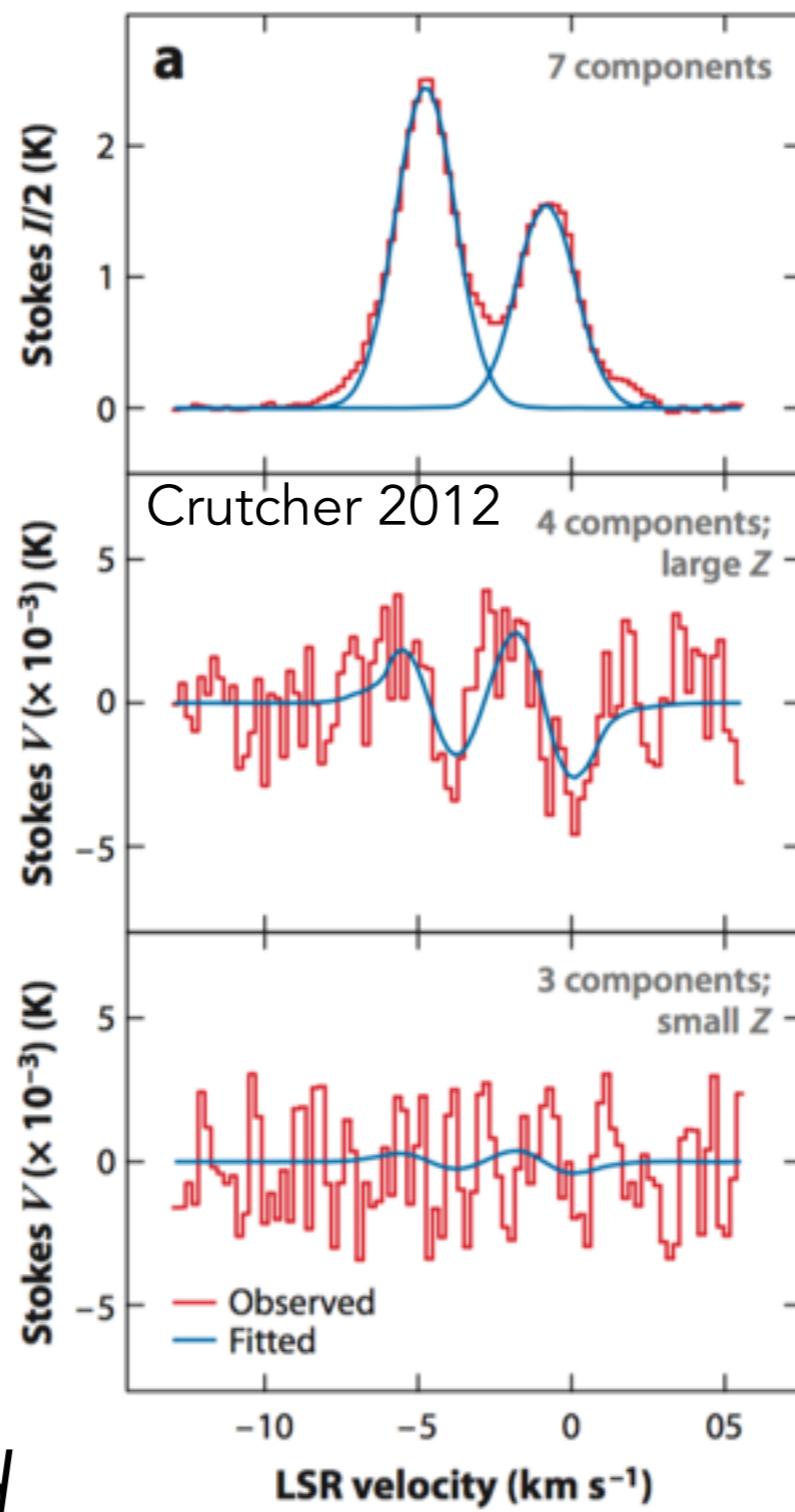
Magnetic Fields in the ISM

Total intensity
7 hyperfine components for
mm rotational lines of CN
two velocity components
along line of sight.

Circularly polarized emission:
4 components with large
Zeeman splitting

Circularly polarized emission:
3 components with small
Zeeman splitting

line-of-sight B-field



Zeeman splitting is largest when there is an unpaired electron in outer shell:

e.g. HI, OH, **CN**, CH,
CCS, SO, and O₂

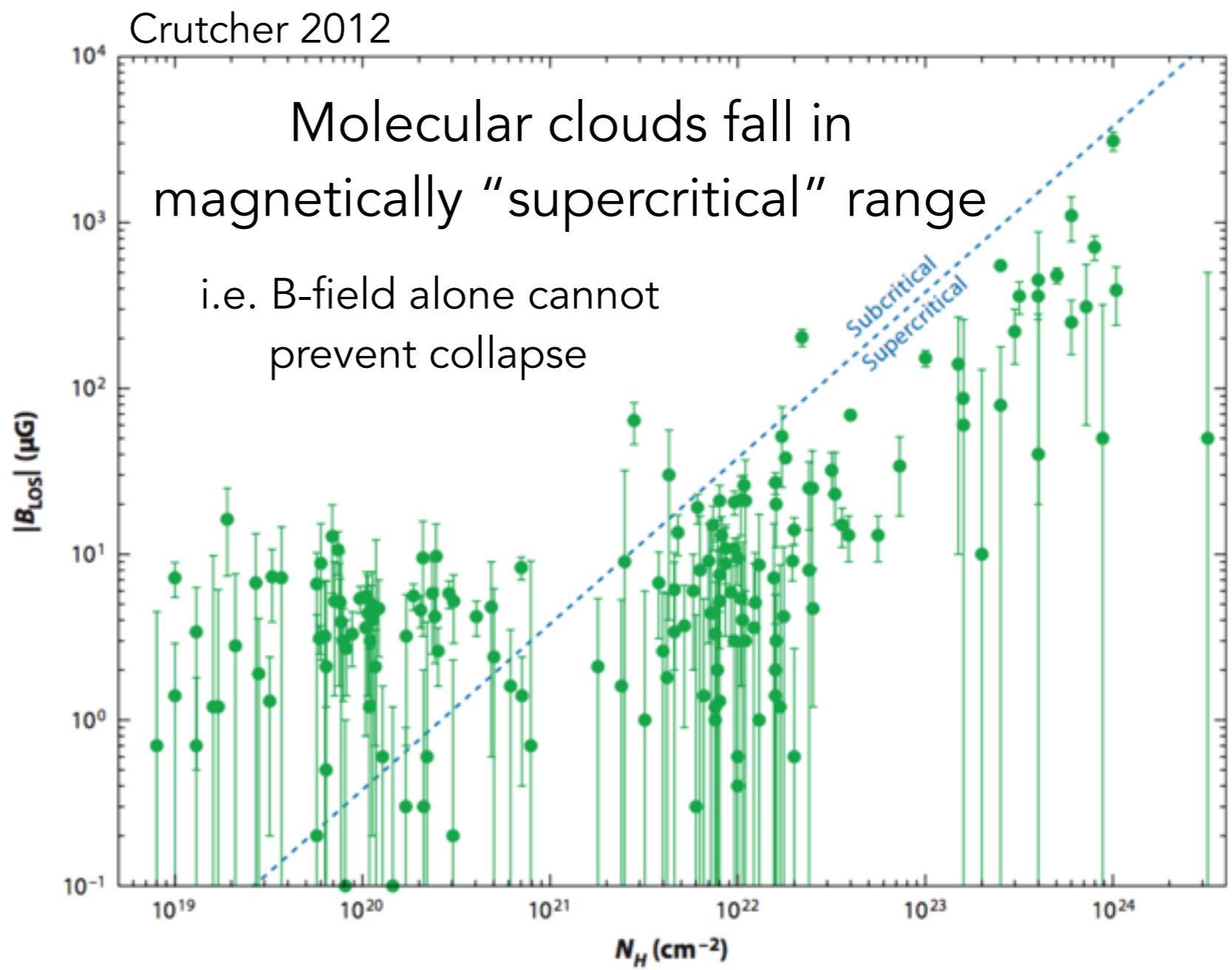
Even then, energy shift is small.

But, shifted levels produce different circular polarizations.

Molecular Clouds

Observed Characteristics

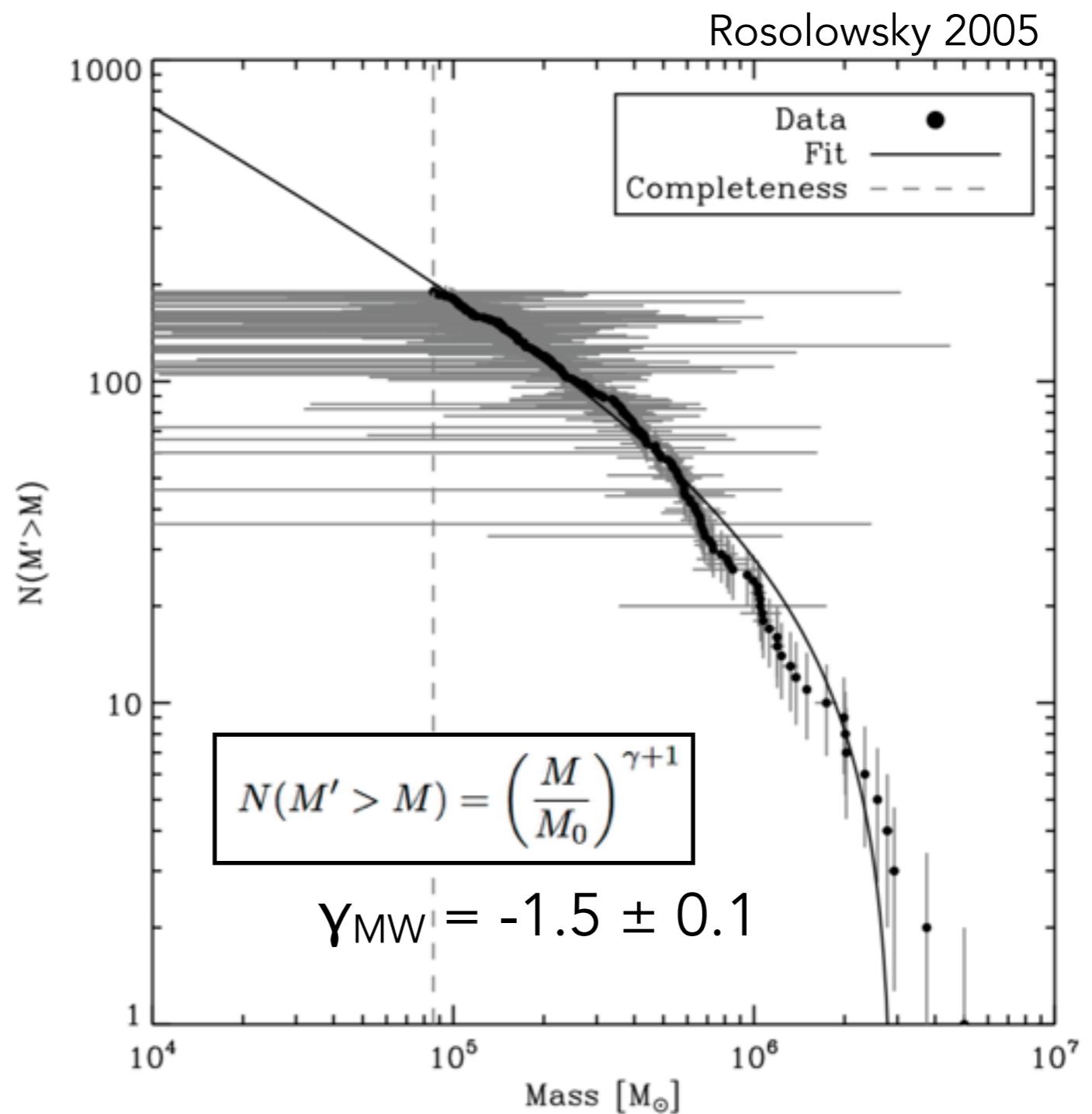
- Self-Gravity
- Turbulence
- Substructure
- **Magnetic Fields**
- Mass Spectrum
- Lifetimes
- Star Formation



Molecular Clouds

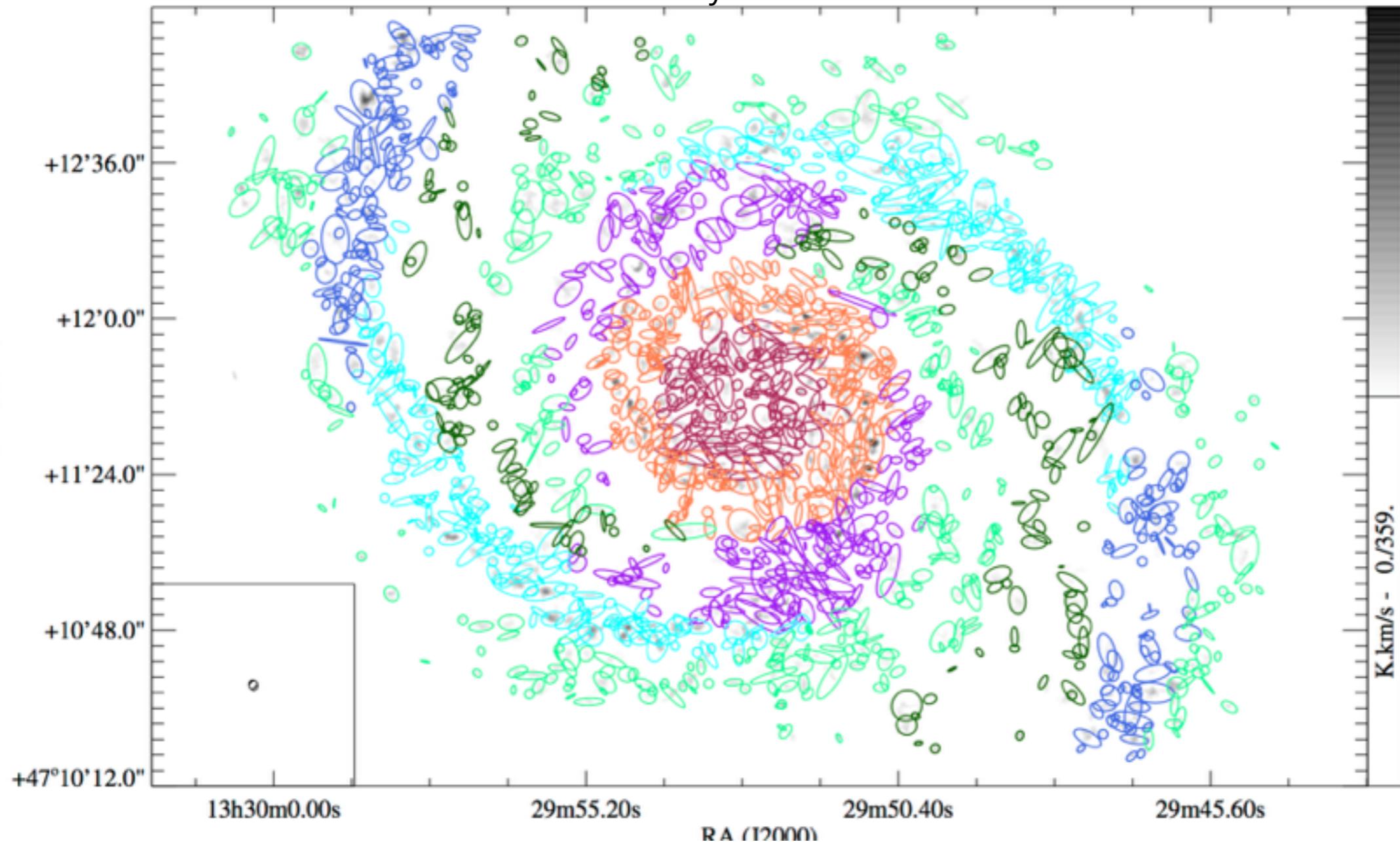
Observed Characteristics

- Self-Gravity
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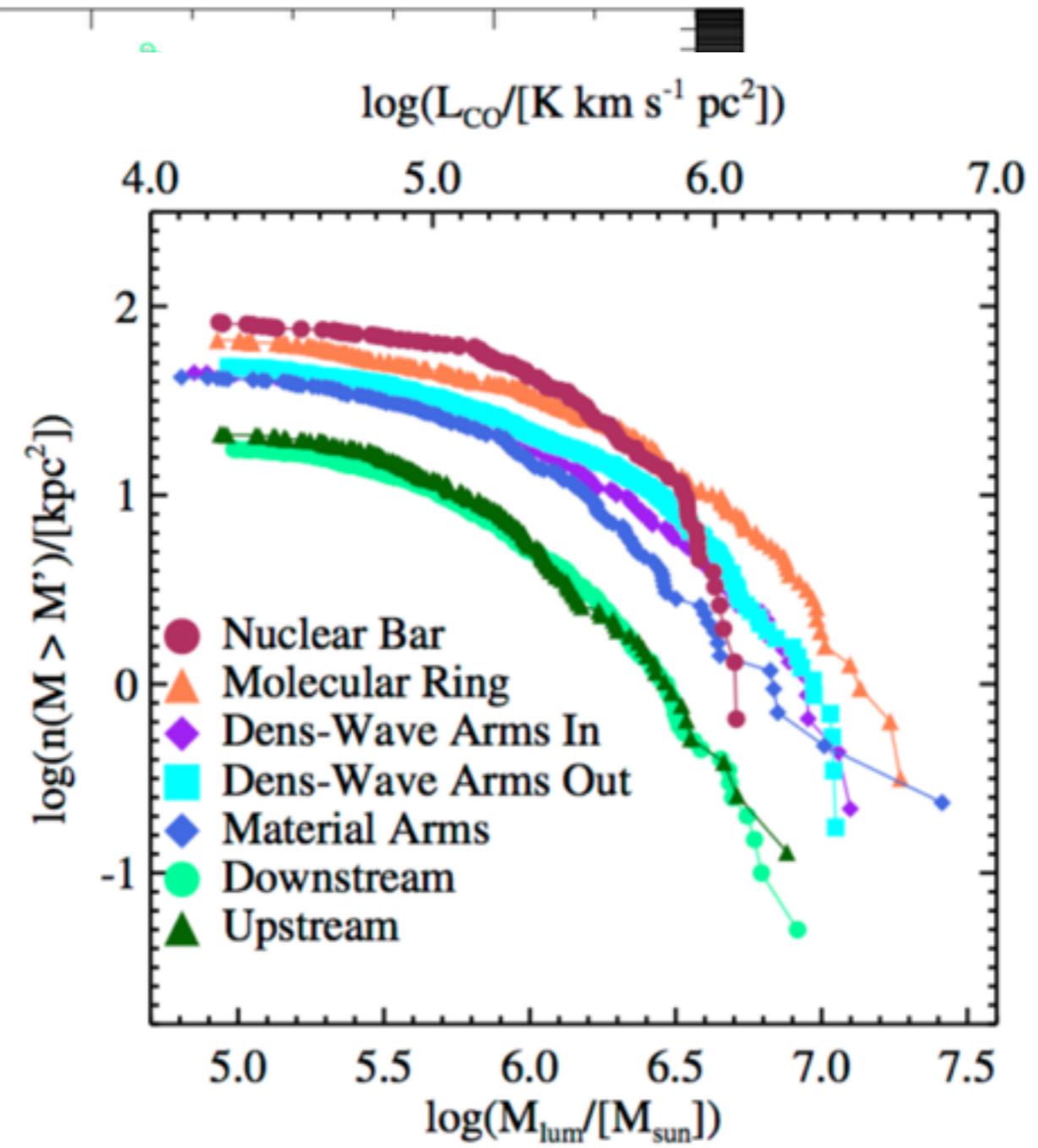
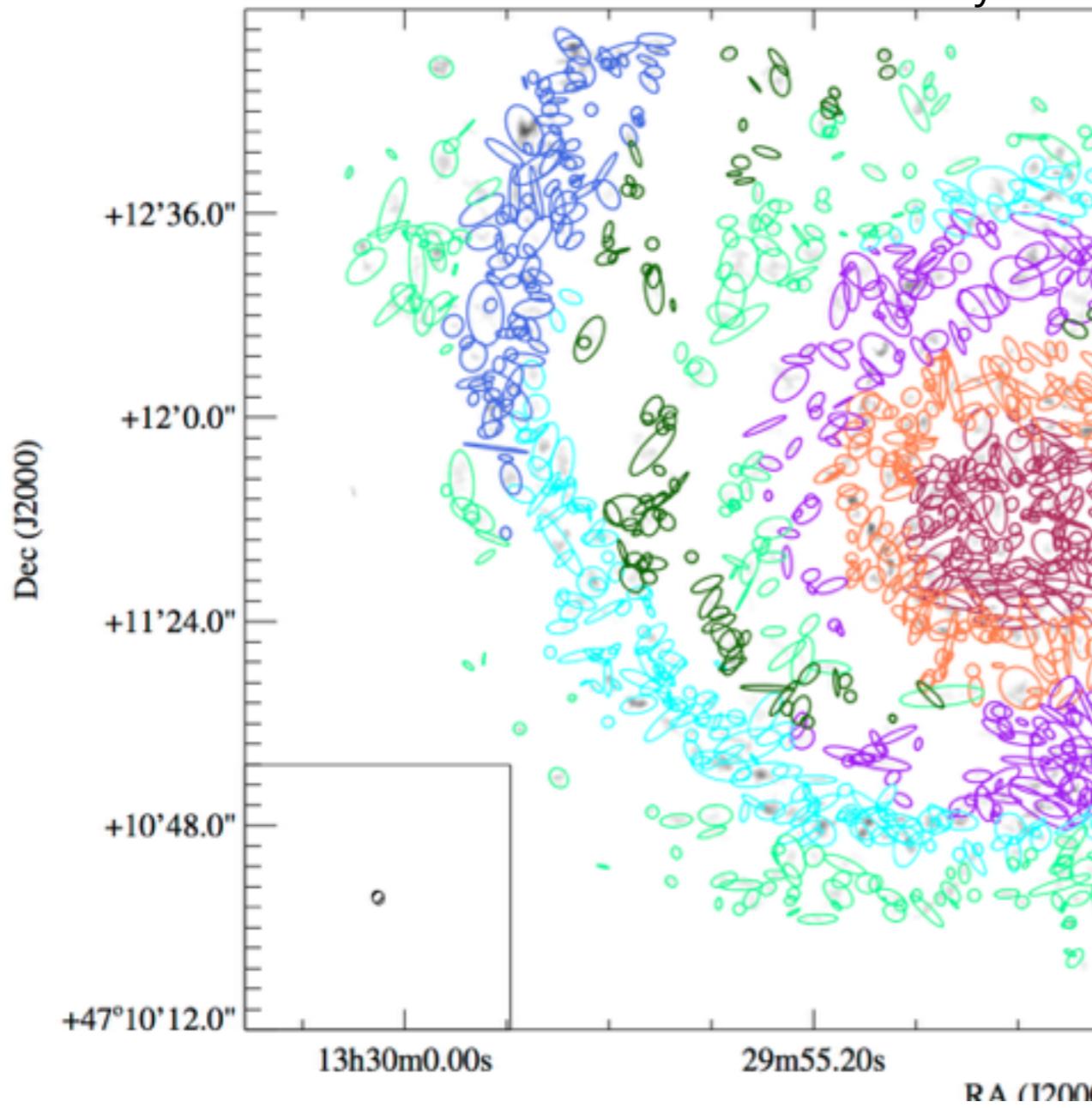
Molecular Clouds

Colombo et al. 2014 - PAWS survey of M51



Molecular Clouds

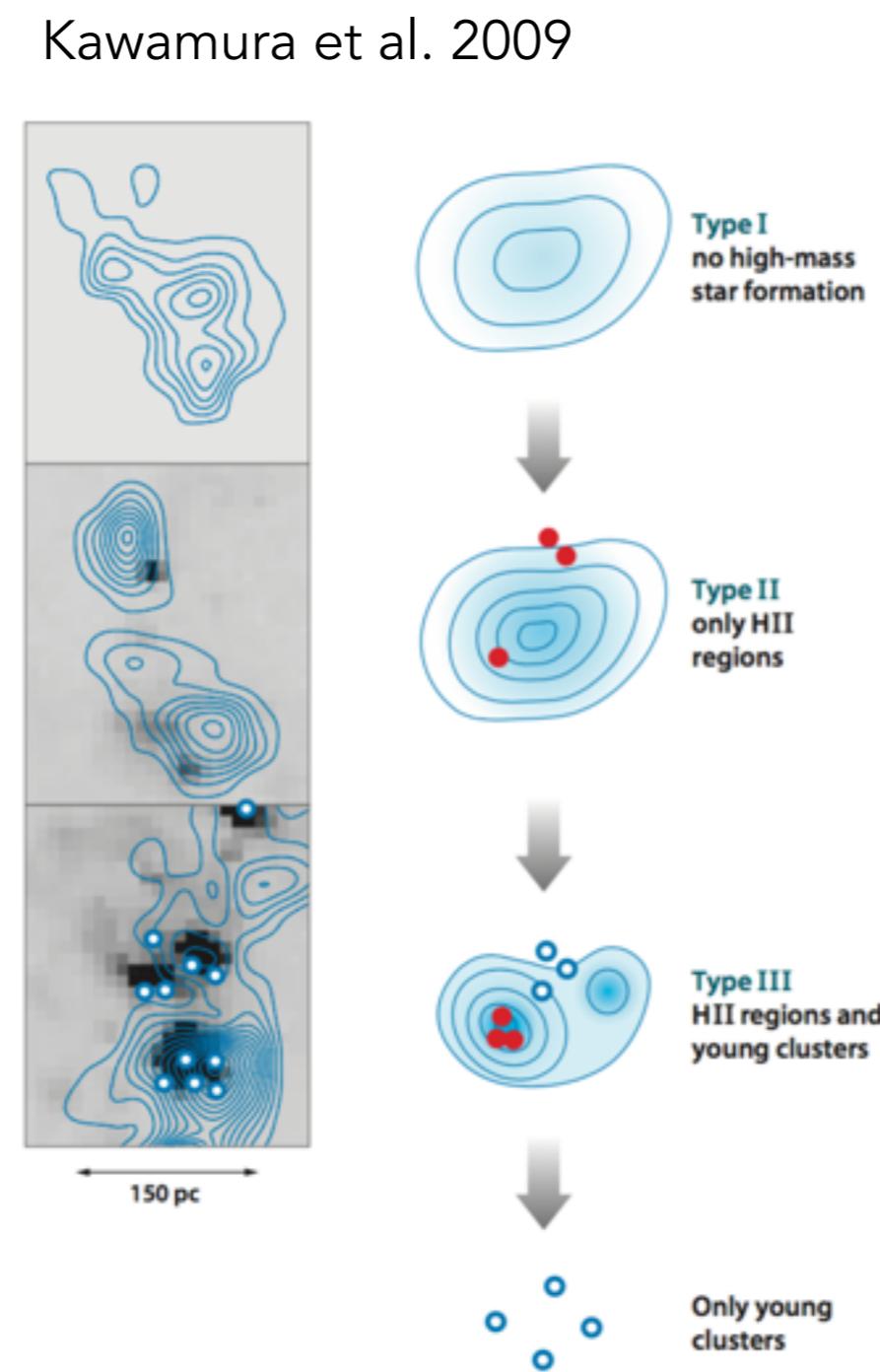
Colombo et al. 2014 - PAWS survey of M51



Molecular Clouds

Observed Characteristics

- Self-Gravity
- Turbulence
- Substructure
- Magnetic Fields
- Mass Spectrum
- **Lifetimes**
- Star Formation



If star formation rate is constant, relative numbers of clouds in each evolutionary state, plus ages of clusters when no molecular gas is around gives you cloud lifetimes.

~20-30 Myr

Star Formation

Table 1 Properties of dark clouds, clumps, and cores

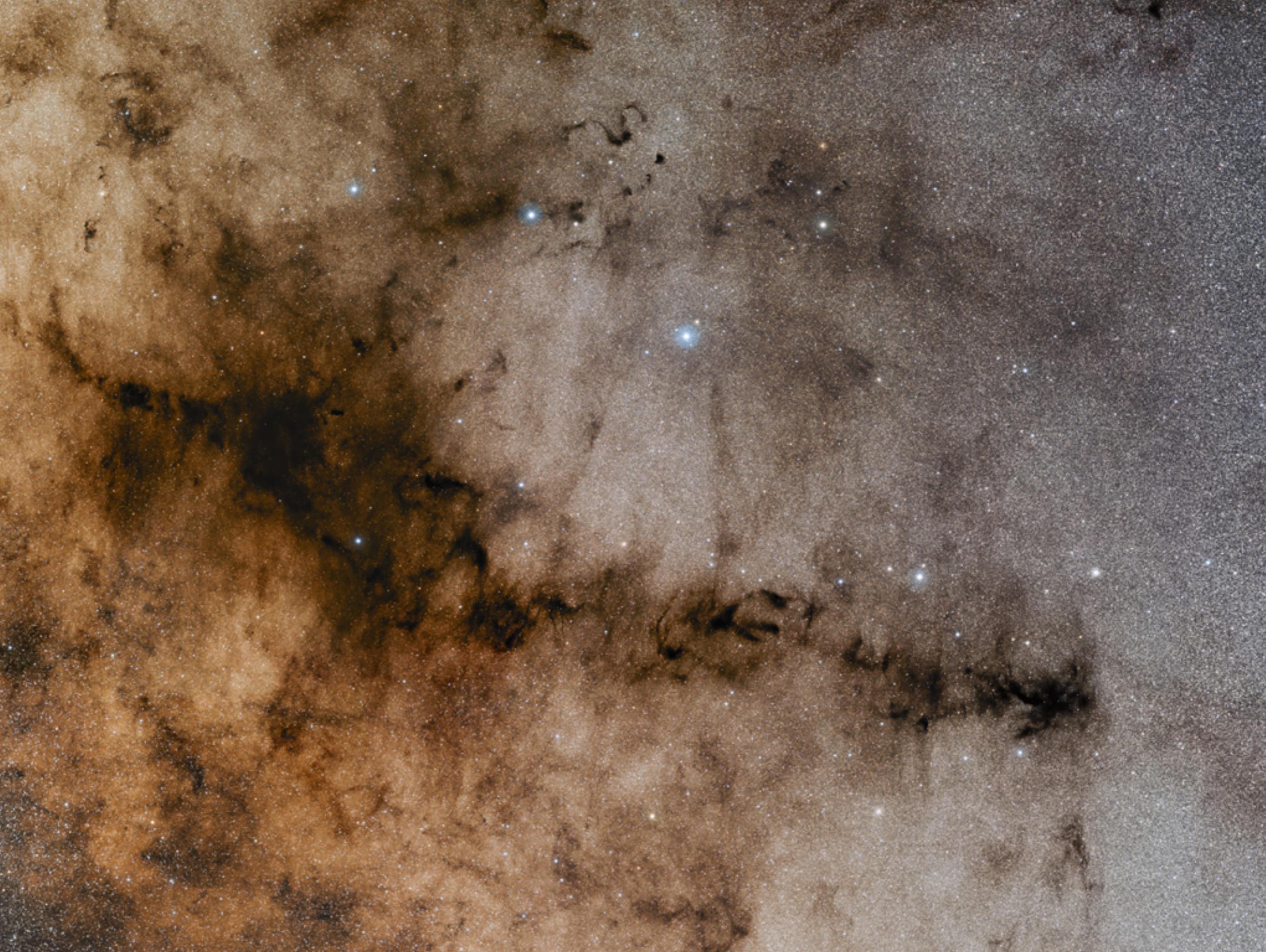
	Clouds ^a	Clumps ^b	Cores ^c
Mass (M_{\odot})	$10^3 - 10^4$	50–500	0.5–5
Size (pc)	2–15	0.3–3	0.03–0.2
Mean density (cm^{-3})	50–500	10^3 – 10^4	10^4 – 10^5
Velocity extent (km s^{-1})	2–5	0.3–3	0.1–0.3
Crossing time (Myr)	2–4	≈ 1	0.5–1
Gas temperature (K)	≈ 10	10–20	8–12
Examples	Taurus, Oph, Musca	B213, L1709	L1544, L1498, B68

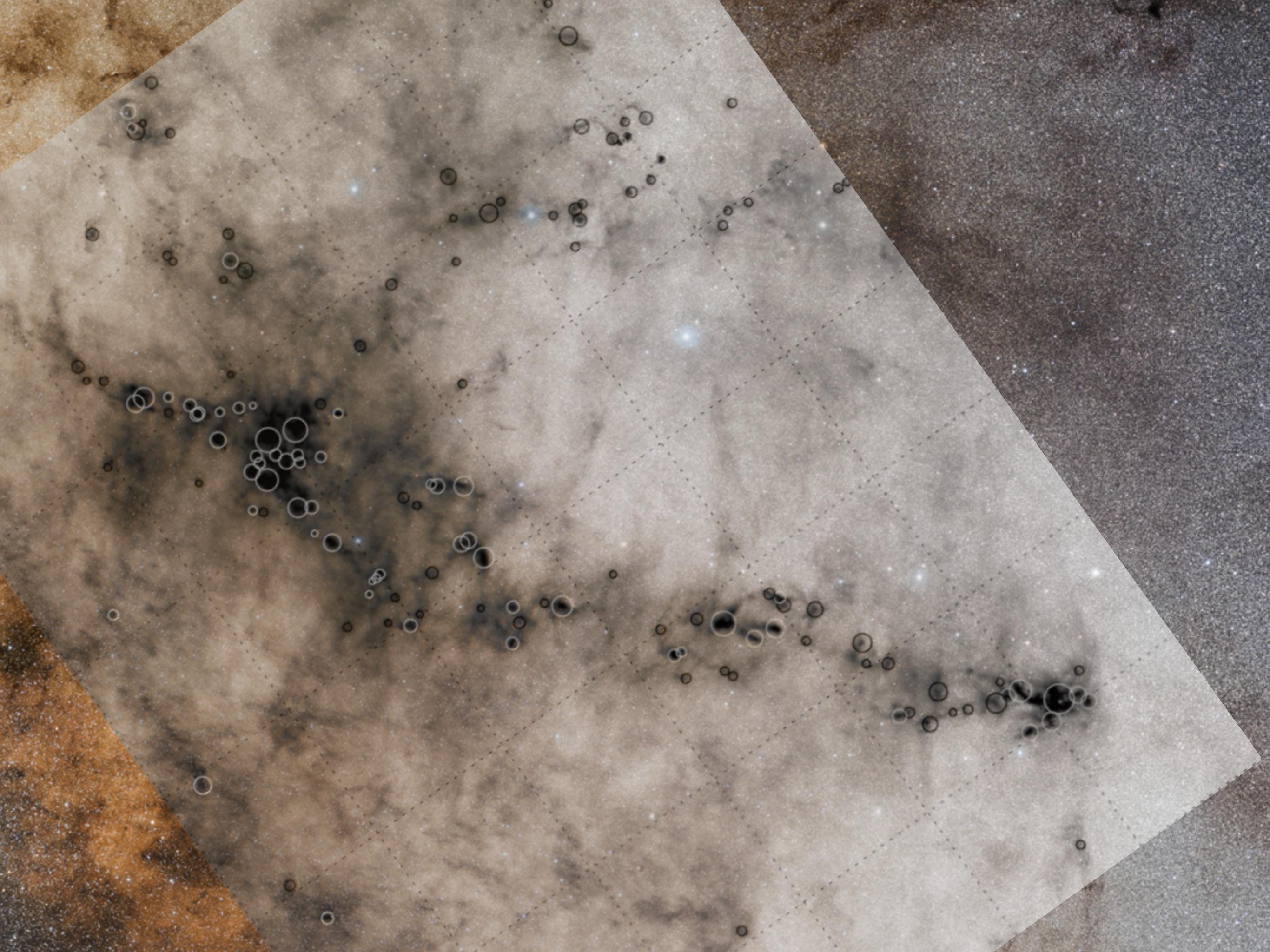
^aCloud masses and sizes from the extinction maps by Cambrésy (1999), velocities and temperatures from individual cloud CO studies.

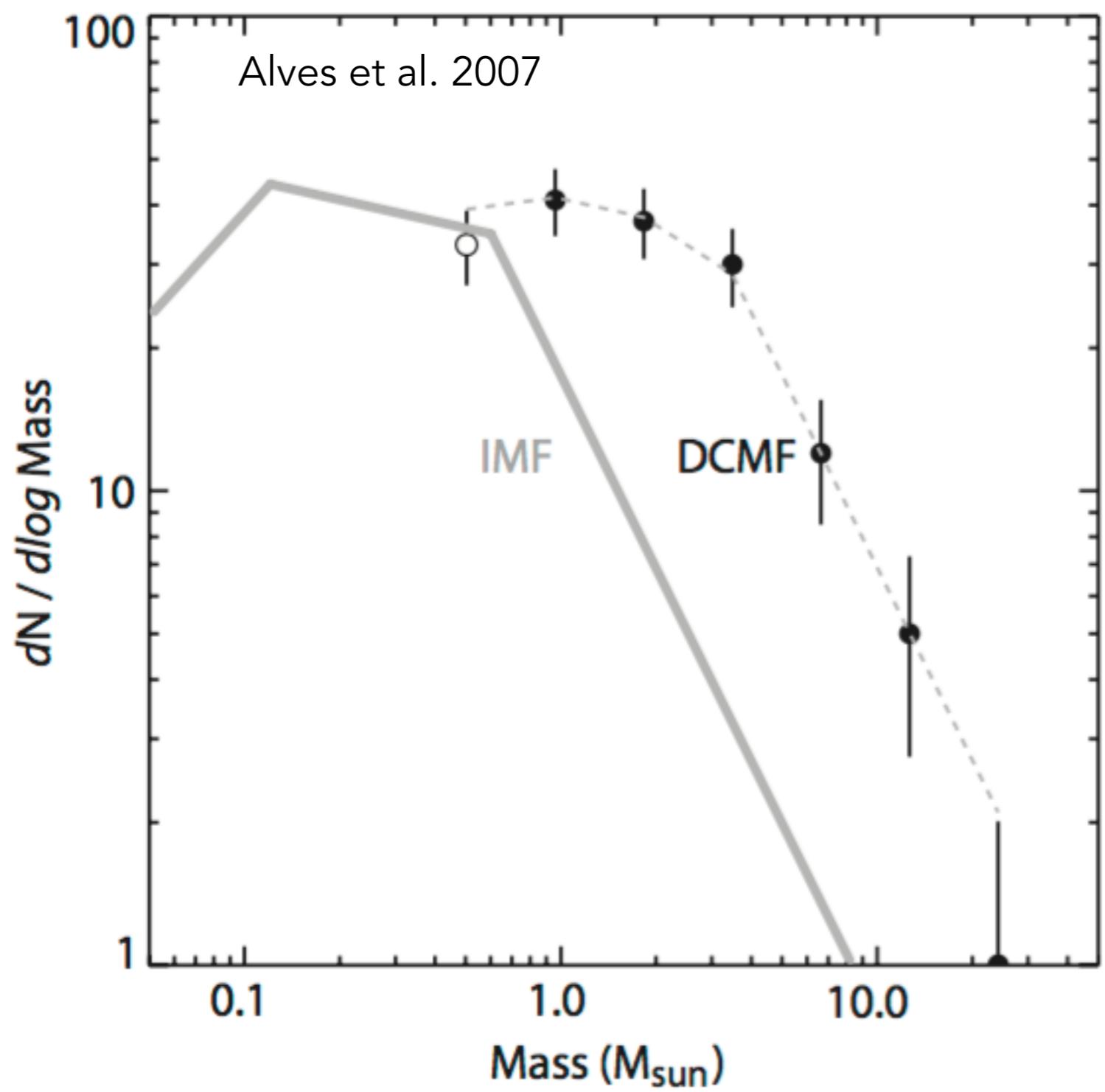
^bClump properties from Loren (1989) (^{13}CO data) and Williams, de Geus & Blitz (1994) (CO data).

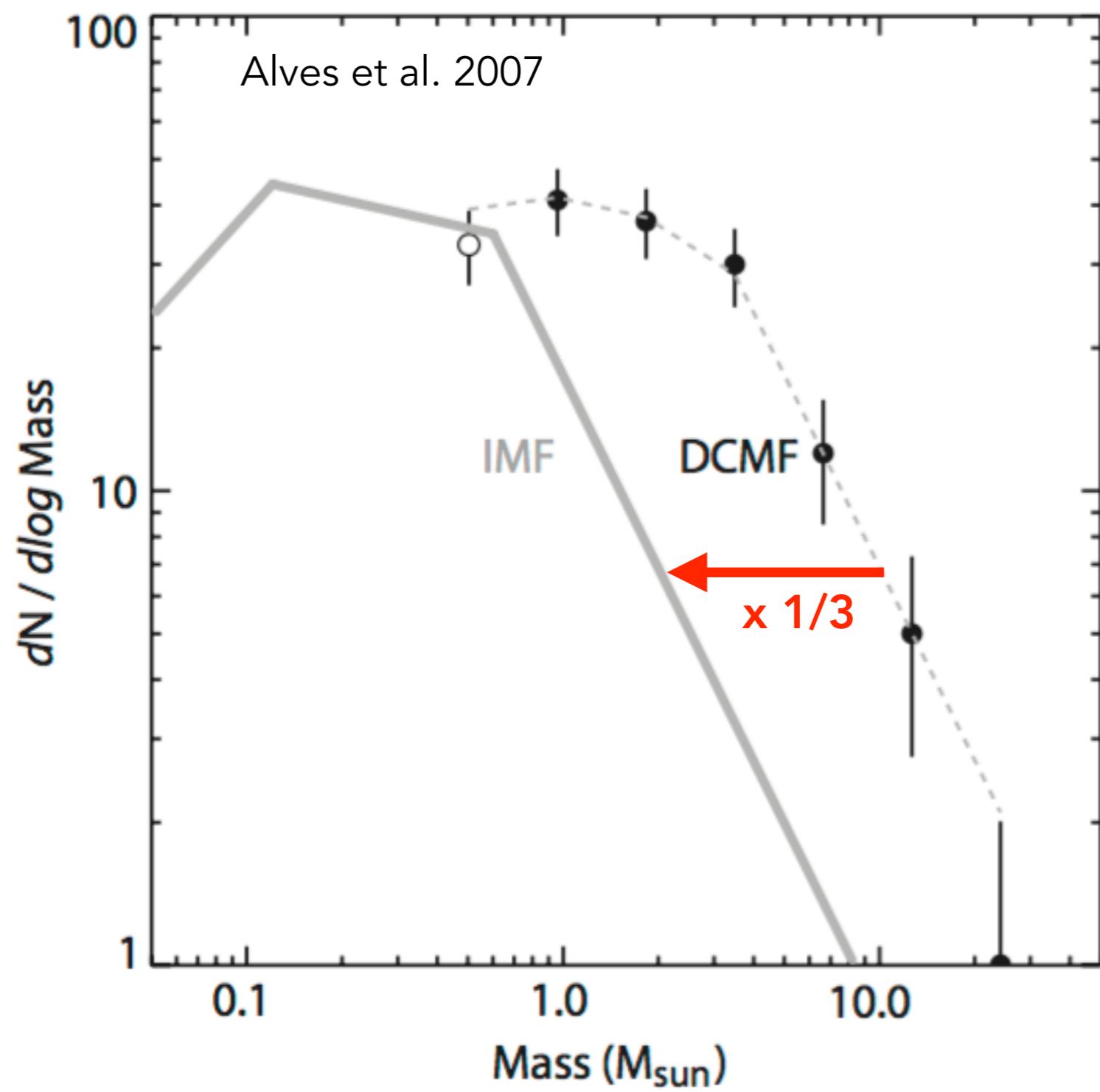
^cCore properties from Jijina, Myers & Adams (1999), Caselli et al. (2002a), Motte, André & Neri (1998), and individual studies using NH_3 and N_2H^+ .

Bergin & Tafalla 2007









The Initial Mass Function

Number of stars per unit $\log(M)$ that are formed.

Controversy persists over whether it is the same everywhere.

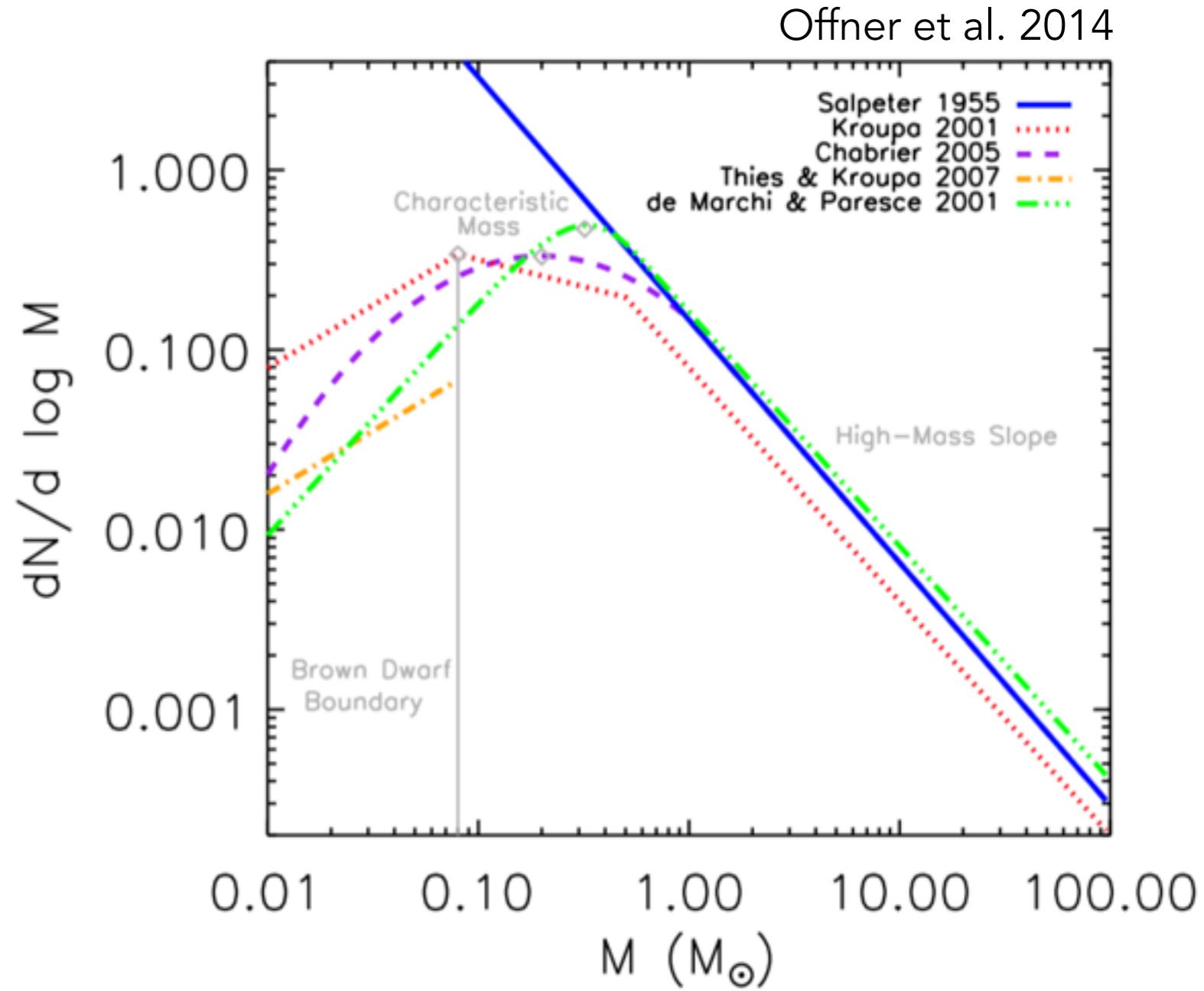
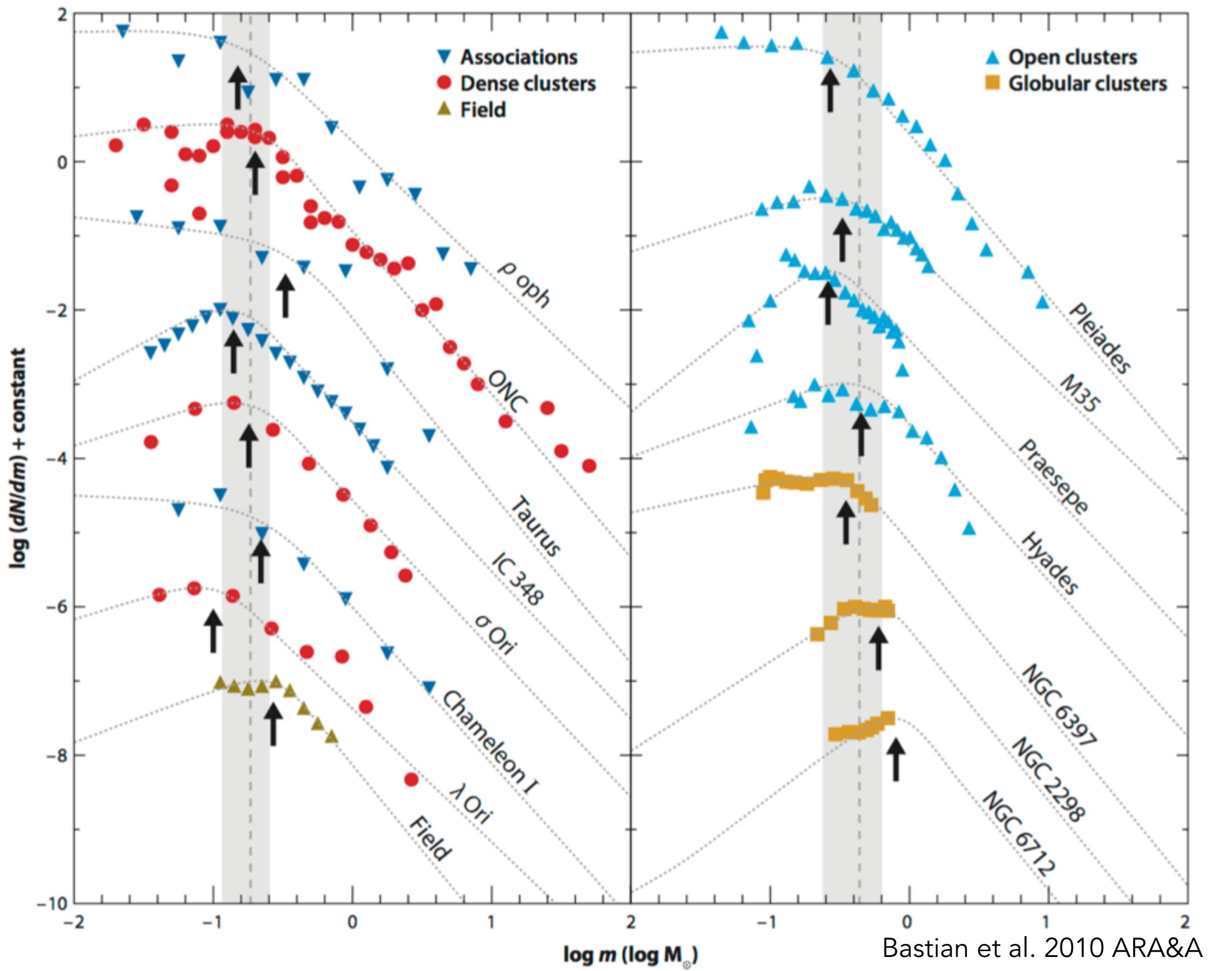
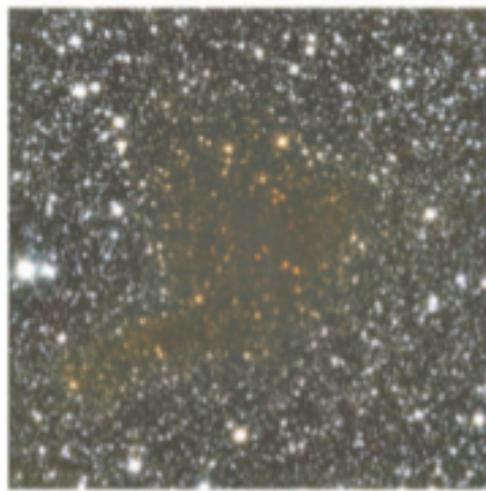


Fig. 1.— IMF functional forms proposed by various authors from fits to Galactic stellar data. With the exception of the Salpeter slope, the curves are normalized such that the integral over mass is unity. When comparing with observational data, the normalization is set by the total number of objects as shown in Figure 2.



Cores in Molecular Clouds

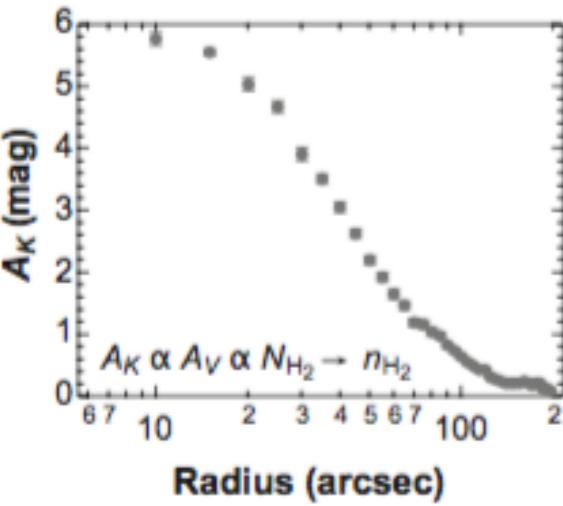
a Barnard 68 K band



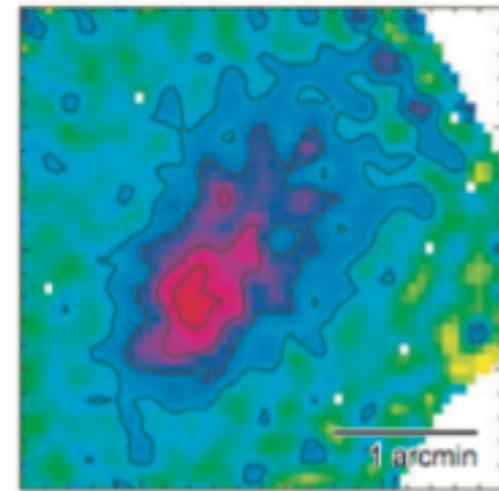
$$A_V = r_V^{H,K} E(H - K)$$

$$A_V = f N_H$$

$$N_H = (r_V^{H,K} f^{-1}) \cdot E(H - K)$$



b L1544 1.2 mm continuum



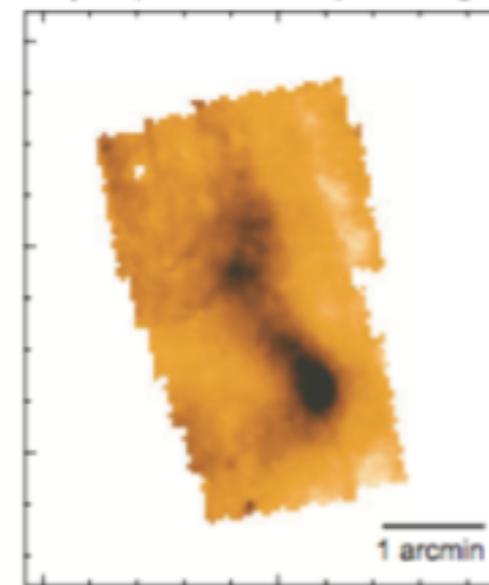
For optically thin emission:

$$I_\nu = \int \kappa_\nu \rho B_\nu(T_d) dl$$

$$I_\nu = m \langle \kappa_\nu B_\nu(T_d) \rangle N_H$$

$$N_H = I_\nu / \langle m \kappa_\nu B_\nu(T_d) \rangle^{-1}$$

c ρ Oph core D 7 μm image

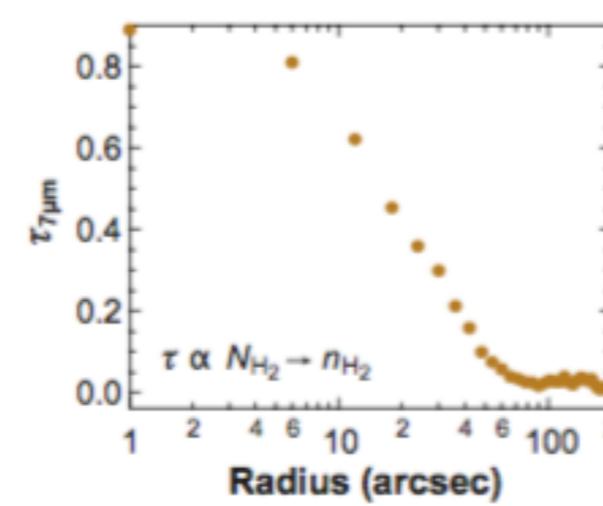
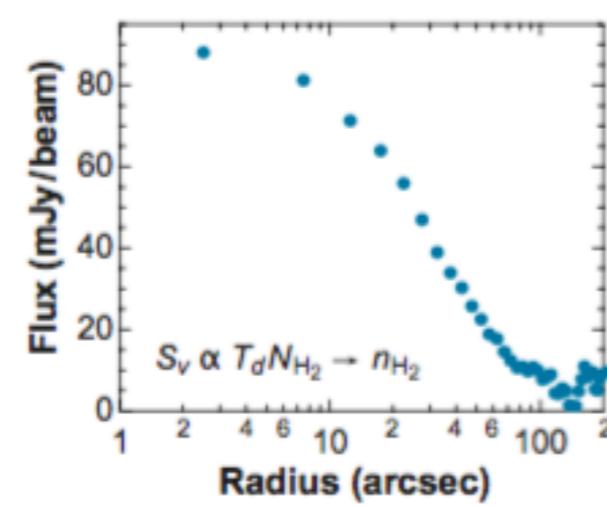


$$I_\nu = I_\nu^{bg} \exp(-\tau_\lambda) + I_\nu^{fg}$$

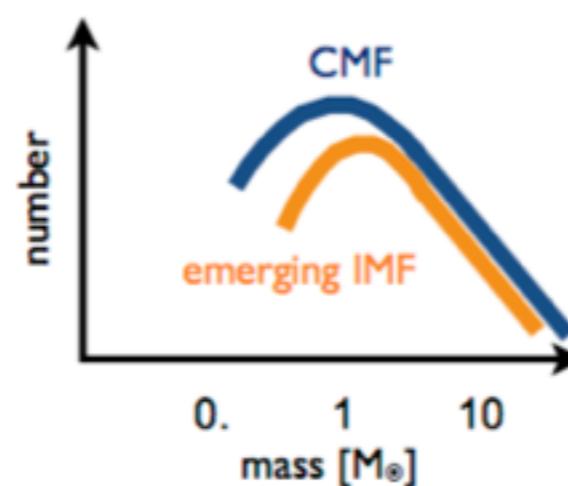
$$\tau_\lambda = \sigma_\lambda N_H$$

$$N_H = \frac{1}{\sigma_\lambda} \ln \left[\frac{I_\nu^{bg}}{I_\nu - I_\nu^{fg}} \right]$$

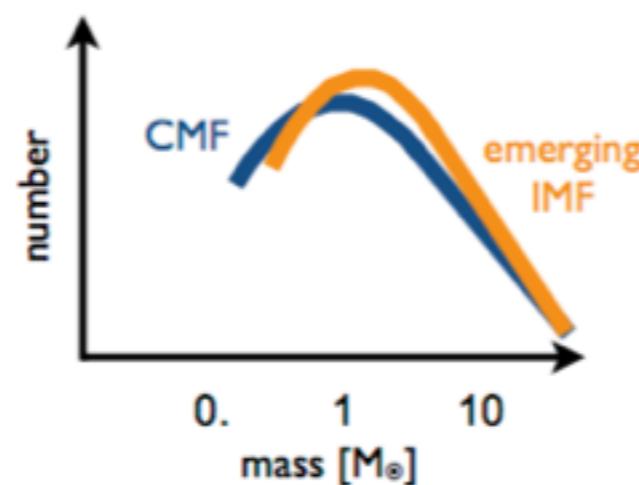
Column density profiles of dense cores are similar to Bonnor-Ebert profile (isothermal, marginally stable spherical cloud)



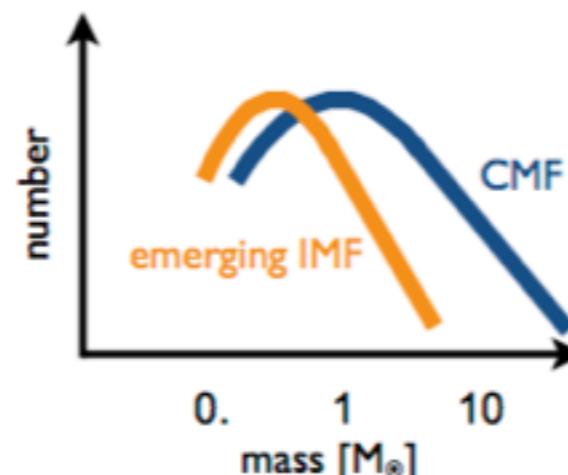
i) Not all cores are 'prestellar'. Here we show the emerging IMF that could arise if the low-mass cores in the CMF are transient 'fluff'.



ii) Core growth is not self-similar. Here we show the emerging IMF that could arise if, say, only the low-mass cores in the CMF are still accreting.



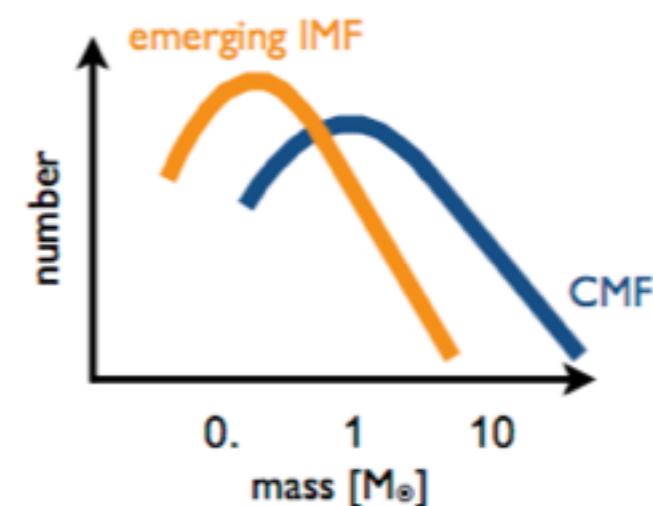
iii) Varying star formation efficiency (SFE). Here we show the emerging IMF that could arise if the high-mass cores in the CMF have a lower SFE than their low-mass siblings.



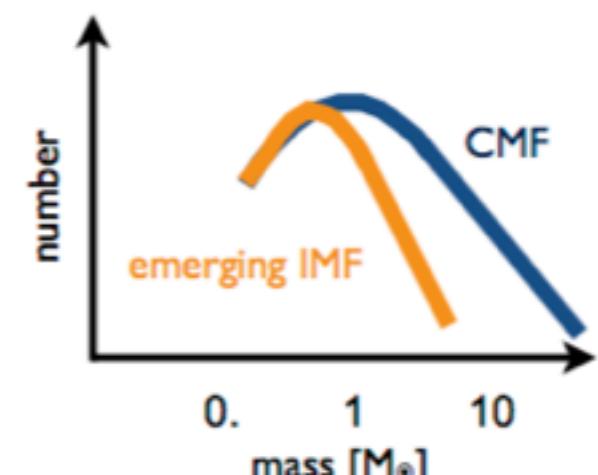
When does the CMF map to the IMF?

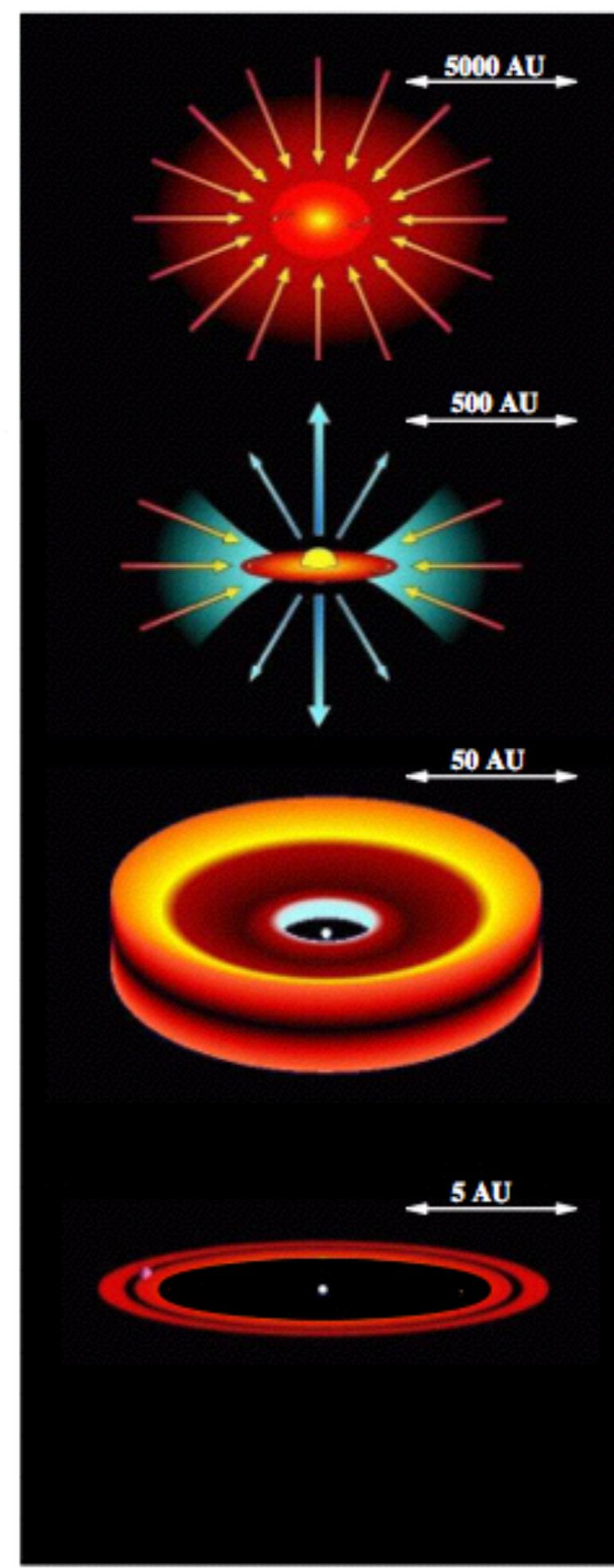
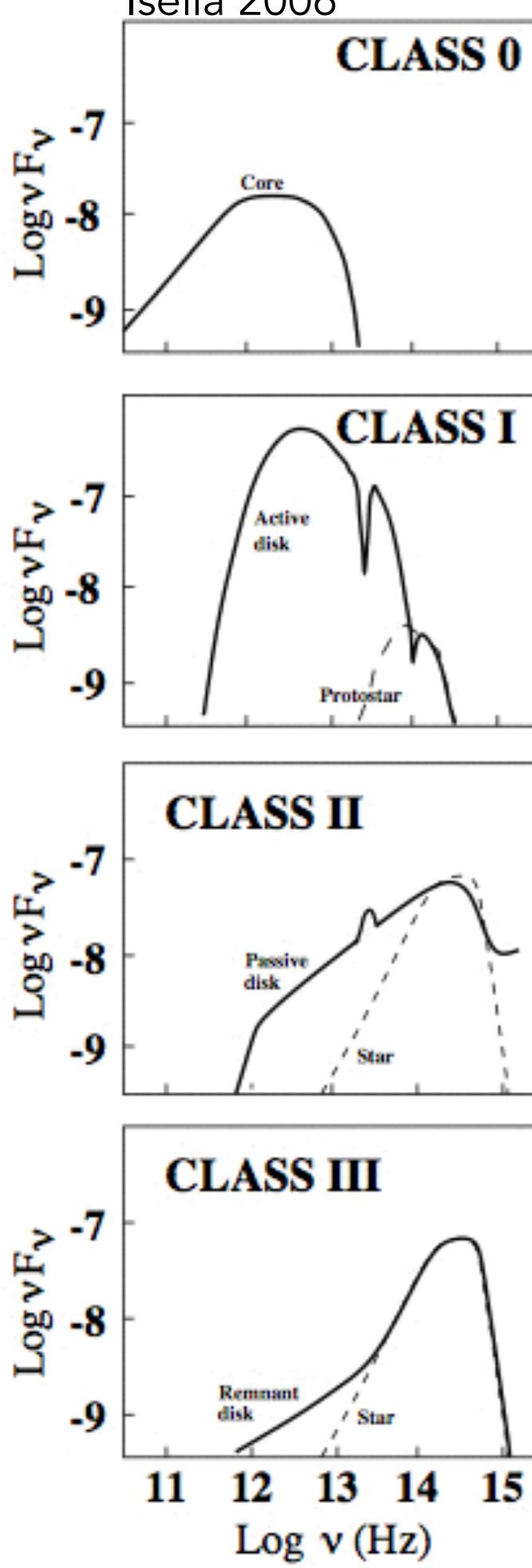
Offner et al. 2014

iv) Fragmentation is not self-similar. Here we show the emerging IMF that could arise if the cores in the CMF fragment based on the number of initial Jeans masses they contain.



v) Varying embedded phase timescale. Here we show the emerging IMF that could arise if the low-mass cores in the CMF finish before the high-mass cores.





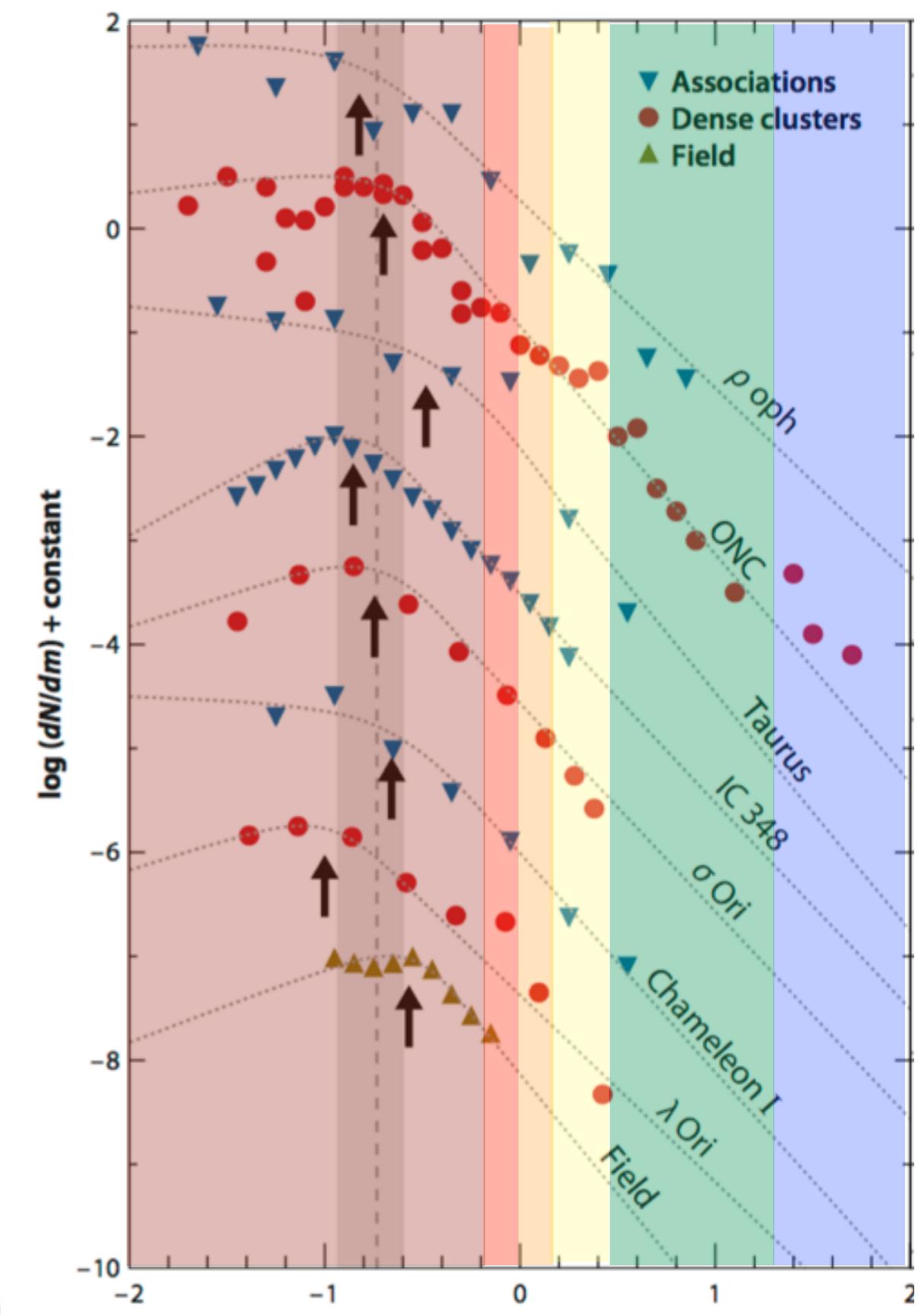
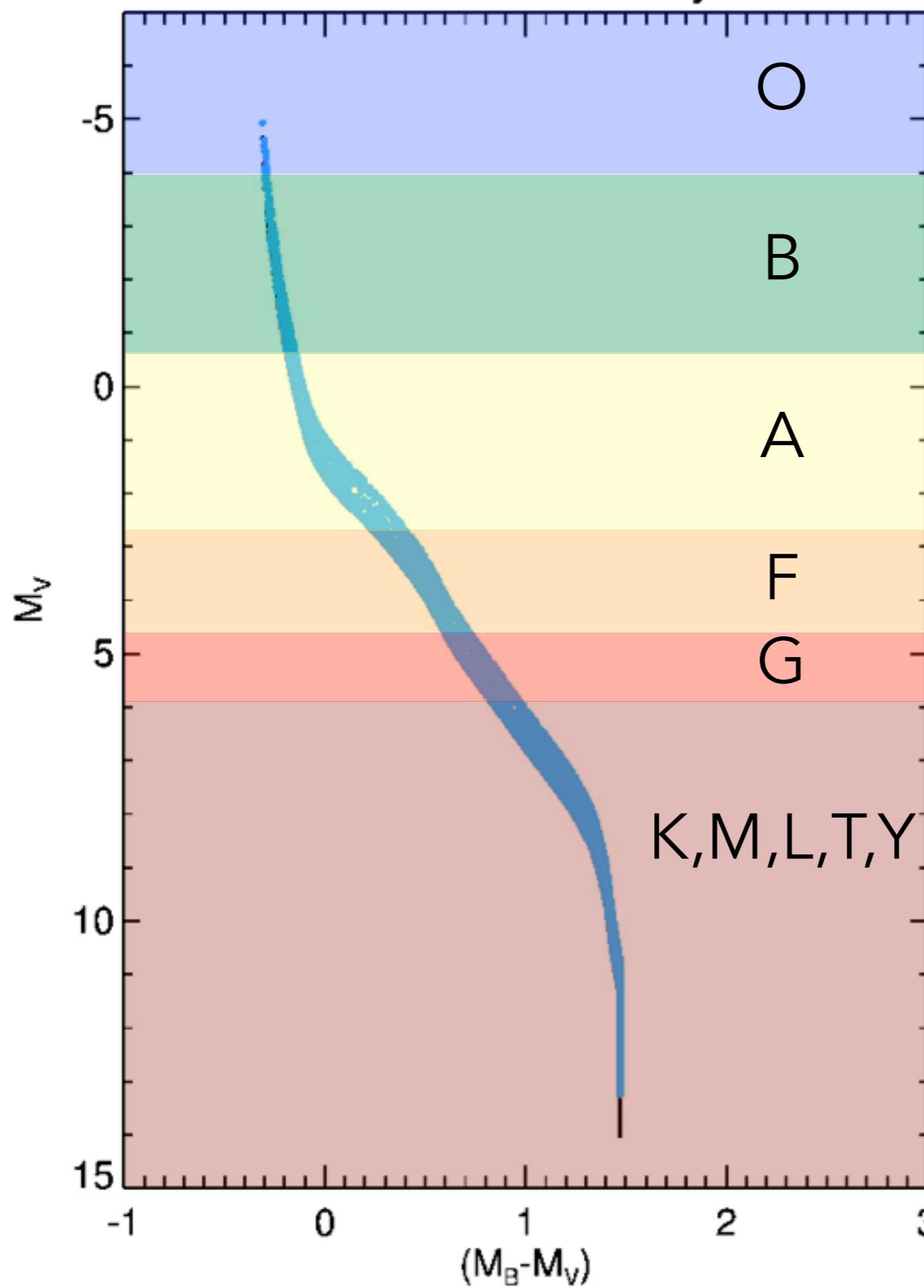
Gravitational collapse

Angular momentum
-> disk formation
-> outflows & jets

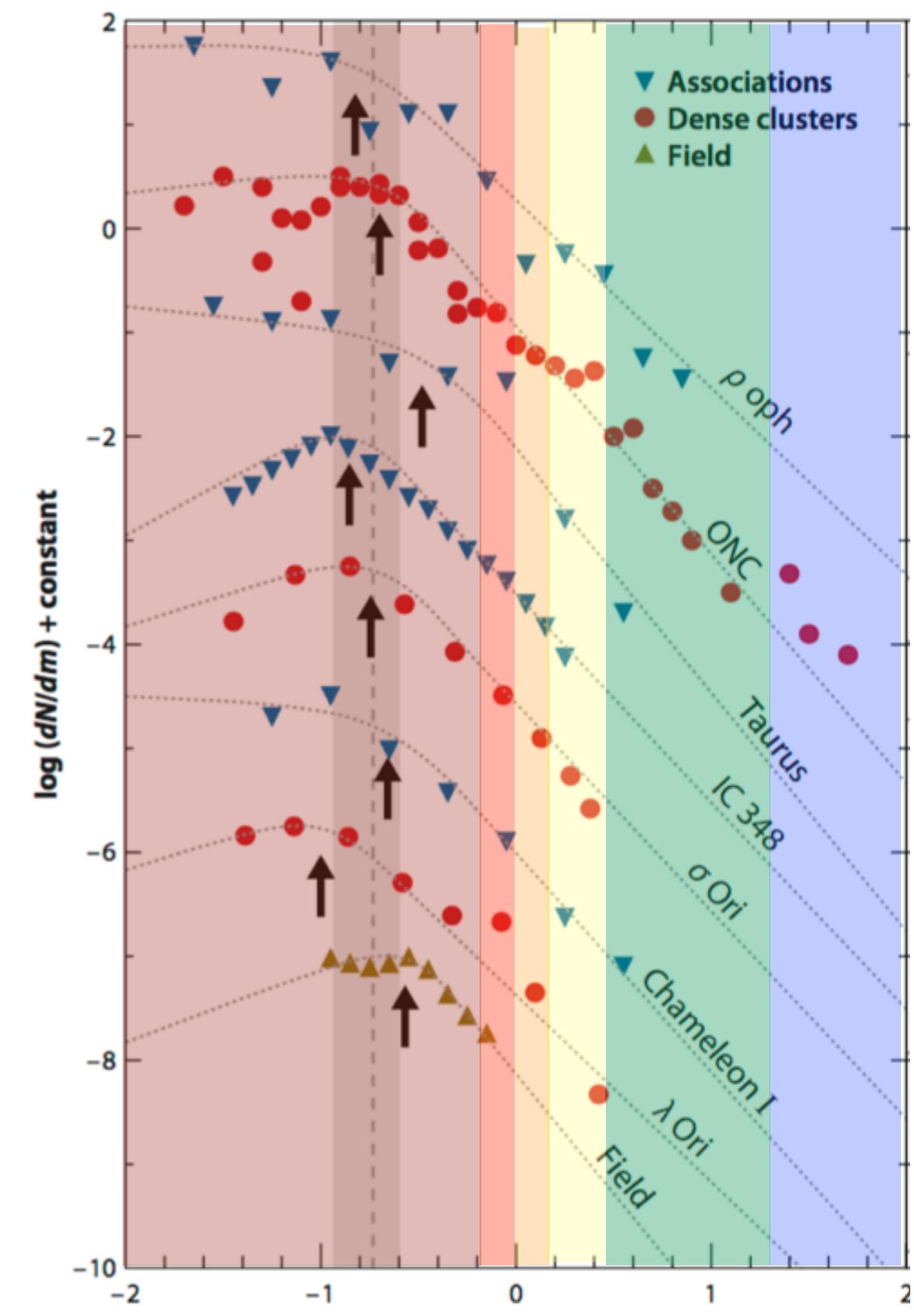
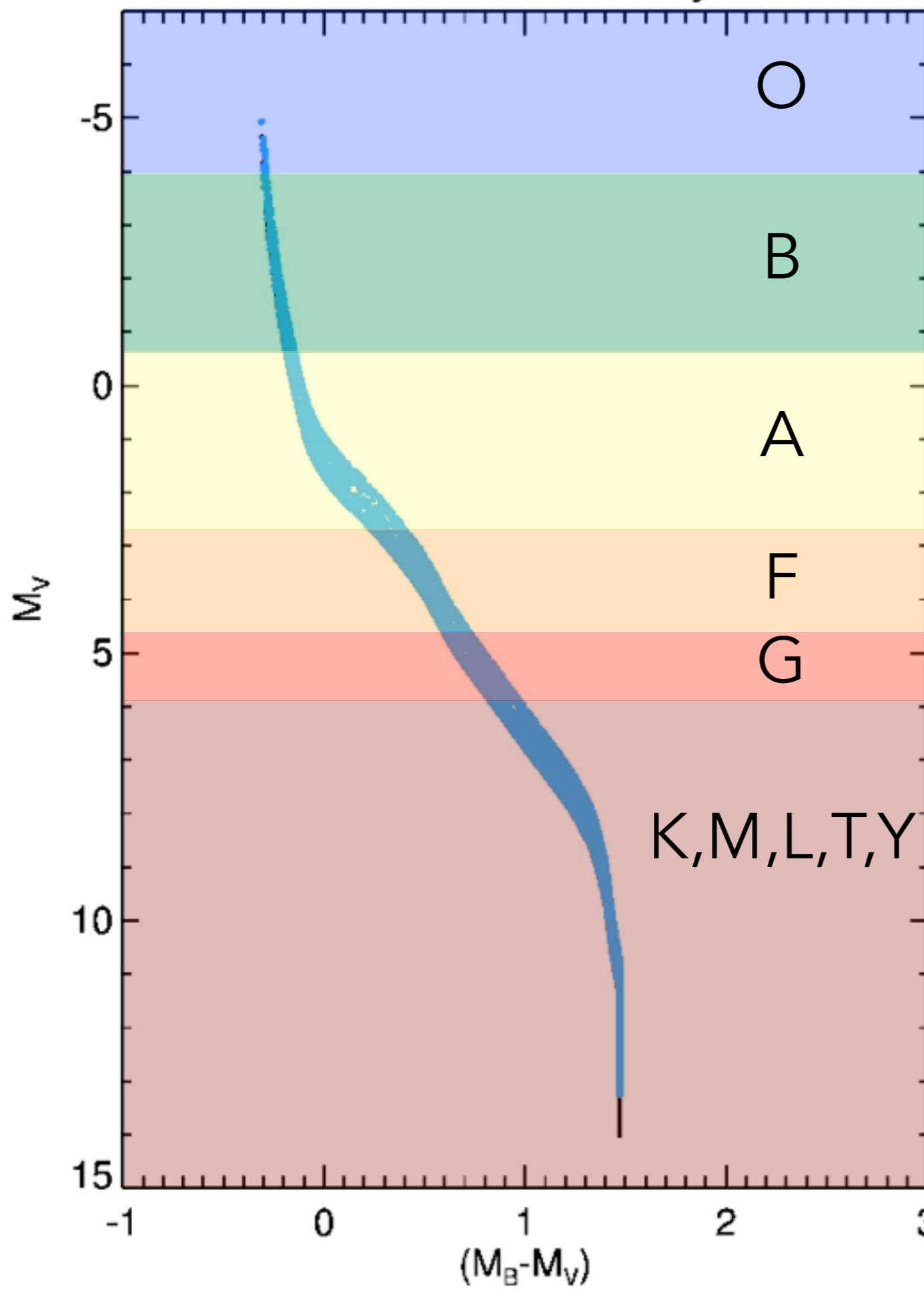
Most material is in a disk,
accretion onto protostar
through disk.

Most material accreted,
remnant disk.

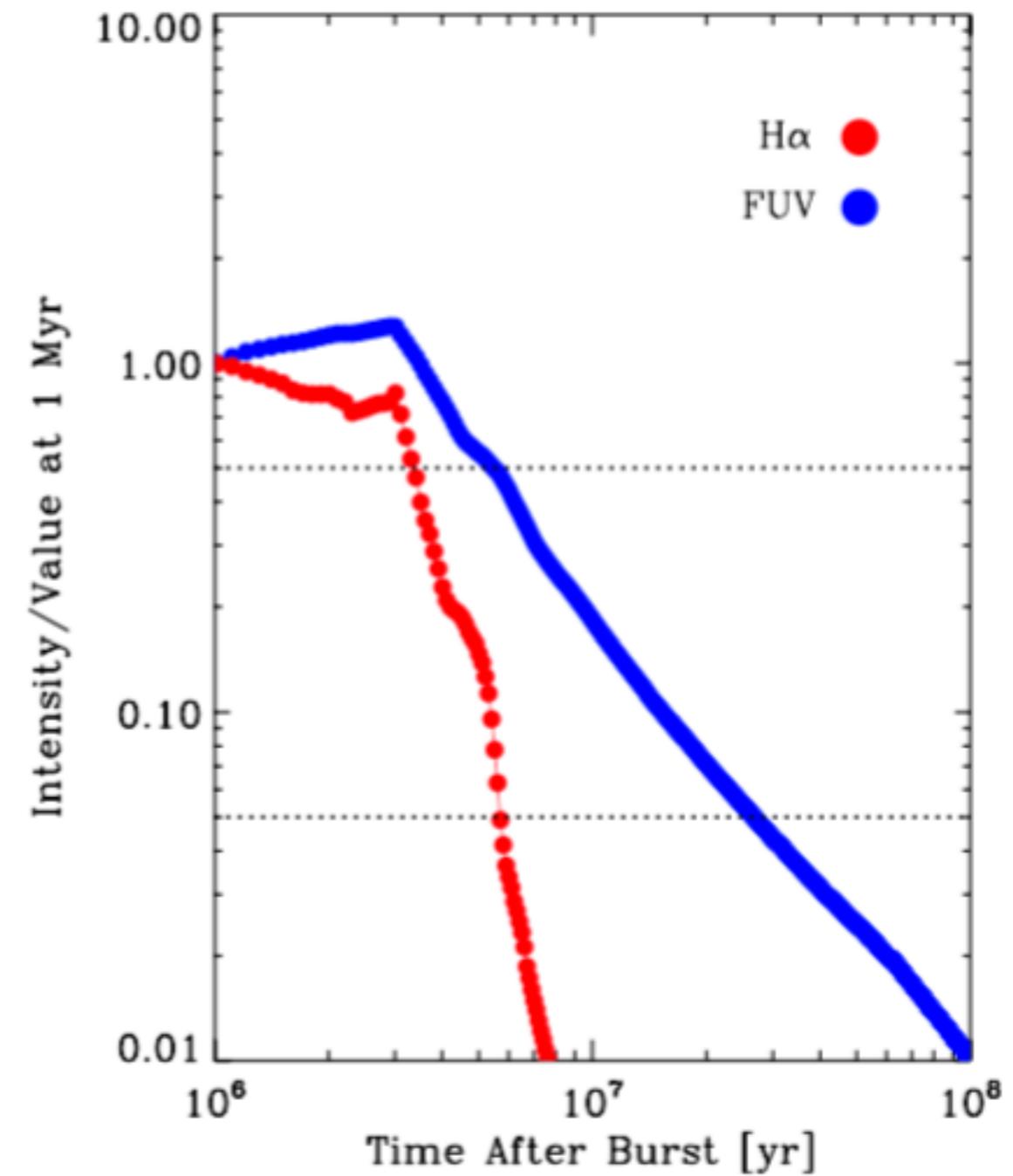
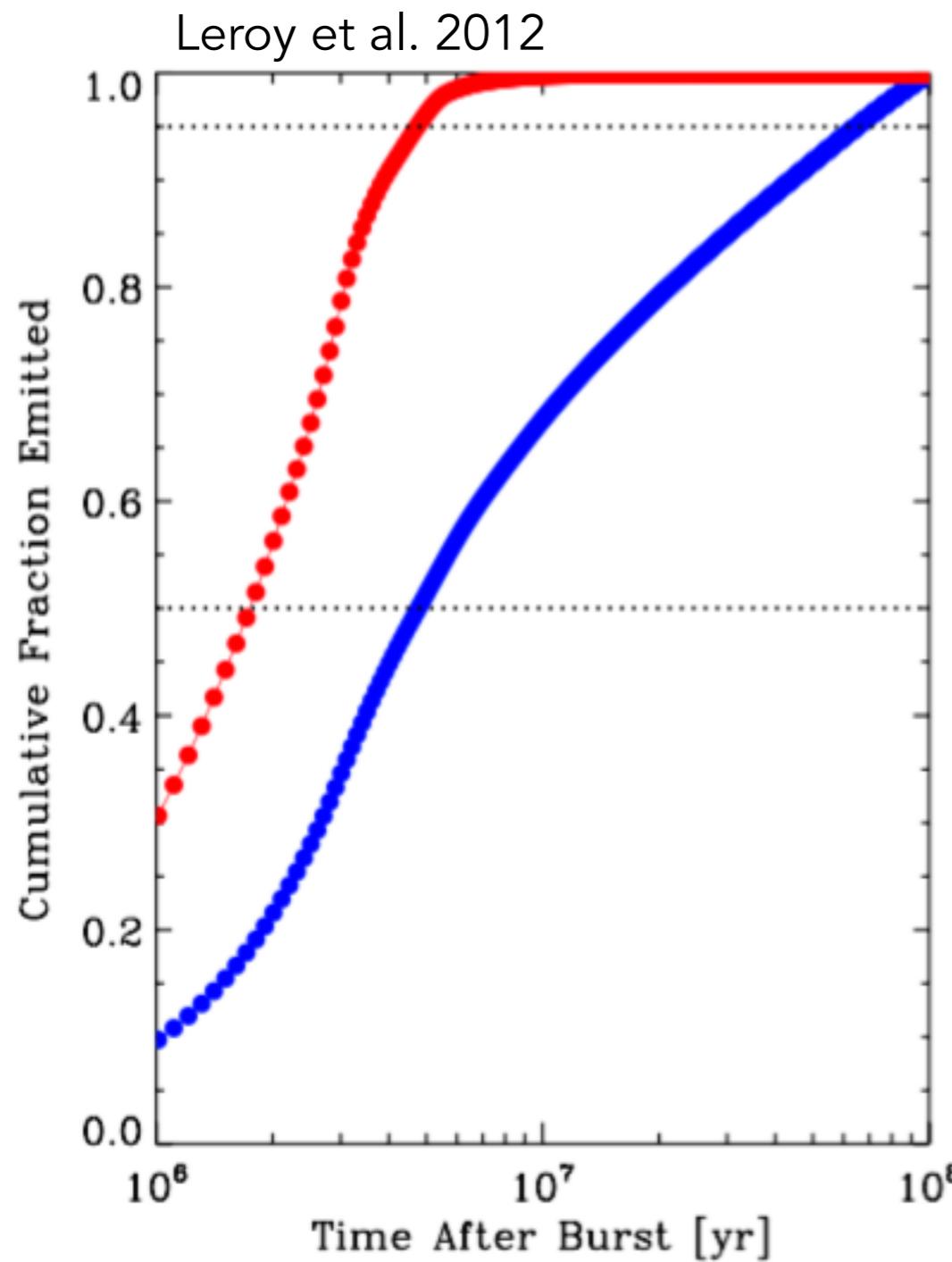
Time = 0.00 Myr



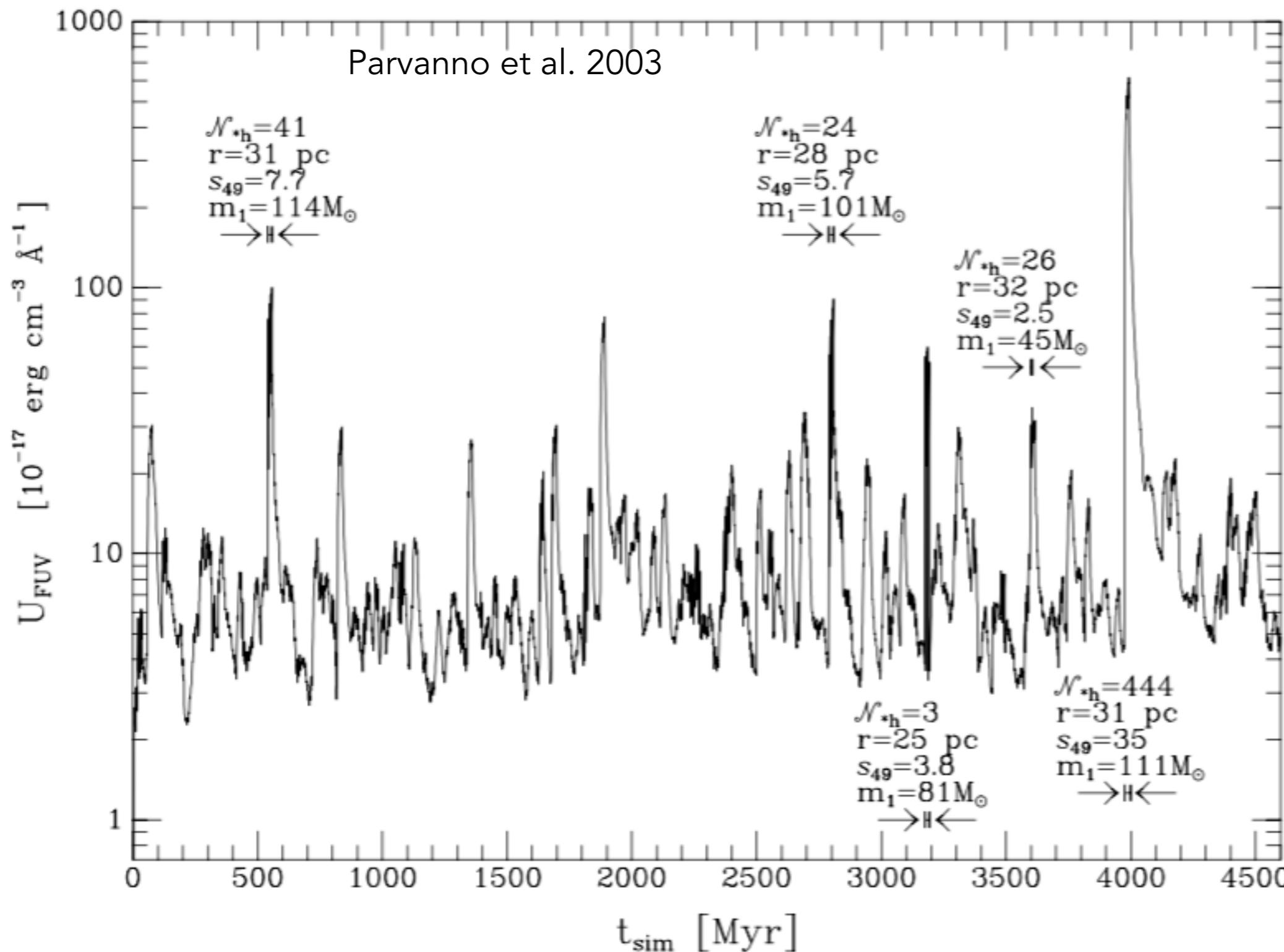
Time = 0.00 Myr



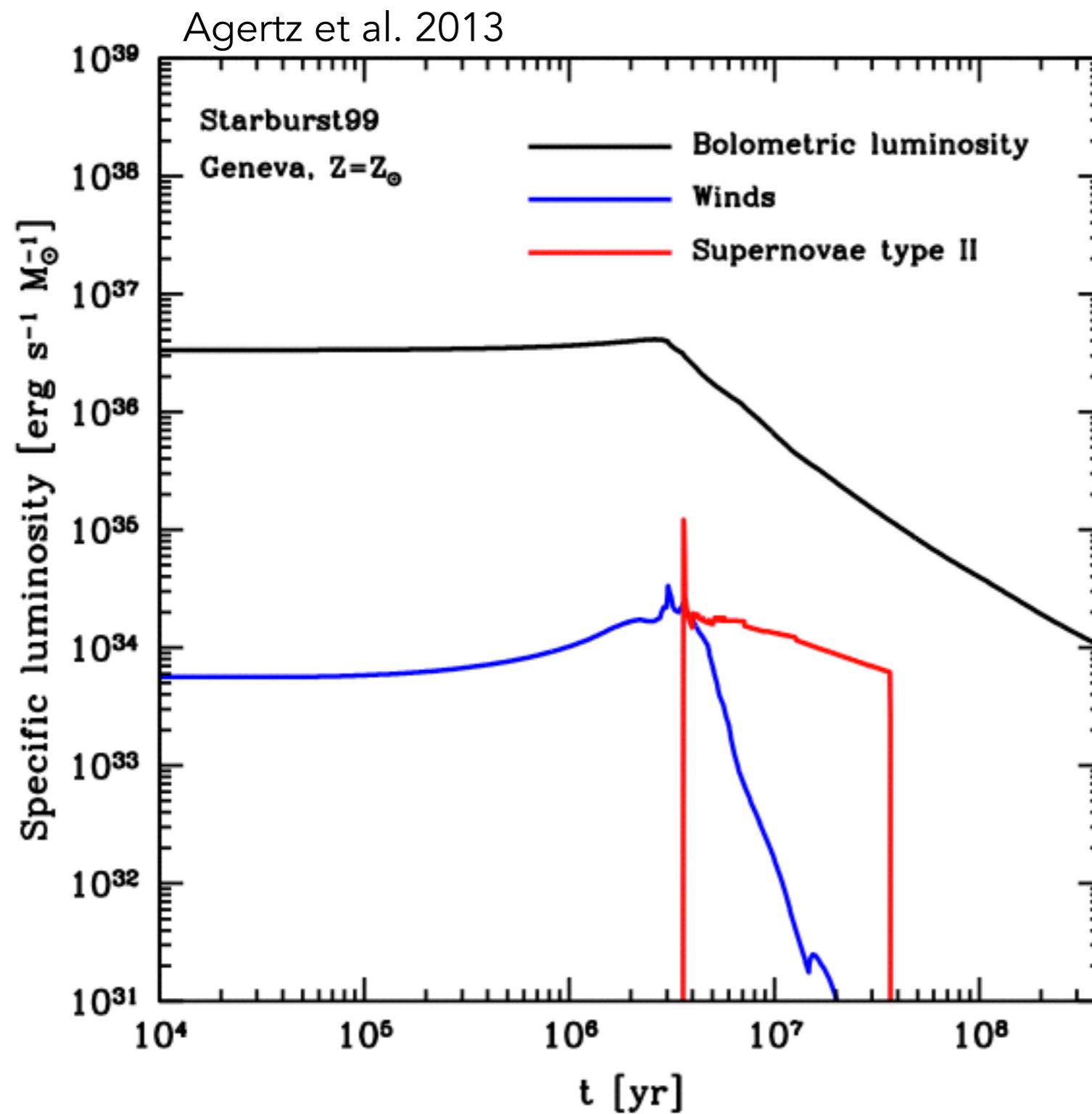
Radiative Feedback



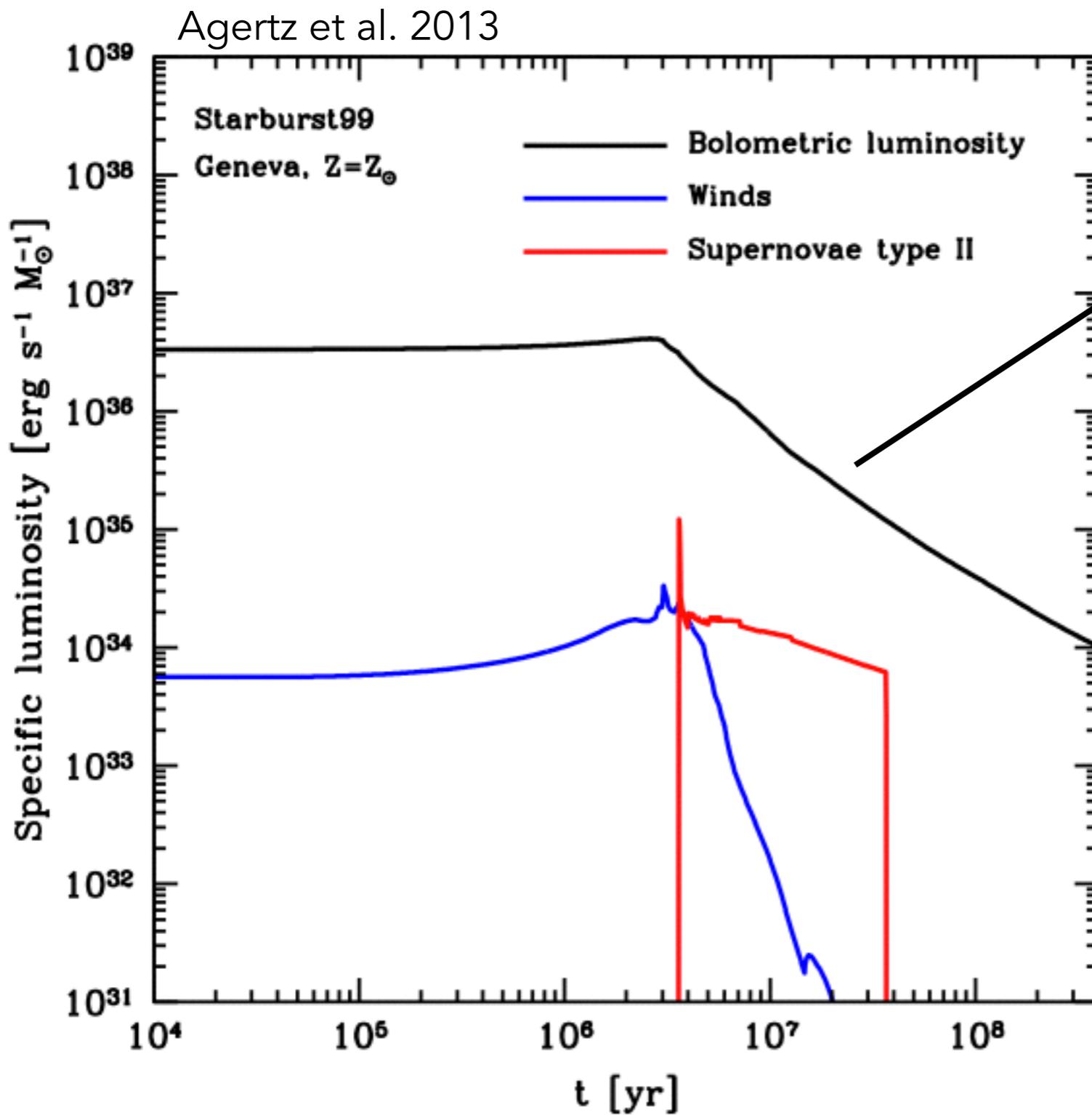
Radiative Feedback



Mechanical Feedback



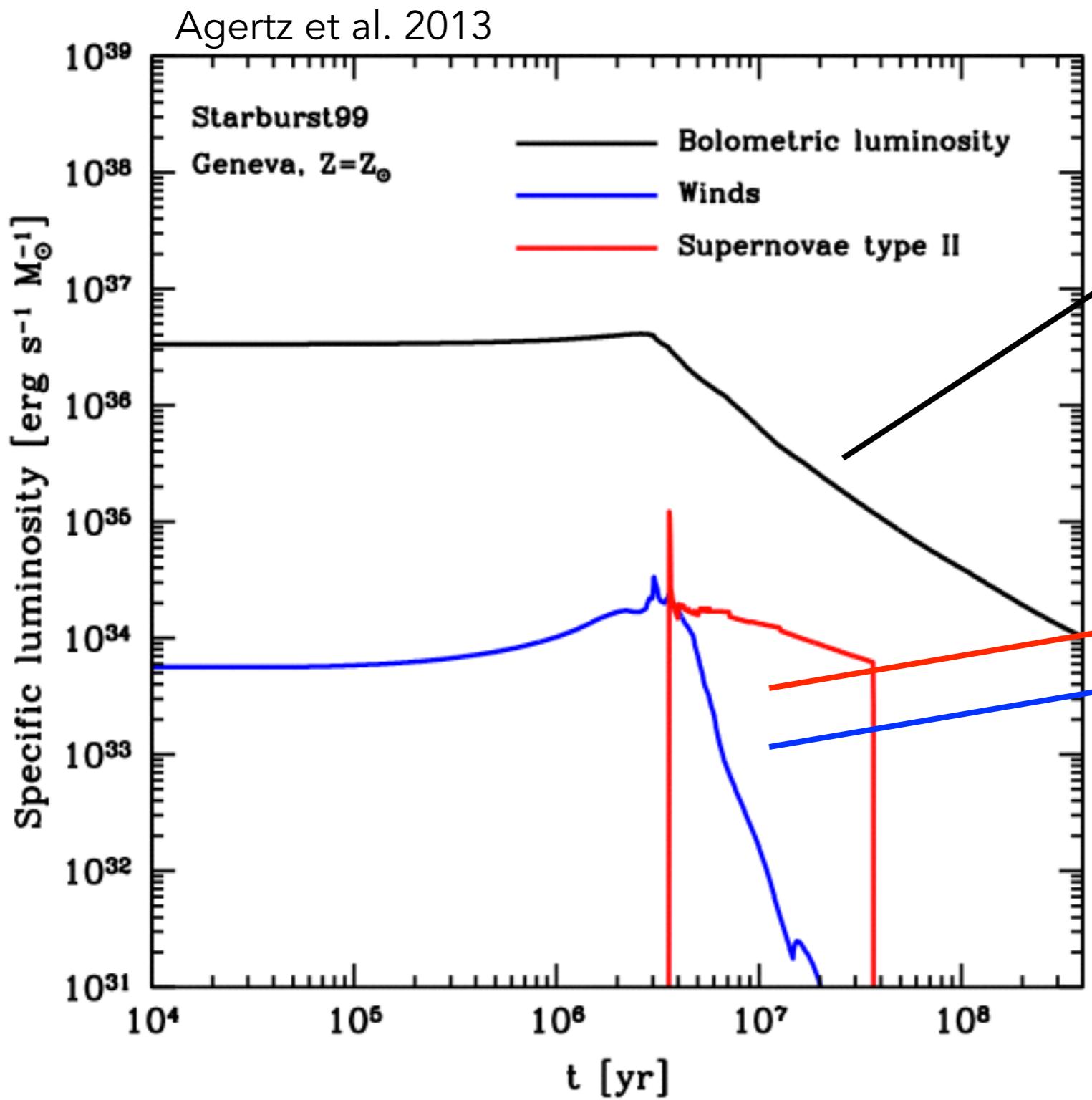
Mechanical Feedback



fraction into ionization,
photoelectric heating,
dust heating?

deposited where?

Mechanical Feedback



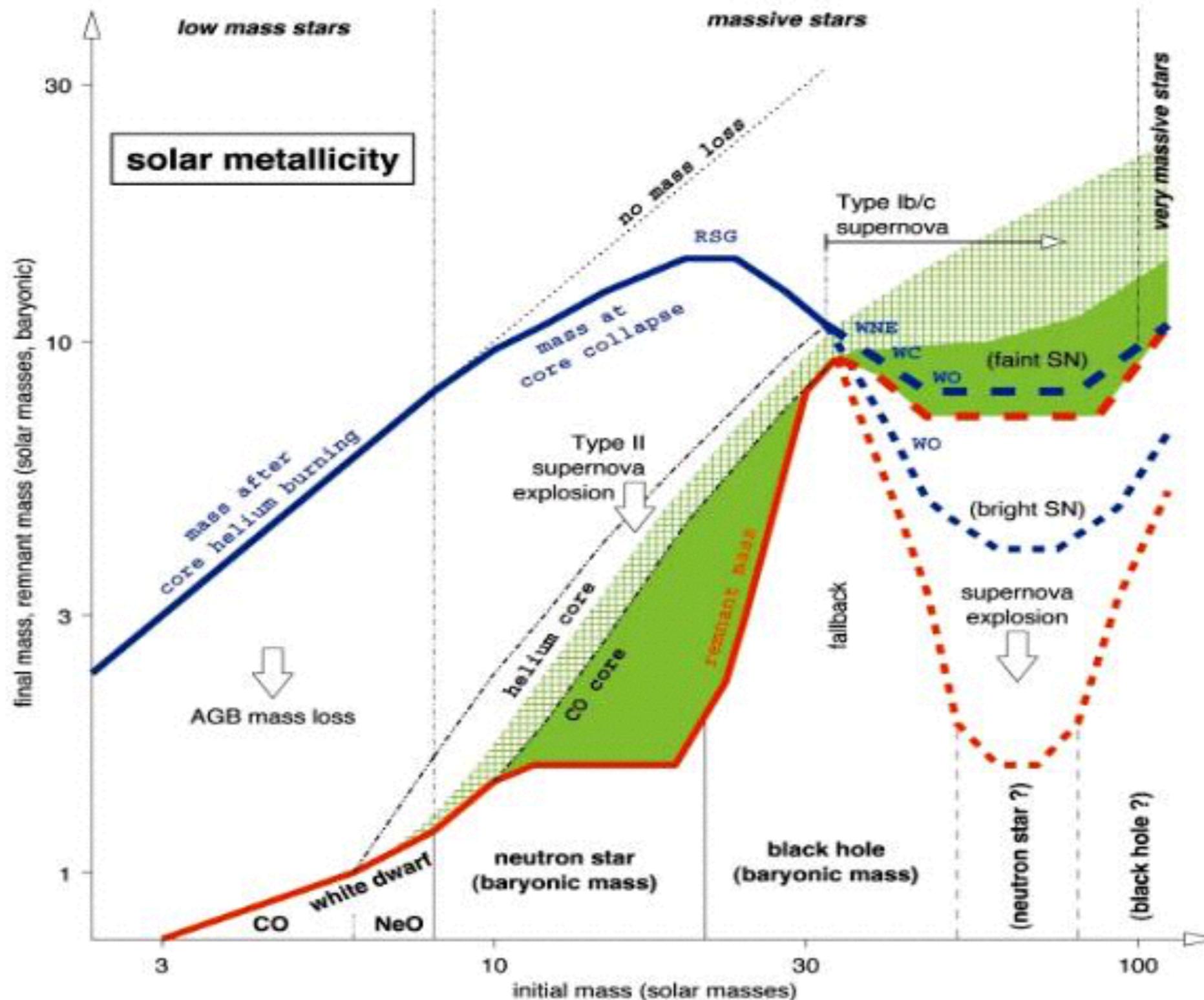
fraction into ionization,
photoelectric heating,
dust heating?

deposited where?

which phases is this
deposited into, at what
distances from stars?

Supernovae

Woosley et al. 2002



Stars with
masses $> 8 M_{\odot}$
explode.

Supernovae
produce
 $\sim 10^{53}$ ergs in
neutrinos
 $\sim 10^{51}$ ergs in
kinetic energy

Supernovae

Initially: $M_{\text{ejecta}} \sim \text{few } M_{\odot}$, $v_{\text{ejecta}} \sim 10^4 \text{ km/s}$

Phase	Characteristics	Ends when...	Radius at end
Free Expansion	ballistic expansion, shock wave into ISM/CSM, ejecta cools due to adiabatic expansion, reverse shock when $P_{\text{shocked ISM}} > P_{\text{ej}}$	$M_{\text{swept}} > M_{\text{ej}}$	$R \sim t$
Sedov-Taylor	ejecta is very hot, $P_{\text{ej}} > P_{\text{ISM}}$ expansion driven by hot gas, radiation losses are unimportant	radiative losses become important	$R \sim t^{2/5}$
Snow Plow	pressure driven expansion with radiative loss, then momentum driven	shock becomes subsonic	$R \sim t^{2/7}$ $R \sim t^{1/4}$
Fadeaway	turbulence dissipates remnant structure and merges with ISM	-	-

SILCC: SImulating the LifeCycle of molecular Clouds



Stefanie Walch
Philipp Girichidis
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Thomas Peters
Dominik Derigs
Christian Baczynski

Walch et al., MNRAS 454, 238 (2015)
Girichidis et al., arXiv:1508.06646

KS SN rate, mixed driving

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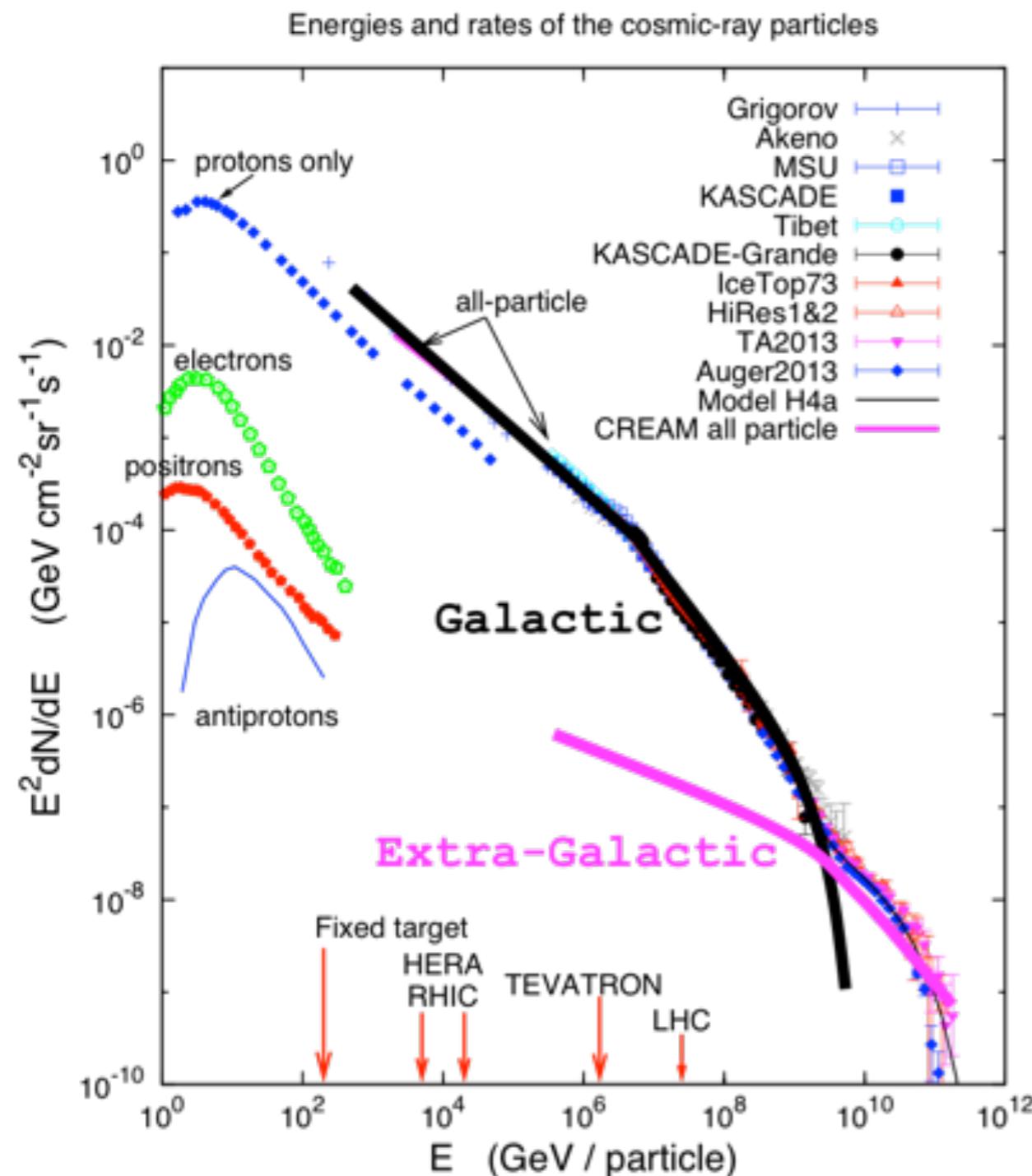


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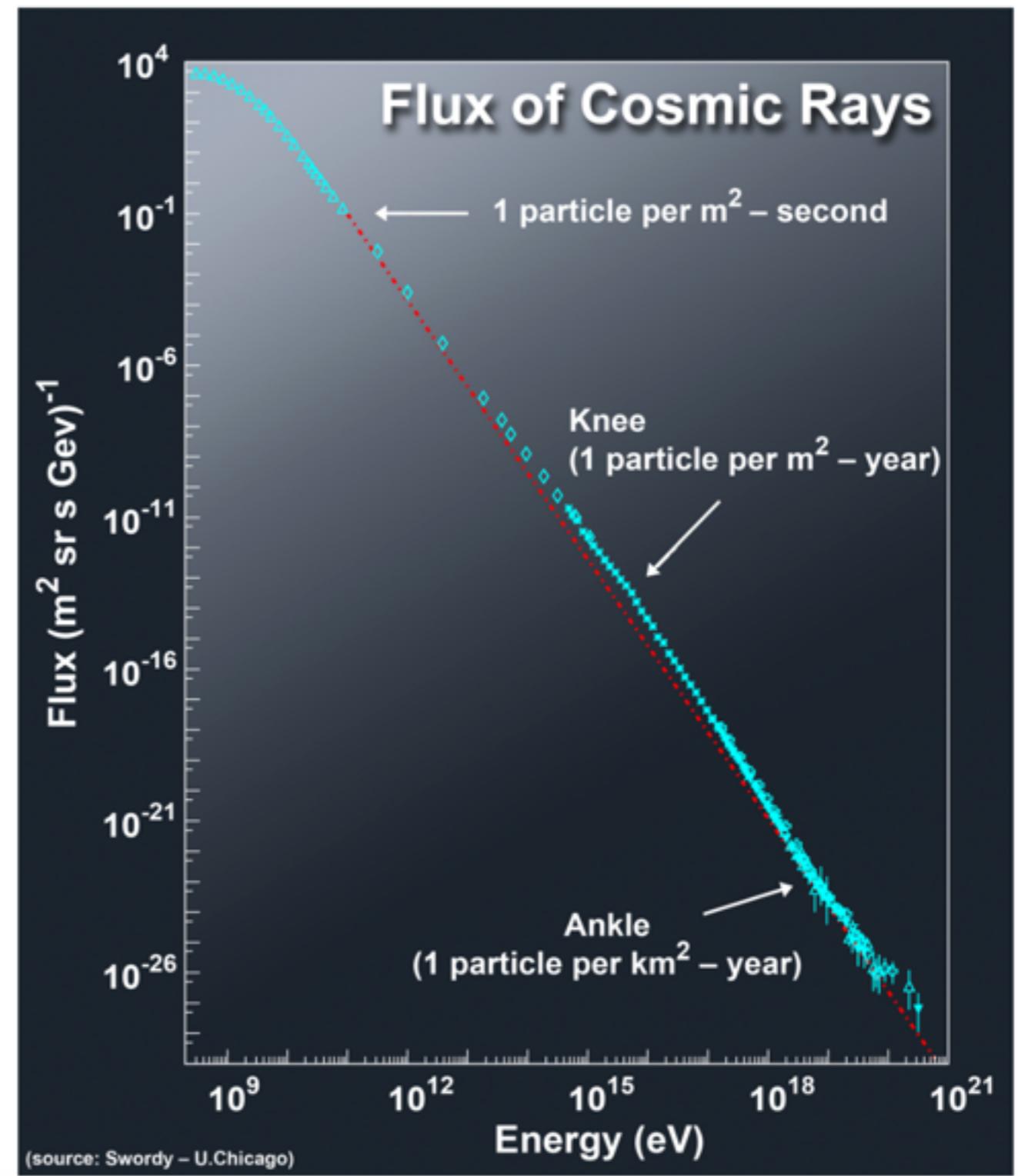
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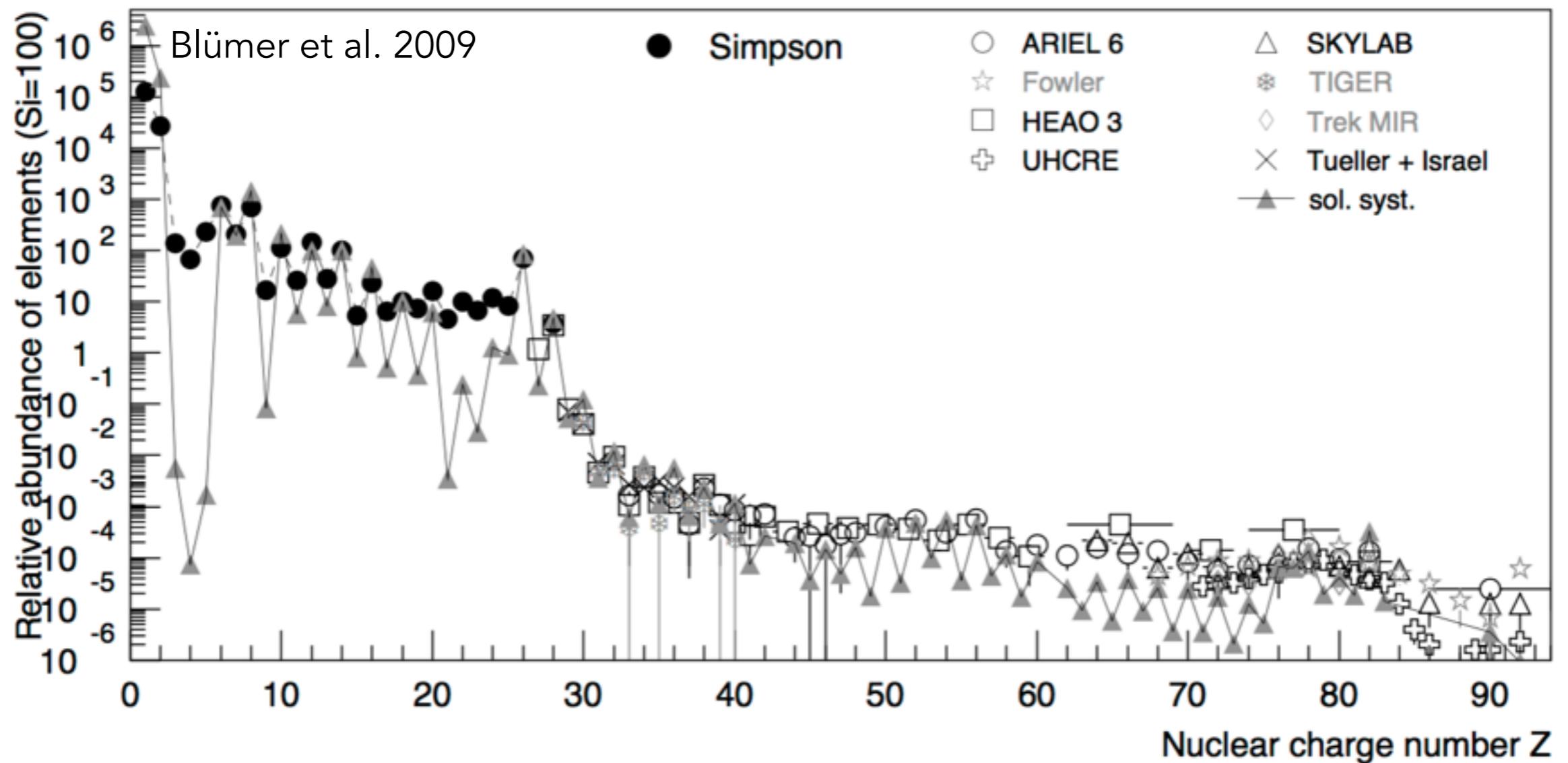
Cosmic Rays



<https://masterclass.icecube.wisc.edu/en/analyses/cosmic-ray-energy-spectrum>



Cosmic Rays

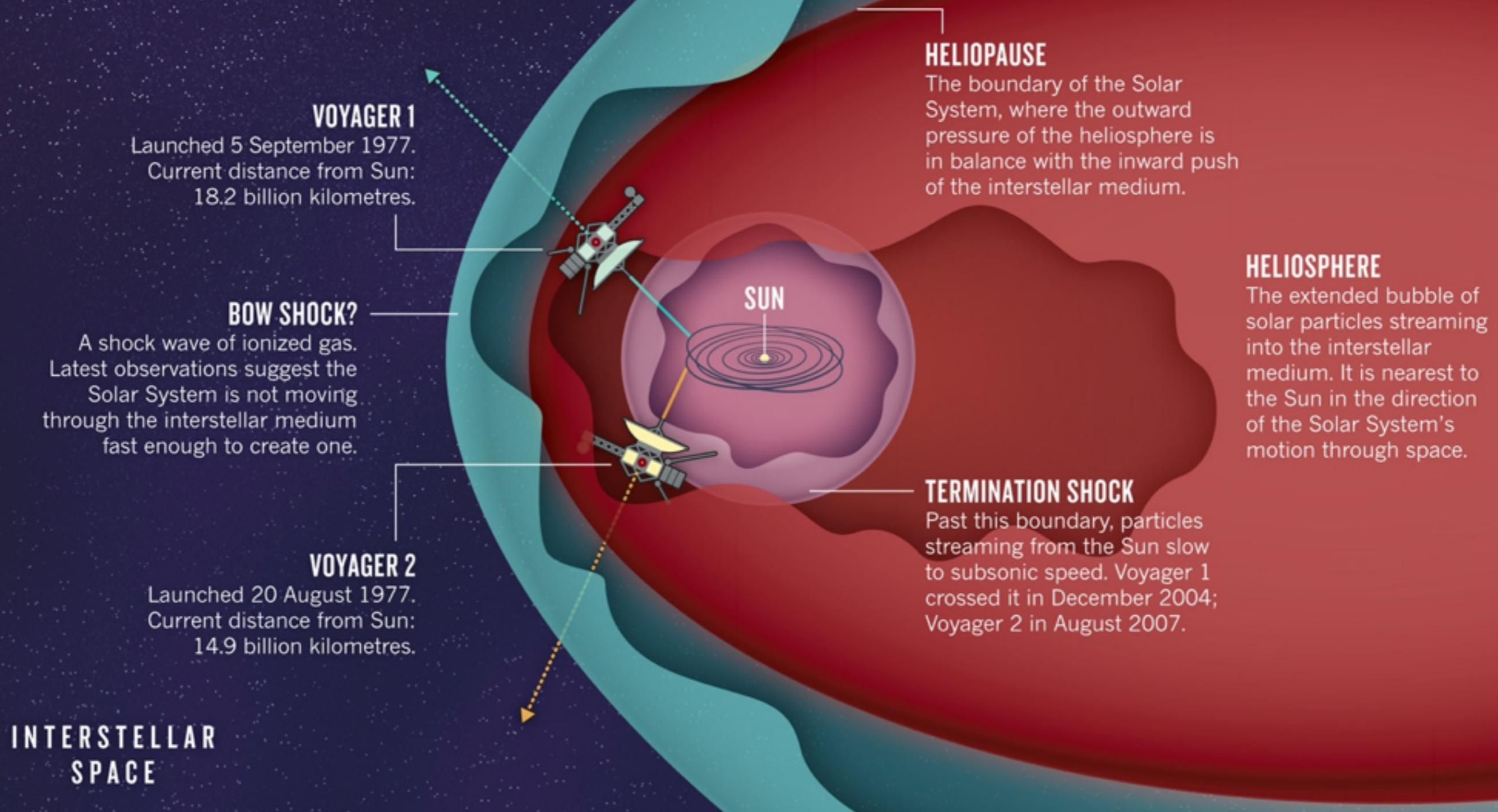


Protons are most abundant, He next.

${}^6\text{Li}$, ${}^9\text{Be}$, ${}^{10}\text{B}$, ${}^{11}\text{B}$ from “spallation”.

EDGING INTO THE UNKNOWN

After 35 years, the Voyager 1 spacecraft may finally be nearing the edge of the Solar System — the heliopause — but the probe's readings are proving difficult to interpret. Its sister craft, Voyager 2, is probably a few years away from reaching the milestone.



Cowen, *Nature*, 2012

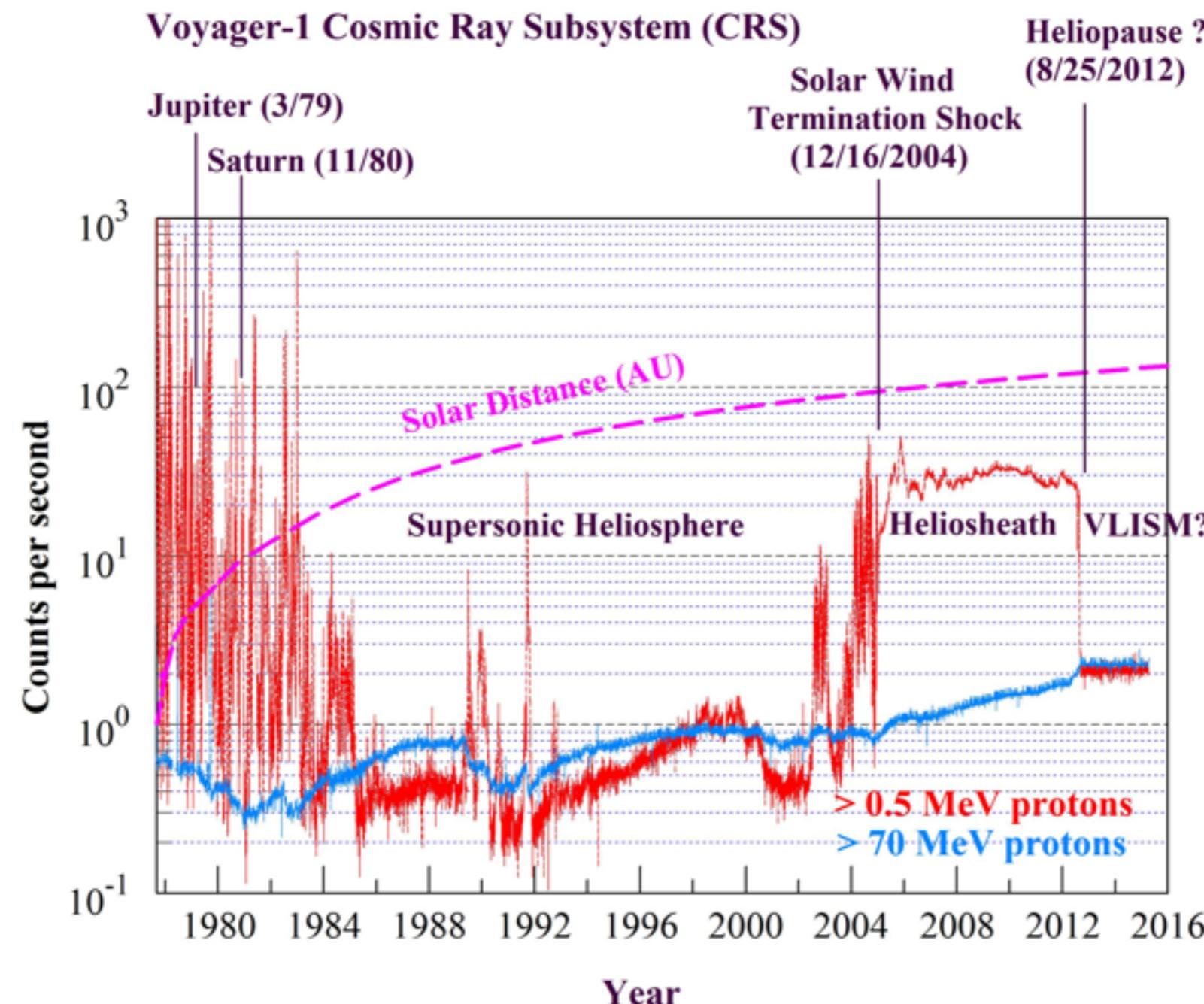
EDGING INTO THE UNKNOWN

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Launched
Curiosity

A shock wave
Latest observations
Solar System
through the interstellar
fast enough

Launched
Curiosity

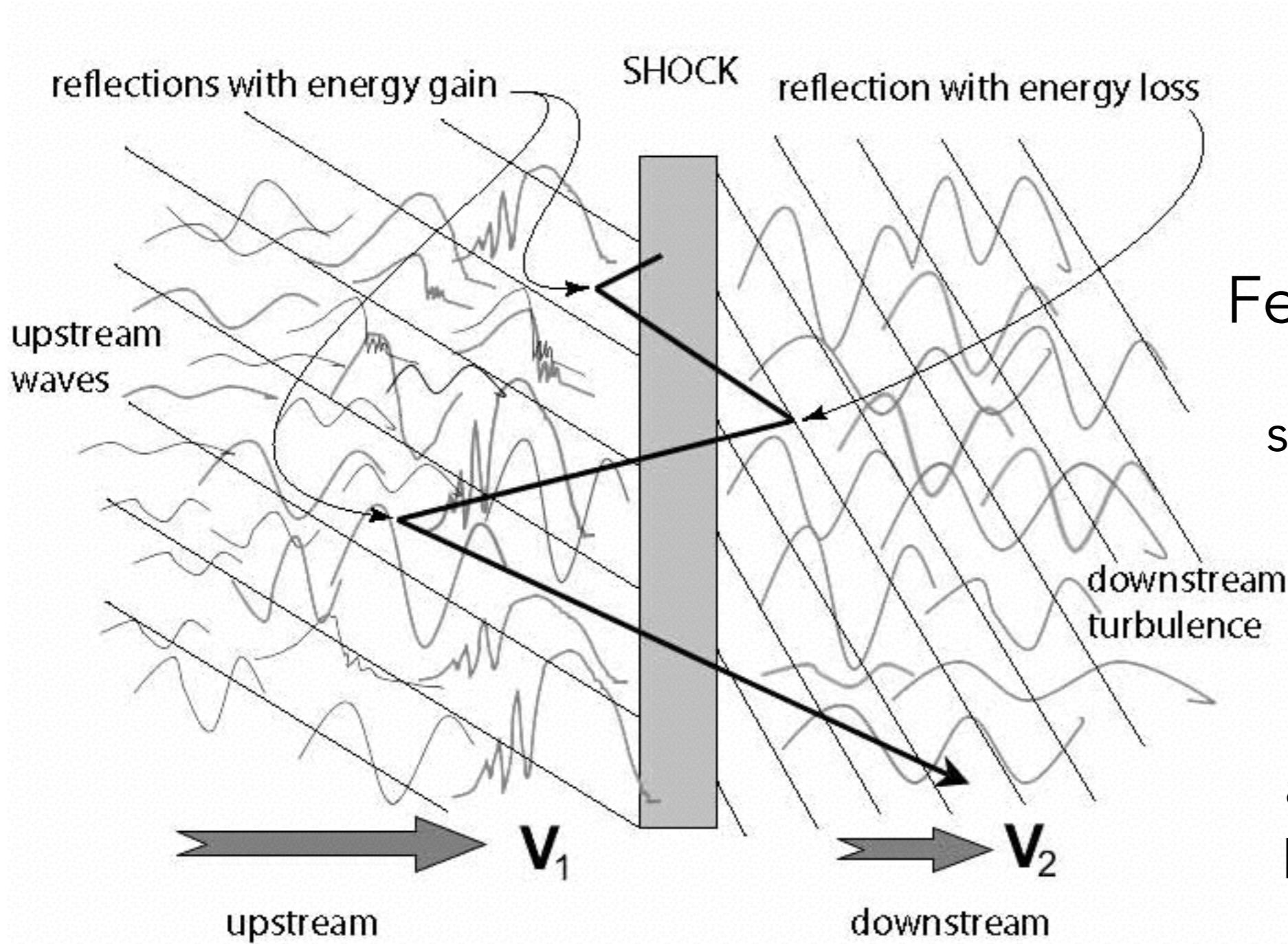


http://vepo.gsfc.nasa.gov/Voyager_heliopause.html

INTERSTELLAR
SPACE

Cowen, Nature, 2012

Acceleration of Cosmic Rays



Diffusive shock
acceleration
(first order
Fermi acceleration)
scattering off B-field
fluctuations
analogy to particle
bouncing between
converging walls

Acceleration of Cosmic Rays

Maximum energy attainable by diffusive shock acceleration is set by when B-fields can no longer confine CR.

gyroradius > scale of system

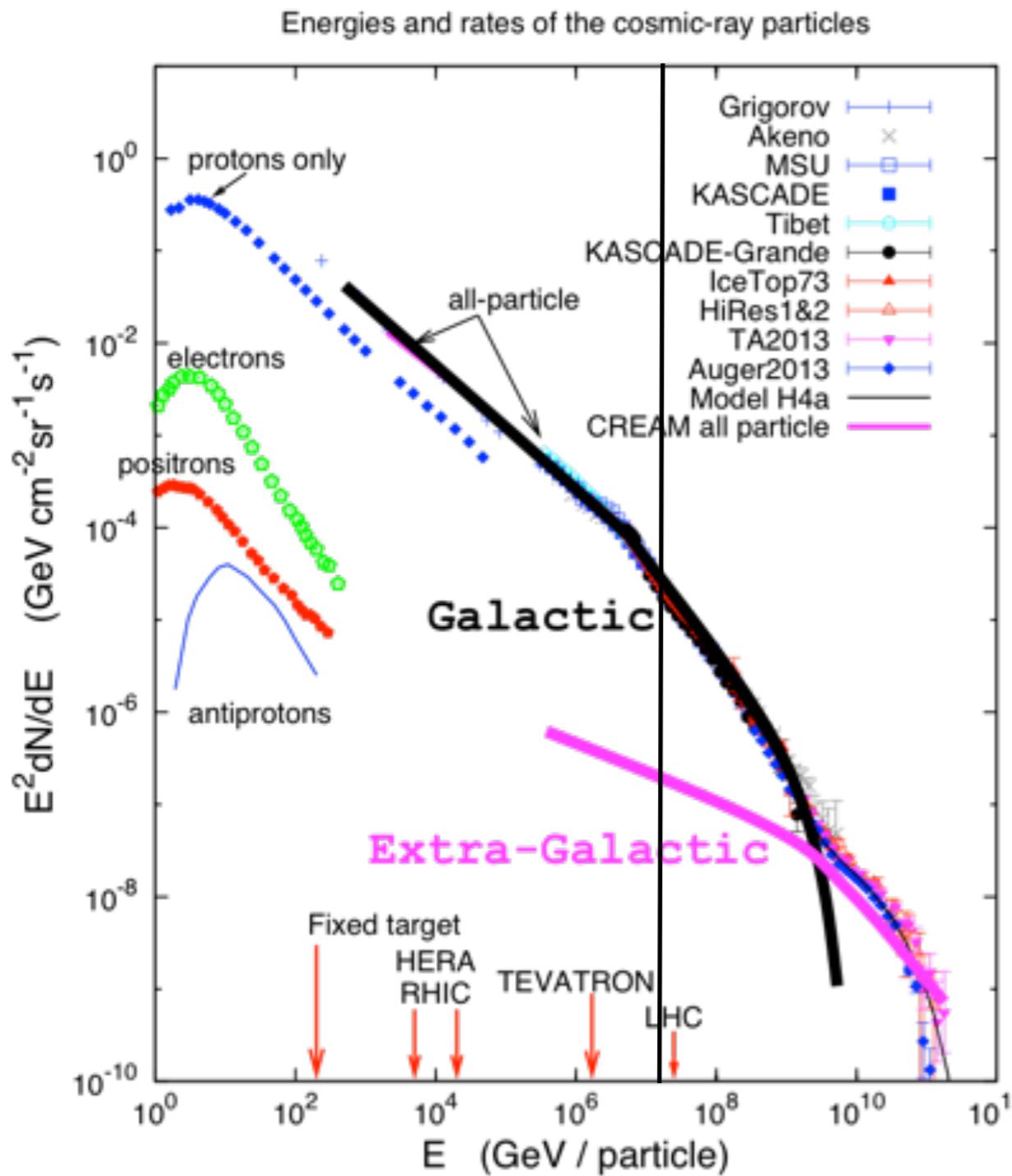
$$R_{\text{gyro}} = \frac{pc}{eB_{\perp}}$$

p = momentum

$$E_{\text{max}} = eB_{\text{SNR}}L$$

$$E_{\text{max}} \approx 10^{7.0} \text{ GeV} \left(\frac{L}{23 \text{ pc}} \right) \left(\frac{B_{\text{SNR}}}{10 \mu\text{G}} \right)$$

Acceleration of Cosmic Rays

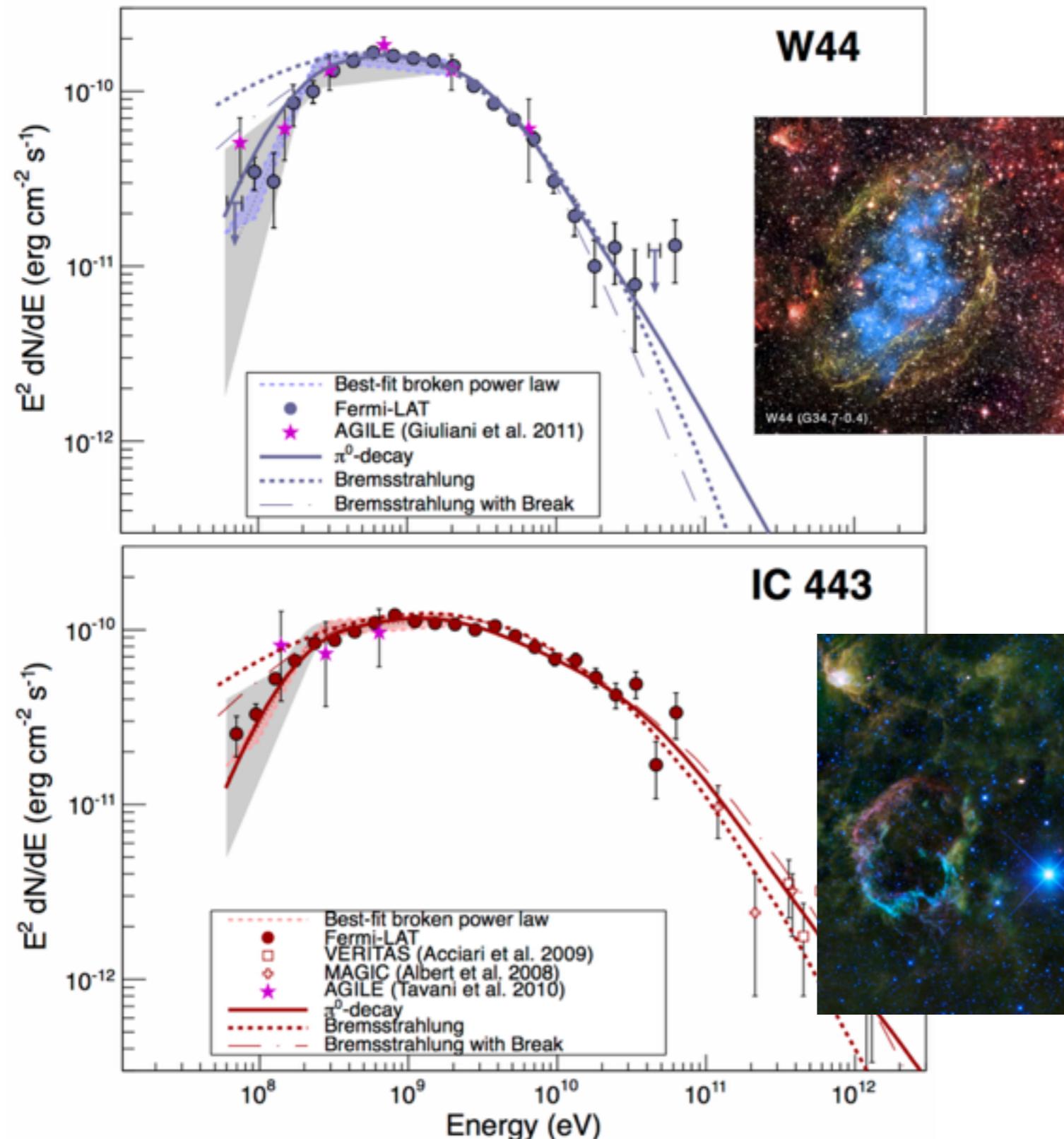


$$E_{\max} \approx 10^{7.0} \text{ GeV} \left(\frac{L}{23 \text{ pc}} \right) \left(\frac{B_{\text{SNR}}}{10 \mu\text{G}} \right)$$

Supernova shocks have long been thought to be the best candidate for CR acceleration.

Recently, first direct evidence...

Acceleration of Cosmic Rays



Accelerated protons create pions when they run into the surrounding ISM. Pions decay and produce gamma rays.

Fermi confirmation of gamma-ray spectrum following pion decay prediction for some SNRs in the MW.
(Ackermann et al. 2013)