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## Principle of Minimum Complementary Energy

- \* In analogy w/ the total potential energy we may define the complementary potential energy

$$\pi' = A' + \underline{\Phi}'$$

Potential at prescribed displacements

Complementary internal strain energy (strain energy)

Related to virtual quantities

$$\delta W_E' = \Delta \cdot \delta F$$

- Assume  $\Delta$  is given through a potential

$$\Delta = - \frac{\partial \underline{\Phi}'}{\partial F}$$

$$* \delta W_E' = - \frac{\partial \underline{\Phi}'}{\partial F} \delta F = - \delta \underline{\Phi}'$$

- Assume internal forces are conservative

$$* \delta W_I' = - \delta A'$$

Recall P.C.V.W

$$\delta W' = \delta W_E' + \delta W_I' = 0$$

$$= - \delta A' - \delta \underline{\Phi}' = 0$$

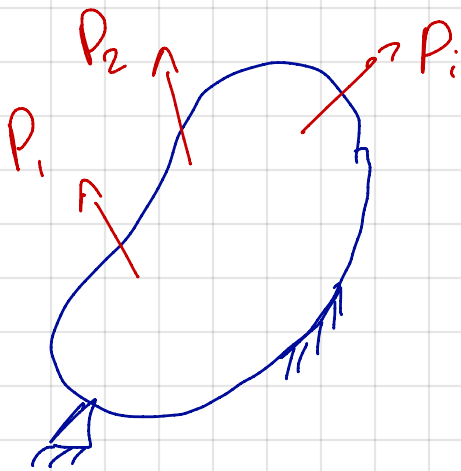
$$\delta W' = - \delta (A' + \underline{\Phi}') = 0$$

$$\rightarrow \boxed{\delta \pi' = 0} \xrightarrow{\pi'} \boxed{d\pi' = 0}$$

$$\rightarrow d\pi' = 0$$

- \* A conservative system undergoes compatible deformations if and only if the variation in the total complementary energy vanishes for all statically admissible forces.

### Castiglione's 1st theorem



$$\begin{aligned}\pi &= A + \bar{\Phi} \\ &= A - \sum_{i=1}^n P_i \cdot \Delta_i\end{aligned}$$

$$\begin{aligned}\frac{\partial \pi}{\partial \Delta_f} &= \frac{\partial A}{\partial \Delta_f} - \frac{\partial}{\partial \Delta_f} \sum_{i=1}^n P_i \Delta_i = 0 \\ &= \frac{\partial A}{\partial \Delta_f} - P_f = 0\end{aligned}$$

$$P_f = \frac{\partial A}{\partial \Delta_f}$$

- \* For an elastic body, the magnitude at the load at a point is equal to the derivative at  $A$  with respect to the projected displacement.

## Castigliano's 2nd Theorem

$$\Pi' = A' + \bar{\Phi}' = A' - \sum_{i=1}^N P_i \cdot \Delta_i$$

$$\frac{\partial \Pi'}{\partial P_j} = \frac{\partial A'}{\partial P_j} - \frac{\partial}{\partial P_j} \sum_{i=1}^N P_i \cdot \Delta_i$$

$$= \frac{\partial A'}{\partial P_j} - \Delta_j = 0$$

$$\Delta_j = \frac{\partial A'}{\partial P_j}$$

Linear Elastic

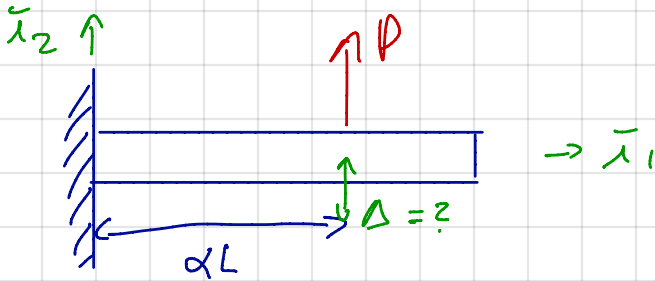
$$A = A'$$

$$\Delta_j = \frac{\partial A}{\partial P_j}$$

\* For an elastic body, the deflection at a point is given by the partial derivative of the complementary strain energy with respect to the driving force.

$$H_{23}^C = 0$$

Example



$$\eta = \frac{x_1}{L}$$

$$u_3 = \begin{cases} PL(\alpha - \eta) & 0 \leq \eta \leq \alpha \\ 0 & \alpha \leq \eta < 1 \end{cases}$$

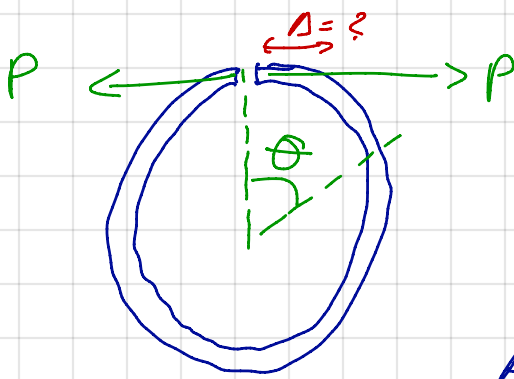
$$A' = \int_0^L \frac{1}{2} \frac{u_3^2}{H_{33}^C} dx_1$$

$$\eta = \frac{x_1}{L} \quad d\eta L = dx_1$$

$$A' = \frac{1}{2H_{33}^C} \int_0^\alpha (PL)^2 (\alpha - \eta)^2 \cdot L d\eta$$

$$A' = \frac{P^2 (\alpha L)^3}{6 H_{33}^C}$$

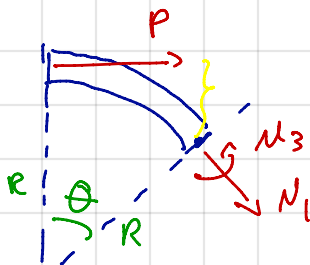
$$\Delta = \frac{\partial A'}{\partial P} = \frac{P(\alpha L)^3}{3 H_{33}^C}$$



Ring is subjected to both bending & axial loads

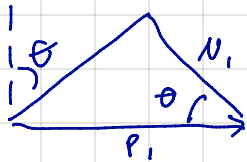
$$A' = \int_0^{2\pi} \frac{1}{2} \frac{M^2}{H_{33}^C} R d\theta + \int_0^{2\pi} \frac{1}{2} \frac{N^2}{S} R d\theta$$

$$\Delta = \frac{\partial A'}{\partial P} = \int_0^{2\pi} \frac{M}{H_{33}^C} \frac{\partial M}{\partial P} R d\theta + \int_0^{2\pi} \frac{N}{S} \frac{\partial N}{\partial P} R d\theta$$



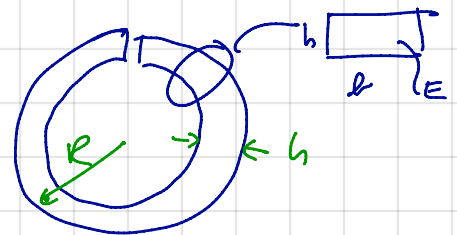
$$\Sigma U \rightarrow M_3 - P(R - R \cos(\theta)) = 0$$

$$\begin{cases} M_3 = PR(1 - \cos(\theta)) \\ N_1 = P \cos(\theta) \end{cases}$$



$$\Delta = \int_0^{2\pi} \frac{PR(1 - \cos(\theta)) R (1 - \cos(\theta))}{H_{33}^C} R d\theta + \int_0^{2\pi} \frac{P}{S} \cos(\theta) \cos(\theta) R d\theta$$

$$\Delta = \frac{PR^3}{H_{33}^C} (3\pi) + \frac{PR}{S} (\pi)$$



$$H_{33}^C = E \frac{bh^3}{12}, \quad S = Ebh$$

$$\Delta = \frac{36\pi}{E} \frac{PR^3}{bh^3} + \frac{\pi}{E} \frac{PR}{bh} = \frac{\pi}{E} \frac{PR}{bh} \left( \underbrace{36 \left( \frac{R}{h} \right)^2}_{\text{Bending}} + \underbrace{1}_{\text{Axial}} \right) = \Delta$$

\* For a thin ring,  $R \gg h$ , bending is the main contributor to displacement!

\* Stiffness is mainly due to axial extension!