

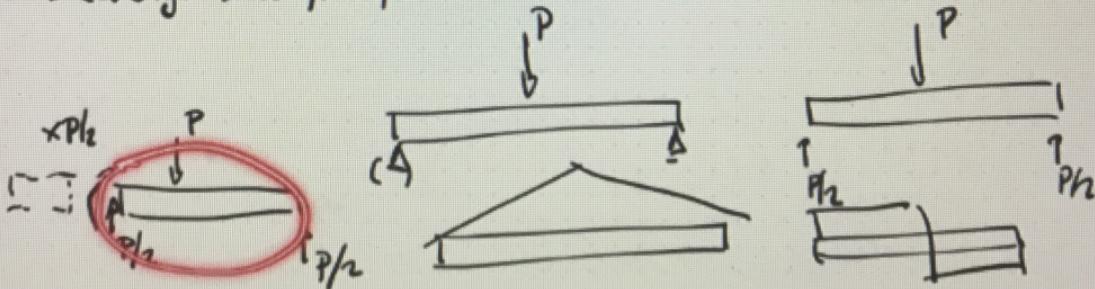
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Recap:

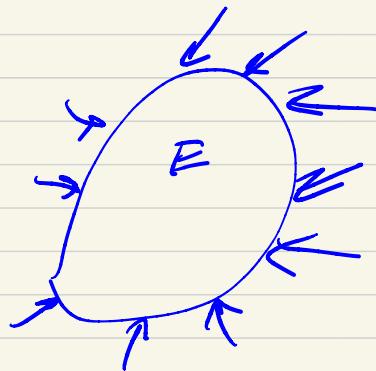
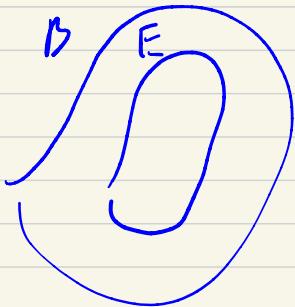
Balance of linear momentum

Recap:

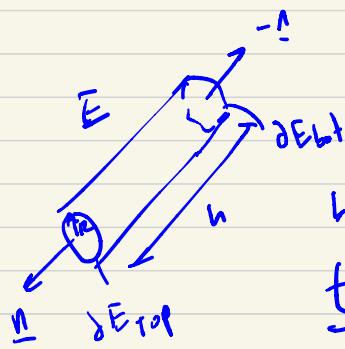
- * System of particles: $\frac{D}{Dt}(\underline{L}) = \underline{F}_{ext}$; $\underline{L} = \sum_{n=1}^N m_n^2 \dot{r}_n^{\perp}$; $\underline{F}_{ext} = \sum_{n=1}^N \underline{f}_n$
- * Continuum system: $\frac{D}{Dt} \int_B \underline{\rho} \underline{\ddot{v}} d\underline{r} = \int_B \underline{\underline{\sigma}} \underline{\dot{v}} d\underline{r} + \int_{\partial B} \underline{\underline{\tau}} \underline{v} dA$
- Using cons. of mass & kinematics: $\int_B \underline{\rho} \underline{\ddot{v}} d\underline{r} = \int_B \underline{\rho} \underline{\dot{v}} d\underline{r} + \int_{\partial B} \underline{\underline{\tau}} \underline{v} dA$
- Cauchy's stress principle



Cauchy stress principle:



$$\int_E \rho \dot{\underline{x}} dv = \int_E \rho b dv + \int_{\partial E} f dA$$

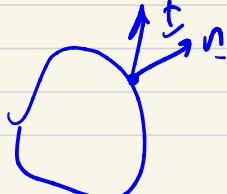


by limits,

$$\underline{t}(x) \Big|_{\partial E_{top}} = - \underline{f}(x) \Big|_{\partial E_{bkt}}$$

represent the
rest of the body or my
subbody E

(action = reaction) traction

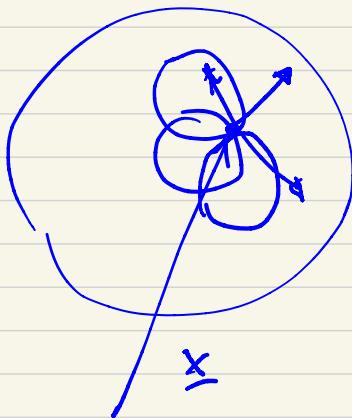


* As $R \rightarrow 0$, the only characteristic from the surface that remains is its normal.

Consequently, we expect $\underline{t}(x, n)$.

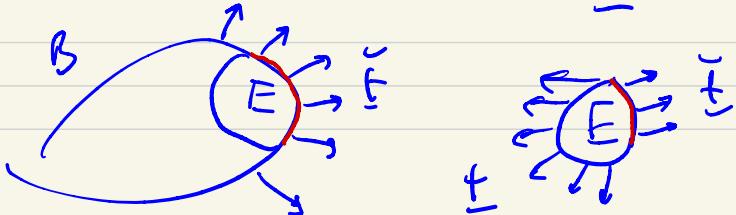
In this way, we can write:

$$\underline{t}(x, n) = -\underline{t}(x, -n)$$



* This means there are an infinite number of tractions at each point!

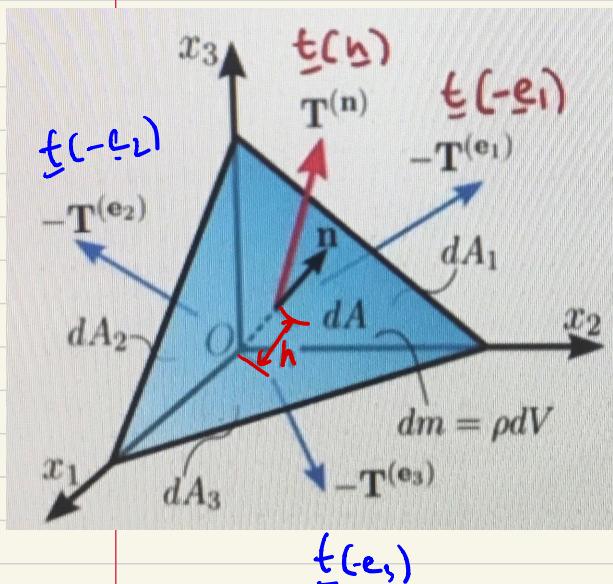
* The totality of these, called the stress state, characterizes the internal forces at x .



* If one side of our subbody coincides w/
an external surface, we can relate the
boundary conditions to the stress state:

$$\underline{f}(\underline{x}, \underline{n}) = \underline{\underline{t}}(\underline{x}) \text{ on } \partial B$$

(not ∂E)



Cauchy shear tensor

By simple geometric projection:

$$\Delta A_1 = \Delta A_n (\underline{n} \cdot \underline{e}_1) \Rightarrow \Delta A_n n_1$$

$$\Delta A_2 = \Delta A_n (\underline{n} \cdot \underline{e}_2) = \Delta A_n n_2$$

$$\Delta A_3 = \Delta A_n (\underline{n} \cdot \underline{e}_3) = \Delta A_n n_3$$

Balance of linear momentum for the tetrahedron is:

$$\int_T p(\underline{x} - \underline{b}) \, dv$$

$$= \int_{\partial T} \underline{f} \cdot \underline{dA} = \int_{\partial T_1} \underline{f} \cdot \underline{dA} + \int_{\partial T_2} \underline{f} \cdot \underline{dA} + \int_{\partial T_3} \underline{f} \cdot \underline{dA} + \int_{\partial T_n} \underline{f} \cdot \underline{dA}$$

By mean value theorem, there's a point
that the function gives the value of the integral



Applying the mean value theorem:

$$\int^*(x^* - \underline{x}^*) \left(\underbrace{\frac{1}{3} h \Delta A_n} \right)$$

ΔV

$$= \underline{f}^*(-\underline{c}_1) \Delta A_1 + \underline{f}^*(-\underline{c}_2) \Delta A_2 \\ + \underline{f}^*(-\underline{c}_3) \Delta A_3 + \underline{f}^*(1) \Delta A_n$$

Substituting all ΔA_i and dividing by ΔA_n
we get:

$$\frac{1}{3} \int^*(x^* - \underline{x}^*) = \underline{f}^*(-\underline{c}_1) n_1 + \underline{f}^*(-\underline{c}_2) n_2 \\ + \underline{f}^*(-\underline{c}_3) n_3 + \underline{f}^*(1)$$

Substituting $\underline{t}^*(\underline{\epsilon}_i) = -\underline{t}^*(\underline{\epsilon}_i)$

& taking $\lim_{h \rightarrow 0}$:

$$\boxed{\underline{t}(\underline{n}) = \underline{t}(\underline{\epsilon}_1)n_1 + \underline{t}(\underline{\epsilon}_2)n_2 + \underline{t}(\underline{\epsilon}_3)n_3}$$

(points collapsed to same point)

Note: $\underline{t}(\underline{\epsilon}_i) = \left\{ \begin{array}{l} t_1(\underline{\epsilon}_i) \\ t_2(\underline{\epsilon}_i) \\ t_3(\underline{\epsilon}_i) \end{array} \right\}$

- full vector \underline{t} , function of direction looking at
- not \underline{t} in direction of $\underline{\epsilon}_i$
- This is \underline{t} in the face of $\underline{\epsilon}_i$
(doesn't have to be perpendicular to face.)

full vector!

Let us define the vectors $\underline{f}_1 = \underline{t}(\underline{e}_1)$

$$\underline{f}_2 = \underline{t}(\underline{e}_2)$$

$$\underline{f}_3 = \underline{t}(\underline{e}_3)$$

The previous expression becomes:

$$\underline{t}(\underline{\sigma}) = \underline{f}_1 n_1 + \underline{f}_2 n_2 + \underline{f}_3 n_3$$

we can see that $\underline{t}(\underline{\sigma})$ is a vector w/ components

$$\begin{aligned} f_i(\underline{\sigma}) &= (\underline{f}_1)_i n_1 + (\underline{f}_2)_i n_2 + (\underline{f}_3)_i n_3 \\ &= \underbrace{(\underline{e}_i \cdot \underline{f}_1)}_{\delta_{i1}} n_1 + \underbrace{(\underline{e}_i \cdot \underline{f}_2)}_{\delta_{i2}} n_2 + \underbrace{(\underline{e}_i \cdot \underline{f}_3)}_{\delta_{i3}} n_3 \end{aligned}$$

$f_i(\underline{\sigma}) = \delta_{ij} n_j$

(Cauchy's relation)

we claim that δ_{ij} are the components of a rank 2 tensor $\underline{\Sigma}$ called Cauchy stress Tensor:

$$\underline{f}(\underline{\sigma}) = \underline{\Sigma} \cdot \underline{n}$$