Computational Physics – Lecture 15

(8 April 2020)

Recap

In the previous lecture, we defined the Fourier Transform, the Fourier Series, and the Discrete Fourier Transform.

Fourier Transform

Function \longleftrightarrow Function

$$\tilde{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx \cdot f(x) \cdot \exp(-ikx)$$
 (1)

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dk \cdot \tilde{f}(k) \cdot \exp(ikx)$$
 (2)



Fourier Series

Periodic Function \longleftrightarrow Numbers

$$\gamma_k = \frac{1}{\sqrt{2\pi}} \int_0^{2\pi} dx \cdot f(x) \cdot \exp(-ikx)$$
 (3)

$$f(x) = \frac{1}{\sqrt{2\pi}} \sum_{k=-\infty}^{\infty} \gamma_k \cdot \exp(ikx)$$
 (4)

where $f(x+2\pi) = f(x)$ and $k = \dots, -3, -2, -1, 0, 1, 2, 3, \dots$



$Numbers \longleftrightarrow Numbers$

$$\tilde{f}(k_q) = \frac{1}{\sqrt{n}} \sum_{p=0}^{n-1} f(x_p) \cdot \exp(-ik_q x_p)$$
 (5)

$$f(x_p) = \frac{1}{\sqrt{n}} \sum_{q=0}^{n-1} \tilde{f}(k_q) \cdot \exp(ik_q x_p)$$
 (6)

where
$$x_p = p\Delta$$
 for $p = 0, 1, \dots, n-1$ and

$$k_q = 2\pi q / n\Delta \text{ for } q = 0, 1, \dots, n-1$$



Today's plan

In today's class, we want to go deeper into the Discrete Fourier Transform.

nth roots of unity

We know that there n complex nth roots of unity, given by

$$\exp\left(\frac{i \cdot 2\pi \cdot k}{n}\right) \tag{7}$$

for k = 0, ..., n - 1.

nth roots of unity

We know that there n complex nth roots of unity, given by

$$\exp\left(\frac{i \cdot 2\pi \cdot k}{n}\right) \tag{7}$$

for k = 0, ..., n - 1.

Notice that

$$\left[\exp\left(\frac{i\cdot 2\pi\cdot k}{n}\right)\right]^n\tag{8}$$

$$= \exp\left(i \cdot 2\pi \cdot k\right) \tag{9}$$

$$= \cos(2\pi \cdot k) + i \cdot \sin(2\pi \cdot k) \tag{10}$$

$$= 1 + i \cdot 0 \tag{11}$$

$$=1 \tag{12}$$



nth roots of unity

For example, the four fourth roots of unity are given by

$$\exp\left(\frac{i \cdot 2\pi \cdot 0}{4}\right) = 1\tag{13}$$

$$\exp\left(\frac{i\cdot 2\pi\cdot 1}{4}\right) = e^{i\pi/2} \tag{14}$$

$$\exp\left(\frac{i\cdot 2\pi\cdot 2}{4}\right) = e^{i\pi} \tag{15}$$

$$\exp\left(\frac{i \cdot 2\pi \cdot 3}{4}\right) = e^{i3\pi/2} \tag{16}$$



Remarkably, the n nth roots of unity can be used to create a basis for the finite-dimensional complex linear vector space \mathbb{C}^n .

Remarkably, the n nth roots of unity can be used to create a basis for the finite-dimensional complex linear vector space \mathbb{C}^n .

(We are assuming the inner product

$$\langle \mathbf{a}, \mathbf{b} \rangle = \mathbf{a} \cdot \mathbf{b}^* \tag{17}$$

where $\mathbf{a}, \mathbf{b} \in \mathbb{C}^n$ and \mathbf{b}^* is the complex conjugate of \mathbf{b} .)



For n = 4, for instance, the four vectors

$$\begin{bmatrix}
\exp\left(\frac{i\cdot2\pi\cdot0}{4}\cdot0\right) \\
\exp\left(\frac{i\cdot2\pi\cdot0}{4}\cdot1\right) \\
\exp\left(\frac{i\cdot2\pi\cdot0}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot0}{4}\cdot3\right)
\end{bmatrix}, \begin{bmatrix}
\exp\left(\frac{i\cdot2\pi\cdot1}{4}\cdot0\right) \\
\exp\left(\frac{i\cdot2\pi\cdot1}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot1}{4}\cdot3\right)
\end{bmatrix}, \\
\begin{bmatrix}
\exp\left(\frac{i\cdot2\pi\cdot2}{4}\cdot0\right) \\
\exp\left(\frac{i\cdot2\pi\cdot2}{4}\cdot0\right) \\
\exp\left(\frac{i\cdot2\pi\cdot2}{4}\cdot1\right) \\
\exp\left(\frac{i\cdot2\pi\cdot2}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot2}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot2\right) \\
\exp\left(\frac{i\cdot2\pi\cdot3}{4}\cdot3\right)
\end{bmatrix}$$
(18)

form a basis for \mathbb{C}^4 .

Or in other words, the four vectors,

$$\begin{bmatrix} 1\\1\\1\\1 \end{bmatrix}, \begin{bmatrix} 1\\e^{i\pi/2}\\e^{i\pi}\\e^{i3\pi/2} \end{bmatrix}, \begin{bmatrix} 1\\e^{i\pi}\\1\\e^{i\pi} \end{bmatrix}, \begin{bmatrix} 1\\e^{i3\pi/2}\\e^{i\pi}\\e^{i\pi/2} \end{bmatrix}$$
(19)

form a basis for \mathbb{C}^4 .

In general, the n n-dimensional vectors

$$\mathbf{v}_k = \left[\exp\left(\frac{i \cdot 2\pi \cdot k}{n} \cdot m\right) \mid m = 0, 1, \dots, n - 1 \right]^T$$
 (20)

for k = 0, 1, ..., n - 1 form a basis for \mathbb{C}^n .



In general, the n n-dimensional vectors

$$\mathbf{v}_k = \left[\exp\left(\frac{i \cdot 2\pi \cdot k}{n} \cdot m\right) \mid m = 0, 1, \dots, n - 1 \right]^T$$
 (20)

for k = 0, 1, ..., n - 1 form a basis for \mathbb{C}^n .

(Here the index k labels the vectors and the index m labels the vector components.)

To show that the \mathbf{v}_k really do form a basis of \mathbb{C}^n , we first notice that these vectors are n in number.

To show that the \mathbf{v}_k really do form a basis of \mathbb{C}^n , we first notice that these vectors are n in number.

Now we show that these vectors are orthogonal.

If
$$k_1 = k_2$$

$$\langle \mathbf{v}_{k_1}, \mathbf{v}_{k_2} \rangle \tag{21}$$

$$= \mathbf{v}_{k_1} \cdot \mathbf{v}_{k_2}^* \tag{22}$$

$$= \sum_{m=0}^{n-1} v_{k_1,m} \cdot v_{k_2,m}^* \tag{23}$$

$$= \sum_{m=0}^{n-1} \exp\left(\frac{i \cdot 2\pi \cdot k_1}{n} \cdot m\right) \cdot \exp\left(\frac{-i \cdot 2\pi \cdot k_2}{n} \cdot m\right) \tag{24}$$

$$= \sum_{m=0}^{n-1} 1 \qquad \text{(for } k_1 = k_2\text{)} \tag{25}$$

$$= n > 0. (26)$$

And if $k_1 \neq k_2$

$$\langle \mathbf{v}_{k_{1}}, \mathbf{v}_{k_{2}} \rangle$$
 (27)
$$= \mathbf{v}_{k_{1}} \cdot \mathbf{v}_{k_{2}}^{*}$$
 (28)
$$= \sum_{m=0}^{n-1} v_{k_{1},m} \cdot v_{k_{2},m}^{*}$$
 (29)
$$= \sum_{m=0}^{n-1} \left[\exp\left(\frac{i \cdot 2\pi \cdot (k_{1} - k_{2})}{n}\right) \right]^{m}$$
 (30)
$$= \frac{1 - \left[\exp\left(\frac{i \cdot 2\pi \cdot (k_{1} - k_{2})}{n}\right) \right]^{n} }{1 - \left[\exp\left(\frac{i \cdot 2\pi \cdot (k_{1} - k_{2})}{n}\right) \right]}$$
 (31)
$$= \frac{1 - 1}{1 - \left[\exp\left(\frac{i \cdot 2\pi \cdot (k_{1} - k_{2})}{n}\right) \right]}$$
 (32)

So we have shown that

$$\langle \mathbf{v}_{k_1}, \mathbf{v}_{k_2} \rangle = n \delta_{k_1, k_2}. \tag{33}$$

So we have shown that

$$\langle \mathbf{v}_{k_1}, \mathbf{v}_{k_2} \rangle = n \delta_{k_1, k_2}. \tag{33}$$

The \mathbf{v}_k are orthogonal and are indeed a basis for \mathbb{C}^n .

So we have shown that

$$\langle \mathbf{v}_{k_1}, \mathbf{v}_{k_2} \rangle = n \delta_{k_1, k_2}. \tag{33}$$

The \mathbf{v}_k are orthogonal and are indeed a basis for \mathbb{C}^n .

That means, any vector $\mathbf{w} \in \mathbb{C}^n$ can now be expanded in this basis to get the components of \mathbf{w} as

$$\tilde{w}_q = \langle \mathbf{w}, \mathbf{v}_q \rangle \tag{34}$$

$$= \mathbf{w} \cdot \mathbf{v}_q^* \tag{35}$$

$$= \sum_{p=0}^{n-1} w_p \cdot v_{q,p}^* \tag{36}$$

$$= \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right) \tag{37}$$

So we have

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{38}$$

where the w_p are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in the standard basis and the \tilde{w}_q are a representation of the same vector in our new basis.



So we have

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{38}$$

where the w_p are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in the standard basis and the \tilde{w}_q are a representation of the same vector in our new basis.

This is the Discrete Fourier Transform.

So we have

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{38}$$

where the w_p are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in the standard basis and the \tilde{w}_q are a representation of the same vector in our new basis.

This is the Discrete Fourier Transform.

We have transformed n numbers w_p into n other numbers \tilde{w}_q .



We can also go from the new basis to the old basis.

We can also go from the new basis to the old basis.

The *n* basis vectors of the standard basis of \mathbb{C}^n are simply

$$\mathbf{u}_x = \left[\delta_{x,m} \mid m = 0, 1, \dots, n - 1\right]^T \tag{39}$$

for $x = 0, 1, \dots, n - 1$.



We can also go from the new basis to the old basis.

The *n* basis vectors of the standard basis of \mathbb{C}^n are simply

$$\mathbf{u}_x = \left[\delta_{x,m} \mid m = 0, 1, \dots, n - 1\right]^T \tag{39}$$

for $x = 0, 1, \dots, n - 1$.

For example, for n = 4, the Standard Basis of \mathbb{C}^4 is

$$\begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \\ 1 \end{bmatrix}$$
 (40)



Now any vector $\mathbf{w} \in \mathbb{C}^n$ can be expanded in the standard basis to get the components of \mathbf{w} as

$$w_p = \langle \mathbf{w}, \mathbf{u}_p \rangle \tag{41}$$

$$= \mathbf{w} \cdot \mathbf{u}_p^* \tag{42}$$



Now any vector $\mathbf{w} \in \mathbb{C}^n$ can be expanded in the standard basis to get the components of \mathbf{w} as

$$w_p = \langle \mathbf{w}, \mathbf{u}_p \rangle \tag{41}$$

$$= \mathbf{w} \cdot \mathbf{u}_p^* \tag{42}$$

But now we only know \mathbf{w} in the new basis.



Now any vector $\mathbf{w} \in \mathbb{C}^n$ can be expanded in the standard basis to get the components of \mathbf{w} as

$$w_p = \langle \mathbf{w}, \mathbf{u}_p \rangle \tag{41}$$

$$= \mathbf{w} \cdot \mathbf{u}_p^* \tag{42}$$

But now we only know \mathbf{w} in the new basis.

So while calculating these inner products we must use the new-basis representation of the standard basis vectors \mathbf{u}_p .



Now any vector $\mathbf{w} \in \mathbb{C}^n$ can be expanded in the standard basis to get the components of \mathbf{w} as

$$w_p = \langle \mathbf{w}, \mathbf{u}_p \rangle \tag{41}$$

$$= \mathbf{w} \cdot \mathbf{u}_p^* \tag{42}$$

But now we only know \mathbf{w} in the new basis.

So while calculating these inner products we must use the new-basis representation of the standard basis vectors \mathbf{u}_p .

What is the new-basis representation of the standard basis vectors?



The new-basis representation of the standard basis vectors is just their Discrete Fourier Transform, that is

$$\tilde{u}_{p,a} = \sum_{b=0}^{n-1} u_{p,b} \cdot \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot b\right) \tag{43}$$

$$= \sum_{b=0}^{n-1} \delta_{p,b} \cdot \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot b\right) \tag{44}$$

$$= \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot p\right) \tag{45}$$



The new-basis representation of the standard basis vectors is just their Discrete Fourier Transform, that is

$$\tilde{u}_{p,a} = \sum_{b=0}^{n-1} u_{p,b} \cdot \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot b\right) \tag{43}$$

$$= \sum_{b=0}^{n-1} \delta_{p,b} \cdot \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot b\right) \tag{44}$$

$$= \exp\left(\frac{-i \cdot 2\pi \cdot a}{n} \cdot p\right) \tag{45}$$

We can now use this representation of the standard basis vectors to calculate the w_p 's.



We have

$$w_p = \langle \mathbf{w}, \mathbf{u}_p \rangle \tag{46}$$

$$= \mathbf{w} \cdot \mathbf{u}_p^* \tag{47}$$

$$= \sum_{q=0}^{n-1} \tilde{w}_q \cdot \tilde{u}_{p,q}^* \tag{48}$$

$$= \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right) \tag{49}$$

So we have

$$w_p = \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right) \tag{50}$$

where the \tilde{w}_q are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in our new basis and the w_p are a representation of the same vector in the standard basis basis.



Inverse Discrete Fourier Transform

So we have

$$w_p = \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right) \tag{50}$$

where the \tilde{w}_q are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in our new basis and the w_p are a representation of the same vector in the standard basis basis.

This is the Inverse Discrete Fourier Transform.



Inverse Discrete Fourier Transform

So we have

$$w_p = \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right) \tag{50}$$

where the \tilde{w}_q are a representation of the vector $\mathbf{w} \in \mathbb{C}^n$ in our new basis and the w_p are a representation of the same vector in the standard basis basis.

This is the Inverse Discrete Fourier Transform.

We have transformed n numbers \tilde{w}_p into n other numbers w_q .



So the Discrete Fourier Transform is given by

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{51}$$

and its inverse is given by

$$w_p = \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right). \tag{52}$$



Now we would expect that if we calculated the DFT of n numbers and then calculated the inverse DFT we would recover the n numbers.

Now we would expect that if we calculated the DFT of n numbers and then calculated the inverse DFT we would recover the n numbers.

Let us check this:

$$\sum_{q=0}^{n-1} \tilde{w}_q \exp\left(\frac{i2\pi qp}{n}\right) \tag{53}$$

$$= \sum_{q=0}^{n-1} \left[\sum_{r=0}^{n-1} w_r \exp\left(\frac{-i2\pi qr}{n}\right) \right] \exp\left(\frac{i2\pi qp}{n}\right)$$
 (54)

$$= \sum_{q=0}^{n-1} \sum_{r=0}^{n-1} w_r \exp\left(\frac{i2\pi q(p-r)}{n}\right)$$
 (55)

$$=\sum_{r=0}^{n-1} w_r \cdot n \cdot \delta_{r,p} = n \cdot w_p \neq w_p. \tag{56}$$



So our DFT does not preserve norm. It is not normalised.

So our DFT does not preserve norm. It is not normalised.

We can make it normalised and unitary by doing

$$\tilde{w}_q = \frac{1}{\sqrt{n}} \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{57}$$

and the inverse is

$$w_p = \frac{1}{\sqrt{n}} \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right). \tag{58}$$



So our DFT does not preserve norm. It is not normalised.

We can make it normalised and unitary by doing

$$\tilde{w}_q = \frac{1}{\sqrt{n}} \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i \cdot 2\pi \cdot q}{n} \cdot p\right),\tag{57}$$

and the inverse is

$$w_p = \frac{1}{\sqrt{n}} \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i \cdot 2\pi \cdot q}{n} \cdot p\right). \tag{58}$$

Now we can freely go back and forth between the Fourier-space representation and the configuration-space representation.



In Numpy, the numpy.fft module contains various functions for Fourier analysis.

In Numpy, the numpy.fft module contains various functions for Fourier analysis.

The function numpy.fft.fft computes the DFT and the function numpy.fft.ifft computes the inverse DFT.

Numpy's DFT is not unitary, so that

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i2\pi qp}{n}\right),\tag{59}$$

and

$$w_p = \frac{1}{n} \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i2\pi qp}{n}\right). \tag{60}$$

Numpy's DFT is not unitary, so that

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i2\pi qp}{n}\right),\tag{59}$$

and

$$w_p = \frac{1}{n} \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i2\pi qp}{n}\right). \tag{60}$$

For example:

```
>>> import numpy as np
>>> w_p = (1,0,0,0)
>>> tw_q = np.fft.fft(w_p)
>>> tw_q
array([1.+0.j, 1.+0.j, 1.+0.j, 1.+0.j])
>>> np.fft.ifft(tw_q)
array([1.+0.j, 0.+0.j, 0.+0.j, 0.+0.j])
```



FFTW is a famous C library written by Matteo Frigo (MIT) and Steven Johnson (MIT).

FFTW is a famous C library written by Matteo Frigo (MIT) and Steven Johnson (MIT).

FFTW's DFT is not even normalised, so that

$$\tilde{w}_q = \sum_{p=0}^{n-1} w_p \cdot \exp\left(\frac{-i2\pi qp}{n}\right),\tag{61}$$

and

$$w_p = \sum_{q=0}^{n-1} \tilde{w}_q \cdot \exp\left(\frac{i2\pi qp}{n}\right). \tag{62}$$



```
#include <stdio.h>
#include <stdlib.h>
#include <fftw3.h>
void main () {
  fftw_complex w_p[4], tw_q[4];
  fftw_plan p;
  w_p[0][0] = 1.0; w_p[0][1] = 0.0;
  w_p[1][0] = 0.0; w_p[1][1] = 0.0;
  w_p[2][0] = 0.0; w_p[2][1] = 0.0;
  w_p[3][0] = 0.0; w_p[3][1] = 0.0;
  int n = 4:
  p = fftw_plan_dft_1d(n, w_p, tw_q,
      FFTW FORWARD, FFTW ESTIMATE):
```

```
fftw_execute(p);
int i;
for(i=0; i<4; i++) {
  printf("%f %f\n", tw_q[i][0], tw_q[i][1]);
}
printf("\n");
p = fftw_plan_dft_1d(n, tw_q, w_p,
    FFTW_BACKWARD, FFTW_ESTIMATE);
fftw_execute(p);
```

```
for(i=0; i<4; i++) {
    printf("%f %f\n", w_p[i][0], w_p[i][1]);
}

fftw_destroy_plan(p);
}</pre>
```

This code, when compiled as

```
$ gcc -o demo demo.c -lfftw3
and run, produces
```

```
$ ./demo
1.000000 0.000000
1.000000 0.000000
1.000000 0.000000
4.000000 0.000000
```

0.00000 0.00000 0.00000 0.00000 0.00000 0.00000

(Notice the unnormalised output.)



Activity

Compute the DFT of $[0,1,0,0]^T$ first manually and then using numpy.fft.fft. Then compute the inverse DFT of the result using numpy.fft.ifft.

We can also change our notation to think of the w's and \tilde{w} 's as functions of p and q (respectively).

We can also change our notation to think of the w's and \tilde{w} 's as functions of p and q (respectively).

So we can write

$$\tilde{f}(q) = \frac{1}{\sqrt{n}} \sum_{p=0}^{n-1} f(p) \cdot \exp\left(\frac{-i2\pi qp}{n}\right),\tag{63}$$

and the inverse as

$$f(p) = \frac{1}{\sqrt{n}} \sum_{q=0}^{n-1} \tilde{f}(q) \cdot \exp\left(\frac{i2\pi qp}{n}\right),\tag{64}$$

for p = 0, 1, ..., n - 1 and q = 0, 1, ..., n - 1.



We can also introduce new variables

$$x_p = p\Delta \text{ for } p = 0, 1, \dots, n-1,$$
 (65)

and

$$k_q = 2\pi q/n\Delta \text{ for } q = 0, 1, \dots, n-1.$$
 (66)

We can also introduce new variables

$$x_p = p\Delta \text{ for } p = 0, 1, \dots, n-1,$$
 (65)

and

$$k_q = 2\pi q/n\Delta \text{ for } q = 0, 1, \dots, n-1.$$
 (66)

so that

$$\tilde{f}(k_q) = \frac{1}{\sqrt{n}} \sum_{n=0}^{n-1} f(x_p) \cdot \exp(-ik_q x_p),$$
 (67)

and the inverse as

$$f(x_p) = \frac{1}{\sqrt{n}} \sum_{q=0}^{n-1} \tilde{f}(k_q) \cdot \exp(ik_q x_p),$$
 (68)

for p = 0, 1, ..., n - 1 and q = 0, 1, ..., n - 1.

