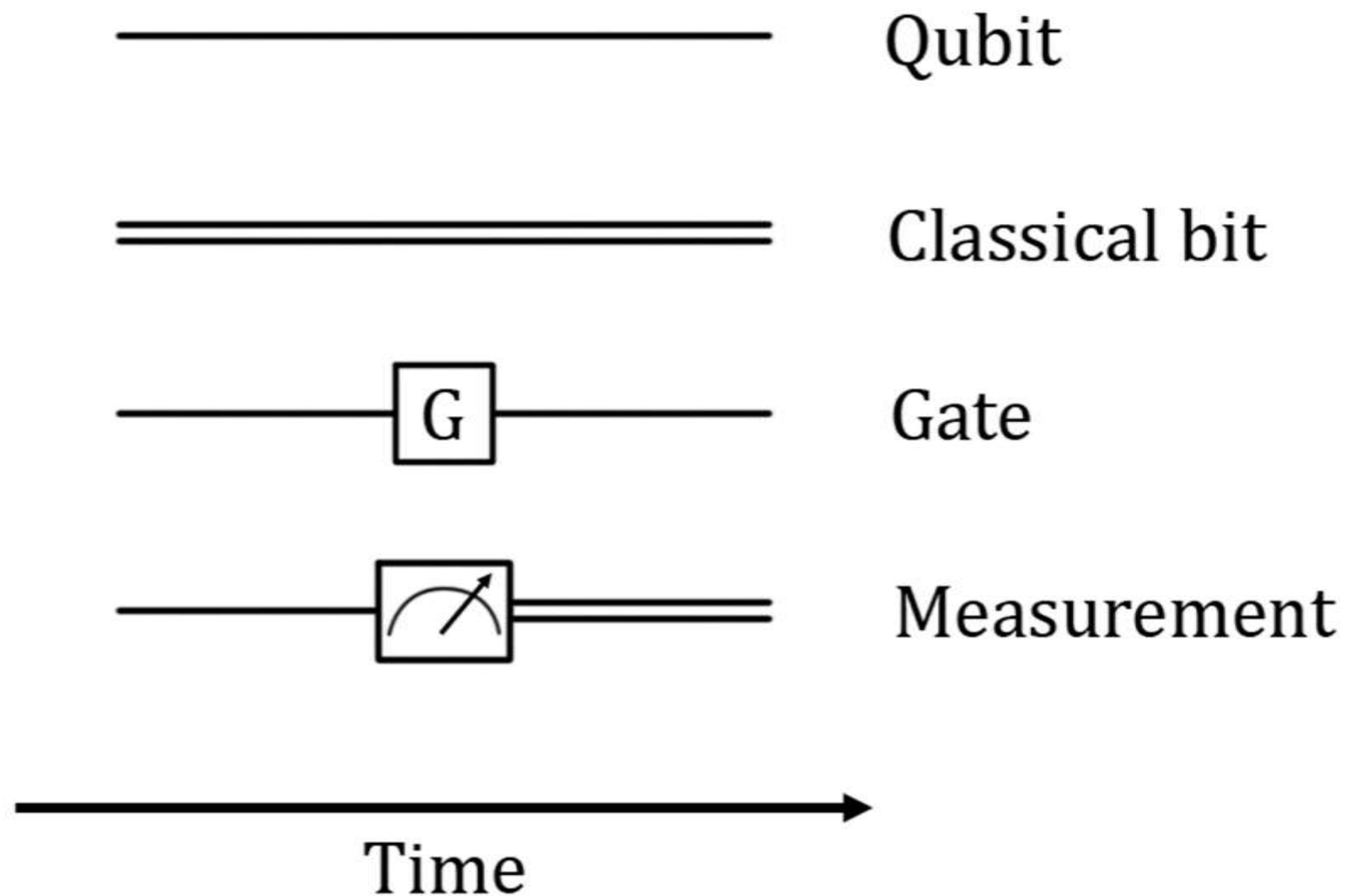


# Multi-Qubit Quantum Systems

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# Lecture outline

Multi-qubit quantum states

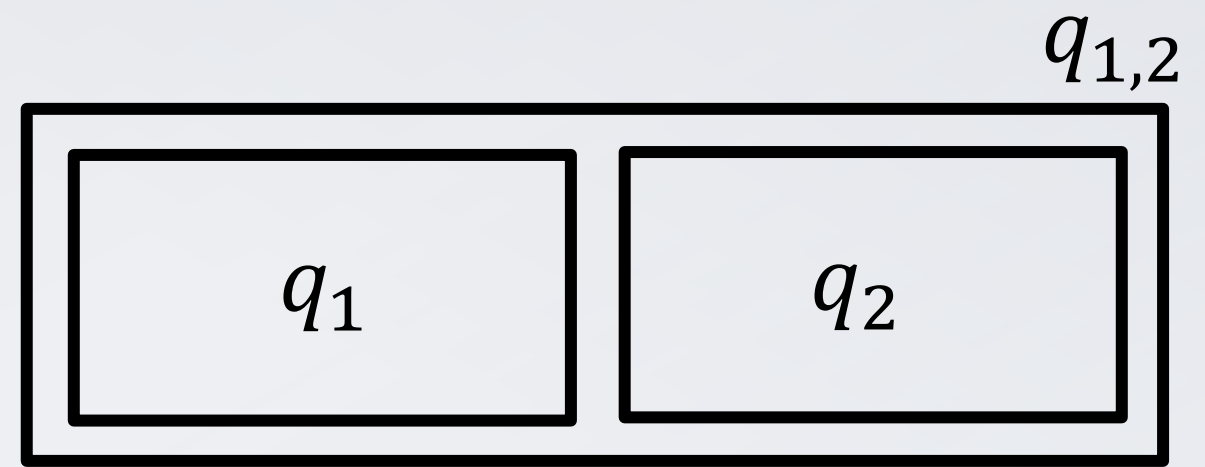
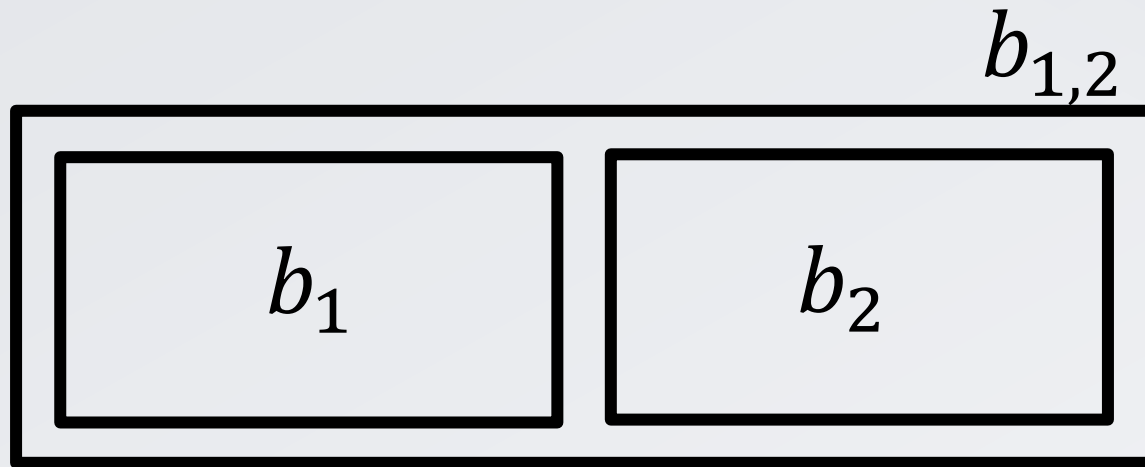
Multi-qubit quantum gates

Notations for quantum algorithms

Measuring multi-qubit systems

# Multi-qubit quantum states

## Scaling up: multiple qubits



$b_{1,2} = 00 \text{ or } 01 \text{ or } 10 \text{ or } 11$

$$\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \text{ or } \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} \begin{matrix} \text{--- "00"} \\ \text{--- "01"} \\ \text{--- "10"} \\ \text{--- "11"} \end{matrix}$$

$$q_{1,2} = \begin{pmatrix} c_{00} \\ c_{01} \\ c_{10} \\ c_{11} \end{pmatrix} \begin{matrix} \text{--- "00"} \\ \text{--- "01"} \\ \text{--- "10"} \\ \text{--- "11"} \end{matrix}$$

$$|c_{00}|^2 + |c_{01}|^2 + |c_{10}|^2 + |c_{11}|^2 = 1$$

# Dirac notation

“Ket” notation:  $|\cdot\rangle$  denotes a column vector

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}, |c\rangle = \begin{pmatrix} c_0 \\ c_1 \end{pmatrix} = c_0|0\rangle + c_1|1\rangle$$

More generally, let's consider an n-qubit quantum state:  
 $|\psi\rangle$  denotes a column vector of size  $2^n$ , where  $c_i \in \mathbb{C}$ :

$$|\psi\rangle = \begin{pmatrix} c_0 \\ c_1 \\ \vdots \\ c_{2^n-2} \\ c_{2^n-1} \end{pmatrix} = \sum_{i=0}^{2^n-1} c_i |i\rangle$$

We will sometimes denote the  $i$ -th basis state in integer form for compactness, rather than in binary

## Dirac notation: examples

$$|\psi\rangle = \begin{pmatrix} c_{00} \\ c_{01} \\ c_{10} \\ c_{11} \end{pmatrix} \begin{matrix} \text{-- "00"} \\ \text{-- "01"} \\ \text{-- "10"} \\ \text{-- "11"} \end{matrix} \quad \begin{aligned} &= c_{00} |00\rangle + c_{01} |01\rangle + c_{10} |10\rangle + c_{11} |11\rangle \\ &= c_0 |0\rangle + c_1 |1\rangle + c_2 |2\rangle + c_3 |3\rangle \end{aligned}$$

### Ket notation:

Popular shorthand notation for sparse column vectors.

$$|\varphi\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} |000\rangle + \frac{1}{\sqrt{2}} |110\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |6\rangle)$$

When working with integer basis states, double-check if the notation is big-endian or little-endian

# Tensor product of vectors

Denoted  $|p\rangle \otimes |q\rangle$ ,  $|p\rangle|q\rangle$ ,  $|pq\rangle$ , or  $|p, q\rangle$

## Input:

Vectors  $|p\rangle, |q\rangle$  with dimensions  $m, n$  respectively:  $|p\rangle = \begin{pmatrix} p_0 \\ \vdots \\ p_{m-1} \end{pmatrix}, |q\rangle = \begin{pmatrix} q_0 \\ \vdots \\ q_{n-1} \end{pmatrix}$

## Output:

Vector  $|p\rangle \otimes |q\rangle$  with dimension  $mn$ :

$$|p\rangle \otimes |q\rangle = \begin{pmatrix} p_0 q_0 \\ p_0 q_1 \\ \vdots \\ p_0 q_{n-1} \\ \vdots \\ p_{m-1} q_0 \\ \vdots \\ p_{m-1} q_{n-1} \end{pmatrix}$$

2 qubits cover a state space of dimension 4,  
3 qubits - dimension 8,  
 $n$  qubits - dimension  $2^n$

# Tensor product of vectors

## Examples:

$$|0\rangle \otimes |0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = |00\rangle = |0\rangle$$

$$|1\rangle \otimes |0\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} = |10\rangle = |2\rangle \quad \text{big-endian}$$

**A (tensor) product state (a.k.a. “separable state”):**

$$(a_0|0\rangle + a_1|1\rangle) \otimes (c_0|0\rangle + c_1|1\rangle) = a_0c_0|00\rangle + a_0c_1|01\rangle + a_1c_0|10\rangle + a_1c_1|11\rangle$$

In ket notation, just open the brackets to calculate the tensor product!



# Entanglement

## Unentangled

Alice

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Bob

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

## Entangled

Alice

Bob

cannot write separately  
(as a product state)

$$\frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$$

# Superposition vs entanglement

## Superposition is relative to the basis, entanglement is absolute

- A state  $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$  is in superposition with respect to the basis  $\{|0\rangle, |1\rangle\}$ , but is a basis state with respect to the basis  $\{|+\rangle, |-\rangle\}$
- Any state that is in superposition with respect to some basis is a basis state in some other basis
- A state that is entangled (cannot be represented as a tensor product of two states) is entangled in all bases

## Both superposition and entanglement are properties of quantum states, not specific states

There are lots of states in superposition and lots of entangled states

## States can be separable even if they're not "neatly" separable

State  $\frac{1}{\sqrt{2}}(|010\rangle + |111\rangle)$  is separable: qubit 2 in state  $|1\rangle$  and qubits 1 and 3 in state  $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ .

# Multi-qubit quantum gates

## Review: matrices and vectors

$$n \times m \text{ matrix } A = \begin{pmatrix} a_{00} & a_{01} & \cdots & a_{0,m-1} \\ a_{10} & a_{11} & \cdots & a_{1,m-1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n-1,0} & a_{n-1,1} & \cdots & a_{n-1,m-1} \end{pmatrix} \quad \text{Vector } x = \begin{pmatrix} x_0 \\ x_1 \\ \vdots \\ x_{m-1} \end{pmatrix}$$

Multiplying a vector by a matrix

$$\begin{pmatrix} a_{00} & a_{01} & \cdots & a_{0,m-1} \\ a_{10} & a_{11} & \cdots & a_{1,m-1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n-1,0} & a_{n-1,1} & \cdots & a_{n-1,m-1} \end{pmatrix} \begin{pmatrix} x_0 \\ x_1 \\ \vdots \\ x_{m-1} \end{pmatrix} = \begin{pmatrix} a_{00}x_0 + a_{01}x_1 + \cdots + a_{0,m-1}x_{m-1} \\ a_{10}x_0 + a_{11}x_1 + \cdots + a_{1,m-1}x_{m-1} \\ \vdots \\ a_{n-1,0}x_0 + \cdots + a_{n-1,m-1}x_{m-1} \end{pmatrix}$$

# Multi-qubit quantum gates

A **quantum gate** that acts on  $n$  qubits is a  $2^n \times 2^n$  unitary matrix  $U$ :

$$U = \begin{pmatrix} a_{00} & a_{01} & \cdots & a_{0,2^n-1} \\ a_{10} & a_{11} & \cdots & a_{1,2^n-1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{2^n-1,0} & a_{2^n-1,1} & \cdots & a_{2^n-1,2^n-1} \end{pmatrix}$$

*Reminder: unitary matrix means that  $U^\dagger U = U U^\dagger = I$*

$n$ -qubit state is a vector of size  $2^n$

To apply a gate to  $n$  qubits, multiply the state vector by the matrix

# Dirac notation: bra

“Bra” notation:  $\langle \cdot |$  denotes a row vector (adjoint of corresponding ket vector  $|\cdot\rangle$ )

$$\langle 0| = (1 \quad 0) = |0\rangle^\dagger$$

$$\langle 1| = (0 \quad 1) = |1\rangle^\dagger$$

$$\langle c| = (c_0 \quad c_1) = c_0\langle 0| + c_1\langle 1| = |c\rangle^\dagger$$

**Inner product** of vectors  $|\varphi\rangle$  and  $|\psi\rangle$ : **bra-ket**

$$\langle \varphi | \psi \rangle = (\varphi_0^* \quad \dots \quad \varphi_{n-1}^*) \begin{pmatrix} \psi_0 \\ \vdots \\ \psi_{n-1} \end{pmatrix} = \varphi_0^* \psi_0 + \dots + \varphi_{n-1}^* \psi_{n-1}$$

Bra-ket notation is used to denote orthogonal vectors (their bra-ket will be zero) and to calculate the probability of measurement outcomes (the probability of getting 0 when measuring a state  $|c\rangle$  is  $\langle 0|c\rangle$ ).

# Dirac notation: gates

Outer product of vectors  $|\varphi\rangle$  and  $|\psi\rangle$ : **ket-bra**

$$|\psi\rangle\langle\varphi| = \begin{pmatrix} \psi_0 \\ \vdots \\ \psi_{n-1} \end{pmatrix} (\varphi_0^* \quad \dots \quad \varphi_{n-1}^*) = \begin{pmatrix} \psi_0\varphi_0^* & \dots & \psi_0\varphi_{n-1}^* \\ \vdots & \ddots & \vdots \\ \psi_{n-1}\varphi_0^* & \dots & \psi_{n-1}\varphi_{n-1}^* \end{pmatrix}$$
$$|0\rangle\langle 1| = \begin{pmatrix} 1 \\ 0 \end{pmatrix} (0 \quad 1) = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

We can use ket-bra notation to write gates

$$\begin{pmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{pmatrix} = a_{00}|0\rangle\langle 0| + a_{01}|0\rangle\langle 1| + a_{10}|1\rangle\langle 0| + a_{11}|1\rangle\langle 1|$$

Convenient for sparse gates (gates with a lot of zero elements in the matrix)

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} = |1\rangle\langle 0| + |0\rangle\langle 1|$$

(read:  $|0\rangle$  becomes  $|1\rangle$ , and  $|1\rangle$  becomes  $|0\rangle$ )

# Controlled NOT gate

$$CX = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \quad CX|\psi\rangle = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \\ c_2 \\ c_3 \end{pmatrix} = \begin{pmatrix} c_0 \\ c_1 \\ c_3 \\ c_2 \end{pmatrix}$$

In ket-bra notation:

$$CX = |00\rangle\langle 00| + |01\rangle\langle 01| + |11\rangle\langle 10| + |10\rangle\langle 11|$$

$$CX = |0\rangle\langle 0| \otimes I + |1\rangle\langle 1| \otimes X$$

**“Controlled X” – the first qubit is “control”, the second qubit is “target”**

- Apply gate X on target qubit if control qubit is in  $|1\rangle$  state
- Otherwise, do nothing



# Controlled NOT gate

Quantum equivalent of classical XOR gate

$$|a, b\rangle \rightarrow |a, a \oplus b\rangle$$

$$00 \rightarrow 00$$

$$01 \rightarrow 01$$

$$10 \rightarrow 11$$

$$11 \rightarrow 10$$

Can entangle qubits (does not always!)

$$CX |+, 0\rangle = CX \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle) = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

$$CX |1, +\rangle = CX \frac{1}{\sqrt{2}} (|10\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (|11\rangle + |10\rangle) = |1, +\rangle$$

**“Phase kickback”:**

Apply the CX gate with control qubit in any superposition and target qubit in state  $|-\rangle$ .

Keeps qubits not entangled but propagates the  $-1$  phase to control in state  $|1\rangle$ .

$$CX |+, -\rangle = CX \frac{1}{2} (|00\rangle + |10\rangle - |01\rangle - |11\rangle) = \frac{1}{2} (|00\rangle + |11\rangle - |01\rangle - |10\rangle) = |-, -\rangle$$

# Endianness in multi-qubit gates

$$\begin{array}{c} \text{00} \\ \text{01} \\ \text{10} \\ \text{11} \end{array} \begin{array}{c} \text{00} \text{ 01} \text{ 10} \text{ 11} \\ \left( \begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{array} \right) \end{array}$$

First (most significant bit) as control

$$\begin{array}{c} \text{00} \\ \text{01} \\ \text{10} \\ \text{11} \end{array} \begin{array}{c} \text{00} \text{ 01} \text{ 10} \text{ 11} \\ \left( \begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{array} \right) \end{array}$$

Second (least significant bit) as control

# SWAP gate

$$SWAP = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$SWAP|\psi\rangle = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \\ c_2 \\ c_3 \end{pmatrix} = \begin{pmatrix} c_0 \\ c_2 \\ c_1 \\ c_3 \end{pmatrix}$$

In ket-bra notation:

$$SWAP = |00\rangle\langle 00| + |01\rangle\langle 10| + |10\rangle\langle 01| + |11\rangle\langle 11|$$

Swaps the states of the first and the second qubits

# Toffoli gate (double-controlled NOT)

**Toffoli gate** ( $CCX$ ) flips the state of the third qubit if and only if the first two qubits are both in the  $|1\rangle$  state (the quantum equivalent of AND gate)

$$|a, b, c\rangle \rightarrow |a, b, (a \wedge b) \oplus c\rangle$$

In ket-bra notation:

$$CCX = (I_2 - |11\rangle\langle 11|) \otimes I_1 + |11\rangle\langle 11| \otimes X$$

$$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix}$$

# Controlled gates

If you have a quantum gate, you can always define its **controlled** variant using one or several qubits as controls

The gate is applied if all control qubits are in the  $|1\rangle$  state, and nothing happens otherwise

$$C_n U_m = (I_n - |1 \dots 1\rangle\langle 1 \dots 1|) \otimes I_m + |1 \dots 1\rangle\langle 1 \dots 1| \otimes U_m$$

**Example: controlled Z gate**

$$CZ |+, +\rangle = CZ \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) = \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle - |11\rangle)$$

# Gates, controlled on patterns

You can also define controlled variants of gates with patterns other than  $|1 \dots 1\rangle$  as controls!

The gate is applied if the control qubits are in the given state  $|ctrl\rangle$ , and nothing happens otherwise

$$C_n U_m = (I_n - |ctrl\rangle\langle ctrl|) \otimes I_m + |ctrl\rangle\langle ctrl| \otimes U_m$$

**Example: zero-controlled Z gate**

$$C_0 Z |+, +\rangle = \frac{1}{2} (|00\rangle - |01\rangle + |10\rangle + |11\rangle)$$

**How to implement? (ControlledOnInt and ControlledOnBitstring functions in Q#)**

- Apply X gates to each qubit in the control register that is in the  $|0\rangle$  state in the  $|ctrl\rangle$  state
- Apply the regular controlled gate
- Apply X gates to each qubit in the control register that is in the  $|0\rangle$  state in the  $|ctrl\rangle$  state again

# Tricks for applying gates to states in Dirac notation

Apply a gate to the relevant qubits of each basis state independently, then regroup terms if needed

Example: Hadamard gate acting on  $\alpha|01\rangle + \beta|10\rangle$

$$H(\alpha|01\rangle + \beta|10\rangle) = \alpha(H|0\rangle)|1\rangle + \beta(H|1\rangle)|0\rangle = \frac{1}{\sqrt{2}}(\alpha|01\rangle + \alpha|11\rangle + \beta|00\rangle - \beta|10\rangle)$$

If the gate acts on multiple qubits, don't spell out the math for it, think about effect

Example: CNOT gate with 3<sup>rd</sup> qubit as control and 1<sup>st</sup> qubit as target

$$\begin{aligned} CNOT_{3,1} \frac{1}{2}(|000\rangle + |001\rangle + |110\rangle + |101\rangle) &= \\ &= \frac{1}{2}(|000\rangle + |101\rangle + |110\rangle + |001\rangle) \end{aligned}$$

Think of controlled gates as acting on basis states under classical conditions

Example: Controlled-on-zero Ry gate with 2<sup>nd</sup> qubit as control and 1<sup>st</sup> qubit as target

$$CRy_{2,1} \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle) = \frac{1}{\sqrt{2}}(|01\rangle + (Ry|1\rangle)|0\rangle)$$

# Universal sets of quantum gates

**What set of gates allows to express an arbitrary gate?**

- Any  $n$ -qubit gate can be represented using 2-qubit gates
- Any 2-qubit gate can be represented using CNOT and single-qubit gates (Krauss-Cirac decomposition)
- A limited set of single-qubit gates can **approximate** all single-qubit gates with arbitrary precision (unitary gate synthesis)

**Example: {H, T, CNOT} is a universal set**

Other sets exist; the choice of the best universal set depends on hardware architecture



# Notations for quantum algorithms

# Matrix notation

Quantum state on  $n$  qubits: a vector of  $2^n$  numbers

$$\begin{pmatrix} c_0 \\ \vdots \\ c_{2^n-1} \end{pmatrix}$$

Quantum gate acting on  $n$  qubits: a  $2^n \times 2^n$  matrix

$$\begin{pmatrix} a_{00} & a_{01} & \cdots & a_{0,2^n-1} \\ a_{10} & a_{11} & \cdots & a_{1,2^n-1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{2^n-1,0} & a_{2^n-1,1} & \cdots & a_{2^n-1,2^n-1} \end{pmatrix}$$

**Grows very inconvenient very quickly!**

# Dirac notation

Quantum state on  $n$  qubits: a sum of up to  $2^n$  ket vectors:

$$|\psi\rangle = \sum_{i=0}^{2^n-1} c_i |i\rangle$$

Quantum gate acting on  $n$  qubits a sum of up to  $4^n$  ket-bra terms:

$$A = \sum_{i,j=0}^{2^n-1} a_{ij} |i\rangle\langle j|$$

**Very convenient for sparse states and matrices and for orderly states and matrices (i.e., ones that allow to compress the sum into a formula)**

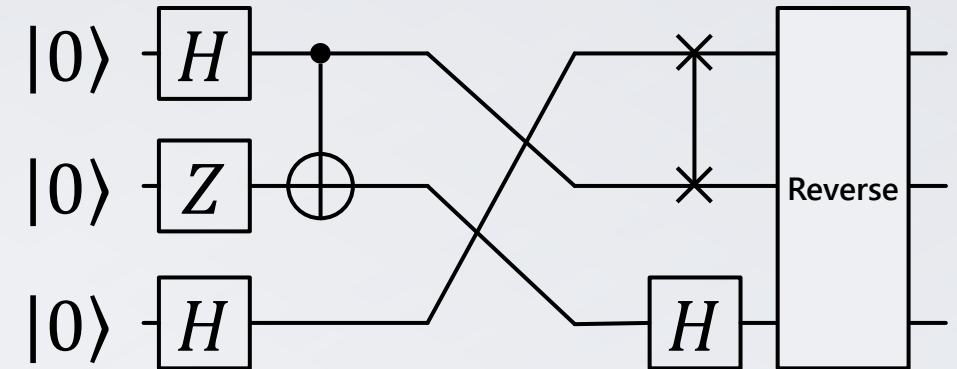
**Not very convenient for large dense notations**

# Circuit notation

Quantum state: a horizontal line (wire)

Quantum gate: an annotated box or symbol

Quantum measurement: an annotated box



Very popular in books and papers

Doesn't represent quantum states (closer to real life)

Supports "procedures" but not loops or classical conditions

Grows quite unreadable very fast!



# Quantum code (Q#)

Quantum state: hidden state of Qubit objects

Quantum gate: an operation on Qubit arrays

Quantum programming languages don't have native quantum state representation

But they can support loops, classical flow control, classical computations, and other capabilities

```
operation WState_PowerOfTwo_Reference (qs : Qubit[]) : Unit is Adj {  
    let N = Length(qs);  
    if (N == 1) {  
        // base of recursion: |1>  
        X(qs[0]);  
    } else {  
        let K = N / 2;  
        // create W state on the first K qubits  
        WState_PowerOfTwo_Reference(qs[0 .. K - 1]);  
        // the next K qubits are in |0...0> state  
        // allocate ancilla in |+> state  
        use anc = Qubit();  
        H(anc);  
  
        for i in 0 .. K - 1 {  
            Controlled SWAP([anc], (qs[i], qs[i + K]));  
        }  
        for i in K .. N - 1 {  
            CNOT(qs[i], anc);  
        }  
    }  
}
```

# Measuring multi-qubit systems

# Measuring multiple qubits

Consider a system of  $n$  qubits in state  $|q\rangle$ ; the  $2^n$  basis states of this system are  $|b_0\rangle, \dots, |b_{2^n-1}\rangle$ .

If we measure the system in the basis  $\{|b_0\rangle, \dots, |b_{2^n-1}\rangle\}$ ,

- we'll get outcome  $b_i$  with probability  $|\langle b_i | q \rangle|^2$
- and the system state will collapse to  $|b_i\rangle$ .

**Example:** measure  $\frac{1}{\sqrt{3}}(|00\rangle + |01\rangle + |10\rangle)$  in the computational basis.

- outcome 00 with probability  $|\langle 00 | q \rangle|^2 = \frac{1}{3}$
- outcome 01 with probability  $|\langle 01 | q \rangle|^2 = \frac{1}{3}$
- outcome 10 with probability  $|\langle 10 | q \rangle|^2 = \frac{1}{3}$

# Partial measurement

Consider a system of  $n$  qubits in the state  $|q\rangle$ ;  
measure the *first* qubit of this system in computational basis.

Same as measurements in Dirac notation: the basis of the measured part of the system is  $\{|b_i\rangle\}$

- We'll get outcome  $b_i$  with probability  $|\langle b_i | q \rangle|^2$ ,  
but this inner product will be a vector!

Proper inner product is defined for vectors of the matching dimensions, so you need to split the basis vectors of  $|q\rangle$  into tensor products of the first qubit and the rest of the system

For example,  $\langle 1 | 10 \rangle = \langle 1 |_1 | 1 \rangle_1 | 0 \rangle_2 = \langle 1 | 1 \rangle_1 | 0 \rangle_2 = | 0 \rangle_2$

- The system state will collapse to  $|b_i\rangle\langle b_i || q \rangle$ , renormalized,  
but this expression will be a vector on  $n$  qubits with norm 1



## Example: partial measurement

Consider the state  $|q\rangle = \frac{1}{\sqrt{3}}(|00\rangle + |01\rangle + |10\rangle)$ . Measure the first qubit of  $|q\rangle$  in the  $\{|0\rangle, |1\rangle\}$  basis.

$$\begin{aligned}\langle 0|q\rangle &= \frac{1}{\sqrt{3}}(\langle 0|_1|0\rangle_1|0\rangle_2 + \langle 0|_1|0\rangle_1|1\rangle_2 + \langle 0|_1|1\rangle_1|0\rangle_2) = \frac{1}{\sqrt{3}}(|0\rangle_2 + |1\rangle_2) \\ \langle 1|q\rangle &= \frac{1}{\sqrt{3}}(\langle 1|00\rangle + \langle 1|01\rangle + \langle 1|10\rangle) = \frac{1}{\sqrt{3}}|0\rangle_2\end{aligned}$$

Not renormalized yet!

Probability of measuring 0 is  $|\langle 0|q\rangle|^2 = \frac{2}{3}$ , and if we measure 0, the system collapses to

$$|0\rangle \langle 0|q\rangle = |0\rangle_1 \otimes \frac{1}{\sqrt{2}}(|0\rangle_2 + |1\rangle_2)$$

Probability of measuring 1 is  $|\langle 1|q\rangle|^2 = \frac{1}{3}$ , and if we measure 1, the system collapses to

$$|1\rangle \langle 1|q\rangle = |1\rangle_1 \otimes |0\rangle_2$$