



PHY-765 SS18 Gravitational Lensing Week 15

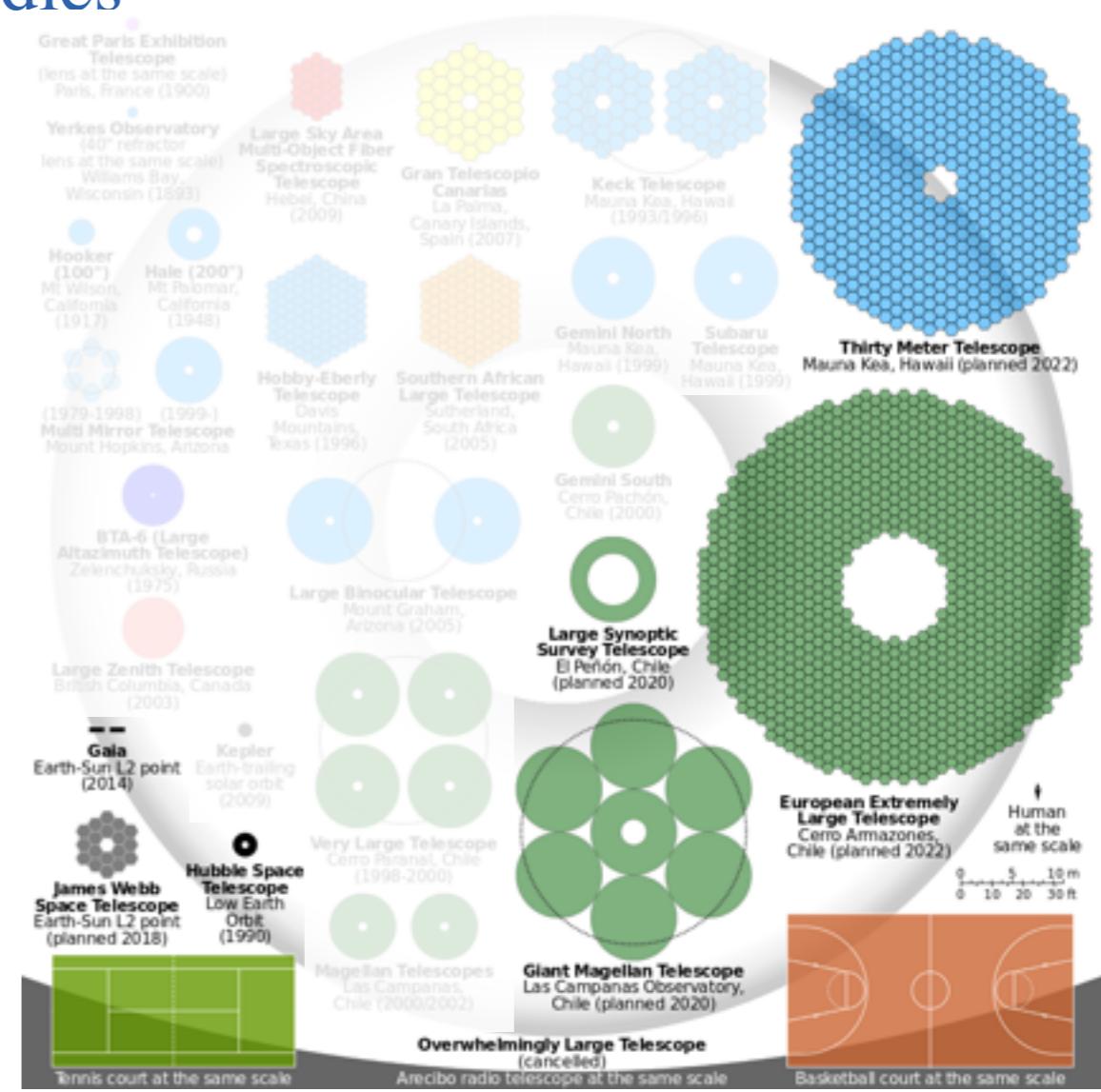
Course Summary
and
Q&A

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Last week

- (Incomplete) Overview of The Future of GL:
 - HST: Source follow-up and lensing clusters
 - OGLE/MicroFUN: Monitoring campaign of microlensing events
 - Gaia: Billions of points source; QSO lens ‘contaminants’
 - SDSS: Spec surveys incl. BAO studies
 - DES: Large-area imaging survey
 - LSST: Large-area imaging survey
 - JWST: Individual objects
 - WFIRST: Large-area survey
 - ELTs: “HST from the ground”

- Cluster lensing
- QSO lensing
- Galaxy-Galaxy lensing
- Star-Star microlensing
- Exoplanet searches with microlensing
- Wide-field weak lensing
- Power Spectrum lensing analysis



The aim of today

Summarize:

Course Topics &

(some of) The Course Essentials

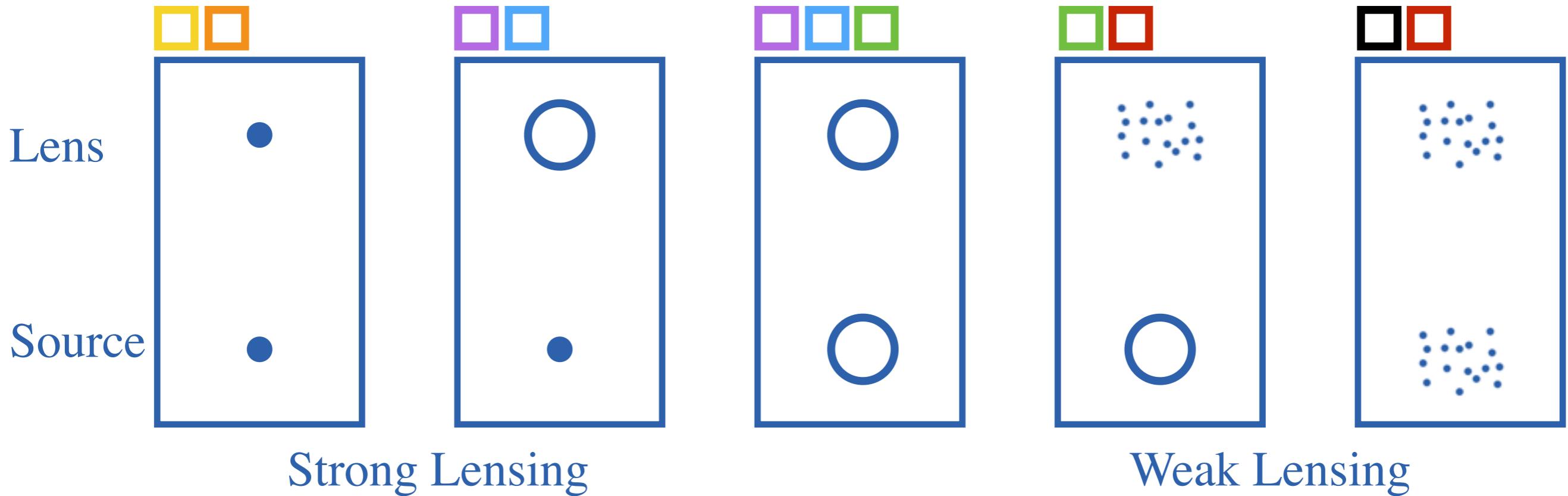
The Course Topics

Week 1: Introduction

- Cluster lensing
- QSO lensing
- Galaxy-Galaxy lensing
- Star-Star microlensing
- Exoplanet searches with microlensing
- Wide-field weak lensing
- Power Spectrum lensing analysis

Week 2: Basic Lens Geometry

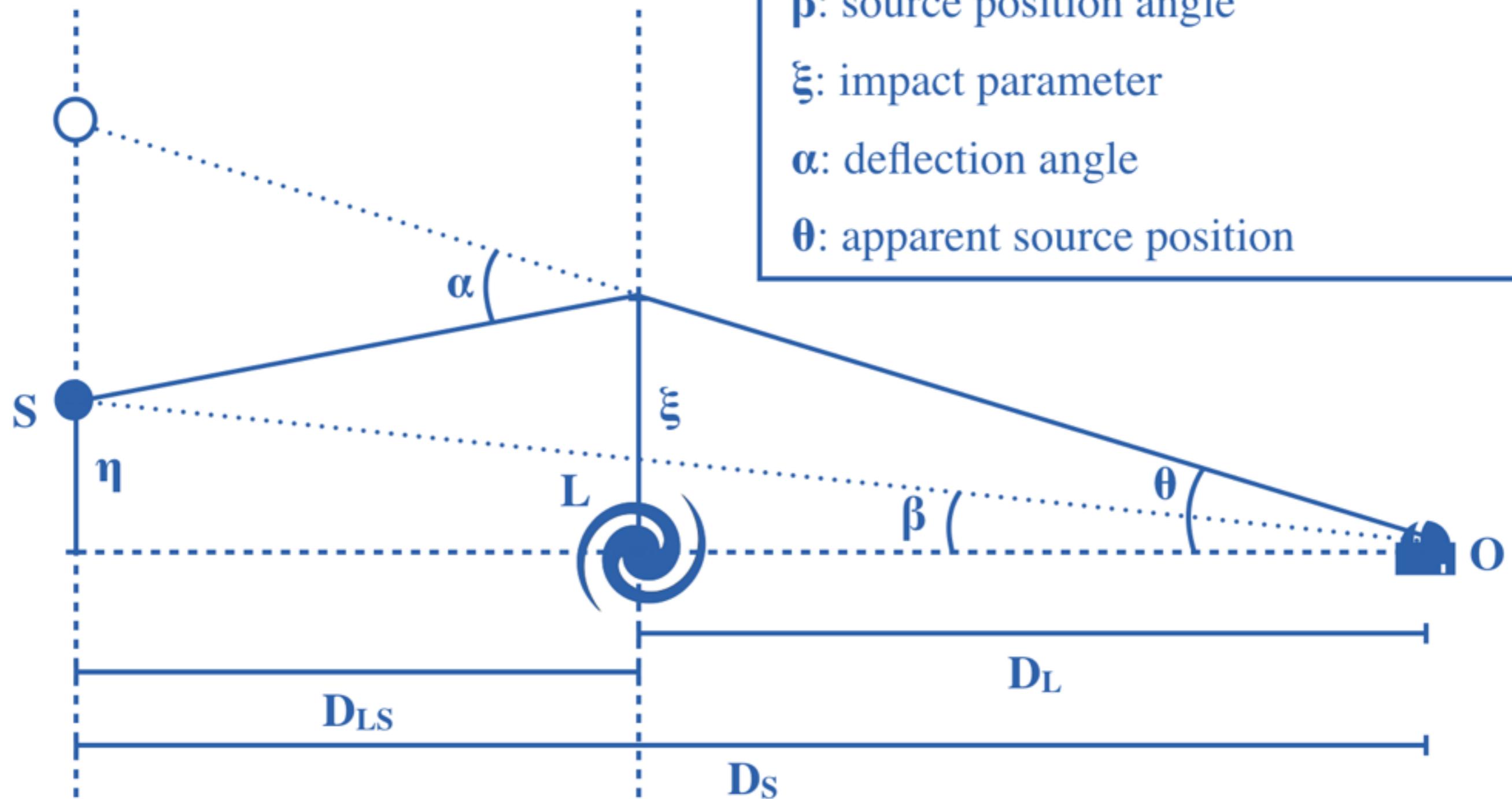
- Week: 4, 5, 6, 10, 14
- Week: 3, 4, 5, 6, 10, 14
- Week: 3, 4, 6, 10, 14
- Week: 7, 8, 14
- Week: 7, 8, 14
- Week: 6, 12, 13, 14
- Week: 12, 13, 14



Lens Geometry

Source
Plane

Lens
Plane



The Lens Equation

$$\beta = \theta - \alpha(\theta)$$

- Obtained from geometrical consideration of GL (deflection angles)
- Provides (non-linear) mapping from source plane to lens/image plane
- The deflection angle, α , is governed by the lens' surface mass distribution

$$\kappa(\theta) \equiv \frac{\Sigma(D_L\theta)}{\Sigma_{\text{cr}}}$$

$$\Sigma_{\text{cr}} \equiv \frac{c^2}{4\pi G} \frac{D_S}{D_L D_{LS}}$$

- The point mass lens (PML):

$$\beta = \theta - \frac{4MG D_{LS}}{c^2 D_S D_L} \frac{\theta}{|\theta|^2}$$

$$\theta_E \equiv \sqrt{\frac{4MG}{c^2} \frac{D_{LS}}{D_S D_L}}$$

LE Consequences 1: Multiple Images

- The PML:

$$\theta_{\pm} = \frac{\beta}{2} \left[1 \pm \sqrt{1 + \frac{4\theta_E^2}{\beta^2}} \right]$$



- Considered Isothermal Sphere (IS) - both spherical and cored version

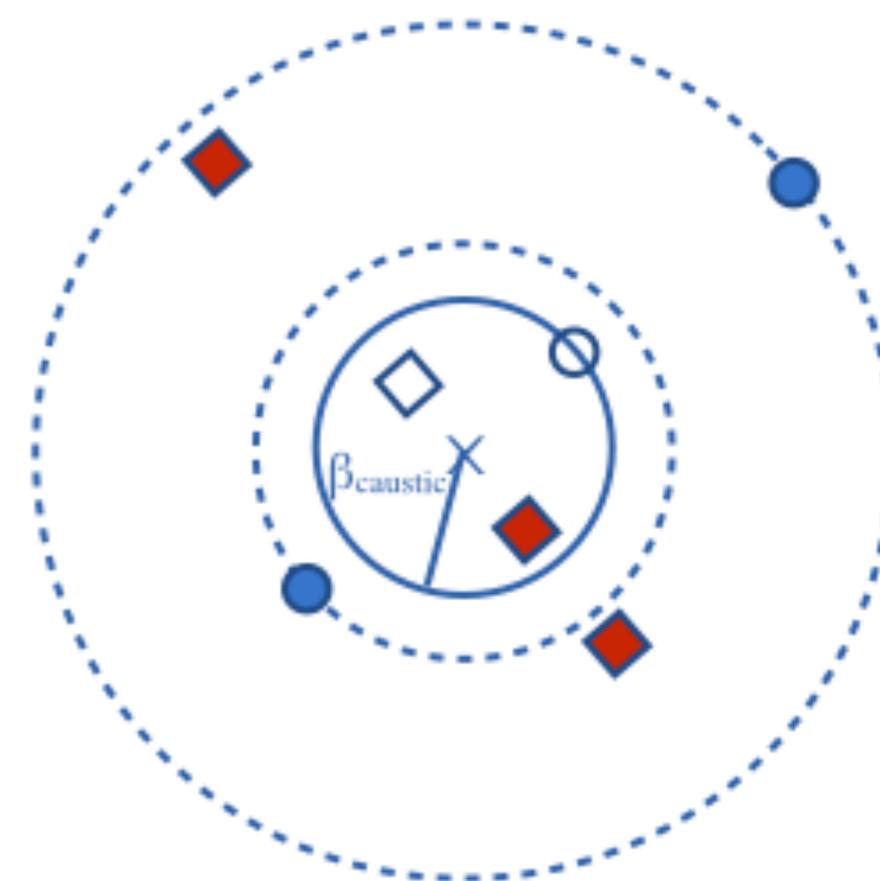
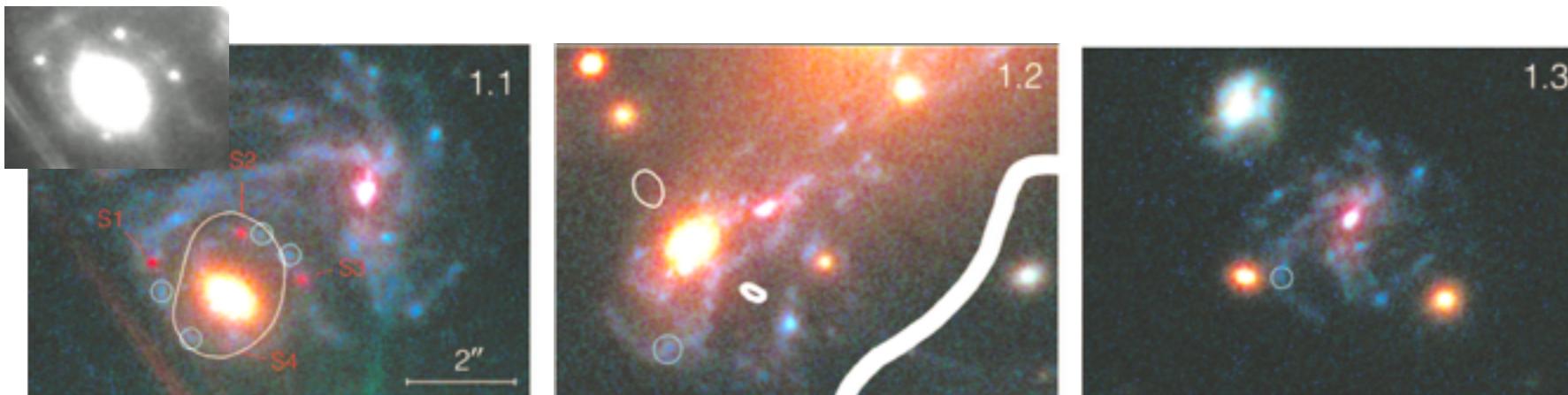
$$\beta = \theta - \frac{\theta_0}{\theta^2} \left[\sqrt{\theta^2 + \theta_{\text{core}}^2} - \theta_{\text{core}} \right] \theta$$

$$\theta_E = \theta_0 \sqrt{1 - 2 \frac{\theta_{\text{core}}}{\theta_0}}$$

- Defined caustics and critical curves

- Critical curves are where images fall if source is on the caustic
- Where pairs of images are created/destroyed
- Caustic is where multiple images appear/disappear

- SN Refsdal: Multiple Images at it's best



LE Consequences 2: Time Delays

- The arrival (travel) time of light from multiple images differ due to
 - change in the gravitational potential (The Shapiro time delay)
 - geometry as the light travels along different paths

$$\Delta t = \frac{D_L D_S}{c D_{LS}} \left[\frac{(\theta - \beta)^2}{2} - \frac{\Phi(\theta)}{c^2} \right]$$

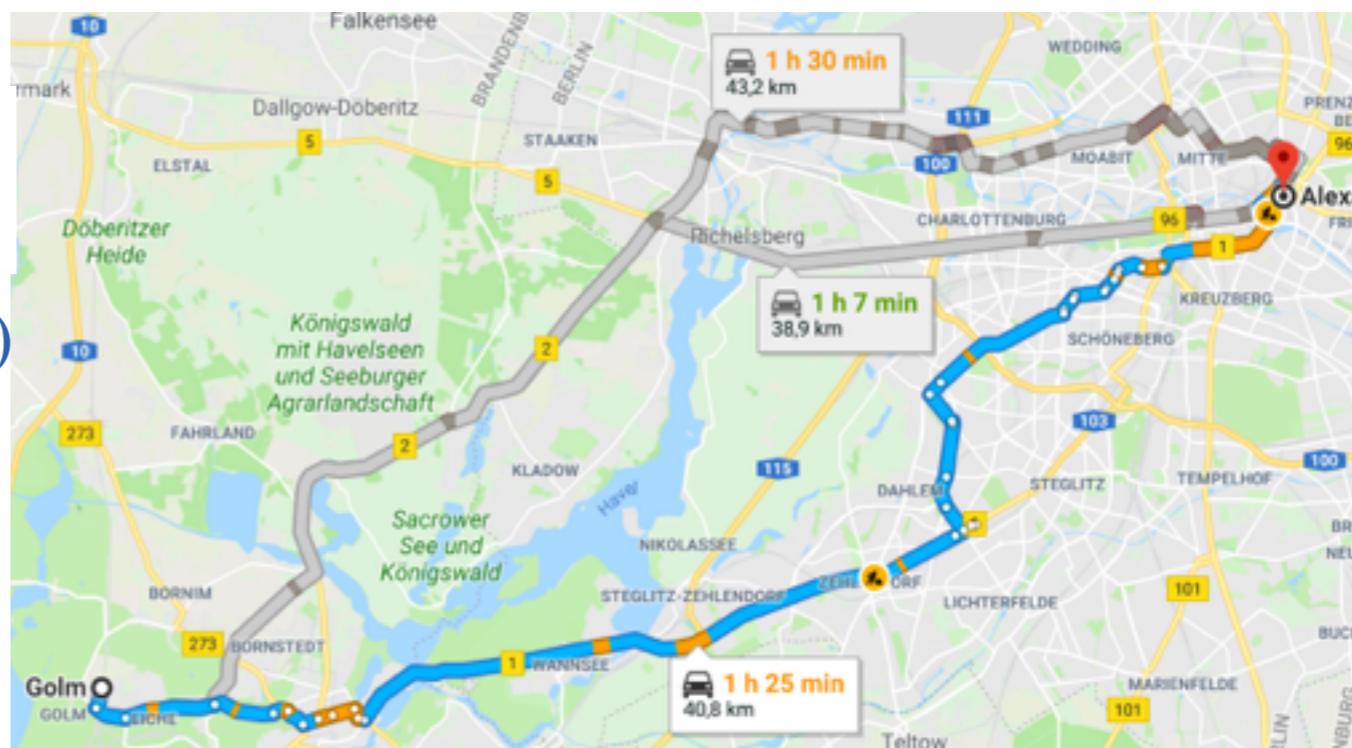
Only depends on distances; no lens details

Only depends on lens mass distribution

- For the PML:

$$t_+ - t_- \simeq -(1 + z_L) \frac{D_L D_S}{c D_{LS}} 2c^2 \theta_E \beta$$

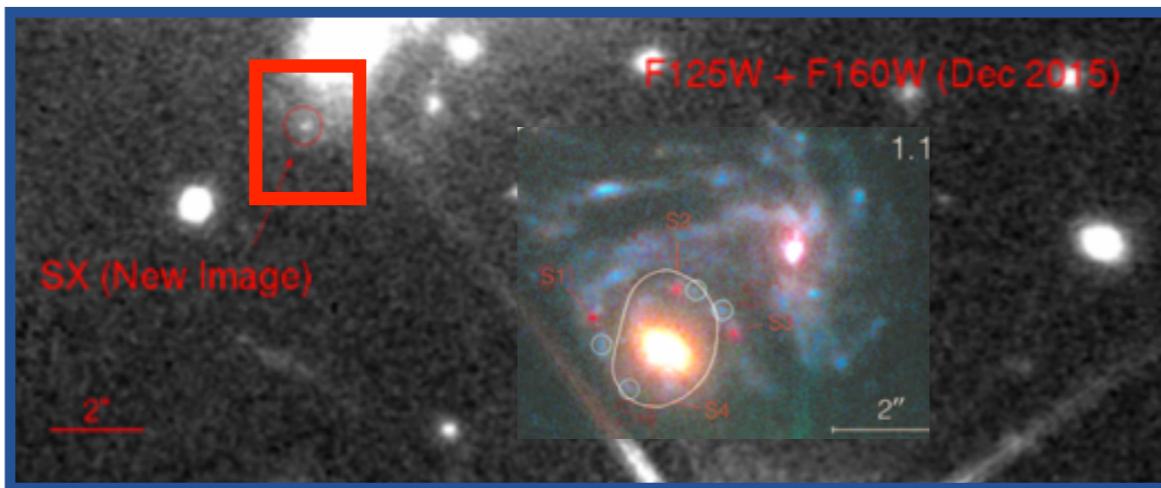
- Light passing closest to the lens (t_-) is delayed the most
- Light from image θ_+ arrive first



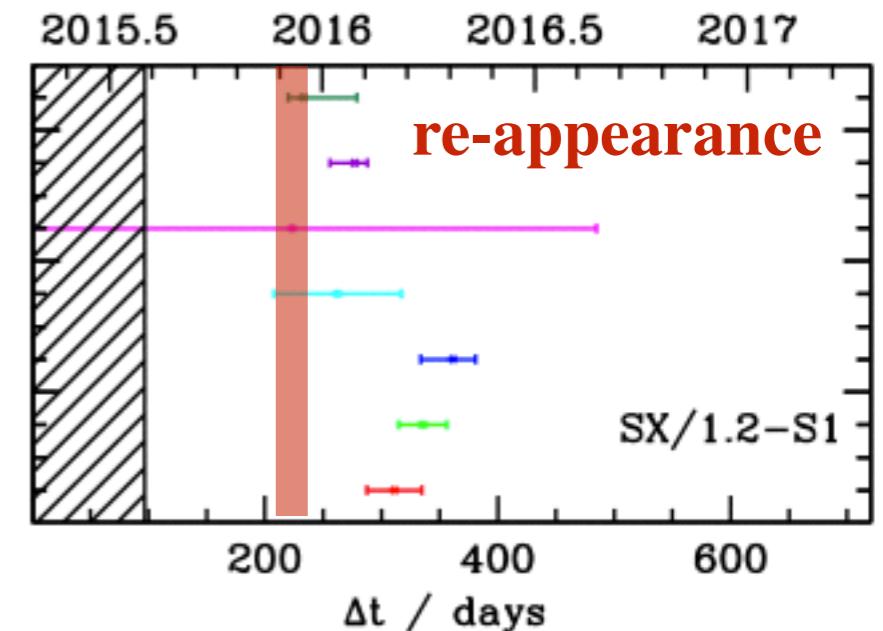
Geometry \sim route taken
Gravitational potential \sim traffic along route

LE Consequences 2: Time Delays

- SN Refsdal Reappearance December 11 2015

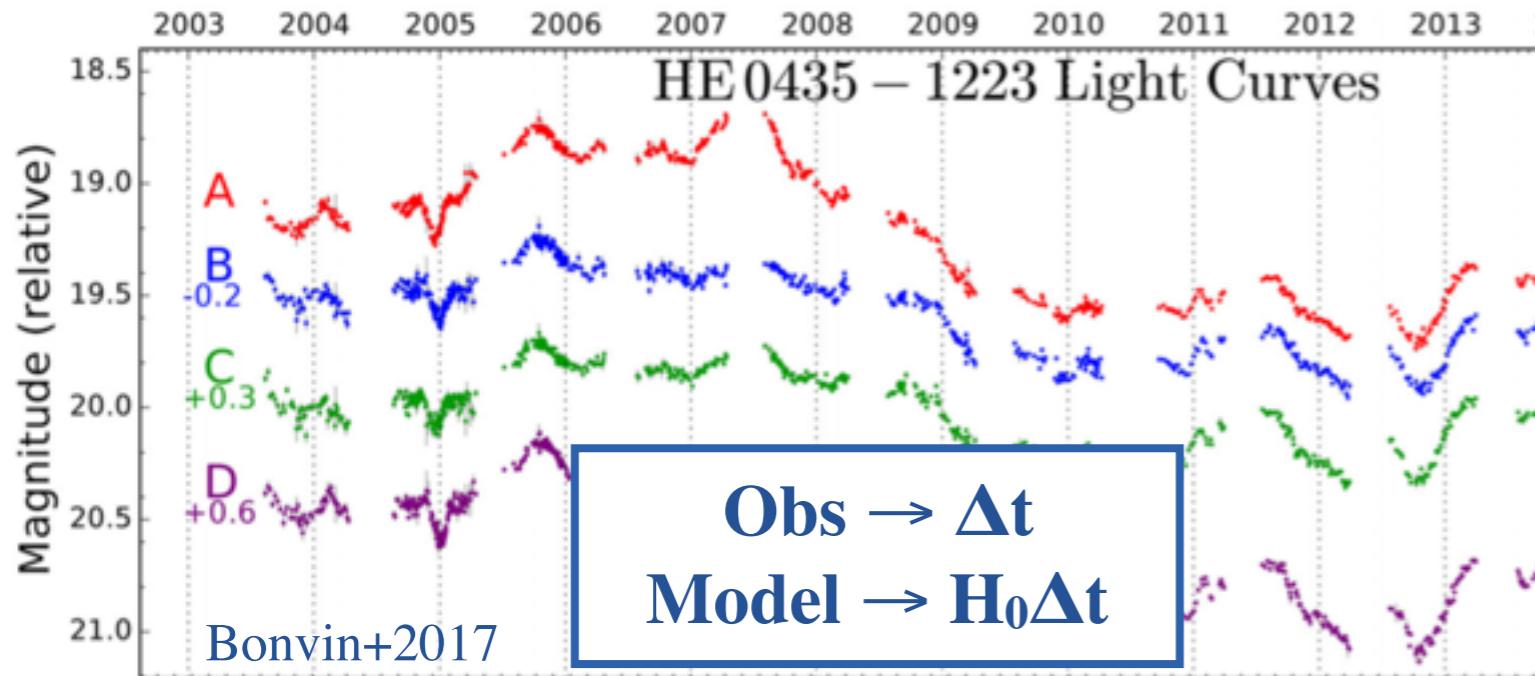
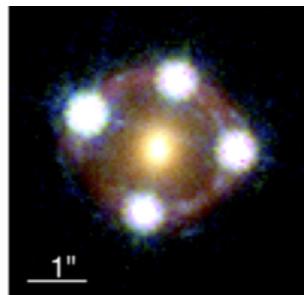


HE 0435–1223

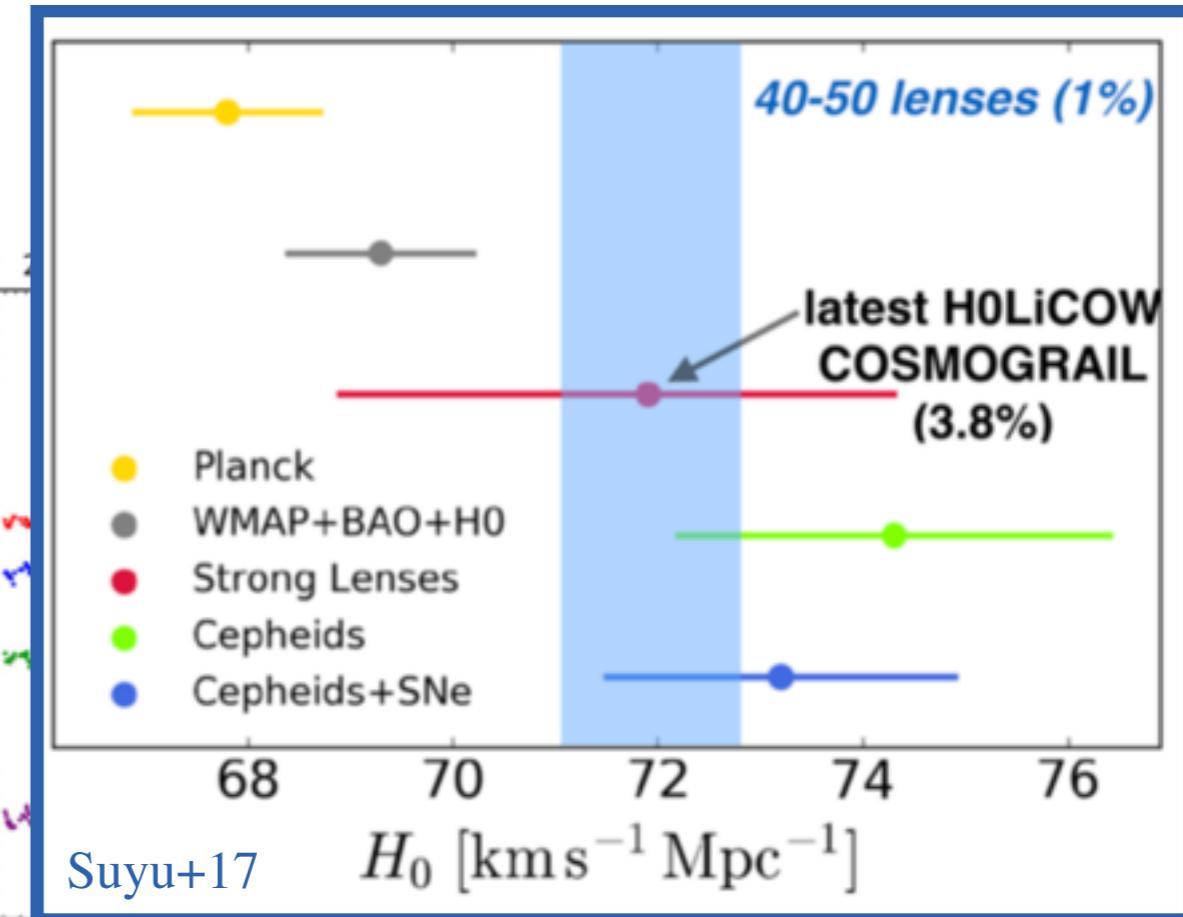


Treu+2016

- COSMOGRAIL & H0LiCOW



Bonvin+2017



Suyu+17

LE Consequences 3: Magnification

- Introduced the Jacobian Matrix

$$\mathcal{A}(\boldsymbol{\theta}) = \frac{\partial \boldsymbol{\beta}}{\partial \boldsymbol{\theta}} = \begin{pmatrix} \frac{\partial \beta_i}{\partial \theta_i} & \frac{\partial \beta_i}{\partial \theta_j} \\ \frac{\partial \beta_j}{\partial \theta_i} & \frac{\partial \beta_j}{\partial \theta_j} \end{pmatrix}$$

- Related this to the deflection angles and the gravitational potential

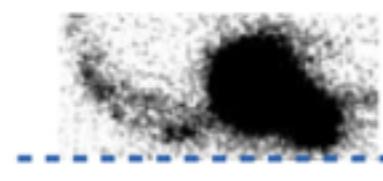
$$\mathcal{A}(\boldsymbol{\theta}) = \begin{pmatrix} 1 - \frac{\partial \alpha_i}{\partial \theta_i} & -\frac{\partial \alpha_i}{\partial \theta_j} \\ -\frac{\partial \alpha_j}{\partial \theta_i} & 1 - \frac{\partial \alpha_j}{\partial \theta_j} \end{pmatrix} = \begin{pmatrix} 1 - \frac{\partial^2 \psi}{\partial \theta_i^2} & -\frac{\partial^2 \psi}{\partial \theta_i \partial \theta_j} \\ -\frac{\partial^2 \psi}{\partial \theta_j \partial \theta_i} & 1 - \frac{\partial^2 \psi}{\partial \theta_j^2} \end{pmatrix} \equiv (\delta_{ij} - \Psi_{ij})$$

- Defined the distortion tensor (Ψ_{ij}), convergence (κ) and shear (γ)
- And from that the magnification of

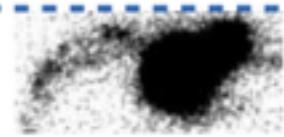
$$\mu \equiv \frac{1}{\det \mathcal{A}(\boldsymbol{\theta})} = \frac{1}{(1 - \kappa)^2 - \gamma^2} ; \quad \gamma^2 \equiv \gamma_1^2 + \gamma_2^2$$

- Considered the magnification and parity of the IS case

$\mu > 0$: positive parity



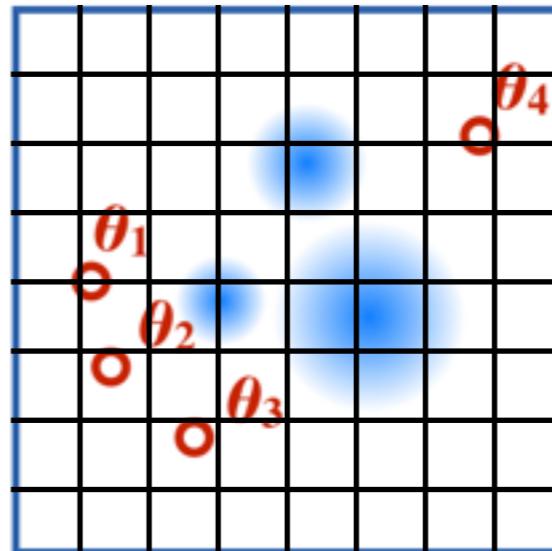
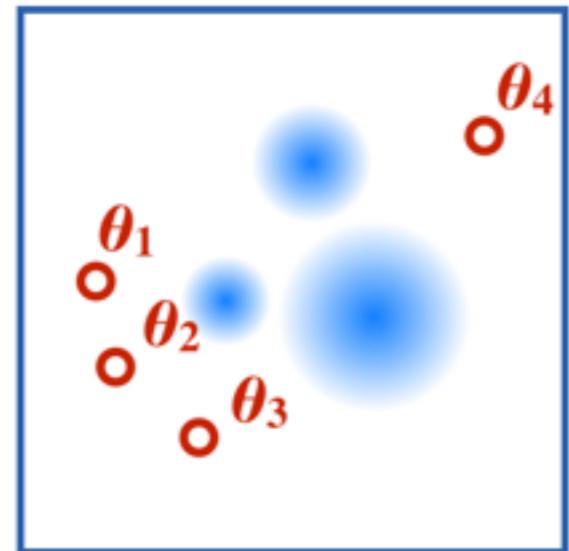
$\mu < 0$: negative parity



Modeling Lenses for Scientific Purposes

- Two approaches: Parametric & non-parametric

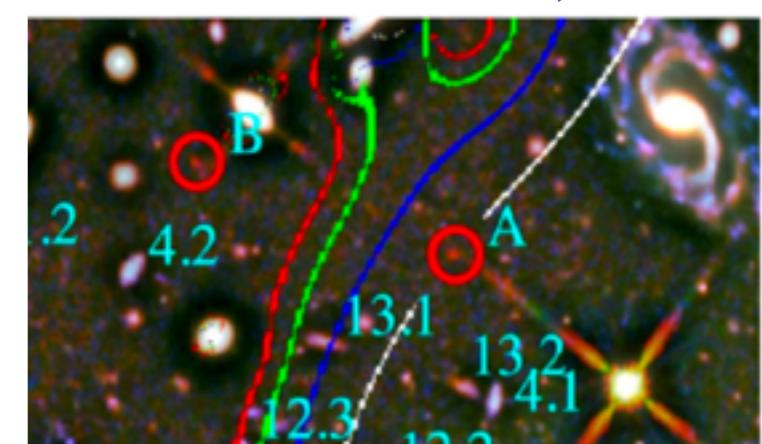
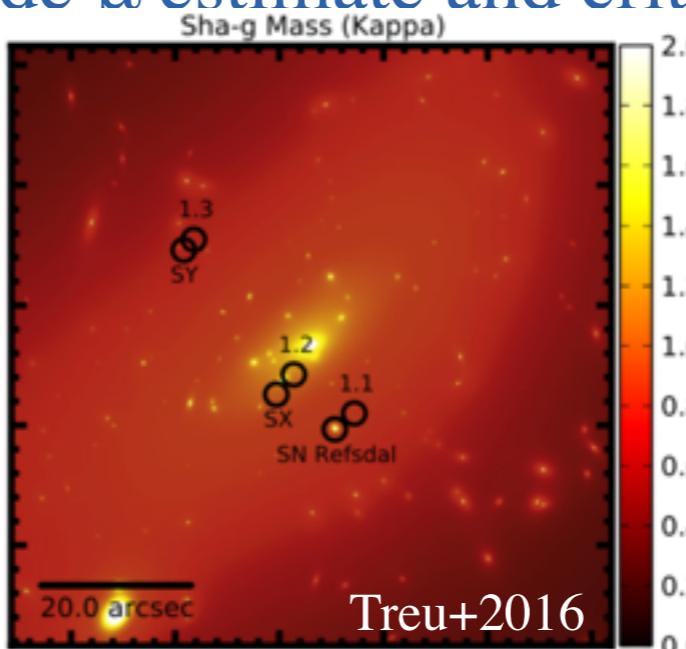
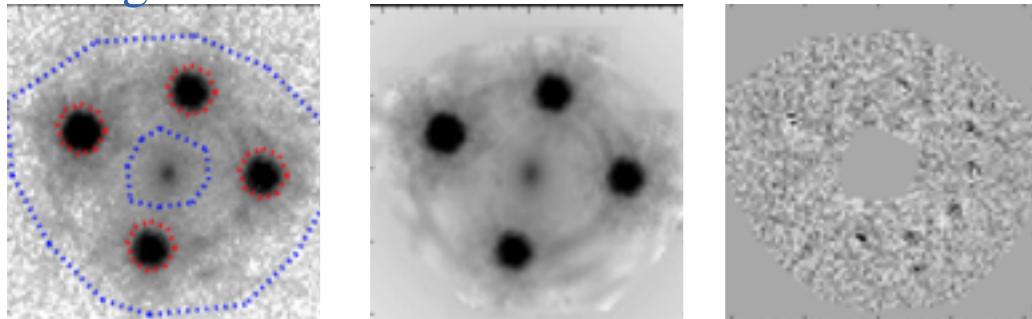
Parametrize matter profile/distribution (e.g. NFW or IS) and re-produce source positions



Model pixelated surface brightness distribution iteratively by “molding” lens mass distribution

- Lens models and their mass (κ) maps are useful for, e.g.
- Finding high- z galaxies (provide u estimate and critical curve locations)
- Determining lens masses
- Predict $H_0 \Delta t$ time-delays

Wang+2017



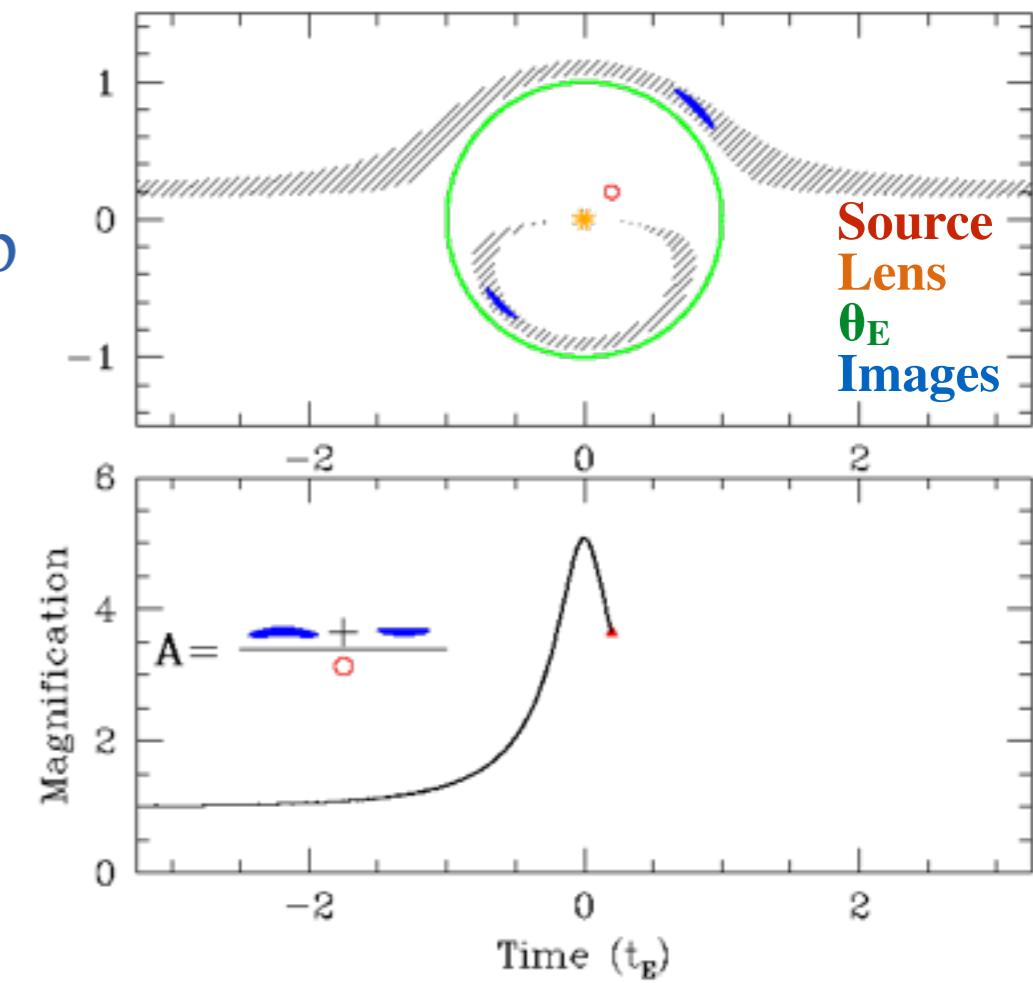
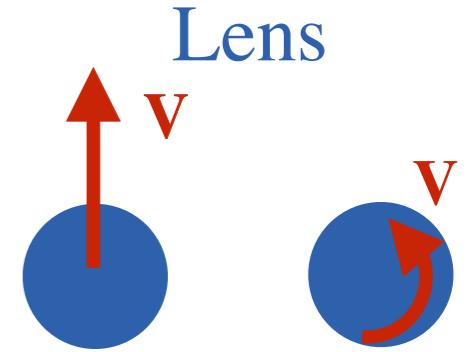
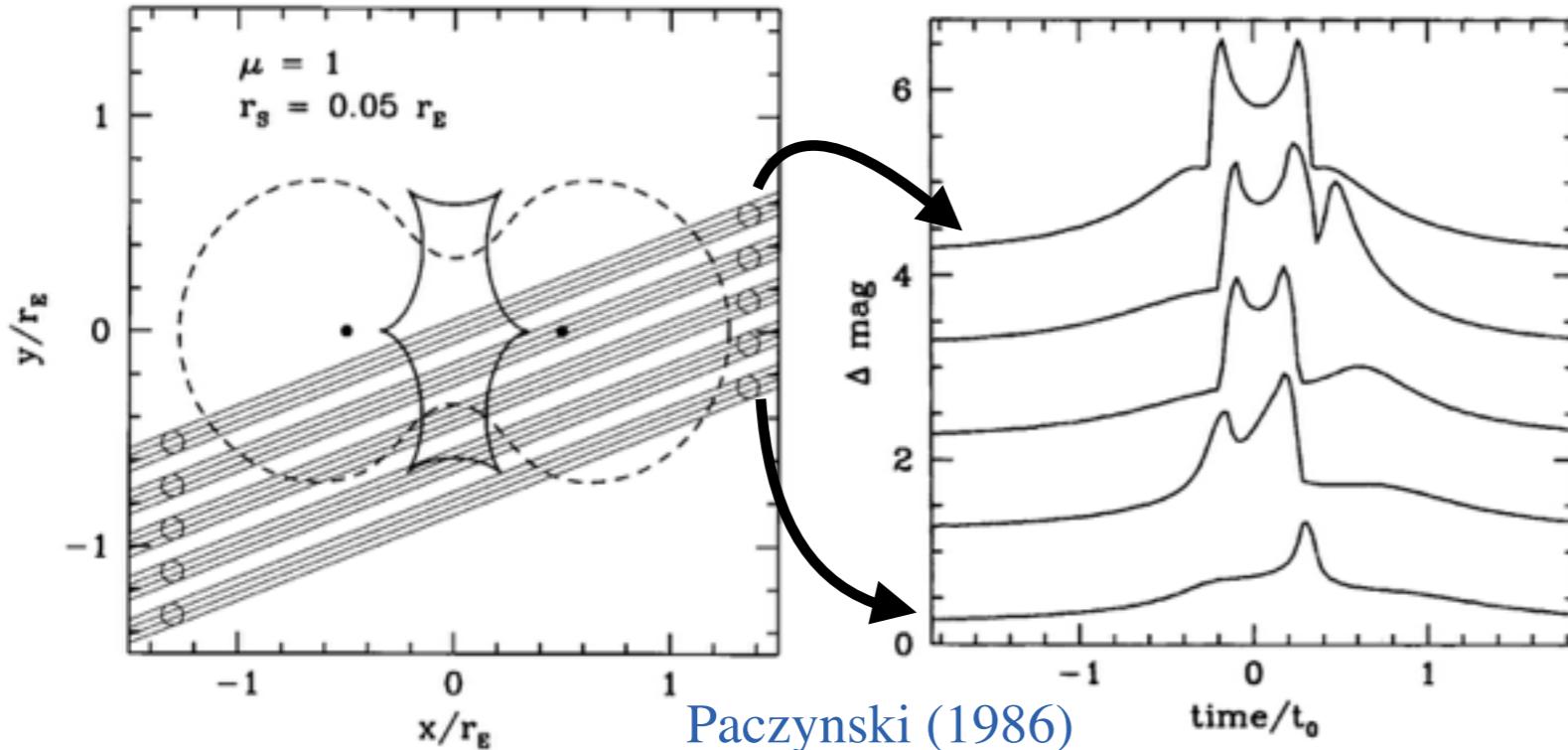
Microlensing

- Galactic and extra galactic micro lensing
- Can be thought of as “unresolved strong lensing”
- Observe variations in μ , i.e., the source brightness
- For the PML we have

$$\mu_{\pm} = \frac{1}{1 - (\theta_E/\theta_{\pm})^4}$$

$$\mu = \frac{y^2 + 2}{y\sqrt{y^2 + 4}} \quad \text{where} \quad y = \frac{\beta}{\theta_E}$$

- Caustics (patterns) provide magnification map



Movie credit: B. Scott Gaudi, OSU

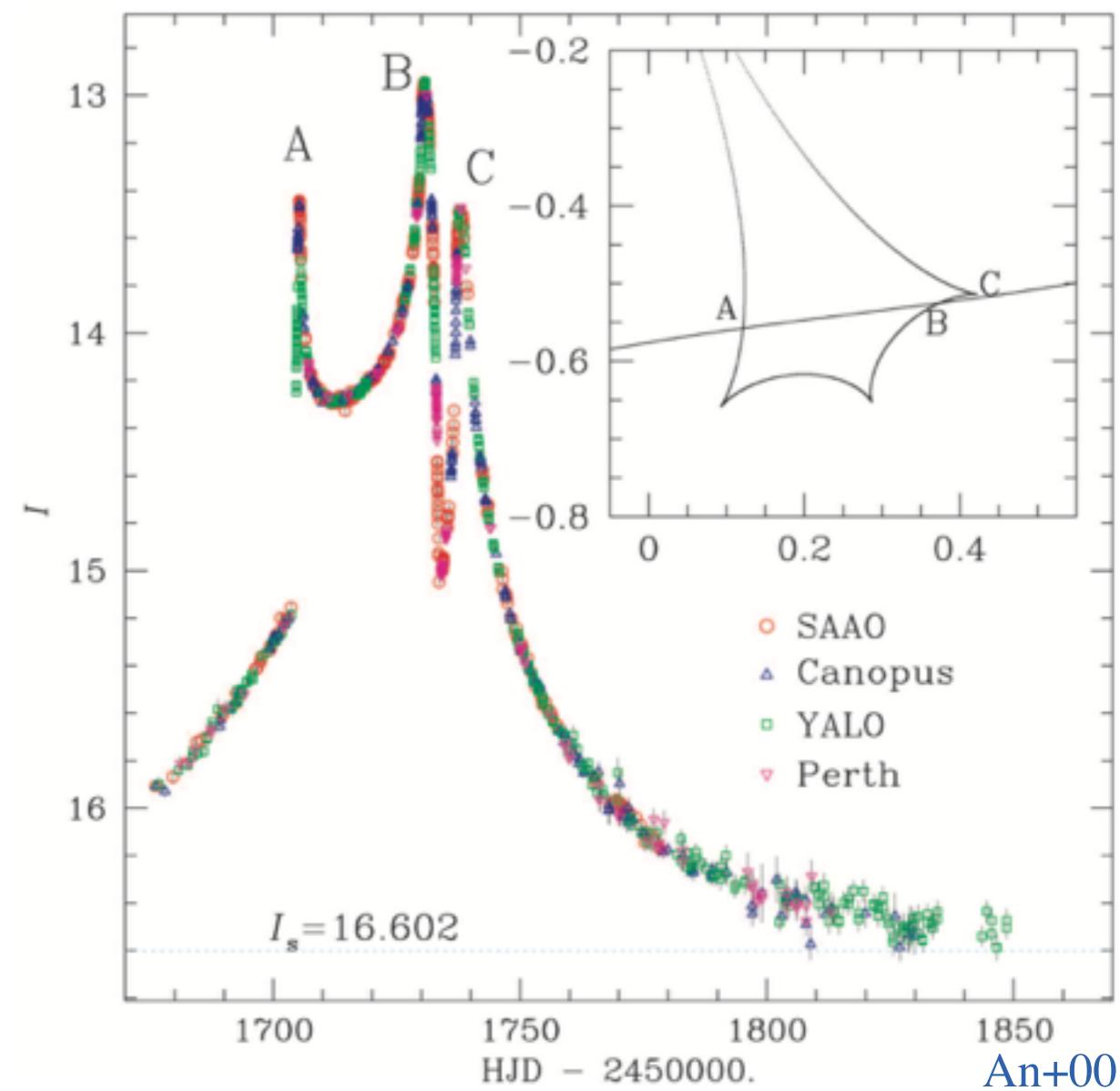
Microlensing

- Particular changes to the magnification curve can indicate double PML
- For an ‘x-axis aligned’ double PML the magnification is

$$\mu = \frac{1}{1 - \frac{1}{x^4} \left(1 + \frac{qx^2}{(x-x_p)^2} \right)^2}$$

- Here the subscript ‘p’ refers to **perturber**
- ... or **planet**
- So if a planet (around lens) is close to lensed image position, $\Delta\mu_p$ is large
- Looked at a range of such events:

$$\Delta\mu_p \simeq \frac{2\mu_0^2 q}{x^2(x - x_p)^2} \quad q = \frac{m_p}{M_\star}$$



Weak Lensing

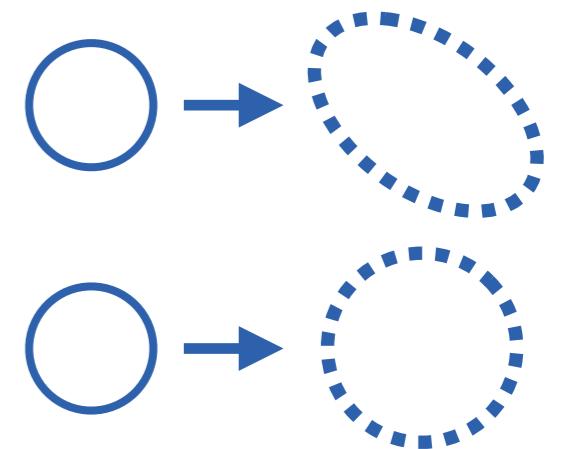
- Characterizing source image distortions from lensing
 - Shearing (γ)
 - Scaling (κ)
- Idealized weak lensing ($\kappa = 0, \gamma_2 = 0$) reveals ellipse nature of distortion

$$1 = \frac{(1 - \gamma_1)^2}{\beta_0^2} \theta_1^2 + \frac{(1 + \gamma_1)^2}{\beta_0^2} \theta_2^2$$

- Ellipticity can be described by 2nd order surface brightness moments

$$q_{ij} \equiv \int d^2\theta \, S^{\text{obs}}(\theta) \theta_i \theta_j$$

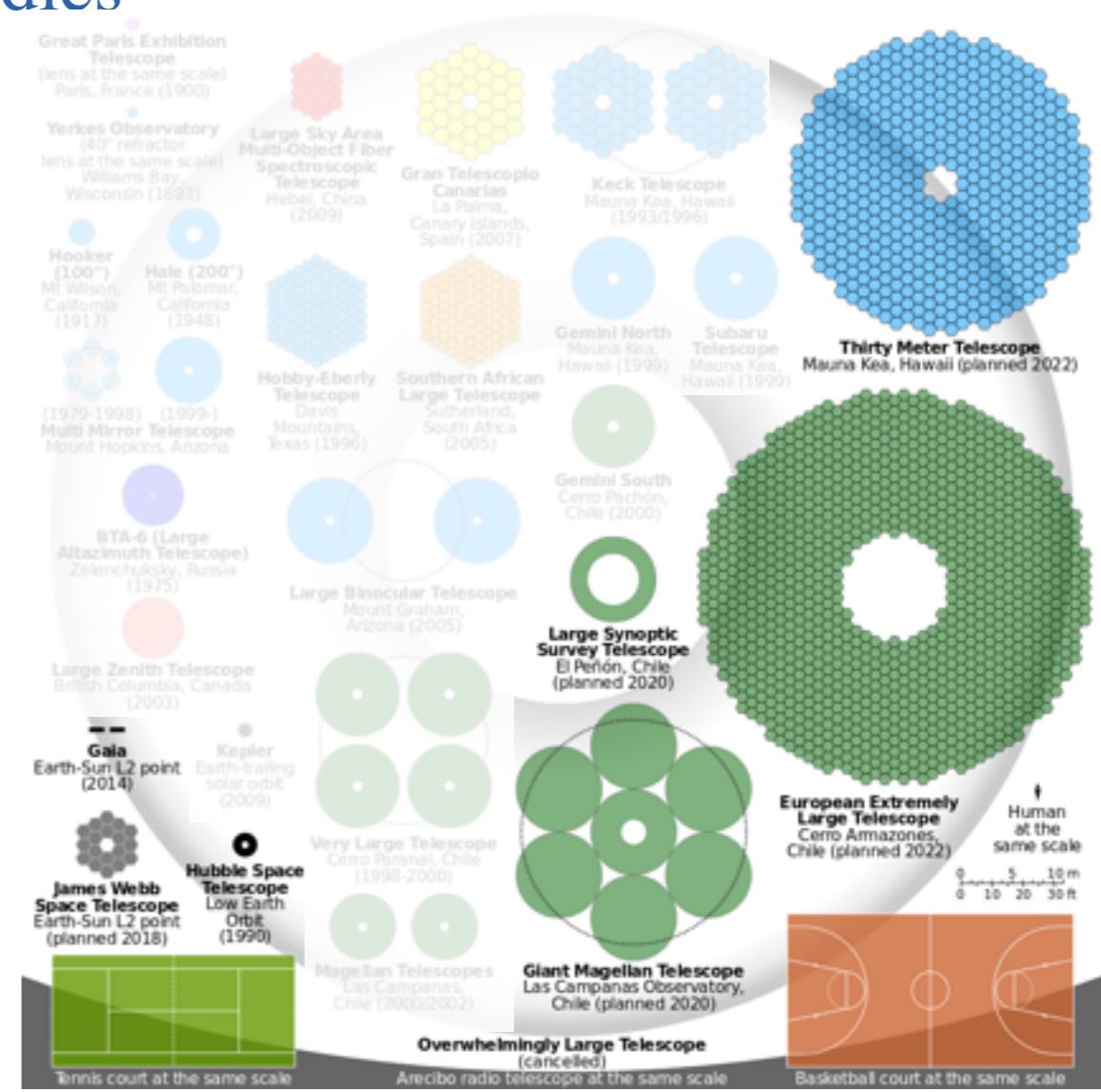
- Useful for large scale mass determination
- To measure this ellipticity in the weak lensing regime, needs statistics
 - Cross correlation & power spectrum on density contrasts $\delta(\mathbf{x}, t) \equiv \frac{\rho_m(\mathbf{x}, t) - \bar{\rho}_m(t)}{\bar{\rho}(t)}$
- The extreme: Cosmic shearing of CMB by integrated foreground mass
 - Probes cosmological parameters and lensing characteristics

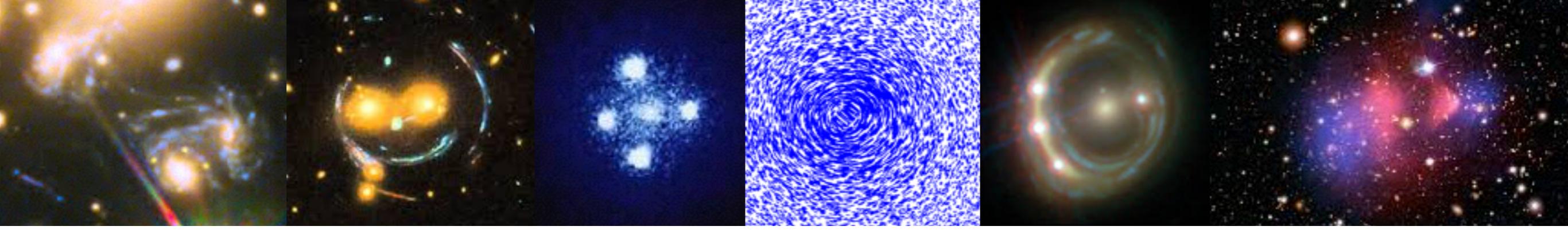


Outlook to the (near) future of lensing

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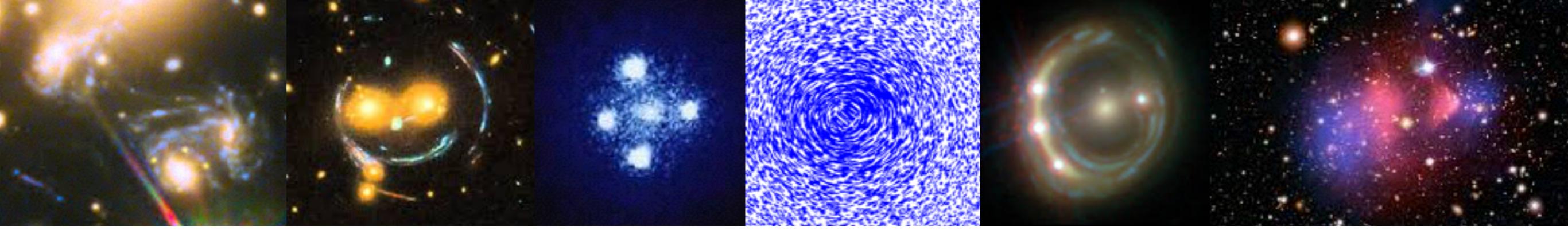
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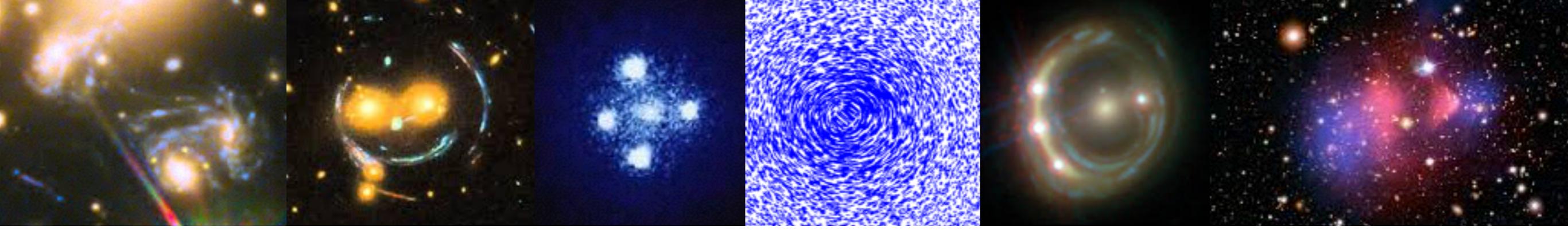
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Questions?



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Last Week's Worksheet



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This Week's Worksheet