



# Geometry of Morse Theory

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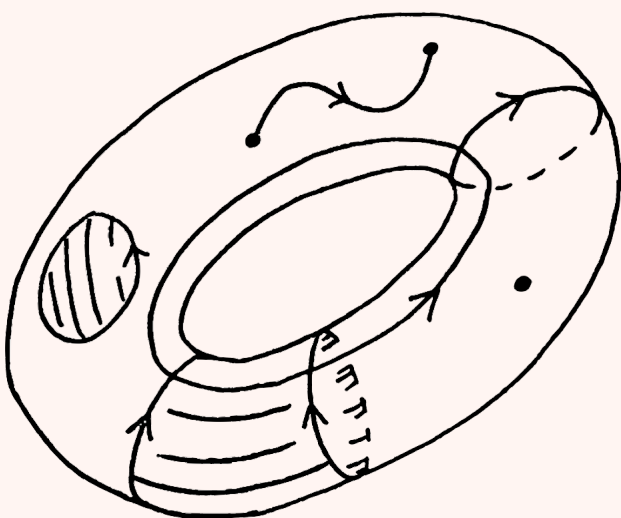
## Motivation

One goal of mathematical physics is to study *topological invariants* of *manifolds* like **Betti numbers** which count the **number of holes in each dimension** and the **Euler characteristic**.

**Definition (Homology groups).** For a given manifold  $M$ , the  $n$ -th Homology group is defined:

$$H_n(M) = \frac{n\text{-dim submanifolds without boundary}}{n\text{-dim boundaries to } (n+1)\text{-dim submanifolds}}$$

Torus:  $H_0(T) = \mathbb{Z}$   $H_1(T) = \mathbb{Z}^2$   $H_2(T) = \mathbb{Z}$



DIM	SUBMANIFOLD	BOUNDARY
0	•	∅
1	↪	{•, •}
1	⊙	∅
2	⊙	⊙
2	⊙	{∅, ∅}
2	⊙	∅

**Definition (Betti numbers).** For a given manifold  $M$ , the  $n$ -th Betti number, denoted  $b_n$ , is the rank of the  $n$ -th Homology group.

**Definition (Euler characteristic).** The *Euler characteristic* for a manifold  $M$  is

$$\chi(M) = \sum_{i=0}^n (-1)^i b_i$$

Torus:  $b_0 = 1$   $b_1 = 2$   $b_2 = 1$

$$\chi(T) = 1 - 2 + 1 = 0$$

## Morse Theory

Morse theory studies the **critical points** on a manifold and how to **reconstruct the manifold** given the critical points.

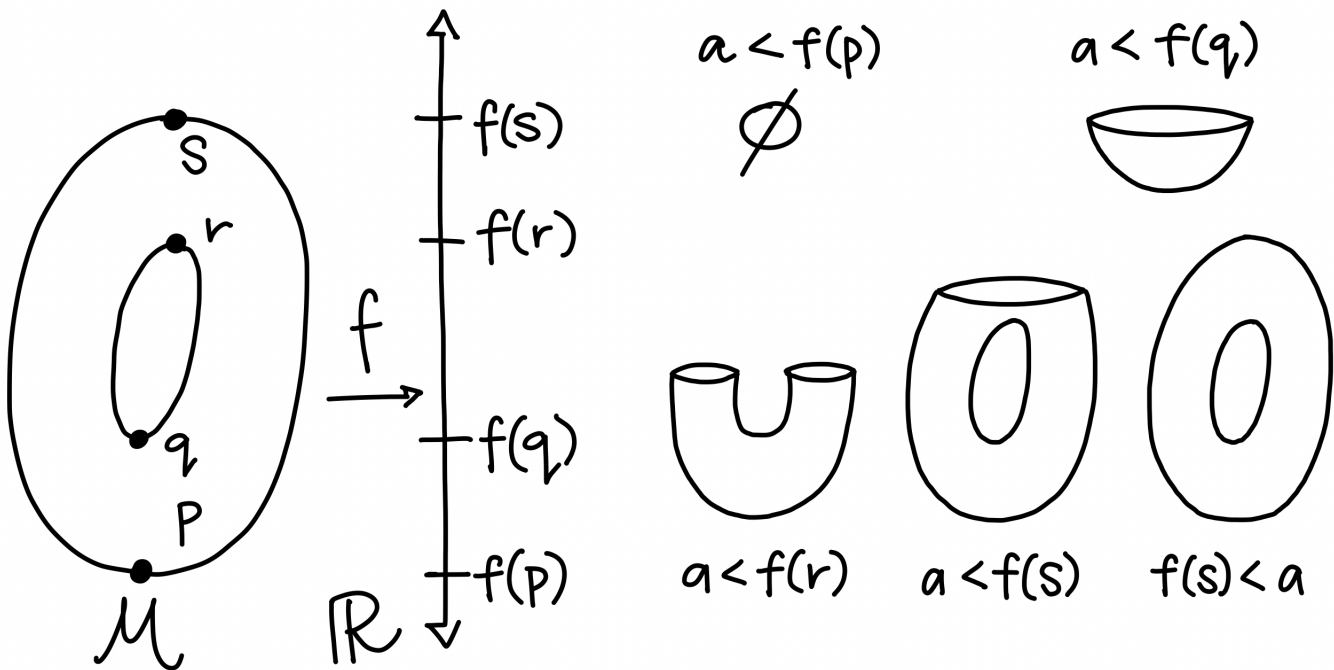
**Definition (Morse function).** Let  $f : M \rightarrow \mathbb{R}$ . We call  $f$  a *Morse function* if each critical point has a non-zero second derivative:

determinant of Hessian matrix  $= \det(H_f) = \det \left[ \frac{\partial^2 f}{\partial x_i \partial x_j} \right] \neq 0$

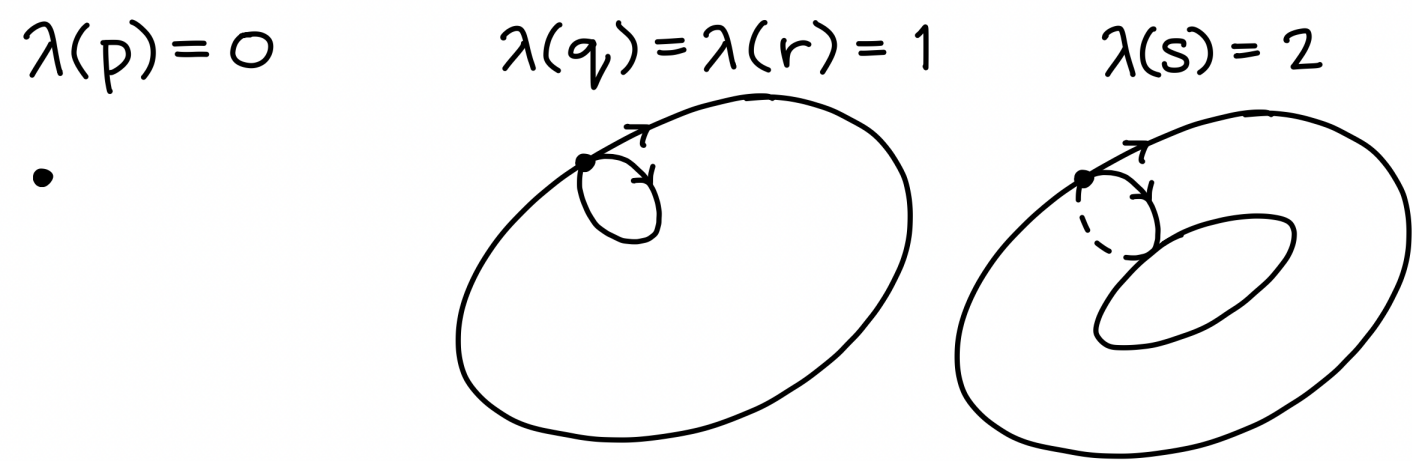
**Definition (index).** The (*Morse*) *index* of a Morse function  $f$  is the number of negative eigenvalues of the function's Hessian matrix.

Suppose that  $M$  is a compact manifold and  $f : M \rightarrow \mathbb{R}$  is a Morse function. Define *compact submanifolds* of  $M$ :

$$M^a = \{x \in M \mid f(x) \leq a\}$$



When  $f$  passes through a critical point, the *homology* of the compact submanifolds change. In fact, when  $M^a$  passes through a critical point  $p$  and “becomes”  $M^b$  **one adds a  $\lambda$ -cell to  $M^a$**  where  $\lambda$  is the index of  $p$  (ie.  $M^b \simeq M^a \cup e^\lambda$ ).



**Theorem (Morse Inequalities).** Let  $c_\lambda$  denote the number of critical points of index  $\lambda$  in  $M$ . Then,

$$b_\lambda \leq c_\lambda$$

$$\sum_{\lambda=0}^n (-1)^\lambda c_\lambda = \chi(M)$$

Torus:  $c_0 = 1$   $c_1 = 2$   $c_2 = 1$

$$\chi(T) = 1 - 2 + 1 = 0$$

## Hodge Theory

Hodge theory counts the number of holes in a manifold by studying **differential forms** and **cohomology classes**, the *dual* to Homology groups.

**Definition (differential form).** Given an  $n$ -dimensional manifold, a (*differential*)  $p$ -form  $\omega \in \Omega^p(M)$  is a family of  $(n - p)$ -dimensional surfaces.

**Definition (exterior derivative).** The *exterior derivative* is a map that turns a  $p$ -form into a  $(p + 1)$ -form. (**generalized gradient**)

$$d : \Omega^p(M) \rightarrow \Omega^{p+1}(M)$$

**Definition (codifferential).** The *codifferential* is a map that turns a  $p$ -form into a  $(p - 1)$ -form. (**generalized divergence**)

$$\delta : \Omega^p(M) \rightarrow \Omega^{p-1}(M)$$

**Definition (Laplacian).** The *Laplacian* is a map from  $p$ -form to themselves.

$$\Delta = d\delta + \delta d$$

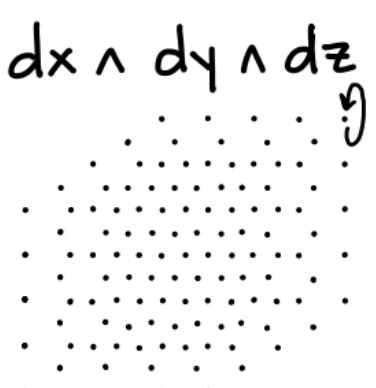
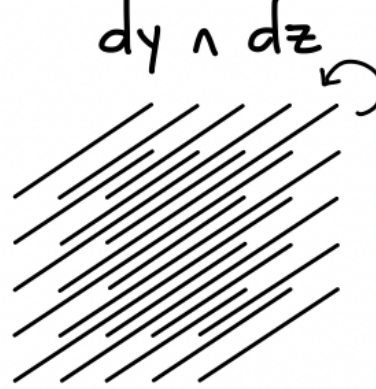
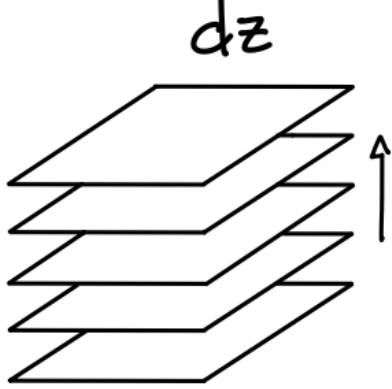
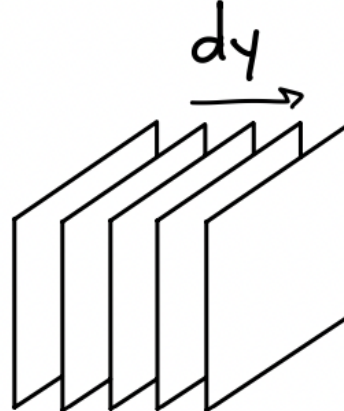
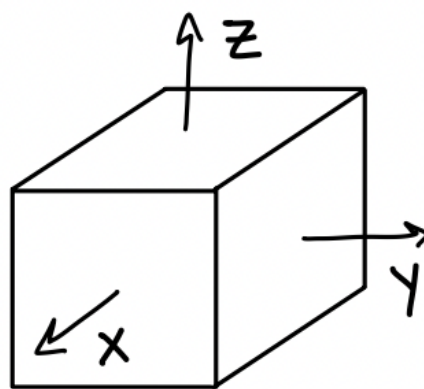
$$\Delta : \Omega^p(M) \rightarrow \Omega^p(M)$$

There are three types of differential forms:

- closed forms ( $d\omega = 0$ )
- exact forms ( $\omega = d\alpha$ )
- harmonic forms ( $\Delta\omega = 0$ )

**Definition (de Rham cohomology classes).** For a given manifold  $M$ , the  $n$ -th *cohomology class* is defined:

$$H^n(M) = \frac{\text{closed } n\text{-forms}}{\text{exact } n\text{-forms}} = [\text{harmonic forms}]$$



## Witten's Insights

Witten used *supersymmetry* and *Hodge theory* in order to **connect Morse theory to cohomology classes**.

Let  $h : M \rightarrow \mathbb{R}$  and  $t \in \mathbb{R}$ .

$$d_t = e^{-ht} d e^{ht} \quad \delta_t = e^{ht} \delta e^{-ht}$$

$$\Omega^0, \Omega^2, \dots, \Omega^n \quad \Omega^1, \Omega^3, \dots, \Omega^{n-1}$$

bosonic states fermionic states

**Definition (Witten symmetry operators).** A supersymmetry theory is one with (*Hermitian*) *symmetry operators*.

$$Q_{1t} = d_t + \delta_t \quad Q_{2t} = i(d_t - \delta_t) \quad H_t = d_t \delta_t + \delta_t d_t$$

**Proof (of First Morse Inequality).**

- $e^{ht}$  is invertible  $\implies$  statements hold for all  $t \in \mathbb{R}$
- number of zero eigenvalues of  $H_t$  acting on  $\Omega^p = b_p(t) = b_p$
- $(t \rightarrow \infty)$  eigenvalues of  $H_t$  can be calculated in terms of the local data at critical points
- around  $A$  (critical point of index  $p$ ),  $H_t$  has **one** zero eigenvalue which is a  $p$ -form, denoted  $|a\rangle$

$\implies$  at most the number of zero energy  $p$ -forms equals the number of critical points with index  $p \implies b_p \leq c_p \square$

**Proof (of Second Morse Inequality).** (*goal: define a coboundary operator using critical points*)

- potential energy has minimum at each critical point  $\implies$  tunneling from one critical point  $A$  to another  $B$
- tunneling paths  $\Lambda : A \rightarrow B$  and  $\Gamma : B \rightarrow A$  are of steepest descent  $\implies A$  and  $B$  have index difference of 1
- paths induce orientation  $\implies$  if  $\Gamma$  and  $\Lambda$  agree then  $n_\Gamma = +1$ , and if not then  $n_\Gamma = -1 \implies$  define  $n(a, b) = \sum_\Gamma n_\Gamma$

$\implies$  define a coboundary operator such that:

$$\partial |a\rangle = \sum_b n(a, b) |b\rangle$$

$\implies$  Betti numbers of  $\partial$  are the same Betti numbers of  $M \square$

## Acknowledgements

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