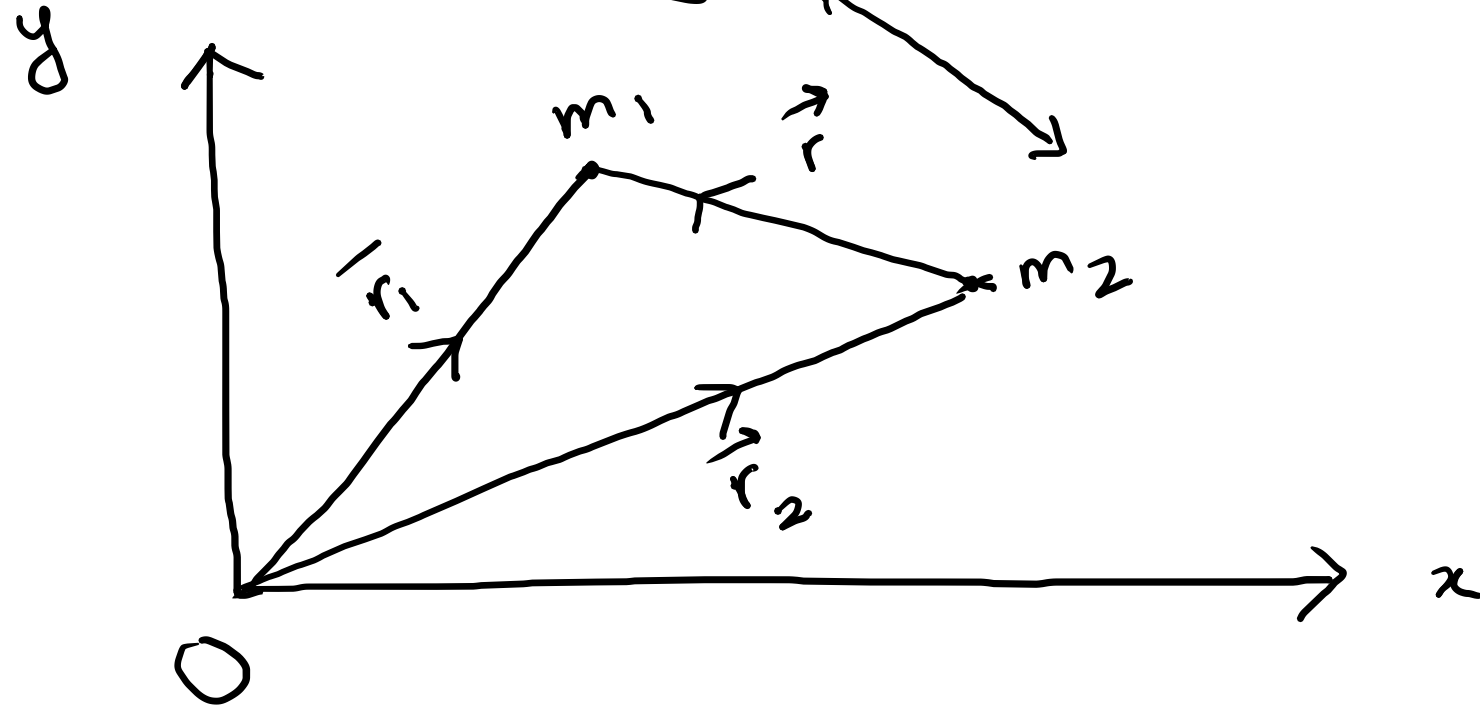


# LECTURE 17

## Two-BODY PROBLEM

Two self-gravitating bodies



$$\begin{aligned}\vec{F} &= - \frac{G m_1 m_2}{r^2} \hat{r} \\ &= - \frac{G m_1 m_2}{r^3} \vec{r}\end{aligned}$$

So far, have considered  
only central force directed  
towards the origin.

EOM :-

$$m_1 \ddot{\vec{r}}_1 = \vec{F}_{12}$$

$$m_2 \ddot{\vec{r}}_2 = \vec{F}_{21}$$

$\vec{F}_{ij} \equiv$  force on body  $i$  due to body  $j$

$$\vec{F}_{12} = -\vec{F}_{21}$$

Aim:- to reduce this problem to a "one-body" problem, so far considered.

Define,  $\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$   $\longrightarrow$  Centre of mass

$\vec{r} = \vec{r}_1 - \vec{r}_2$   $\longrightarrow$  relative separation.

Helpful to know the rel<sup>n</sup>s between  $(R, r)$  and  $(\vec{r}_1, \vec{r}_2)$

$$\boxed{\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}}$$

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

$$= \frac{m_1 \vec{r}_1 + m_2 (\vec{r}_1 - \vec{r})}{m_1 + m_2}$$

$$= \vec{r}_1 - \frac{m_2 \vec{r}}{m_1 + m_2}$$

$$\Rightarrow \vec{r}_1 = \vec{R} + \frac{m_2 \vec{r}}{m_1 + m_2} \quad //$$

$$\vec{r}_2 = \vec{r}_1 - \vec{r} = \vec{R} + \frac{m_2 \vec{r}}{m_1 + m_2} - \vec{r}$$

$$= \vec{R} - \frac{m_1 \vec{r}}{m_1 + m_2} \quad //$$

- EOM in the new variables,

$$m_1 \ddot{\vec{r}}_1 = \vec{F}_{12} \quad \text{--- (1)}$$

$$m_2 \ddot{\vec{r}}_2 = \vec{F}_{21} \quad \text{--- (2)}$$

$$\vec{F}_{12} = -\vec{F}_{21}$$

$$m_2 \times (1) - m_1 \times (2),$$

$$m_1 m_2 (\ddot{\vec{r}}_1 - \ddot{\vec{r}}_2) = m_2 \vec{F}_{12} - m_1 \vec{F}_{21} = (m_1 + m_2) \vec{F}_{12}$$

$$\Rightarrow \frac{m_1 m_2}{m_1 + m_2} \ddot{\vec{r}} = \vec{F}_{12}$$

Effectively represents a "one-body problem" if  
and  $\vec{r}$  is the position vector.

$$\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$$

$$\boxed{\mu \ddot{\vec{r}} = \vec{F}_{12}}$$

$$\left. \begin{aligned} m_1 \ddot{\vec{r}}_1 &= \vec{F}_{12} \\ m_2 \ddot{\vec{r}}_2 &= \vec{F}_{21} \end{aligned} \right\} \begin{aligned} \frac{m_1 \ddot{\vec{r}}_1}{m_1 + m_2} + \frac{m_2 \ddot{\vec{r}}_2}{m_1 + m_2} &= \frac{\vec{F}_{12}}{m_1 + m_2} + \frac{\vec{F}_{21}}{m_1 + m_2} \\ \Rightarrow \frac{m_1 \ddot{\vec{r}}_1 + m_2 \ddot{\vec{r}}_2}{m_1 + m_2} &= 0 \\ \Rightarrow \ddot{\vec{R}} = 0 \quad \Rightarrow \quad \dot{\vec{R}} = \text{const} . \end{aligned}$$

$$\begin{aligned} T &= \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 \\ &= \frac{1}{2} m_1 \left[ \dot{\vec{R}} + \frac{m_2 \dot{\vec{r}}}{m_1 + m_2} \right]^2 + \frac{1}{2} m_2 \left[ \dot{\vec{R}} - \frac{m_1 \dot{\vec{r}}}{m_1 + m_2} \right]^2 \\ &= \frac{1}{2} (m_1 + m_2) \dot{\vec{R}}^2 + \frac{1}{2} m_1 \frac{m_2^2 \dot{\vec{r}}^2}{(m_1 + m_2)^2} + \frac{1}{2} m_2 \frac{m_1^2 \dot{\vec{r}}^2}{(m_1 + m_2)^2} \\ &= \frac{1}{2} (m_1 + m_2) \dot{\vec{R}}^2 + \frac{1}{2} \frac{m_1 m_2}{(m_1 + m_2)} \dot{\vec{r}}^2 \quad (\text{since } \frac{m_1^2}{m_1 + m_2} + \frac{m_2^2}{m_1 + m_2} = \frac{m_1 m_2}{m_1 + m_2}) \end{aligned}$$

$$= \underbrace{\frac{1}{2}(m_1 + m_2)}_{\text{const.}} \dot{R}^2 + \frac{1}{2} \frac{m_1 m_2}{(m_1 + m_2)} \dot{r}^2$$

$$= \text{const} + \underbrace{\frac{1}{2} \mu \dot{r}^2}$$

$$E = \frac{1}{2} \mu \dot{r}^2 + U(|\vec{r}|) \rightarrow \text{expression similar to} \\ \text{conv. eqn. for one-body} \\ \text{problem.}$$

Small oscillations about extrema of a potential.

$$U(x) = U_0(-ax^2 + bx^4), \quad U_0, a, b > 0.$$

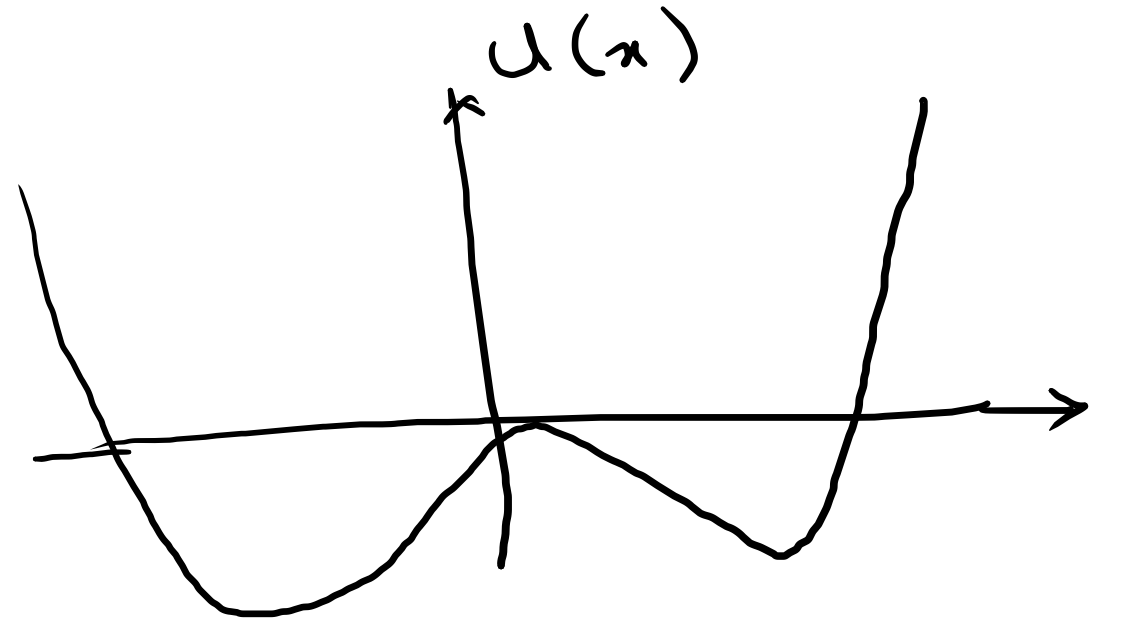
$$\frac{dU(x)}{dx} = U_0(-2ax + 4bx^3) = 0$$

$$\Rightarrow 2x(a - 2bx^2) = 0$$

$$x = 0, \pm \sqrt{\frac{a}{2b}}$$

$$U(0) = 0$$

$$U(x = \pm \sqrt{\frac{a}{2b}}) = U_0 \left[ -a \frac{a}{2b} + b \cdot \frac{a^2}{4b^2} \right] = U_0 \left( -\frac{a^2}{4b} \right)$$



$$\frac{d^2 U}{dx^2} = U_0 (-2a + 12bx^2)$$

$$\left. \frac{d^2 U}{dx^2} \right|_{x=0} = -2aU_0 < 0.$$

$$\left. \frac{d^2 U}{dx^2} \right|_{x = \pm \sqrt{\frac{a}{2b}}} = U_0 \left( -2a + \overset{6}{12} \cancel{b} \cdot \frac{a}{\cancel{2b}} \right) = 4aU_0 > 0$$

$$U(x) = \underline{U(x_0)} + \cancel{\left. \frac{dU}{dx} \right|_{x_0} (x-x_0)} + \underbrace{\frac{1}{2} \left. \frac{d^2 U}{dx^2} \right|_{x_0} (x-x_0)^2}$$

$$F(x) = - \frac{dU(x)}{dx} = - \underbrace{\left. \frac{d^2 U}{dx^2} \right|_{x_0}}_{k} (x-x_0).$$

$$\omega = \sqrt{4aU_0}.$$

$$m \ddot{x} = -k(x-x_0)$$

$$\Rightarrow \ddot{x} + \frac{k}{m}(x-x_0) = 0$$

$$\Rightarrow \ddot{x} + \omega^2(x-x_0) = 0$$



▣ Different type of force.

$$\vec{F} = \frac{A}{r} \hat{\theta}$$

$$d\vec{r} = dr \hat{r} + r d\theta \hat{\theta}$$

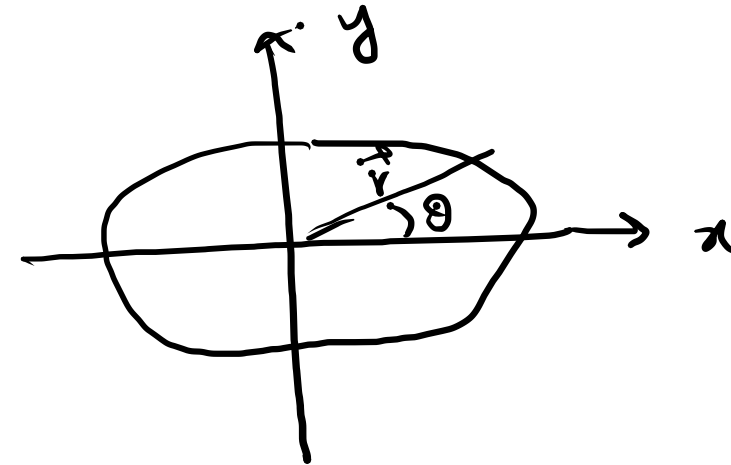
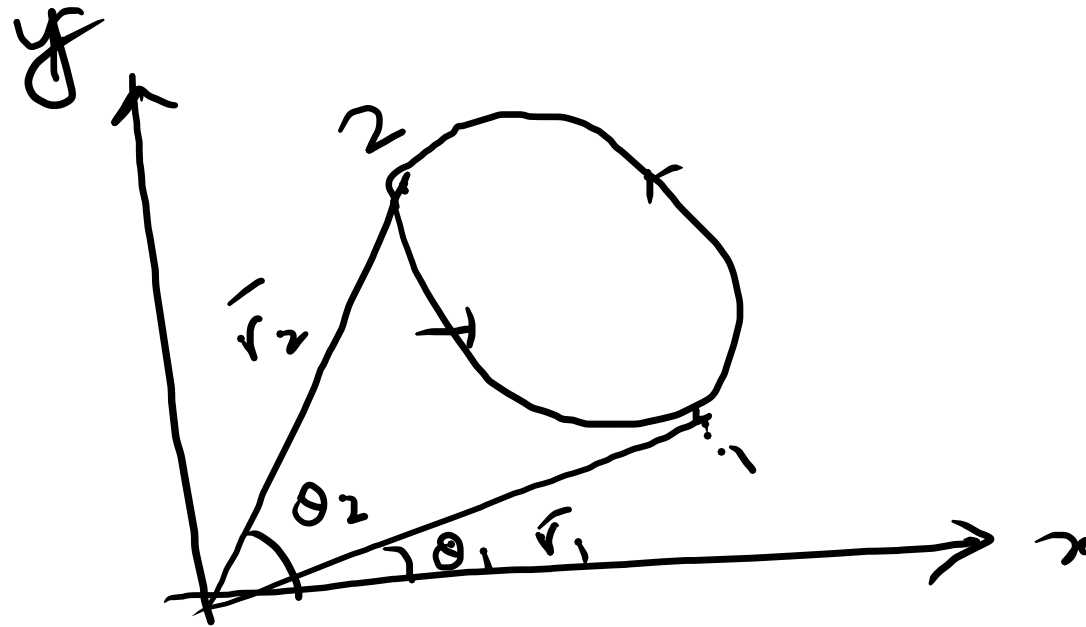
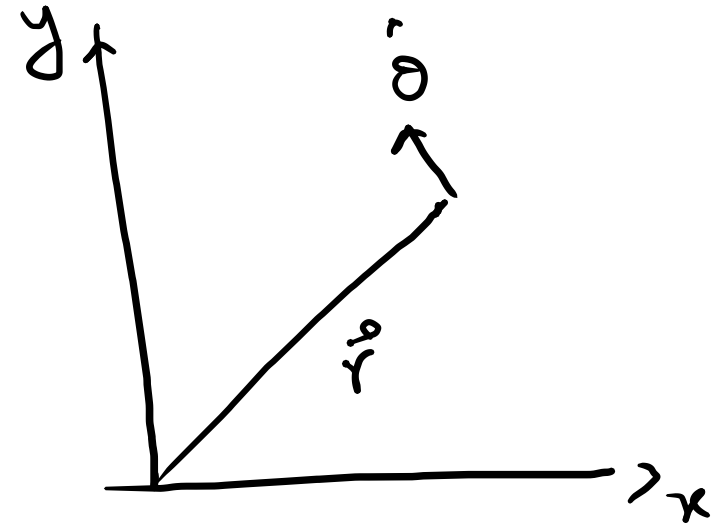
$$dW = \vec{F} \cdot d\vec{r}$$

$$= \frac{A}{r} r d\theta$$

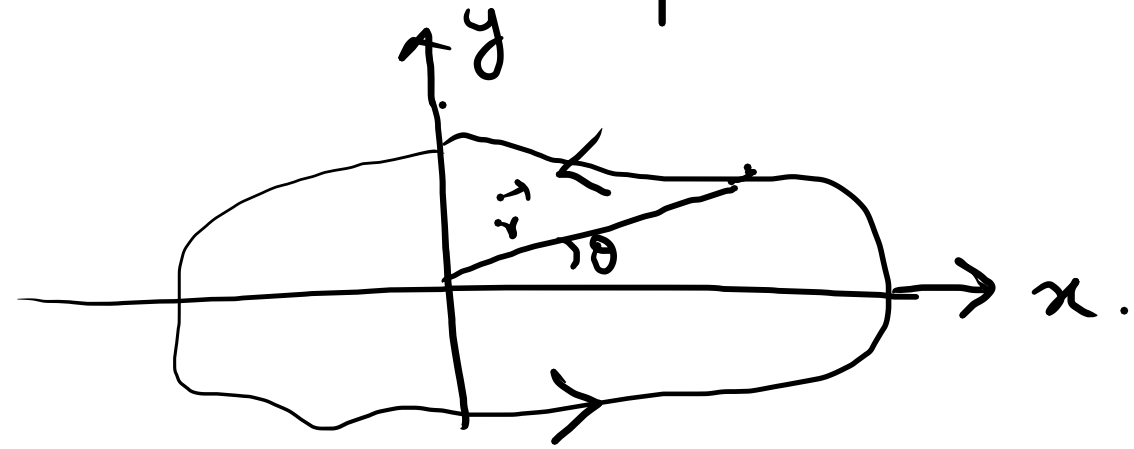
$$= A d\theta$$

$$W = \int_1^2 dW + \int_2^1 dW$$

$$= A(\theta_2 - \theta_1) + A(\theta_1 - \theta_2) = 0$$



Another closed path

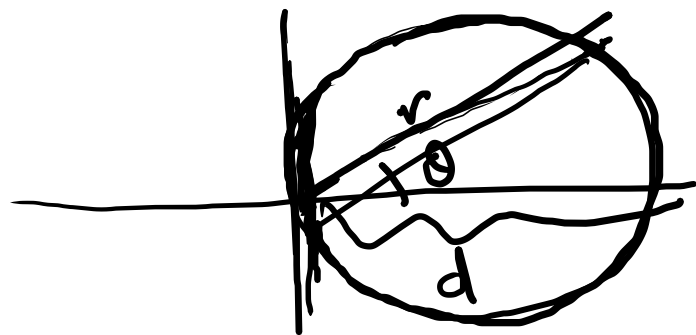


$$W = A \int_0^{2\pi} d\theta = 2\pi A \neq 0.$$

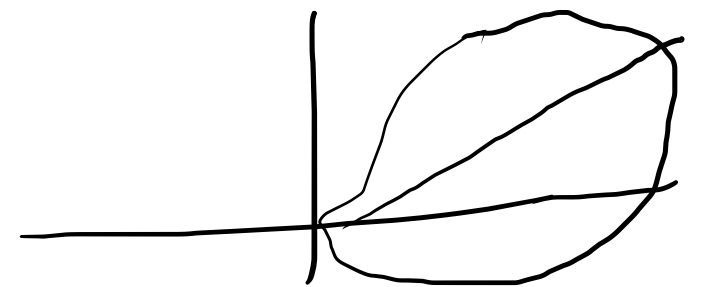
$$\vec{F} = \frac{A}{r} \hat{\theta}$$

$\hookrightarrow$  not defined at the origin.

So, any closed path containing the origin is treated differently from closed path not enclosing the origin.



$$r = d \cos \theta$$



# OSCILLATORY MOTION

$$\ddot{x} + \omega^2 x = 0$$

$$x(t) = C_1 e^{i\omega t} + C_2 e^{-i\omega t}$$

$$= B_1 \cos \omega t + B_2 \sin \omega t$$

$$= A \left( \frac{B_1}{A} \cos \omega t + \frac{B_2}{A} \sin \omega t \right)$$

$$= A (\cos \omega t \cos \delta + \sin \omega t \sin \delta)$$

$$= A \cos(\omega t - \delta)$$

↪ phase.

