

# LECTURE 15

## Recap:

- Solved the eqn. for the path/orbit under a central force with few convenient forms of  $V(r)$  /  $F(r)$ .

Forms of  $V(r)$  /  $F(r)$  dictated by whether the eqn. for orbit is integrable or not.

- Consider general  $V(r)$  with extremum at  $r = r_0$ .

$$V(r) \approx \underbrace{V(r_0)}_{\text{const}} + (r-r_0) \underbrace{\frac{dV}{dr} \Big|_{r=r_0}}_0 + \frac{1}{2} (r-r_0)^2 \underbrace{\frac{d^2 V}{dr^2} \Big|_{r=r_0}} + \dots$$

Near  $r = r_0$ ,

$$V(r) \approx \frac{1}{2} k (r - r_0)^2$$

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$$k = \left. \frac{d^2 V}{dr^2} \right|_{r=r_0}.$$

☐ Try to work out some facts about orbit under a general force field having an extremum.

$$\vec{F} = f(r) \hat{r}$$

→ In the context of a central force  $V \equiv V_{\text{eff}}(r)$ .

$$\omega_r = \sqrt{\frac{k}{m}} = \sqrt{V_{\text{eff}}''(r_0)}$$

EOM :-

$$\ddot{r} - r \dot{\theta}^2 = f(r)$$

$$\Rightarrow \ddot{r} - r \frac{L^2}{r^4} = f(r)$$

$$\Rightarrow \ddot{r} - \frac{L^2}{r^3} = f(r).$$

Sub :-

$$-L^2 u^2 \frac{d^2 u}{d\theta^2} - L^2 u^3 = f(1/u)$$

$$\Rightarrow \boxed{\frac{d^2 u}{d\theta^2} + u = -\frac{f(1/u)}{L^2 u^2}}$$

$$L = r^2 \dot{\theta} \Rightarrow L u^2 = \dot{\theta}$$

$$r = \frac{1}{u}$$

$$\begin{aligned} \dot{r} &= -\frac{1}{u^2} \dot{u} = -\frac{\dot{\theta}}{u^2} \frac{du}{d\theta} \\ &= -L \frac{du}{d\theta} \end{aligned}$$

$$\ddot{r} = -L^2 u^2 \frac{d^2 u}{d\theta^2}$$

$\Rightarrow$  2<sup>nd</sup> order ODE for orbit.

— Physically, presence of extremum in the context of central force implies presence of circular orbit: Since, small deviations from the extremum  $O[(r-r_0)^2]$  are considered, geometrically, this corresponds to small deviations/perturbations from circular orbit.

— Particle moving in a circular orbit of radius  $a$  under some attractive central force.  $\Rightarrow -f(r)$ .

$$\frac{v^2}{a} = f(a)$$

$$L = va \Rightarrow L^2 = a^3 f(a).$$

Radius of the circular orbit =  $a$ .  
 $r \equiv a + \epsilon(t)$  where  $\epsilon$  is small.

$$u = \frac{1}{a} (1 + \epsilon(\theta)), \quad \epsilon \text{ is small.}$$

Eqn. for orbit,

$$\frac{d^2 u}{d\theta^2} + u = - \frac{f(1/u)}{L^2 u^2}$$

$$\text{Sub: } u = \frac{1}{a} (1 + \epsilon(\theta))$$

$$\begin{aligned} \frac{1}{a} \frac{d^2 \epsilon}{d\theta^2} + \frac{1}{a} (1 + \epsilon(\theta)) &= \frac{f\left(\frac{a}{1+\epsilon(\theta)}\right)}{L^2 u^2} \\ \Rightarrow \frac{d^2 \epsilon}{d\theta^2} + 1 + \epsilon &= \frac{a}{L^2 u^2} \frac{f\left(\frac{a}{1+\epsilon}\right)}{f\left(\frac{a}{1+\epsilon}\right)} = \frac{a^3}{L^2} \frac{(1+\epsilon)^{-2} f\left(\frac{a}{1+\epsilon}\right)}{(1+\epsilon)^{-2} f\left(\frac{a}{1+\epsilon}\right)} \\ &= \frac{a^3}{L^2} \frac{f(a)}{f(a)}. \end{aligned}$$

$$\begin{aligned}
 f\left(\frac{a}{1+\xi}\right) &= f\left(a - \frac{a\xi}{1+\xi}\right) \\
 &= f(a) - \frac{a\xi}{1+\xi} f'(a) + O\left(\frac{\xi}{1+\xi}\right)^2 \\
 &\approx f(a) - a f'(a) \xi + O(\xi^2)
 \end{aligned}$$

$$(1+\xi)^{-2} \approx 1 - 2\xi + O(\xi^2).$$

$$\begin{aligned}
 (1+\xi)^{-2} f\left(\frac{a}{1+\xi}\right) &\approx (1-2\xi) \left(f(a) - a f'(a) \xi\right) \\
 &\approx f(a) - a f'(a) \xi - 2\xi f(a).
 \end{aligned}$$

$$\frac{(1+\xi)^{-2} f\left(\frac{a}{1+\xi}\right)}{f(a)} = 1 - a \frac{f'(a)}{f(a)} \xi - 2\xi.$$

Eqn. for path :-

$$\frac{d^2 \xi}{d\theta^2} + 1 + \xi = 1 - a \frac{f'(a)}{f(a)} \xi - 2\xi$$

$$\Rightarrow \frac{d^2 \xi}{d\theta^2} + \left( 3 + \frac{a f'(a)}{f(a)} \right) \xi = 0$$

Case 1 :  $3 + \frac{a f'(a)}{f(a)} < 0$ .

$$\xi = ( ) e^{m_1 \theta} + ( ) e^{-m_1 \theta}$$

$\xi$  grows with time  $\Rightarrow$  orbit does not remain circular.

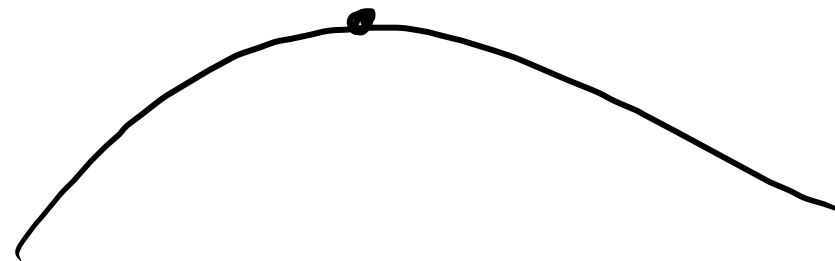
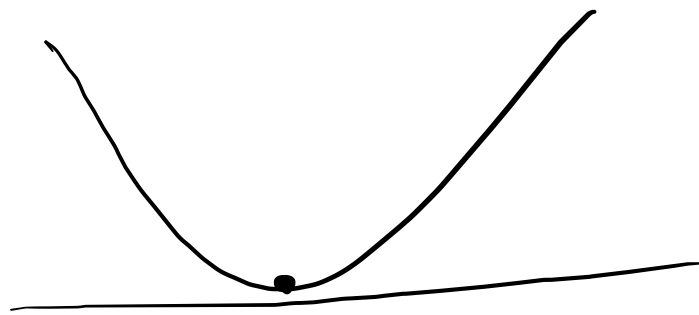
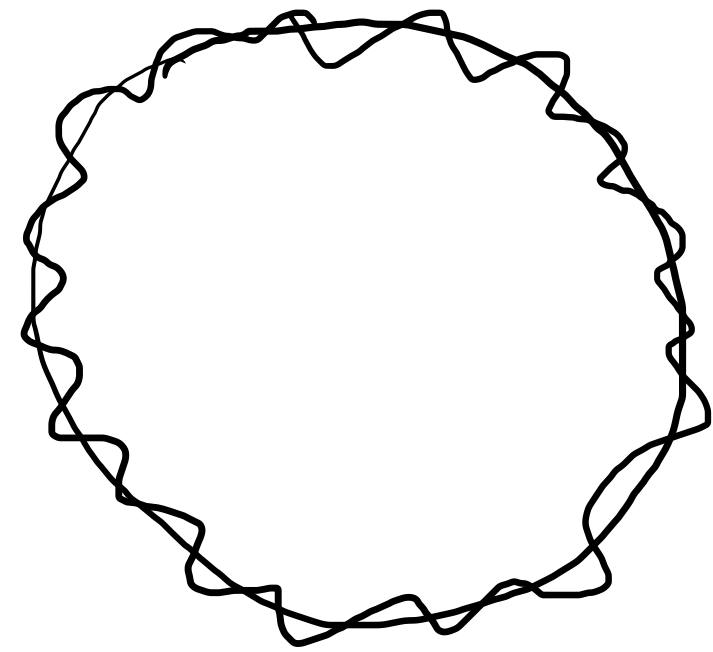
Case 2:-  $\Omega^2 = 3 + \frac{a f'(a)}{f(a)} > 0$ .

$$\frac{d^2 \xi}{d\theta^2} + \Omega^2 \xi = 0.$$

$$\xi = A \cos \Omega \theta + B \sin \Omega \theta$$

Circular orbit retains its shape.

For the perturbed orbit to close:-  $\Omega \rightarrow +ve$  integer.





— Consider  $f(r) = k r^\nu$ .

$$\Omega^2 = 3 + \frac{a f'(a)}{f(a)} = 3 + \frac{\cancel{a} \cancel{k} \nu a^{\nu-1}}{\cancel{k} a^\nu} = \nu + 3 > 0$$
$$\Rightarrow \nu > -3.$$

for closure of orbit.

$$\nu + 3 = m^2 \rightarrow \text{+ve integer,}$$

$$\Rightarrow \nu = m^2 - 3.$$

$$f(r) = k r^\nu$$

$$V(r) = - \int dr f(r).$$

$$V_{\text{eff}}(r) = \frac{L^2}{2r^2} + k r^\nu$$

$$\omega_r = \sqrt{V_{\text{eff}}''(a)}$$

$$\frac{d^2 \phi}{d\theta^2} + \omega_r^2 \phi^2 = 0.$$

Exercise!:- Stability analysis for  $V_{\text{eff}}'(r)$  leads to same results as linearisation of eqn. for orbit.

$$\int \frac{dx}{\sqrt{a^2 - x^2}}$$

$$\int \frac{dx}{\sqrt{x^2 - a^2}} \quad \checkmark$$