

PH 107 :Quantum Physics and Applications

Step potential

Lecture 16: 08-02-2022

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Recap (Finite Potential Well)

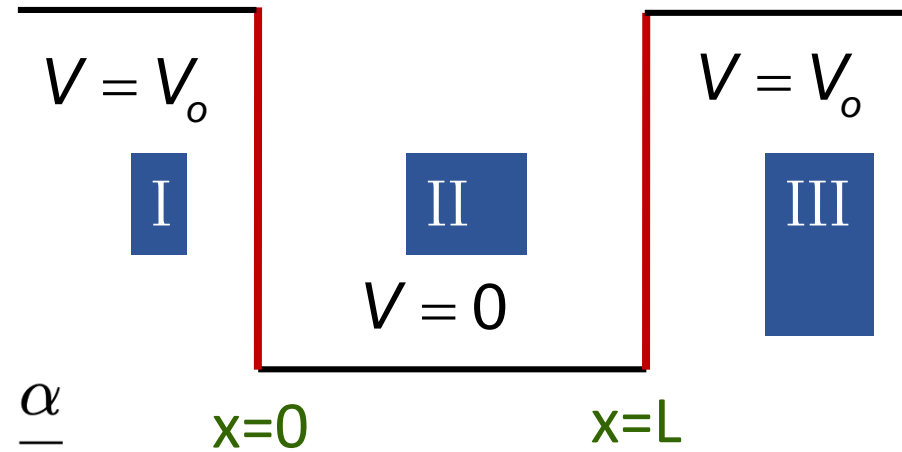
$$V(x) = 0 \quad \text{for } 0 < x < L$$

$$= V_o \quad \text{for } x < 0 \text{ or } x > L$$

$$\tan\left(\frac{kL}{2}\right) = \frac{\alpha}{k} \quad -\cot\left(\frac{kL}{2}\right) = \frac{\alpha}{k}$$

$$\tan\left(\frac{kL}{2}\right) = \sqrt{\left(\frac{k_o L/2}{kL/2}\right)^2 - 1}$$

$$\text{and } -\cot\left(\frac{kL}{2}\right) = \sqrt{\left(\frac{k_o L/2}{kL/2}\right)^2 - 1}$$



$$\text{Using } \frac{2m}{\hbar^2}(V_o - E) = \alpha^2 ;$$

$$\frac{2mE}{\hbar^2} = k^2$$

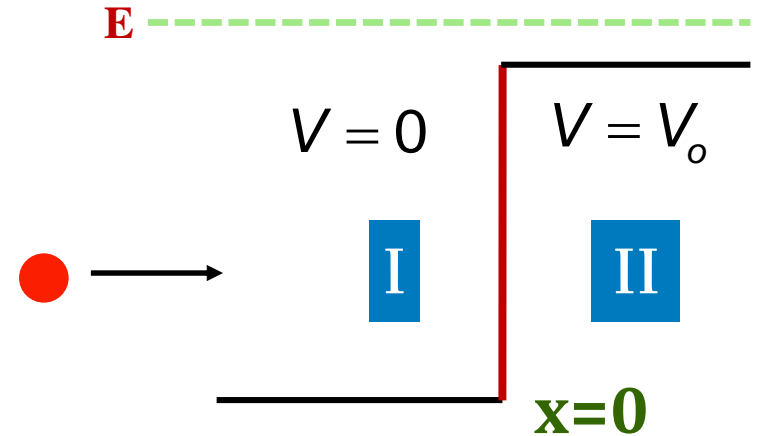
$$\text{and } k_o = \sqrt{\frac{2mV_o}{\hbar^2}}$$

Graphical intersection of LHS and RHS is the estimate of the allowed energy states.

Step potential

$$V(x) = 0 \quad \forall x \leq 0 \\ = V_0 \quad \forall x > 0$$

Consider, $E > V_0$



Classically, there will be total transmission of all particles and particle will have kinetic energy $E - V_0$. This is a simple *scattering* problem in 1-D.

Quantum mechanically, the dynamics is regulated by SE.

$$\varphi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}, \text{ where } k_1^2 = \frac{2mE}{\hbar^2}, \quad x < 0$$

$$\varphi_{II}(x) = Ce^{ik_2x} + De^{-ik_2x}, \text{ where } k_2^2 = \frac{2m(E-V_0)}{\hbar^2}, \quad x > 0$$

1. Since there is no incidence from the right side, in

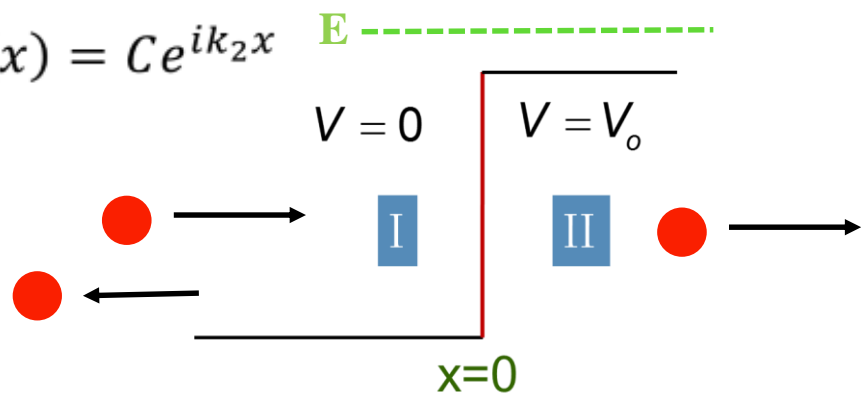
$$\varphi_{II}(x) = Ce^{ik_2x} + De^{-ik_2x} \quad \longrightarrow \quad D = 0$$

$$\longrightarrow \varphi_{II}(x) = Ce^{ik_2x}$$

2. Boundary conditions

$$(a) \varphi_I(0) = \varphi_{II}(0) \quad \Longrightarrow \quad A + B = C$$

$$(b) \varphi'_I(0) = \varphi'_{II}(0) \quad \Longrightarrow \quad ik_1(A - B) = ik_2C$$



Add

$$\Longrightarrow A = \frac{C}{2} \left(1 + \frac{k_2}{k_1} \right)$$

Divide

$$\Longrightarrow \frac{C}{A} = \frac{2k_1}{k_1 + k_2}$$

Subtract

$$\Longrightarrow B = \frac{C}{2} \left(1 - \frac{k_2}{k_1} \right)$$

Divide

$$\Longrightarrow \frac{B}{A} = \frac{k_1 - k_2}{k_1 + k_2}$$

The wave functions; $\varphi_I(x) = A \left(e^{ik_1x} + \left(\frac{k_1-k_2}{k_1+k_2} \right) e^{-ik_1x} \right)$, $x < 0$

in terms of A

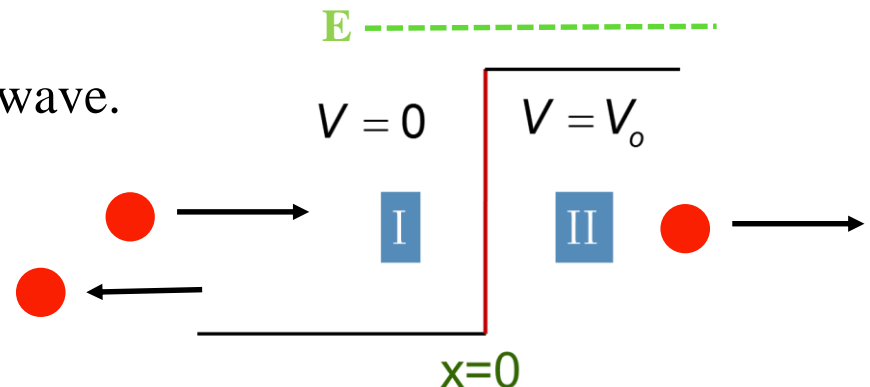
$$\varphi_{II}(x) = A \left(\frac{2k_1}{k_1 + k_2} \right) e^{ik_2x} \quad , x < 0$$

- This implies that the probability of the particle being **reflected** is non-zero. However **classically this is forbidden**.
- This effect is attributed to **wave like behavior of particles**.

Ae^{ik_1x} represents the incident wave

$A \left(\frac{k_1-k_2}{k_1+k_2} \right) e^{-ik_1x}$ represents the reflected wave

$A \left(\frac{2k_1}{k_1+k_2} \right) e^{ik_2x}$ represents the transmitted wave.

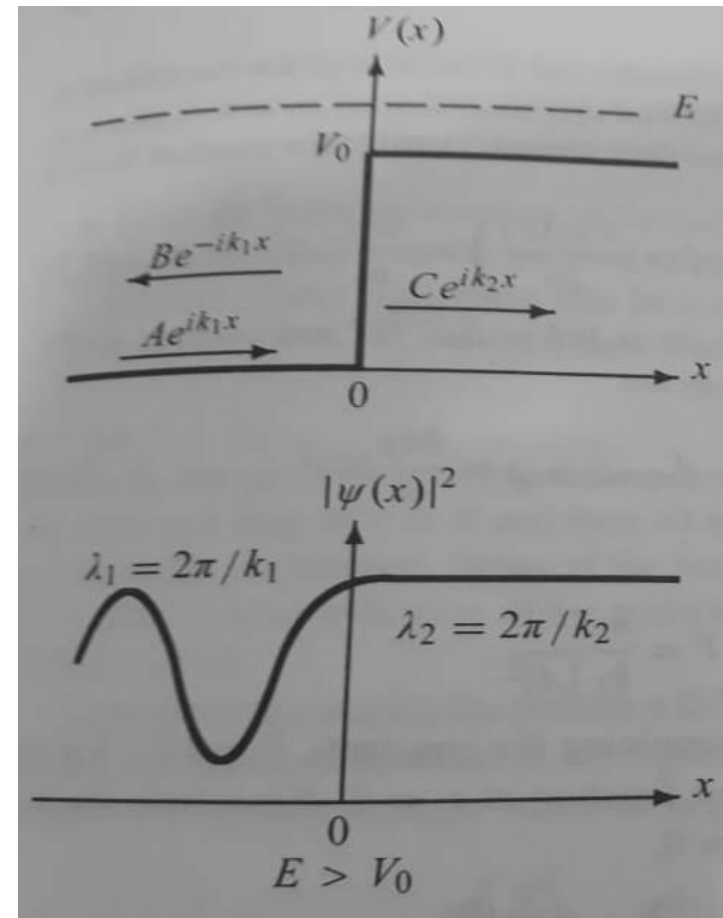
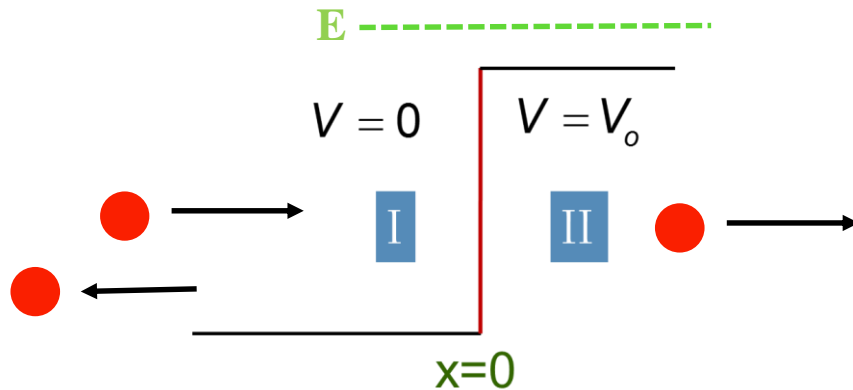


Probability Density, $E > V_0$

Ae^{ik_1x} represents the incident wave

$A \left(\frac{k_1 - k_2}{k_1 + k_2} \right) e^{-ik_1x}$ represents the reflected wave

$A \left(\frac{2k_1}{k_1 + k_2} \right) e^{ik_2x}$ represents the transmitted wave.



Case $E = V_0$

$$\varphi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}, \text{ where } k_1^2 = \frac{2mE}{\hbar^2}, \quad x < 0$$

$$\varphi_{II}(x) = Ce^{ik_2x}, \text{ where } k_2^2 = \frac{2m(E-V_0)}{\hbar^2}, \quad x > 0$$

$$\frac{B}{A} = \frac{k_1 - k_2}{k_1 + k_2}$$

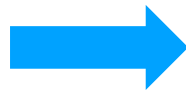


$$B=A$$

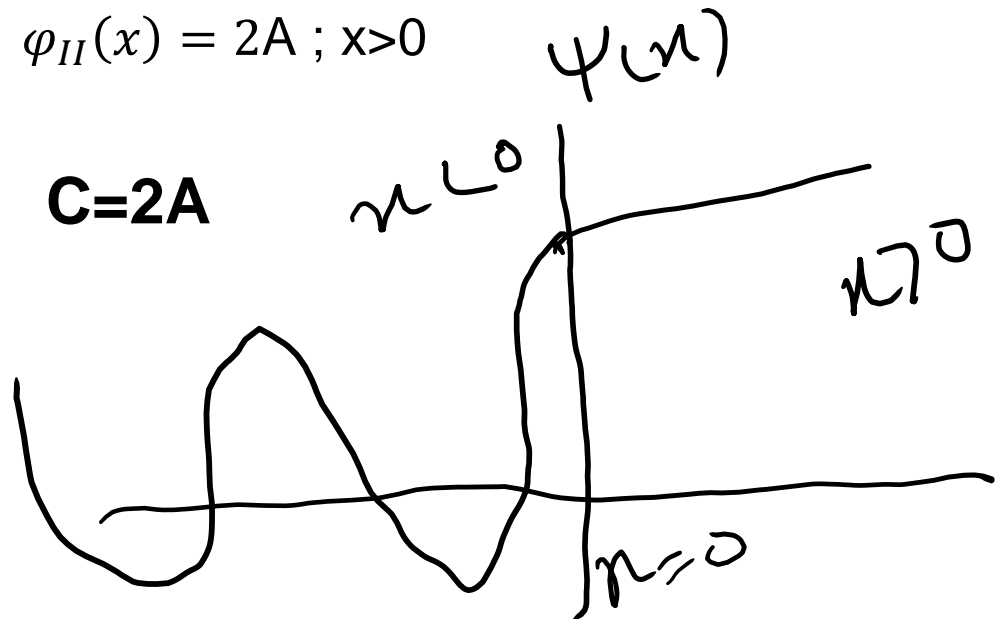
$$\varphi_I(x) = 2A \cos k_1x, \quad x < 0$$

$$\varphi_{II}(x) = 2A; \quad x > 0$$

$$\frac{C}{A} = \frac{2k_1}{k_1 + k_2}$$



$$C=2A$$



Probability current density

$$J(x, t) = \frac{\hbar}{2im} \left[\psi^*(x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right]$$

$J(x, t)$ is the current associated with charge density ρ

Derivation provided in lecture notes

$$\frac{\partial \rho}{\partial t} + \frac{\partial J}{\partial x} = 0$$

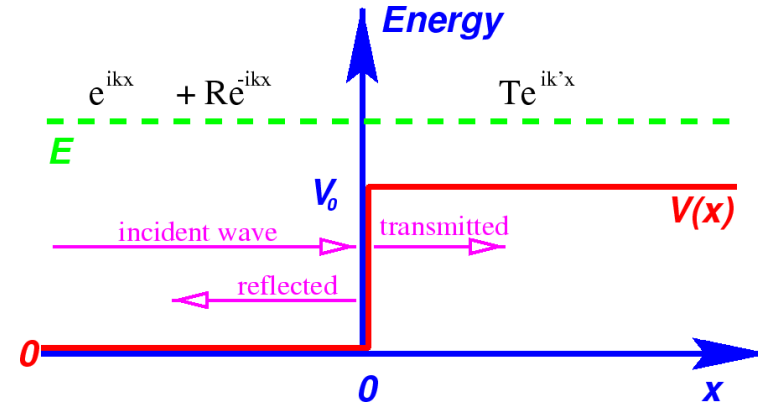
Change in charge is associated with current escaping/entering the volume.

Charge density in electromagnetism is equivalent to probability density in Quantum Mechanics.

Probability Current density

The current density/flux is defined as :

$$J = \frac{\hbar}{2im} \left[\psi^*(x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right]$$



$$\varphi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}, \text{ where } k_1^2 = \frac{2mE}{\hbar^2}, \quad x < 0$$

$$\varphi_{II}(x) = Ce^{ik_2x} + De^{-ik_2x}, \text{ where } k_2^2 = \frac{2m(E-V_0)}{\hbar^2}, \quad x > 0$$

$$j_{\text{incident}} = \frac{\hbar k_1}{m} |A|^2 \qquad j_{\text{reflected}} = \frac{-\hbar k_1}{m} |B|^2$$

$$j_{\text{transmitted}} = \frac{\hbar k_2}{m} |C|^2$$

Probability Current density

The rate at which the incident particles approach the barrier is $(\hbar k_1/m)|A|^2$.

The rate at which they are reflected is $(\hbar k_1/m)|B|^2$

and the rate at which they move forward is $(\hbar k_2/m)|C|^2$.

The factor of $\frac{k_2}{k_1}$ in T due to **different rate** at which the incident and transmitted particles move in region I and II.

$$j_{x<0} = \frac{\hbar k_1}{m} [|A|^2 - |B|^2]$$

$$j_{x>0} = \frac{\hbar k_2}{m} |C|^2$$

Conversation of Probability Current

$$j_{x<0} = \frac{\hbar k_1}{m} [|A|^2 - |B|^2] \qquad j_{x>0} = \frac{\hbar k_2}{m} |C|^2$$

These two currents density should be same.

$$\begin{aligned} \frac{\hbar k_1}{m} \left[1 - \frac{|B|^2}{|A|^2} \right] |A|^2 &= \frac{\hbar k_1}{m} \left[1 - \left(\frac{k_1 - k_2}{k_1 + k_2} \right)^2 \right] |A|^2 \\ &= \frac{\hbar k_1}{m} \left[\frac{4k_1 k_2}{(k_1 + k_2)^2} \right] |A|^2 \\ &= \frac{\hbar k_2}{m} \left[\frac{4k_1^2}{(k_1 + k_2)^2} \right] |A|^2 \\ &= \frac{\hbar k_2}{m} |C|^2 \end{aligned}$$

Conservation of probability current holds

Transmission and Reflection coefficients

Reflection coefficient,

$$R = \left| \frac{\text{reflected current density}}{\text{incident current density}} \right| = \left| \frac{J_{\text{reflected}}}{J_{\text{incident}}} \right|$$

Transmission coefficient,

$$T = \left| \frac{\text{transmitted current density}}{\text{incident current density}} \right| = \left| \frac{J_{\text{transmitted}}}{J_{\text{incident}}} \right|$$

To estimate the reflection and transmission coefficient, we should the wave function in the appropriate region.

$$\varphi_I(x) = A \left(e^{ik_1x} + \left(\frac{k_1 - k_2}{k_1 + k_2} \right) e^{-ik_1x} \right) \quad \varphi_{II}(x) = A \left(\frac{2k_1}{k_1 + k_2} \right) e^{ik_2x}$$

$$j_{incident} = \frac{\hbar k_1}{m} |A|^2 \quad j_{reflected} = \frac{-\hbar k_1}{m} |B|^2 \quad j_{transmitted} = \frac{\hbar k_2}{m} |C|^2$$

Reflection coefficient,

$$R = \left| \frac{B}{A} \right|^2 = \left(\frac{k_1 - k_2}{k_1 + k_2} \right)^2$$

Transmission coefficient,

$$T = \frac{k_2}{k_1} \left| \frac{C}{A} \right|^2 = \frac{4k_1 k_2}{(k_1 + k_2)^2}$$

- In contrast to classical mechanics, which states that none of the particles gets reflected, quantum mechanical reflection coefficient, $R \neq 0$.
- There are particles that get reflected inspite of $E > V_0$. This effect must be attributed to the **wavelike behavior** of the particles.
- The sum of Reflection and Transmission coefficient,

$$R + T = \left(\frac{k_1 - k_2}{k_1 + k_2} \right)^2 + \frac{4k_1 k_2}{(k_1 + k_2)^2} = 1$$

Transmission and Reflection coefficients, $E > V_0$

$$R = \left(\frac{k_1 - k_2}{k_1 + k_2} \right)^2 = \left(\frac{1 - \kappa}{1 + \kappa} \right)^2$$

$$T = \frac{4k_1k_2}{(k_1 + k_2)^2} = \frac{4\kappa}{(1 + \kappa)^2}$$

$$\text{where } \kappa = \frac{k_2}{k_1} = \sqrt{1 - \frac{V_0}{E}}$$

- For small value of E , T decreases.
- For $E = V_0$, $T = 0$ and $R = 1$.
- For $E \gg V_0$, $\kappa \sim 1$, hence $R = 0$ and $T = 1$.
- Particles with very high energies, the potential step is so weak that it produces no noticeable effect on their motion.

Continuity equations

Schrodinger Equation for $\psi(x, t)$

$$i\hbar \frac{\partial \psi(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} + V(x)\psi(x, t) \quad (1)$$

Schrodinger Equation for $\psi^*(x, t)$

$$i\hbar \frac{\partial \psi^*(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi^*(x, t)}{\partial x^2} + V(x)\psi^*(x, t) \quad (2)$$

Multiply 1 by $\psi^*(x, t)$ and 2 by $\psi(x, t)$

$$i\hbar \psi^*(x, t) \frac{\partial \psi(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \psi^*(x, t) \frac{\partial^2 \psi(x, t)}{\partial x^2} + \psi^*(x, t) V(x) \psi(x, t) \quad (3)$$

$$i\hbar\psi(x,t)\frac{\partial\psi^*(x,t)}{\partial t} = -\frac{\hbar^2}{2m}\psi(x,t)\frac{\partial^2\psi^*(x,t)}{\partial x^2} + \psi(x,t)V(x)\psi^*(x,t) \quad (4)$$

Subtract eq 3 from eq 4

$$i\hbar\left[\psi^*(x,t)\frac{\partial\psi(x,t)}{\partial t} + \psi(x,t)\frac{\partial\psi^*(x,t)}{\partial t}\right] = -\frac{\hbar^2}{2m}\left[\psi^*(x,t)\frac{\partial^2\psi(x,t)}{\partial x^2} - \psi(x,t)\frac{\partial^2\psi^*(x,t)}{\partial x^2}\right] \quad (5)$$

Since

$$\frac{\partial\psi^*(x,t)\psi(x,t)}{\partial t} = \psi^*(x,t)\frac{\partial\psi(x,t)}{\partial t} + \frac{\partial\psi^*(x,t)}{\partial t}\psi(x,t)$$

$$\begin{aligned}
& \frac{\partial}{\partial x} \left\{ \psi * (x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right\} \\
&= \frac{\partial \psi^*(x, t)}{\partial x} \frac{\partial \psi(x, t)}{\partial x} + \psi * (x, t) \frac{\partial^2 \psi(x, t)}{\partial x^2} - \frac{\partial \psi(x, t)}{\partial x} \frac{\partial \psi^*(x, t)}{\partial x} - \psi(x, t) \frac{\partial^2 \psi^*(x, t)}{\partial x^2} \\
&= \psi * (x, t) \frac{\partial^2 \psi(x, t)}{\partial x^2} - \psi(x, t) \frac{\partial^2 \psi^*(x, t)}{\partial x^2}
\end{aligned}$$

$$\frac{\partial}{\partial x} \left\{ \psi * (x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right\} = \psi * (x, t) \frac{\partial^2 \psi(x, t)}{\partial x^2} - \psi(x, t) \frac{\partial^2 \psi^*(x, t)}{\partial x^2}$$


(6)

From eqn 5

$$i\hbar \left[\psi^*(x, t) \frac{\partial \psi(x, t)}{\partial t} + \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial t} \right] = -\frac{\hbar^2}{2m} \left[\psi^*(x, t) \frac{\partial^2 \psi(x, t)}{\partial x^2} - \psi(x, t) \frac{\partial^2 \psi^*(x, t)}{\partial x^2} \right]$$



$$i\hbar \left[\frac{\partial \psi(x, t) \psi^*(x, t)}{\partial t} \right] = -\frac{\hbar^2}{2m} \frac{\partial}{\partial x} \left[\psi^*(x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right]$$


$$i\hbar \left[\frac{\partial \psi(x, t) \psi^*(x, t)}{\partial t} \right] + \frac{\hbar^2}{2im} \frac{\partial}{\partial x} \left[\psi^*(x, t) \frac{\partial \psi(x, t)}{\partial x} - \psi(x, t) \frac{\partial \psi^*(x, t)}{\partial x} \right] = 0$$

$\rho = \psi(x, t) \psi^*(x, t) = \text{probability}$

$$\frac{\partial \rho}{\partial t} + \frac{\partial J}{\partial x} = 0$$

$J(x,t)$ is the current associated with charge density ρ

$$J = \frac{\hbar}{2im} \left[\psi^*(x,t) \frac{\partial \psi(x,t)}{\partial x} - \psi(x,t) \frac{\partial \psi^*(x,t)}{\partial x} \right]$$