

## Fluid Mechanics (theory).

### ① Representation of frame

- Eulerian = looks at a point fixed in space

$$u = F_1(x, y, z, t) \quad v = F_2(x, y, z, t) \quad w = F_3(x, y, z, t)$$

[Easier to integrate]

- Lagrangian = looks at each particle in space  
Let  $(x_0, y_0, z_0)$  be coordinates of a chosen particle at  $t_0$ .

At a later time  $t$ ,

$$x = f_1(x_0, y_0, z_0, t) \quad y = f_2(x_0, y_0, z_0, t) \quad z = f_3(x_0, y_0, z_0, t)$$

$$\vec{v} = \text{velocity} = (u, v, w)$$

## ② Acceleration

$$u + du = u(x, y, z, t) + \frac{\partial u}{\partial x} \Delta x + \frac{\partial u}{\partial y} \Delta y + \frac{\partial u}{\partial z} \Delta z + \frac{\partial u}{\partial t} \Delta t$$

$$\lim_{\Delta t \rightarrow 0} \left( \frac{du}{dt} = \frac{\partial u}{\partial x} \frac{\Delta x}{\Delta t} + \frac{\partial u}{\partial y} \frac{\Delta y}{\Delta t} + \frac{\partial u}{\partial z} \frac{\Delta z}{\Delta t} + \frac{\partial u}{\partial t} \right)$$

$$\frac{du}{dt} = \underbrace{U \frac{\partial U}{\partial x} + V \frac{\partial U}{\partial y} + W \frac{\partial U}{\partial z}}_{\text{material derivative}} + \underbrace{\frac{\partial u}{\partial t}}_{\text{temporal acceleration}}$$

connective acceleration

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + U \frac{\partial}{\partial x} + V \frac{\partial}{\partial y} + W \frac{\partial}{\partial z}.$$

## \* Cylindrical

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + V_r \frac{\partial}{\partial r} + \frac{V_\theta}{r} \frac{\partial}{\partial \theta} + V_z \frac{\partial}{\partial z}$$

$$a_r = \frac{D V_r}{Dt} - \frac{V_\theta^2}{r}$$

$$a_\theta = \frac{D V_\theta}{Dt} + \frac{V_r V_\theta}{r}$$

$$a_z = \frac{D V_z}{Dt}$$

\* Spherical

$$a_r = \frac{\partial V_r}{\partial t} + V_r \frac{\partial V_r}{\partial r} + \frac{V_\theta}{r} \frac{\partial V_r}{\partial \theta} + \frac{V_\phi}{r \sin \theta} \frac{\partial V_r}{\partial \phi} - \frac{V_r^2 + V_\theta^2}{r}$$

$$a_\theta = \frac{\partial V_\theta}{\partial t} + V_r \frac{\partial V_\theta}{\partial r} + \frac{V_\theta}{r} \frac{\partial V_\theta}{\partial \theta} + \frac{V_\phi}{r \sin \theta} \frac{\partial V_\theta}{\partial \phi} + \frac{V_r V_\theta}{r} - \frac{V_\theta^2 \cot \theta}{r}$$

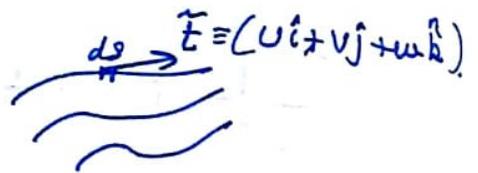
$$a_\phi = \frac{\partial V_\phi}{\partial t} + V_r \frac{\partial V_\phi}{\partial r} + \frac{V_\theta}{r} \frac{\partial V_\phi}{\partial \theta} + \frac{V_\phi}{r \sin \theta} \frac{\partial V_\phi}{\partial \phi} + \frac{V_r V_\phi}{r} + \frac{V_\theta V_\phi \cot \theta}{r}$$

$$\nabla \rightarrow \frac{\partial}{\partial r} + \frac{1}{r} \frac{\partial}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}$$

$$\rightarrow \frac{\partial}{\partial r} + \frac{1}{r} \frac{\partial}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}$$

Lines tangent at which, give velocity at a point at any time instant  $t$ .

$$|ds \times \vec{t}| = 0.$$



$$\frac{dx}{U} = \frac{dy}{V} = \frac{dz}{W}$$

$$ds = dx\hat{i} + dy\hat{j} + dz\hat{k}$$

- Cannot intersect  $\equiv$  as can't have 2 velocities at one point at same time

$$\Rightarrow \text{Surface orthogonal to streamlines} \equiv Udx + Vdy + Wdz = 0$$

- Path lines

Path of any particle  $(x_0, y_0, z_0 \text{ at } t_0)$  over entire time

$$\frac{dx}{dt} = U(x, y, z, t)$$

$$\frac{dy}{dt} = V(x, y, z, t)$$

$$\frac{dz}{dt} = W(x, y, z, t)$$

$$\begin{aligned} I.C.S. &\Rightarrow \\ \psi(x_0) &= x_0 \\ y(t_0) &= y_0 \\ z(t_0) &= z_0 \end{aligned}$$



Steady flow  $\equiv$  Pathline = Streamlines

- Velocity potential  
 $\phi$  is velocity potential if  $\vec{q} \rightarrow$  velocity vector  
 $\vec{q} = -\nabla\phi$
- $\hookrightarrow$  Existence if  $\nabla \times \vec{q} = 0$  (irrotational or conservative)
- $\phi(x, y, z, t) = C_i$  are equipotential surfaces for diff. C:s.
- For an irrotational flow of incompressible fluids  
 $\hookrightarrow \nabla^2 \phi = 0$  (Combining  $\nabla \times \vec{q} = 0$  & continuity  $\nabla \cdot \vec{q} = 0$ )

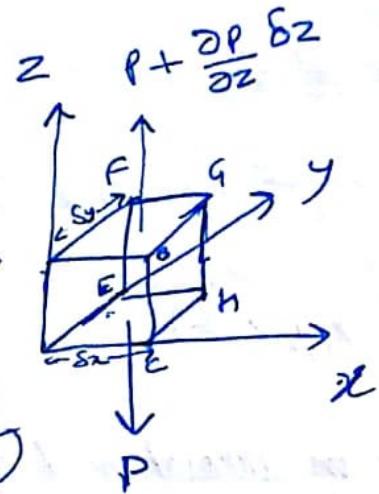
## ④ Euler Equations

$$I\text{-direction} = I_p \delta_x \delta_y \delta_z$$

$$m a_z = \rho \ddot{z} - \frac{\partial p}{\partial z} (\delta_x \delta_y \delta_z)$$

$$\rho \ddot{s}_z \delta_x \delta_y \delta_z = I_p \delta_x \delta_y \delta_z - \frac{\partial p}{\partial z} (\delta_x \delta_y \delta_z)$$

$$\boxed{\frac{Dw}{Dt} = I - \frac{1}{\rho} \frac{\partial p}{\partial z}}$$



## ⑤ Continuity Equation.

$$\text{Along } x = \text{Inflow} - \text{Outflow} = (\rho \delta_y \delta_z) u - (\cancel{\rho \delta_y \delta_z}) \left[ u + \frac{\partial u}{\partial x} \delta_x \right]$$

$$= (\rho \delta_y \delta_z) u - \left[ (\rho \delta_y \delta_z) u + \delta_x \cdot \frac{\partial}{\partial x} (\rho \delta_y \delta_z u) \dots \right]$$

$$= -\delta_x \delta_y \delta_z \frac{\partial (\rho u)}{\partial x}$$

$$\text{Along } y = -\delta_x \delta_y \delta_z \frac{\partial (\rho v)}{\partial y} \quad I = -\delta_x \delta_y \delta_z \frac{\partial (\rho w)}{\partial z}$$

$$x+y+z = \frac{\partial}{\partial t} (\rho \delta_x \delta_y \delta_z)$$

$$\text{Equating} \Rightarrow \boxed{\frac{\partial p}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} = 0}$$

$$= \frac{DP}{Dt} + \rho \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) = 0$$

General  
Continuity  
Equation

For homogeneous, incompressible fluids

$$\boxed{\frac{\partial p}{\partial t} = 0}$$

$$\boxed{\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0}$$

## Problems

(Q)  $\mu = x+y+2t$      $v = 2y+t$ . Determine Lagrange coordinates  $x_0, y_0$  at time  $t$ .

$$\frac{dx}{dt} = x+y+2t \quad \frac{dy}{dt} - 2y = t \Rightarrow e^{-\int 2 dt} = e^{-2t}$$

$$y \cdot e^{-2t} = \int te^{-2t} + C = -\frac{1}{2}te^{-2t} - \frac{1}{4}e^{-2t} + C$$

$$\frac{dx}{dt} - x = Ce^{2t} + \frac{1}{4}(6t+1) \quad y = Ce^{2t} - \frac{(2t+1)}{4}$$

$$x \cdot e^{-t} = C_2 + Ce^t + \int \frac{(6t-1)}{4} e^{-t} dt \Rightarrow$$

$$x = C_2 e^t + Ce^{2t} - \frac{(6t+5)}{4}$$

$$\text{Using } x_0, y_0 \text{ at } t=0 \Rightarrow y_0 = C - V_4 \quad x_0 = C_1 + C - \frac{5}{4}$$

$$C = y_0 + V_4 \quad C_1 = x_0 - y_0 + 1$$

$$x = (x_0 - y_0 + 1)e^t + (y_0 + V_4)e^{2t} - \frac{(6t+5)}{4}$$

$$y = (y_0 + V_4)e^{2t} - \frac{(2t+1)}{4}.$$

(Q) Velocity  $= (\gamma^2 z \cos \theta, \gamma z \sin \theta, z^2 t)$ . Find acceleration.

$$\frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{\gamma} \frac{\partial v_r}{\partial \theta} + v_z \frac{\partial v_r}{\partial z} - \frac{v_\theta^2}{\gamma} = a_r$$

$$a_r = 2\gamma^3 z^2 \cos^2 \theta + z \sin \theta (-\gamma^2 z \sin \theta) + z^2 t (\gamma^2 \cos \theta) - \gamma z^2 \sin^2 \theta$$

$$a_\theta = \gamma^2 z \cos \theta, z \sin \theta + z \sin \theta (\gamma z \cos \theta) + z^2 t (\gamma \sin \theta) + \gamma^2 z^2 \sin \theta \cos \theta.$$

$$a_z = z^2 + z^2 t (2z t)$$

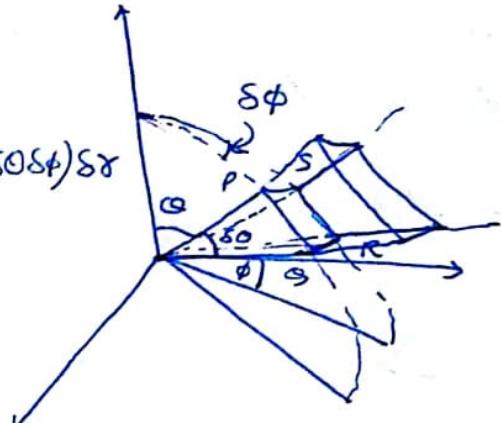
$-\frac{V_0^2}{\gamma}$	$a_r$
$+\frac{V_r V_\theta}{\gamma}$	$a_\theta$

## ④ Continuity for Spherical Coordinates.

$$PQRS \Rightarrow (\rho \sin \theta \sin \phi u_r, \rho \sin \theta \cos \phi u_\theta, \rho \cos \theta u_\phi)$$

$$P'Q'R'S' \Rightarrow (\rho \sin \theta \sin \phi u_r, \rho \sin \theta \cos \phi u_\theta) - \frac{\partial}{\partial r} (\rho^2 \mu_r \sin \theta \cos \phi) \delta r$$

$$\text{Net} \Rightarrow - \frac{\partial}{\partial r} (\rho^2 \sin \theta \mu_r) \delta \theta \delta r$$

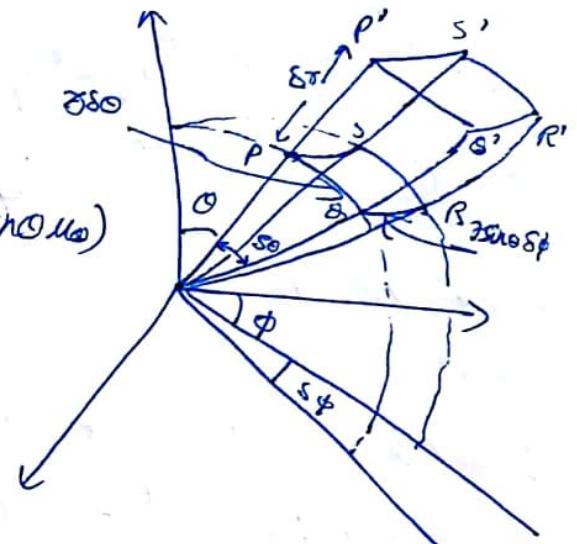


On other faces  $\Rightarrow$

$$\text{Net} = \frac{\partial}{\partial r} (\rho \sin \theta \sin \phi) \delta \theta \delta \phi$$

Equating we get,

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (\rho r^2 u_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\rho \sin \theta u_\theta) \\ + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (\rho u_\phi) = 0. \end{aligned}$$

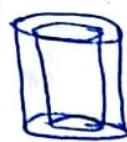


## ⑤ Symmetrical forms of Continuity Equations.

Cylindrical  
(Take  $h=1$ )

$$\frac{\partial}{\partial t} (\rho 2\pi r s_r) = \rho q_r (2\pi r) - \left[ \rho q_r 2\pi r + s_r \frac{\partial}{\partial r} (\rho q_r 2\pi r) \right]$$

$$2\pi r \left( \frac{\partial \rho}{\partial t} \right) s_r = -s_r \frac{\partial}{\partial r} (\rho q_r 2\pi r)$$



$$\text{If } \rho \text{ is constant} \Rightarrow \frac{\partial}{\partial r} (\rho q_r) = 0 \rightarrow \rho q_r = \rho g(t)$$

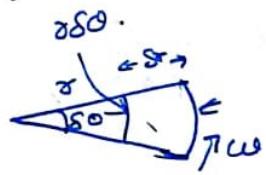
$$\Rightarrow q_r = g(t)$$

$$\text{If steady flow} \Rightarrow q_r = C$$



(2) Fluid moves such that each particle describes circle in a plane moving with  $\omega$ . Find continuity equation.

$$\frac{\partial}{\partial t} (p \tau s \delta s) = - \tau s \frac{\partial (p \tau \omega \delta s)}{\partial s}$$



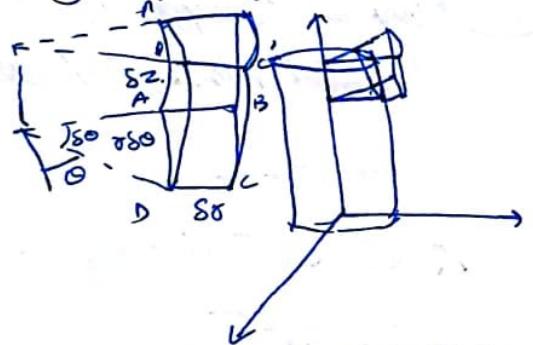
$$\boxed{\frac{\partial p}{\partial t} = - \frac{\partial (p\omega)}{\partial s}} \quad \text{Ans}$$

(3) Mass of fluid so that lines of motion lie on surface of co-axial cylinders. Find Continuity Equation.  
 $v, u$  are parallel to  $z$

No flow across  $A'B'D'A$

$$ABCD = -s_z \frac{\partial}{\partial z} (p v \tau s \delta s)$$

$$D'C'CD = -\tau s \frac{\partial}{\partial s} (p u \delta s \delta z).$$



$$\frac{\partial}{\partial t} (p \tau s \delta s \delta z) + s_z \frac{\partial}{\partial z} (p v \tau s \delta s) + \tau s \frac{\partial}{\partial s} (p u \delta s \delta z) = 0.$$

$$\boxed{\tau \frac{\partial p}{\partial t} + \tau \frac{\partial (p v)}{\partial z} + \frac{\partial (p u)}{\partial s} = 0} \quad \text{Ans}$$



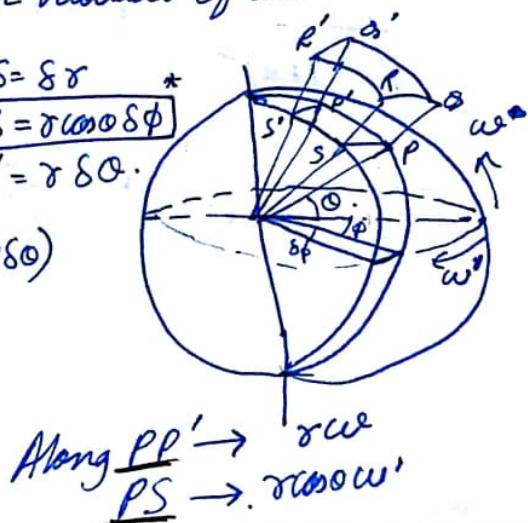
(4) Particle moves on surface of sphere. Given  $p$ ,  $\theta$  (lat.),  $\phi$  (longitude) of any element,  $w$  and  $w'$  angular velocities of element in latitude and longitude.

$$\rightarrow \text{Along } PP' = -\tau s \frac{\partial}{\partial s} (\tau w \cdot \delta r \cdot \tau \cos \phi) \quad \boxed{PS = \tau \cos \phi \delta \phi} \quad PS = \tau \cos \phi \delta \phi$$

$$\text{Along } PS = -\frac{\partial \cos \phi}{\partial s} \delta \phi \quad (P \cdot \tau \cos \phi \cdot \delta r \cdot \tau s)$$

$$\text{Equate as usual} = \frac{\partial}{\partial t} (\dots)$$

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(Q) Show continuity equation is satisfied

$$u = \frac{k(x^2 - y^2)}{(x^2 + y^2)^2} \quad v = \frac{2kxy}{(x^2 + y^2)^2}$$

$$\text{TS: } \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

$$\frac{2kx}{(x^2 + y^2)^2} - 2 \cdot \frac{k(x^2 - y^2) \cdot 2x}{(x^2 + y^2)^3} + \frac{2kx}{(x^2 + y^2)^2} - \frac{8kxy^2}{(x^2 + y^2)^3}$$

$$= 0 \quad \text{Hence } \underline{\text{Possible}}$$

(Q) Flow by a pipe  $y = a + \left(\frac{kx^2}{a}\right)$  from  $-a$  to  $a$ .

If initial velocity  $= V$  at  $x = -a$ , find time taken.

$$A_1 v_1 = A_2 v_2$$

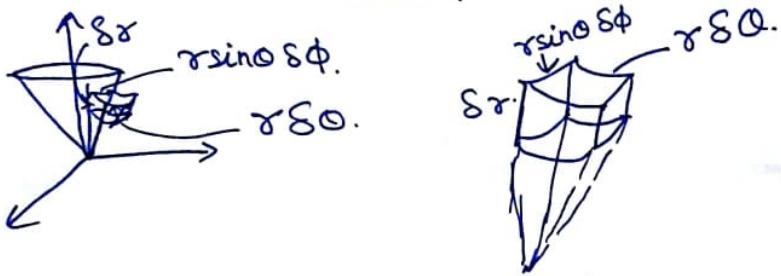
$$\frac{\pi (a+ka)^2 V}{\pi (a+\frac{kx^2}{a})^2} = V_x = \frac{dx}{dt}$$

$$\int \frac{dx}{V(a+ka)^2} \left( a^2 + \frac{k^2 x^4}{a^2} + 2kx^2 \right) = \int dt$$

$$t = \frac{1}{Va^2(1+k)^2} \left[ a^2 x + \frac{k^2 x^5}{5a^2} + \frac{2kx^3}{3} \right]_a^{a}$$

$$= \frac{2}{V(1+k)^2 a^2} \left[ a^3 + \frac{k^2}{5} a^3 + \frac{2}{3} k a^3 \right] = \frac{2a}{V(1+k)^2} \left[ 1 + \frac{2k}{3} + \frac{k^2}{5} \right]$$

Q) Flow on co-axial cones.

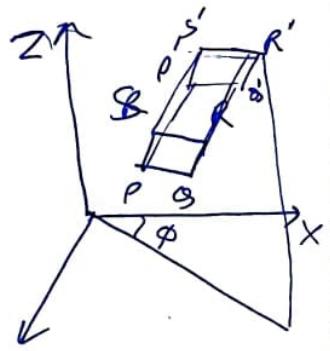


Q) Each particle of liquid ~~does~~ moves in a plane through z-axis.

PQRS lies in given plane.

$$PP' \equiv \text{no fluid along this} = \underline{v_r \sin \theta \phi}$$

$$\frac{PQ}{u} = \delta x, \quad \frac{PS}{v} = \delta y$$



$u$  and  $v$  are velocity along them.

→ Now trivial construction of Continuity Equation

## ⑤ Boundary conditions.

- At any point on a solid surface inside a liquid.  
if normal to surface be  $\hat{n}$  and  $\vec{v}$  be velocity of fluid.  
Then,  $\vec{v} \cdot \hat{n} = 0$ . (at all surface points)

If body is moving with  $\vec{U}$  then  $\vec{v} \cdot \hat{n} = \vec{U} \cdot \hat{n}$  or  $\vec{v} - \vec{U} \cdot \hat{n} = 0$

Condition for a surface to be a boundary surface. (F)

$$\frac{\partial F}{\partial t} + q_v \cdot \nabla F = 0$$

$$q_v = \vec{V}$$

$$\frac{DF}{Dt} = 0.$$

\* Normal velocity of boundary =  $\bar{u} \cdot \hat{n} = \bar{u} \cdot \frac{\nabla F}{|\nabla F|}$ .

IMPORTANT

$$\hat{n} = \frac{\nabla F}{|\nabla F|}$$

## ⑤ Rotational flow.

- Vorticity =  $\boxed{\Omega = \text{curl } \vec{q}}$

L Vortex line = Line at tangent to it at each point  
is in the direction of  $\Omega$  vector

$$\frac{dx}{\Omega_x} = \frac{dy}{\Omega_y} = \frac{dz}{\Omega_z}$$

- Angular velocity =  $\boxed{\omega = \frac{1}{2} \Omega = \frac{1}{2} \text{curl } \vec{q}}$

$$\cdot \frac{\partial \phi}{\partial x} = -\Omega \quad \rightarrow \quad \frac{\partial \phi}{\partial y} = -\Omega.$$

- Streamlines show each particle at an instant
- Pathlines show a particle at all times
- Streaklines look at a particular point in space

**Q** Show that the surface  $\frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} = 1 + kt^2 \left( \frac{y^2}{b^2} + \frac{z^2}{c^2} \right)$  is a possible boundary surface.

$$-\frac{4x^2}{a^2 k^2 t^4} + 2kt \left[ \frac{y^2}{b^2} + \frac{z^2}{c^2} \right] + \mu \left[ \frac{2x}{a^2 k^2 t^4} \right] + v \left[ \frac{2kt^2 y}{b^2} \right] + w \left[ \frac{2kt^2 z}{c^2} \right] = 0$$

$$\text{Putting } u = \frac{-2x}{t}, \quad v = -\frac{y}{t}, \quad w = \frac{-z}{t}$$

the equation is identically satisfied.

Now, check for  $(u, v, w)$  if continuity equation holds

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0. \quad \text{Hence possible}$$

**Q** Show  $\frac{x^2 \tan^2 t}{a^2} + \frac{y^2 \cot^2 t}{b^2} = 1$  is possible boundary and find normal velocity.

$$2x \frac{\tan^2 t \sec^2 t}{a^2} - \frac{2y^2 \cot t \csc^2 t}{b^2} + \mu \left( \frac{2x \tan^2 t}{a^2} \right) + v \left( \frac{2y \cot^2 t}{b^2} \right) = 0$$

$$\mu = -\frac{x \sec^2 t}{\tan^2 t} \quad v = \frac{y \csc^2 t}{\cot t}$$

$$-\frac{\sec^2 t}{\tan^2 t} + \frac{\csc^2 t}{\cot t} \Rightarrow -\frac{\cot t}{\sin t \cos^2 t} + \frac{\sin t}{\cos t \sin^2 t} = 0.$$

Hence possible.

$$\vec{n} = \frac{1}{q} \cdot \left( \frac{2x \tan^2 t}{a^2} \hat{i} + \frac{2y \cot^2 t}{b^2} \hat{j} \right)$$

$$= \frac{1}{q} \left( \left( \frac{2x \tan^2 t}{a^2} \right)^2 + \left( \frac{2y \cot^2 t}{b^2} \right)^2 \right)^{-\frac{1}{2}}$$

$$U = \frac{x}{1+t} \quad V = \frac{y}{2+t} \quad W = \frac{z}{3+t}$$

Stream line = Pathline

$$\frac{dx(1+t)}{x} = \frac{dy(2+t)}{y} \Rightarrow \frac{dx}{x} - \frac{2dy}{y} = t \left( \frac{dy}{y} - \frac{dx}{x} \right)$$

$$\Rightarrow \log x - 2\log y = t(\log y - \log x) + \log c$$

$$\log \left( \frac{x}{y^2} \right) = \log \left( c \frac{y}{x} \right)^t \Rightarrow \boxed{\frac{x}{y^2} = \left( \frac{y}{x} \right)^t} \quad (1)$$

$$\frac{d\ln(x^2+y^2)}{2y} = \frac{d\ln(2+t)}{2} \Rightarrow 2\frac{dx}{y} - \frac{3dx}{2} = t \left( \frac{dy}{2} - \frac{dx}{y} \right)$$

$$\Rightarrow \log(x^2+y^2) = t \log(2+t) + \log c.$$

$$\boxed{\frac{y^2}{x^2} = \left( \frac{y}{x} \right)^t} \Rightarrow \boxed{\frac{y^2}{z^3} = c_2 \left( \frac{y}{z} \right)^t} \quad (2)$$

Interaction of (1) and (2) shows streamlines.

→ Path lines

$$\frac{dx}{dt} = \frac{x}{1+t} \Rightarrow x = c_3(1+t).$$

$$y = c_4(2+t)$$

$$z = c_5(3+t)$$

Streamlines and pathlines are same in case of steady motion.

- Pathlines  $\Rightarrow x = \frac{x_0}{1+t_0}(1+t), y = \frac{y_0}{2+t_0}(2+t), z = \frac{z_0}{3+t_0}(3+t)$ .  
 $(x_0, y_0, z_0)$  at  $t_0$

(A)

- Streak lines → let particle  $x_0, y_0, z_0$  pass through  $x_1, y_1, z_1$  at  $t = S$  where  $t_0 \leq S \leq t$
- Use pathlines.

$$\rightarrow \frac{x_1(1+t_0)}{(1+S)}, \frac{y_1(2+t_0)}{2+S}, \frac{z_1(3+t_0)}{3+S} \equiv x_1, y_1, z_1$$

Substitute in (A)  $\Rightarrow (x_1, y_1, z_1) \equiv \left[ \frac{x_0(1+t)}{1+S}, \frac{y_0(2+t)}{2+S}, \frac{z_0(3+t)}{3+S} \right]$

Ex. 11 (2.5u)

In steady motion of homogeneous liquid in surfaces  $f_1 = a_1, f_2 = a_2$ , define streamlines. Prove that most general values of velocity components  $u, v, w$  are

$$F(f_1, f_2) \frac{\partial(f_1, f_2)}{\partial(y, z)}, F(f_1, f_2) \frac{\partial(f_1, f_2)}{\partial(z, x)}, F(f_1, f_2) \frac{\partial(f_1, f_2)}{\partial(x, y)}.$$

$$\rightarrow f_1 = a_1, df_1 = 0 \Rightarrow \frac{\partial f_1}{\partial x} dx + \frac{\partial f_1}{\partial y} dy + \frac{\partial f_1}{\partial z} dz = 0$$

$$f_2 = a_2, df_2 = 0 \Rightarrow \frac{\partial f_2}{\partial x} dx + \frac{\partial f_2}{\partial y} dy + \frac{\partial f_2}{\partial z} dz = 0$$

$$\frac{dx}{f_{1y}f_{2z} - f_{1z}f_{2y}} = \frac{dy}{f_{1z}f_{2x} - f_{1x}f_{2z}} = \frac{dz}{f_{1x}f_{2y} - f_{1y}f_{2x}}$$

$$\text{or } \frac{dx}{J_1} = \frac{dy}{J_2} = \frac{dz}{J_3}; J_1 = \frac{\partial(f_1, f_2)}{\partial(y, z)}, J_2 = \frac{\partial(f_1, f_2)}{\partial(z, x)}, J_3 = \frac{\partial(f_1, f_2)}{\partial(x, y)}$$

$$- \text{Streamlines} \Rightarrow \frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w} \Rightarrow u = F J_1, v = F J_2, w = F J_3$$

where  $F$  is arbitrary.

$$- \text{Continuity} \equiv \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \Rightarrow F(J_{1x} + J_{2y} + J_{3z}) +$$

$$\boxed{F(J_{1x} + J_{2y} + J_{3z} = 0)}$$

$$J_1 F_x + J_2 F_y + J_3 F_z = 0 \Rightarrow \frac{\partial(F_J, f_1, f_2)}{\partial(x, y, z)} = 0$$

$F$  is dependent  $f_1, f_2$ . Hence, Proved

$\Rightarrow F$  is a function of  $f_1, f_2$  only. Hence, Proved

(Q) For  $\varrho = \frac{k^2(x_j - y_i)}{x^2 + y^2}$ . Show possible motion  
find streamlines  
velocity potential units. find it.

$$\rightarrow \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0.$$

$$\frac{k^2 y \cdot 2x}{(x^2 + y^2)^2} + \frac{-2y \cdot x}{(x^2 + y^2)} = 0. \quad \checkmark$$

$$\rightarrow \frac{dx}{(-y)} = \frac{dy}{x} \Rightarrow x^2 + y^2 = C_2 \quad z = C_1$$

$$\rightarrow \nabla \times \vec{q} = \begin{bmatrix} 0 & 0 & k \\ \frac{\partial}{\partial x} & 0 & -\frac{\partial}{\partial z} \\ -\frac{\partial^2 y}{\partial x^2} & \frac{k^2 x}{x^2 + y^2} & 0 \end{bmatrix} \rightarrow i(0) + j(0) + k \left( \frac{k^2}{x^2 + y^2} - \frac{2k^2 x^2}{(x^2 + y^2)^2} - \frac{k^2 y^2}{x^2 + y^2} + f \right) = 0 \quad \text{Hence potential found}$$

$$\rightarrow \frac{\partial \phi}{\partial x} = -\frac{2k^2 x}{x^2 + y^2} \quad \frac{\partial \phi}{\partial y} = -\frac{2k^2 y}{x^2 + y^2} \quad \frac{\partial \phi}{\partial z} = 0$$

$$\rightarrow \frac{\partial \phi}{\partial x} = \frac{y k^2}{x^2 + y^2} \quad \rightarrow \phi(x, y) = k^2 \tan^{-1}(xy) + f(y).$$

$$\rightarrow \frac{\partial \phi}{\partial y} = f'(y) - \frac{k^2 x}{x^2 + y^2} \quad \rightarrow f'(y) = 0.$$

So, stream function  $\Rightarrow \phi(x, y) = k^2 \tan^{-1}(xy) + C$

$$\text{Q5] } \phi = \frac{Z}{x^3} \tan^{-1} \frac{y}{x}, \text{ then find streamlines.}$$

UPS C  
DO AGAIN

$$U = -\frac{\partial \phi}{\partial x} = +32 \cdot x \tan^{-1} \frac{y}{x} + \frac{Z}{x^3} \frac{x^2}{x^2+y^2} \cdot \frac{y}{x}$$

$$V = -\frac{\partial \phi}{\partial y} = \frac{32 y \tan^{-1} \left( \frac{y}{x} \right)}{x^5} + \frac{Z}{x^3} \frac{x^2}{x^2+y^2} \cdot \frac{1}{x}$$

$$W = -\frac{\partial \phi}{\partial z} = -\frac{\tan^{-1} \frac{y}{x}}{x^3} + 32 \frac{\tan^{-1} \left( \frac{y}{x} \right) z}{x^4} \cdot \frac{1}{x}$$

$$\frac{dx}{x^{-5} \left( 3x^2 \tan^{-1} \frac{y}{x} + yz \right)} = \frac{dy}{x^{-5} \left( 3y^2 \tan^{-1} \frac{y}{x} - xz \right)} = \frac{dz}{x^3 \tan^{-1} \frac{y}{x} \left( \frac{3z^2}{x^2} - 1 \right)}$$

Remember  
A

~~$$= \frac{x dx + y dy + z dz}{x^{-3} \tan^{-1} \left( \frac{y}{x} \right) (3z) \cdot x^2 - 2 x^{-3} \tan^{-1} \frac{y}{x} \cdot 2x^{-5} (3y^2 \tan^{-1} \frac{y}{x})}$$~~

~~$$= \frac{x dy + y dx + 2 dz}{2 z} = \frac{x dx + y dy}{2 x^{-2}}$$~~

~~$$\frac{x dx + y dy + 2 dz}{32 \cdot 2 x^{-3} \tan^{-1} \left( \frac{y}{x} \right)} = \frac{x dy + y dx}{x^3}$$~~

$$\frac{x dx + y dy}{3(x^2 + y^2)} = \frac{x dx + y dy + z dz}{2 x^2}$$

$$\Rightarrow \boxed{\log (x^2 + y^2 + z^2) = \frac{2}{3} \log (x^2 + y^2) + \log c.}$$

$$\frac{x^2 + y^2 + z^2}{x^2 + y^2} = \frac{3x^2 \tan^{-1} \frac{y}{x} + y^2 z^2 + 3y^2 z^2 \tan^{-1} \frac{y}{x}}{x^2 + y^2}$$

$$\frac{x^2 + y^2 + z^2}{x^2 + y^2} = \frac{3x^2 - \tan^{-1} \frac{y}{x}^2}{x^2 + y^2}$$

$$\text{Q} \quad q_r = \left[ 2\frac{M}{r^3} \cos \theta, \frac{M}{r^3} \sin \theta, 0 \right]$$

Show velocity is of potential. Find  $\phi$  and streamlines.

$$\text{TS: } \nabla \times q = 0 \rightarrow \begin{pmatrix} c_r & r \cos \phi & r \sin \phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ r^2 \sin^2 \theta & 2\frac{M}{r^3} \cos \theta & \frac{M}{r^3} \sin \theta \end{pmatrix}$$

$$= c_r(0) + r c_\theta(0) + r \sin \phi \left( -\frac{3M}{r^4} \sin \theta - \frac{2M \sin \theta}{r^3} \right)$$

$\Rightarrow ??$

$$dF = \frac{\partial F}{\partial r} dr + \frac{\partial F}{\partial \theta} d\theta + \frac{\partial F}{\partial \phi} d\phi$$

$$(q_r, q_\theta, q_\phi) = \left( -\frac{\partial F}{\partial r}, -\frac{1}{r} \frac{\partial F}{\partial \theta}, \frac{1}{r \sin \theta} \frac{\partial F}{\partial \phi} \right)$$

Streamlines.

$$\frac{dr}{q_r} = \frac{d\theta}{q_\theta} = \frac{d\phi}{q_\phi}$$

SOLVE

Show all conditions satisfied by

(a)

(i)  $\phi = \alpha x^2 + \beta y^2 + \gamma z^2$ , bounding surface  $F = \alpha x^2 + \beta y^2 + \gamma z^2 - \mathcal{H} = 0$   
 $\boxed{\alpha, \beta, \gamma, a, b, c \text{ are } \frac{1}{2} \text{ of time}}$

TS<sub>o</sub>:

(i)  $\phi$  satisfies Laplace Equation

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = 2\alpha + 2\beta + 2\gamma = 0$$

(ii).  $F$  is bounding surface

$$\frac{\partial F}{\partial t} + \mu \frac{\partial F}{\partial x} + \nu \frac{\partial F}{\partial y} + \omega \frac{\partial F}{\partial z} = 0.$$

$$x^4(\dot{a} - 8ax) + y^4(\dot{b} - 8bx) + z^4(\dot{c} - 8cz) - \dot{F} = 0. \quad (1)$$

All points of surface satisfy (1) and (2). So,

$$\frac{\dot{a} - 8ax}{a} = \frac{\dot{b} - 8bx}{b} = \frac{\dot{c} - 8cz}{c} = \frac{\dot{x}}{\pi}$$

$$\begin{aligned} \log a &= 8 \int \alpha dt + \log \pi \\ \log b &= 8 \int \beta dt + \log \pi \\ \log c &= 8 \int \gamma dt + \log \pi \end{aligned}$$

(Q) Find velocity potential and streamlines [2001]

$$\left( \frac{3xz}{\sigma^5}, \frac{3yz}{\sigma^5}, \frac{3z^2 - x^2}{\sigma^5} \right).$$

$$d\phi = \frac{\partial \phi}{\partial x} dx + \frac{\partial \phi}{\partial y} dy + \frac{\partial \phi}{\partial z} dz$$

$$= -3 \frac{xz}{\sigma^5} dx - 3 \frac{yz}{\sigma^5} dy - \frac{(3z^2 - x^2)}{\sigma^5} dz$$

$$= \left( -3z \frac{(x dx + y dy + z dz)}{\sigma^5} + \sigma^2 dz \right) \sigma$$

$$= -3z \frac{\sigma (\sigma d\sigma) + \sigma^3 dz}{(\sigma^3)^2} = d \left( \frac{z}{\sigma^3} \right)$$

$$\phi = \frac{z}{\sigma^3} \Rightarrow z = \sigma \cos \theta \Rightarrow \boxed{\phi = \frac{\cos \theta}{\sigma^2}} \text{ Ans}$$

$$\text{Streamlines} \rightarrow \frac{dx}{3xz/\sigma^5} = \frac{dy}{3yz/\sigma^5} = \frac{dz}{(3z^2 - x^2)/\sigma^5}$$

$$\frac{xdy + ydy + 2dz}{3z(x^2 + y^2 + z^2) - x^2} = \frac{dx}{3xz} \Rightarrow \frac{2(xdx + ydy + 2dz)}{2\sigma^2} = \frac{2dx}{3\sigma^2}$$

$$\boxed{\frac{d\sigma}{2\sigma^2} = \frac{2}{3} \frac{dx}{\sigma^2}} \rightarrow \text{Ans}$$

• F

③ Show  $\phi = xf(x)$  is a possible potential. Given that liquid speed  $q \rightarrow 0$  as  $r \rightarrow \infty$ , deduce that surfaces of constant speed are  $(r^2 + 3x^2)D^{-\delta} = \text{constant}$

- For  $\phi$  to be possible
  - $\phi$  must satisfy Laplace.

$$\rightarrow \frac{\partial \phi}{\partial x} = f(x) + x f'(x) \frac{x}{r}.$$

$$\frac{\partial \phi}{\partial y} = x f'(x) \frac{y}{r}. \quad \frac{\partial \phi}{\partial z} = \frac{x f'(x) z}{r}.$$

$$(u, v, w) = \left( -f(x) - x^2 \frac{f'(x)}{r}, -xy \frac{f'(x)}{r}, -xz \frac{f'(x)}{r} \right)$$

$$\cdot \frac{\partial^2 \phi}{\partial x^2} \Rightarrow f'(x) \frac{x}{r} + 2x \frac{f''(x)}{r} - \frac{f'(x)x^3 + x^3 f''(x)}{r^3}$$

$$\cdot \frac{\partial^2 \phi}{\partial y^2} \Rightarrow \frac{x f'(x)}{r} + \frac{x f''(x)y^2}{r^2} - \frac{xy^2 f'(x)}{r^3}$$

$$\cdot \frac{\partial^2 \phi}{\partial z^2} \Rightarrow \frac{x f'(x)}{r} + \frac{x f''(x)z^2}{r^2} - \frac{x z^2 f'(x)}{r^3}.$$

$$\text{Adding} \Rightarrow \frac{5xf'(x)}{r} + f'(x) \left[ \frac{5x}{r} - \frac{x}{r} \right] + f''(x) [x].$$

$$\Rightarrow \frac{4f'(x)}{r} + f''(x) = 0 \Rightarrow \log(f'(x)) + 4\log r = \log^4$$

$$\Rightarrow f'_r(x) = C_1 r^{-4} \Rightarrow f(x) = -\frac{C_1}{3} x r^{-3} + C_2$$

$$\text{As } q \rightarrow 0 \text{ as } r \rightarrow \infty \Rightarrow C_2 = 0$$

Surfaces of constant speed.

$\therefore q^2 = \text{constant}$

$$\text{Using } f(x) = -\frac{C}{3x^3}$$

$$q^2 = \frac{C}{3x^3} \left( i - \frac{3x}{x^2} \bar{\tau} \right)$$

$$q^2 = \frac{C}{9x^8} (x^2 + 3x^2)$$

## \* Equations of motion of inviscid fluids

(Sphere)

$$\rho_{\text{avg}} \equiv \text{Continuity} = \sigma^2 v_\sigma = F(t) \quad (\text{Sphere})$$

(Cylinder)

$$\begin{aligned} & \rightarrow \boxed{\frac{\partial v_\sigma}{\partial t} = \bar{F} - \frac{\nabla p}{\rho}} \\ & \sigma v_\sigma = F(t) \\ & \Rightarrow \frac{\partial v_\sigma}{\partial t} + v_\sigma \frac{\partial v_\sigma}{\partial x} = F_{\text{ext}} - \frac{1}{\rho} \frac{dp}{dx}. \\ & \equiv \text{Euler} \end{aligned}$$

### Problems

Type 1 ① Assuming Boyle's law, show radius  $\sqrt{\frac{V}{\pi}}$  of sphere oscillates between a min. and max and find  $n$ .

Step 1  $\Rightarrow$  At any time  $t$ , at  $x'$ ,  $v'$ ,  $p'$  cavity be  $x, v$ .

$$\gamma'^2 v' = \gamma^2 v = F(t). \quad \Rightarrow \quad \frac{\partial v'}{\partial t} = \frac{F'(t)}{\gamma'^2}.$$

Step 2  $\Rightarrow$  Euler.

$$\frac{\partial v'}{\partial t} + v' \frac{\partial v'}{\partial x} = -\frac{1}{\rho} \frac{\partial p'}{\partial x} \Rightarrow -\frac{F(t)}{\gamma'} + \frac{1}{2} v'^2 = C - \frac{P_1}{\rho}$$

$\rightarrow$  If  $\sigma = \infty$ ,  $v' = 0$ ,  $P_1 = \pi \Rightarrow C = \frac{\pi}{\rho}$

$$-\frac{F(t)}{\gamma'} + \frac{1}{2} v'^2 = \frac{\pi - P_1}{\rho}$$

Step 3  $\Rightarrow$  Boyle's law  $P_1 V_1 = P_2 V \Rightarrow m \pi \frac{4}{3} \sigma^3 a^3 = P_0 \frac{4}{3} \pi \sigma^3 \cdot (\rho_0 \text{ at } x)$

$$\boxed{P_0 = \frac{m \pi a^3}{\sigma^3}}$$

Step 4  $\Rightarrow$   $F'(t) = 2 \gamma v^2 + \gamma^2 v \frac{dv}{dx}$

$$\begin{aligned} & \text{Step 5} \Rightarrow \text{at } x' = \sigma, v' = v, P_1 = P_0 \\ & \Rightarrow F'(t) = 2 \gamma v^2 + \gamma^2 v^2 = \frac{\pi - P_0}{\rho} = \frac{\pi}{\rho} \left( 1 - \frac{m a^3}{\sigma^3} \right) \\ & -\frac{2v^2}{\rho} - \sigma v \frac{dv}{dx} + \frac{1}{2} v^2 = \frac{\pi - P_0}{\rho} = \frac{\pi}{\rho} \left( 1 - \frac{m a^3}{\sigma^3} \right) \end{aligned}$$

$$\begin{aligned}
 \textcircled{1} \text{ When } t \text{ not involved } &\equiv F'(r) = d(r^2 v) = 2rv + r \frac{dv}{dr} \frac{dr}{dt} \\
 &= 2rv^2 + r^2 v \frac{dv}{dr}.
 \end{aligned}$$

radius oscillates  $a, na \Rightarrow v=0$  when  $r=a, na$ .

$$\Rightarrow [1 + 3m \log n - n^3 = 0]$$

$$\begin{aligned}
 \uparrow \\
 r=a, v=0 \Rightarrow C' &= \frac{2\pi a^3}{3\rho} - 2a^3 m \pi \log a
 \end{aligned}$$

$$\begin{aligned}
 \uparrow \\
 r^2 v \frac{dv}{dr} \text{ side } &\Rightarrow -3r^2 v^2 - 2r^3 v dv = \frac{\pi}{\rho} \left(1 - \frac{ma^3}{r^3}\right) dr \\
 \Rightarrow \int d(r^3 v^2) &= \frac{2\pi}{\rho} \left(\frac{r^3}{3} - \frac{2ma^3 \log r}{\pi}\right) \\
 \Rightarrow \boxed{r^3 v^2 = -\frac{2\pi}{\rho} \left(\frac{r^3}{3} - 2am \log r\right) + C'}
 \end{aligned}$$

- (2) Liquid between 2 parallel planes. Free Surface is concave cylinder of radius  $a$ , a cylinder of  $b$  is suddenly annihilated. Prove if  $\Pi$  is pressure at outer surface initial pressure at  $r$  from centre is  $\frac{\Pi \log r - \log b}{\log a - \log b}$



$$\rightarrow \sigma v = F(+)$$

$$\frac{\partial v}{\partial t} = \frac{F'(+)v}{\sigma}$$

$$Mt \sigma \rightarrow v, \rho$$

$$\frac{F'(+)v}{\sigma} + v \frac{\partial v}{\partial r} = -\frac{1}{\rho} \frac{\partial P}{\partial r}$$

$$F'(+) \log r + \frac{1}{2} v^2 = C - \frac{P}{\rho}$$

$$At t=0 \Rightarrow$$

$$x=b, v=0, \rho=0 \Rightarrow F'(0) \log b = C$$

$$x=a, v=0, \rho=\Pi \Rightarrow F'(0) \log a = C - \frac{\Pi}{\rho}$$

$$F'(0) = \frac{\Pi}{\rho(\log b - \log a)} \quad C = \frac{\Pi \log b}{\rho(\log b - \log a)}$$

Initially at  $r, v=0$

$$F'(0) \log r = \frac{\Pi \log b}{\rho(\log b - \log a)} - \frac{P}{\rho}$$

$$\frac{\Pi}{\rho} \left( \frac{\log r - \log b}{\log b - \log a} \right) = -\frac{P}{\rho} \Rightarrow$$

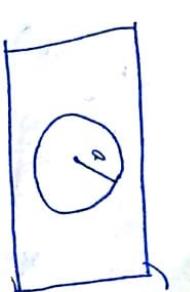
$$\boxed{P = \frac{\Pi(\log r - \log b)}{\log a - \log b}}$$

Put

2

1

(3) Constant pressure at infinity. Pressure at  $r'$  from centre for when radius carry is  $a$ .



fluid

$$r'^2 v' = r^2 v = F(r).$$

Let  $\rho_0$  be  $\rho_i$ .

$$\frac{\partial v'}{\partial r} = \frac{F'(r)}{r'^2}.$$

$$\rightarrow -\frac{F'(r)}{r'} + \frac{1}{2} v'^2 = C - \frac{\rho_i}{\rho}$$

$$C = \frac{\pi}{\rho}$$

$$\textcircled{*} \quad \boxed{-\frac{F'(r)}{r'} + \frac{1}{2} v'^2 = \frac{\pi - \rho_i}{\rho}} \rightarrow$$

$$\boxed{-\frac{F'(r)}{r} + \frac{1}{2} v^2 = \frac{\pi}{\rho}} \quad \text{--- (1)}$$

Now find  $F'(r)$  and  $v'^2$

$$F'(r) = d(r^2 v) = 2r v^2 + r^2 v \frac{dv}{dr}.$$

$$At \quad r' = r, \quad v' = v.$$

$$\frac{1}{r} \left( \frac{1}{2} \cancel{r^2} 2r v^2 - 2r^2 v dv \right) + \frac{1}{2} v^2 dr = \left( \frac{\pi - \rho_i}{\rho} \right) r^2 . dr.$$

$\times$   $\cancel{r}$   $\times$   $\cancel{r}$    
 both sides

$$\Rightarrow -3r v^2 dr - 2r^2 v dv = 2 \left( \frac{\pi - \rho_i}{\rho} \right) r^2 dr$$

$$- \left( d(r^3 v^2) \right) = 2 \left( \frac{\pi - \rho_i}{\rho} \right) \frac{r^3}{3} + C.$$

At  $r = r, \rho = 0$

$$-r^3 v^2 = 2 \left( \frac{\pi - \rho_i}{\rho} \right) \frac{r^3}{3} + C.$$

At  $r = a, v = 0$

$$2 \left[ \frac{\pi a^3}{\rho} \right] + C = 0 \Rightarrow \boxed{r^3 v^2 = \frac{2\pi}{3\rho} (a^3 - r^3)} \quad \text{--- (2)}$$

## TOOLS AVAILABLE

$$\textcircled{1} \quad F'(t) = \frac{d}{dt} (R^2 V) = \frac{d}{dt} \left( R^2 \frac{dR}{dt} \right)$$

$$\textcircled{2} \quad \text{Calculating } F'(t) \rightarrow F(t) = \gamma^2 V$$

$$F'(t) = 2\gamma \frac{d\gamma}{dt} V + \gamma^2 \frac{dV}{dt} = 2\gamma V^2 + \gamma^2 \frac{dV}{dx} \frac{dx}{dt}$$

$$= 2\gamma V^2 + \gamma^2 V \frac{dV}{dx} \xrightarrow{\text{remove}} \text{always } \frac{d}{dt} \text{ term.}$$

$$\textcircled{1} \quad KE = \int_{\frac{1}{2}}^{\infty} \rho 4\pi x^2 dx V^2$$

(Outer surface mass)

$$\Rightarrow V' = \frac{\gamma^2 V}{x}$$

$$\textcircled{2} \quad PE = \int \pi 4\pi x^2 dx \int V dm$$

(Can take inner or outer radii - Just one, look how it changes!)

Type 2 Using energy method.

$$\cdot KE = PE \rightarrow \int p \cdot dv$$

$$\int_{\frac{1}{2}}^1 4\pi r^2 dr \rho v^2$$

① Infinite fluid. force  $\equiv \frac{\mu}{r^{3/2}}$  to origin. Initial cavity (C).

Find time of filling up.

boundary condition

$$\Rightarrow r'^2 v' = r^2 v = F(+)$$

$$At t= \int_{\frac{1}{2}}^{\infty} 4\pi r'^2 dr' \rho v'^2 = 2\pi \rho \int_r^{\infty} r'^2 \cdot \frac{r^4 v^2}{r'^4} dr'$$

$$KE = 2\pi \rho r^4 v^2 \left[ -\frac{1}{3} \right]_r^{\infty} = \boxed{2\rho \pi r^3 v^2}$$

$$PE = \int_{\sigma}^c \frac{2\mu}{r'^2} \cdot 4\pi r'^2 dr' \rho \cdot \left( \int r' dm \right)$$

$$= \frac{16\pi \rho \mu (C^{5/2} - r^{5/2})}{5}$$

$$F = -\frac{\partial V}{\partial r} = F$$

$$\frac{\partial V}{\partial r} = +\frac{\mu}{r^{3/2}}$$

$$\Rightarrow V = \frac{2\mu}{r^{4/3}}$$

(With r decreases with time)

$$\frac{\partial F}{\partial r} \rho \mu (C^{5/2} - r^{5/2}) = 2\pi \rho r^3 v^2$$

$$\frac{\partial F}{\partial r} \rho \mu (C^{5/2} - r^{5/2}) \Rightarrow \frac{d\sigma}{dt} = -\left( \frac{8\mu (C^{5/2} - r^{5/2})}{5 \cdot r^3} \right) t$$

$$r^{5/2} = C^{5/2} \sin^2 \theta$$

$$-\frac{r}{2} \frac{dr}{dt} = dt$$

$$\sqrt{\frac{8\mu}{5}} \sqrt{\frac{2}{C^{5/2} - r^{5/2}}} = dt$$

$$C^{5/2} \left( \frac{2}{5} \right) \cdot 2 e^{\frac{5}{2} \sin \theta} \cos \theta \cdot \sqrt{5} = dt$$

$$\sqrt{8\mu} \cdot \frac{e^{\frac{5}{2} \sin \theta} \cos \theta}{\sqrt{C^{5/2} - r^{5/2}}} dt = \left( \frac{2}{5} \right)^{1/2} C^{5/4}$$

(2) Outside pressure  $\pi$ . Initial pressure inside =  $P_1$

Find velocity of inner surface ( $v$ ) when  
 $R_2, C_2$  is reached.

$\Rightarrow$  gas obeys Boyle's law.

$$\frac{4}{3}\pi(R_1^3 - C_1^3) = \frac{4}{3}\pi(R^3 - C^3)$$

$$x^2 v = \sigma^2 v_1 e^{-\text{inner boundary}}$$

$$-\frac{F'(C)}{\gamma^2} + \frac{1}{2} v^2 = C - \frac{P}{\rho}$$

$$P_1 \cdot \frac{4}{3}\pi C_1^3 = P_2 \cdot \frac{4}{3}\pi C_2^3$$

Let  $(\bar{r}, R)$  be intermediate config

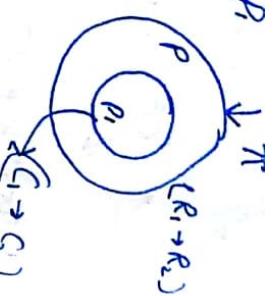
$$KE = \frac{1}{2} \int_{R_1}^R 4\pi x^2 \delta x \rho v^2 = 2\pi \rho \int_0^R x^2 \frac{\sigma^4 v_1^2}{x^4} dx$$

$$= 2\pi \rho \sigma^4 v_1^2 \left[ \frac{1}{\alpha} - \frac{1}{R} \right]$$

$$PE = \int_{R_1}^R \pi \cdot 4\pi x^2 (\delta x) = \underbrace{4\pi \rho \left( \frac{R_1^3}{3} - \frac{R^3}{3} \right)}_{\text{+}} + \underbrace{4\pi \rho C_1^3 \log \frac{C}{C_1}}_{C_1 \int \frac{P_1 C_1^3}{x^3} dx}$$

$$\frac{4}{3}\pi \rho (R_1^3 - R^3) + 4\pi \rho C_1^3 \log \frac{C}{C_1} = 2\pi \rho C_1^4 v_1^2 \left[ \frac{1}{\alpha} - \frac{1}{R} \right]$$

$$\alpha + r - r = R = R_1 \Rightarrow$$



③ Fluid ( $\rho$ , volume =  $\frac{4\pi r^3}{3}$ ) in spherical shell.

$\tau$  on external shell, no  $\rho$  on internal shell.

Initial  $\equiv$  initial radius =  $2c$ .

Find velocity of inner surface when radius is  $c$ .

At any time  $t$ ,  $r$  and  $R$  be radii  
 $v$  and  $V$  be velocity

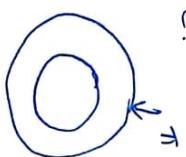
$$r'^2 v' = r v = F(t)$$

$$r'^2 v' = r^2 v = F(t).$$

$$R^3 - r^3 = c^3 \Rightarrow R^3 = r^3 + c^3$$

$$R^3 = \frac{1}{2} \rho (4\pi r^3) dr, v'^2 \Rightarrow 2\pi \rho \int_{r_1}^r \frac{r^2}{r^4} v^2 dr,$$

$$KE = \int_{r_1}^r \frac{1}{2} \rho (4\pi r^3) dr \cdot v'^2 = 2\pi \rho r^4 v^2 \left[ \frac{1}{2} - \frac{1}{R} \right].$$



Done

$$\rho E = \int_{2c}^r -\pi \cdot 4\pi r^2 dr,$$

$$= + \frac{4\pi \rho}{3} (-r^3 + 8c^3).$$

$$KE = \rho E \Rightarrow \frac{4\pi \rho}{3} (8c^3 - r^3) = 2\pi \rho r^4 v^2 \left[ \frac{1}{2} - \frac{1}{8c^3} \right]$$

Set  $r = c$  to get  $v^2$ :

$$\frac{4\pi}{3} \cdot 7\rho^2 = 2\pi \rho c^4 v^2 \left[ \frac{1}{2} - \frac{1}{8c^3} \right]$$

$$v^2 = \left( \frac{14\pi}{3\rho} \frac{2^{4/3}}{2^{10/3}-1} \right)^{1/2}$$

Type 3

- ① A steady inviscid incompressible flow with velocity  $(fx, -fy, 0)$ . Find pressure field  $p(x, y, z)$  if  $p(0, 0, 0) = p_0$  and  $F = -\frac{gx}{\rho} - g_2 z$ .

$\Rightarrow$  Must satisfy equations of motion.

$$fx, f = X^2 \frac{\partial P}{\partial x} - f^2 y = X^2 - \frac{\partial P}{\partial xy}$$

$$0 = -g^2 - \frac{1}{\rho} \frac{\partial P}{\partial z}$$

$$\frac{\partial P}{\partial x} = -f^2 \rho x ; \frac{\partial P}{\partial y} = f^2 \rho y ; \frac{\partial P}{\partial z} = -\rho g z.$$

$$dp = -f^2 \rho x dx + f^2 \rho y dy - \rho g z dz.$$

$$p = \left( -f^2 \rho x^2 + f^2 \rho y^2 - \rho g z^2 \right) \frac{1}{2} + C \rho$$

Q) A fluid of density  $\rho$  bound by concentric spheres outer surface at uniform pressure  $P$  and contracts  $R_1$  to  $R_2$ . Hollow is filled with gas obeying Boyle's law with initial pressure  $P_1$  and contracts  $C_1$  to  $C_2$ . Find velocity of inner surface when  $(R_2, C_2)$  is reached.

$$\rightarrow \pi R_1^2 V_1 = \pi R_2^2 V_2 = F(t).$$

$$\rightarrow P_1 \frac{4}{3} \pi C_1^3 = P_2 \frac{4}{3} \pi C_2^3 \Rightarrow \left( P_2 = \frac{C_1^3}{C_2^3} P_1 \right)$$

$$KE = \int_{C_1}^{R_2} \frac{1}{2} (4\pi \dot{x}^2 d\sigma, P) \frac{C_2^4 V^2}{\dot{x}^4 r^2}$$

$$= 2\pi \rho C_2^4 V^2 \left( \frac{1}{C_2} - \frac{1}{R_2} \right) \quad (\text{No - sign})$$

$$\text{Work done} = \int_{R_1}^{R_2} \pi 4\pi R_2^2 (-dR_2) + \int_{C_1}^{C_2} \frac{\rho C_1^3}{C_2^3} 4\pi C^3 (dc)$$

$$= \pi R_1 \cdot \frac{4}{3} (R_2^3 - R_1^3) + 4\pi \rho C_1^3 \log \frac{C_2}{C_1}$$

By COM  $\Rightarrow \boxed{R_1^3 - R_2^3 = C_1^3 - C_2^3}$

$$\frac{4}{3} \pi R_1^3 (C_1^3 - C_2^3) + 4\pi \rho C_1^3 \log \frac{C_2}{C_1} = 2\pi \rho C_2^3 V^2 \left( 1 - \frac{C_2}{R_2} \right)$$

(Q) Liquid of density between cylinders (a, b). Outer pressure on a is  $\pi$ . Find velocity  $v$  of inner surface when inner radius is  $x$ .

At any time  $t \equiv$  Inner ( $\sigma, v$ )  $p = 0$   
Outer ( $R, V$ )  $p = \pi$ .

USELESS

For any point in fluid  $\Rightarrow x' v' = F(t)$

$$\frac{\partial v'}{\partial t} = \frac{F'(t)}{x'}$$

$$\frac{F'(t)}{x'} + v' \frac{\partial v'}{\partial x'} = -\frac{1}{\rho} \frac{\partial p}{\partial x'}$$

(ENERGY METHOD)

$$\text{Integrating} \rightarrow F'(t) \log x' + \frac{1}{2} v'^2 = C - \frac{p}{\rho}.$$

(A)

$$1. F'(t) \log x + \frac{v^2}{2} = C. \\ 2. F'(t) \log R + \frac{V^2}{2} = C - \frac{\pi}{\rho}.$$

$$\sigma v = R V = F(t) \Rightarrow x \frac{dv}{dt} = R \frac{dR}{dt} = F(t) \Rightarrow \boxed{x dv = R dR = F(t) dt}$$

$$R^2 - x^2 = a^2 - b^2.$$

$$F'(t) = \frac{d}{dt}(\sigma v) = \frac{d}{dx}(\sigma v) \cdot \frac{dx}{dt} = \frac{d}{dx}(\sigma v) \cdot v. \quad (2)$$

But, values in (A) of  $F'(t)$ ,  $V$  from (1), (2)

$$v \frac{d}{dx}(x v)$$

### \*Energy method

when ( $\sigma, R$ )

$$\sigma' v' = \sigma v = F(+)$$

$$R \int_{\sigma}^{\sigma'} \frac{1}{2} (4 \pi \sigma'^2) p v'^2 = 2 \pi p \int_{\sigma}^{\sigma'} \frac{\sigma'^2 v'^2}{\sigma^2} d\sigma'$$

$K E$

$$= 2 \pi p \sigma^2 v^2 (R - \sigma)$$

Potential Energy

~~$$R \int_{\sigma}^{\sigma'} (4 \pi \sigma'^2) \cancel{F} d\sigma' = R \int_{\sigma}^{\sigma'} (-2 \pi \sigma' h) d\sigma' \cancel{\pi}$$~~

$$= \boxed{2 \pi h \cancel{\pi} \cdot (\sigma^2 - R^2)}$$

$$KE = \int_{\sigma}^{R'} \frac{1}{2} (2 \pi \sigma' h) d\sigma' p \cdot \frac{\sigma'^2 v^2}{\sigma'^2}$$

$$= \pi \sigma^2 v^2 p h \log(R/\sigma)$$

$$\boxed{\sigma^2 - b^2 = R^2 - \sigma^2}$$

$$2 \pi K \pi (a^2 - R^2) = \cancel{\pi} \sigma^2 v^2 p \log(R - \sigma)$$

$$v^2 = \frac{\pi (a^2 - R^2)}{\sigma^2 p \log(\sqrt{\sigma^2 + a^2 - b^2} / \sigma)}$$

$$= \frac{2 \pi (b^2 - \cancel{\sigma^2})}{\sigma^2 p \log(\sqrt{\sigma^2 + a^2 - b^2})}$$

Maths

## \* Sources and Sinks

$$(q_{1x} = \frac{1}{2} \frac{\partial \psi}{\partial x}, q_{1y} = \frac{\partial \psi}{\partial y})$$

$w = \phi + i\psi$  Stream function  
 Complex potential Velocity potential.

$$\cdot (U, V) = \left( -\frac{\partial \phi}{\partial x}, -\frac{\partial \phi}{\partial y} \right) = \left( -\frac{\partial \psi}{\partial y}, \frac{\partial \psi}{\partial x} \right) \quad (\phi, \psi \text{ satisfy C-R})$$

also Laplace

$$\cdot q_r = \left| \frac{dw}{dz} \right| = (U^2 + V^2)^{1/2} \quad [\text{Stagnation points } \Rightarrow q_r = 0]$$

• To Show, Velocity potential exist  $\rightarrow$  Show irrotational  
Evaluate  $\zeta$

$\rightarrow$  Complex Potential  $\equiv$  Strength =  $m$ , Rate of output =  $2\pi m$

$\cdot -m \log(z - z_0) = w = \phi + i\psi$   $\downarrow = -m\theta$

$-m \log z$

$$\log(\alpha + i\beta) = \frac{1}{2} \log(\alpha^2 + \beta^2) + i \tan^{-1} \frac{\beta}{\alpha}$$

## → Images

$\cdot$  Plane boundaries at  $\infty \rightarrow n = \frac{\sigma_2}{\sigma_1}$   $\sigma, \alpha \rightarrow \sigma^n, \alpha$

$(B_1, B_2) \rightarrow (nB_1, nB_2)$  Reflect in transform boundary

$w = -\log(z^n - \sigma^n \alpha)$

$$2 \quad z \rightarrow z^n$$

$\cdot$  Circle  $\rightarrow m$  outside  $\Rightarrow [m(a^2/F), -m \text{ at } (0)]$  (Same if  $m$  inside)

$\cdot$  Milne Thomson  $\equiv w$  before cylinder (no boundary flow)  
 After  $|z|=a$  inserted  $\rightarrow [w_{\text{new}} = f(z) + \bar{f}(a^2/z)]$  for  $|z| > a$

- Curves of  $\phi = k_1$ ,  $\psi = k_2$  intersect orthogonally

Then,  $\frac{\partial \phi}{\partial x} \cdot \frac{\partial \psi}{\partial x} + \frac{\partial \phi}{\partial y} \cdot \frac{\partial \psi}{\partial y} = 0$

• 
$$\boxed{\frac{dw}{dz} = -u + iv}$$

- irrotational  $\Rightarrow \nabla \times \mathbf{q} = 0$ . (Not  $\nabla^2 \psi$  or  $\nabla^2 \phi$ )  
If irrotational, then  $\underline{\nabla^2 \phi / \nabla^2 \psi = 0}$ .

- Slope of  $\phi \equiv m_1 = -\frac{\partial \phi / \partial x}{\partial \phi / \partial y}$  (Same for  $\psi$ ).

- Ovals of cassini  $\Rightarrow \gamma \gamma' = \text{constant}$ .

$$\left| \frac{1}{z+a} - \frac{1}{z-a} \right| = \frac{2a}{|z+a||z-a|} = \frac{2a}{\gamma \gamma'}$$

- $X - iY = \frac{1}{2} \pi i \int_C \underbrace{\left( \frac{dw}{dz} \right)^2}_{\text{Residues}} dz$ .

$$M = \operatorname{Re} \left\{ -\frac{1}{2} i \int_C z \left( \frac{dw}{dz} \right)^2 dz \right\}$$

\* For Reflection, Only that boundary on which source/sink does not lie.

① Show  $\phi = \frac{1}{2} \log \frac{(x+a)^2 + y^2}{(x-a)^2 + y^2}$  is possible.

Find streamlines.

$$\bullet \quad \phi = \frac{1}{2} \left[ \log(x+a)^2 + y^2 - \log(x-a)^2 + y^2 \right]$$

$$U = -\frac{\partial \phi}{\partial x} = \frac{-1}{2} \left[ \frac{2(x+a)}{(x+a)^2 + y^2} - \frac{2(x-a)}{(x-a)^2 + y^2} \right]$$

$$V = -\frac{\partial \phi}{\partial y} = -\frac{1}{2} \left[ \frac{2y}{(x+a)^2 + y^2} - \frac{2y}{(x-a)^2 + y^2} \right]$$

$$\frac{\partial U}{\partial x} = -\frac{1}{2} \left[ \frac{2((x+a)^2 + y^2) - 2(x+a)}{(x+a)^2 + y^2} - \frac{2((x-a)^2 + y^2) - 2(x-a)}{(x-a)^2 + y^2} \right]$$

$$\frac{\partial V}{\partial y} = -\frac{1}{2} \left[ \frac{2((x+a)^2 + y^2) - 2y \cdot 2y}{((x+a)^2 + y^2)^2} - \frac{2((x-a)^2 + y^2) - 2y \cdot 2y}{((x-a)^2 + y^2)^2} \right]$$

$$\text{So, } \frac{\partial U}{\partial x} + \frac{\partial V}{\partial y} = 0$$

② Sources ( $m$ ) at  $(-a, 0), (a, 0)$ , sink ( $-2m$ ) at  $(0, 0)$ .  
 Find streamlines and fluid speed at any point.

$$w = -m \log(z+a) - m \log(z-a) + 2m \log z.$$

$$= -m \log(x^2 - y^2 - a^2 + i2xy) + m \log(x^2 - y^2 + i2xy)$$

$$\psi = -m \left( \tan^{-1} \left( \frac{2xy}{x^2 - y^2 - a^2} \right) + \tan^{-1} \left( \frac{2xy}{x^2 - y^2} \right) \right)$$

$$= -m \tan^{-1} \left[ \frac{\frac{2xy}{x^2 - y^2 - a^2} + \frac{2xy}{x^2 - y^2}}{(x^2 - y^2 - a^2)(x^2 - y^2) + 4x^2y^2} \right]$$

$$= -m \tan^{-1} \left[ \frac{2x^3y - 2xy^3 + 2x^2y + 2xy^3 + 2a^2xy}{(x^2 + y^2)^2 - a^2(x^2 - y^2)} \right]$$

$$= m \tan^{-1} \left[ \frac{-2a^2xy}{(x^2 + y^2)^2 - a^2(x^2 - y^2)} \right] = \text{constant} = \cancel{m \tan^{-1}}$$

$$= m \tan^{-1} \left( \frac{2}{\lambda} \right)$$

$$\text{Fluid speed} = \left| -\frac{m}{z+a} - \frac{m}{z-a} + \frac{2m}{z} \right|$$

$$= \left| \frac{-mz^2 + amz - mz^2 - amz + 2mz^2 - 2ma^2}{2(z+a)(z-a)} \right|$$

$$= \left| \frac{2ma^2}{8r_1 r_2} \right| \quad \underline{\underline{\Delta m}}$$

(3) Boundaries  $\{-\frac{\pi}{4}, \frac{\pi}{4}\}$ . Source  $m(a, 0)$ ,  
 Sink  $-m(b, 0)$ .

Find stream function.

$$\rightarrow \frac{\nabla \psi_2}{\nabla \psi_4} = 2. \quad \begin{array}{c} \uparrow \\ \downarrow \end{array} \quad \begin{array}{ll} m(a^2, 0) & m(-a^2, 0) \\ -m(b^2, 0) & -m(-b^2, 0). \end{array}$$

$$\text{Boundary} = \left\{-\frac{\pi}{2}, \frac{\pi}{2}\right\}$$

$$W = -m \log(z^2 - a^2) - m \log(z^2 + a^2) + m \log(z^2 - b^2) \\ + m \log(z^2 + b^2).$$

$$= -m \log(\gamma^4 e^{i40^\circ} - a^4) + m \log(\gamma^4 e^{i40^\circ} - b^4)$$

$$\psi = -m \tan^{-1} \left( \frac{\gamma^4 \sin 40^\circ}{\gamma^4 \cos 40^\circ - a^4} \right) + m \tan^{-1} \left( \frac{\gamma^4 \sin 40^\circ}{\gamma^4 \cos 40^\circ - b^4} \right)$$

$$\psi = -m \tan^{-1} \left( \frac{\gamma^8 \sin 40^\circ \cos 40^\circ - b^4 \gamma^4 \sin 40^\circ - \gamma^8 + a^4 \gamma^4 \sin 40^\circ}{\gamma^8 \cos^2 40^\circ - a^4 \gamma^4 \cos 40^\circ - b^4 \gamma^4 \cos 40^\circ + a^4 b^2} \right)$$

$$\boxed{\psi = -m \tan^{-1} \left( \frac{\gamma^4 \sin 40^\circ (a^4 - b^4)}{\gamma^8 - \gamma^4 a^4 \cos 40^\circ - b^4 \cos 40^\circ + a^4 b^4} \right)} \quad \underline{\text{Ans}}$$

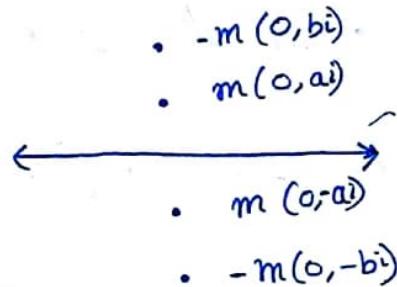
If Boundary at  $\theta = 0, \frac{\pi}{3}$ .  $\Rightarrow \frac{\frac{\pi}{2}}{\frac{\pi}{6}} = 3$

$\Rightarrow \boxed{\theta = 0, \pi} \leftarrow \text{Transformed}$

① Fluid fills space on the side of  $x$  axis which is a rigid boundary and source  $m(0, a)$ , sink  $(-m)(0, b)$  and if pressure on negative side be same as  $P_0$  at  $\infty$ . Show resultant pressure on boundary is  $\frac{\pi P m^2 (a-b)^2}{2ab(a+b)}$

To find pressure at boundary, find  $q$  at boundary

$$\frac{P}{P_0} + \frac{q^2}{2} = \text{constant} = \frac{P_0}{P} \quad \text{at } \infty$$



$$w = -m \log(z^2 + a^2) + m \log(z^2 + b^2)$$

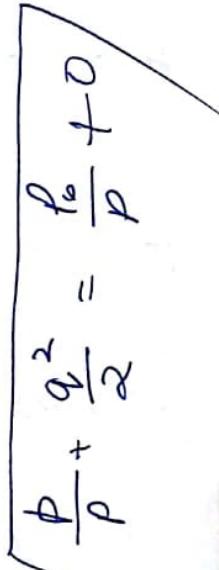
~~$w = -m \log(z^2 + a^2)$~~

$$\text{Find } \left[ \frac{dw}{dz} \right]_{y=0} = \left| -\frac{m \cdot 2z}{(z^2 + a^2)} + \frac{m \cdot 2z}{(z^2 + b^2)} \right|$$

$$= \left| 2ym \left( \frac{1}{z^2 + b^2} - \frac{1}{z^2 + a^2} \right) \right|$$

$$q_y = \frac{2mn(a^2 - b^2)}{(x^2 + b^2)(x^2 + a^2)}$$

$$\frac{q_y^2}{2} = \frac{P_0 - P}{P_0} \Rightarrow P_0 - P \Rightarrow \boxed{\frac{1}{2} P q_y^2} \rightarrow \left( \frac{x}{x^2 + b^2} - \frac{x}{x^2 + a^2} \right)$$



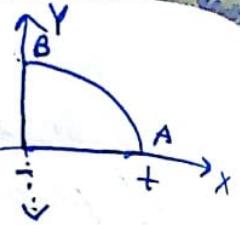
$$\text{Resultant pressure} \Rightarrow \frac{1}{2} P \int_0^\infty \frac{(2m)^2 (a^2 - b^2)^2 x^2}{((x^2 + b^2)(x^2 + a^2))^2} dx$$

$$= 2pm^2 \int_0^\infty \frac{x^2 (a^2 - b^2)^2 dx}{(x^2 + b^2)^2 (x^2 + a^2)^2}$$

$$= \frac{\pi P m^2 (a-b)^2}{2ab(a+b)} \underline{\underline{dx}}$$

$$\int \text{Assume } \frac{A}{x^2 + a^2} + \frac{B}{(x^2 + a^2)^2} + \frac{C}{(x^2 + b^2)} + \dots$$

⑤ Show that the streamline leaving either end at  $\angle \alpha$  with radius is  $\frac{r \sin(\alpha + \theta)}{a^2 \sin(\alpha - \theta)}$



⑥ Saw Sha

Equivalent Image system

• m at  $Z=a$ ,  $Z=-a$       . -m at  $Z=0$ .

$$\begin{aligned} w &= -m \log \left( \frac{z^2 - a^2}{z} \right) = -m \log (z - a^2 z^{-1}) \\ &= -m \log (r \cos \theta + i r \sin \theta - a^2 \bar{z} \cos \theta + i a^2 \bar{z} \sin \theta) \\ &= -m \log \left( r \cos \theta \left( 1 - \frac{a^2}{r} \right) + i \left( r \sin \theta + \frac{a^2}{r} \sin \theta \right) \right) \end{aligned}$$

$$\psi = -m \tan^{-1} \left( \frac{\sin \theta \left( 1 - \frac{a^2}{r} \right)}{\cos \theta \left( 1 - \frac{a^2}{r} \right)} \right)$$

$$\psi = -m \tan^{-1} \left\{ \frac{r^2 + a^2}{r^2 - a^2} \tan \theta \right\}$$

Streamline at angle  $\alpha = -m(\pi - \alpha)$

$$\tan \tan^{-1} \left\{ \frac{r^2 + a^2}{r^2 - a^2} \tan \theta \right\} = \tan(\pi - \alpha) = -\tan \alpha$$

$$\boxed{(r^2 + a^2) \tan \theta + (r^2 - a^2) \tan \alpha = 0} \quad \text{Ans}$$

IMP  
Takeaway

After,

w =

=

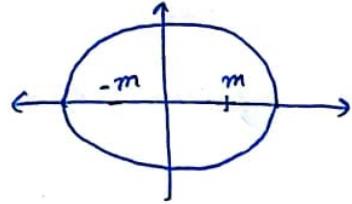
$\Rightarrow$  As u use

- W/C

$w = f_1$

⑥ Source and sink at  $(\pm a/2, 0)$  within  $|z|=a$ .  
 Show streamlines are  $\left(\gamma^2 = \frac{a^2}{4}\right) \left(\gamma^2 = 4a^2\right) - 4a^2y^2 = k\gamma(\gamma^2 - a^2)$

$\Rightarrow$  As circular boundary, so we can  
 use Milne Thomson theorem



- W/o boundary

$$w=f(z) = -m \log(z - a/2) + m \log(z + a/2)$$

$$= -m \log \left( \frac{z - a/2}{z + a/2} \right)$$

$$\begin{aligned} & -m \log(z - a/2) + m \log(z + a/2) \\ & -m \log(2a^2 - az) + m \log(2z) \\ & + m \log(2a^2 + az) - m \log(2z) \end{aligned}$$

\* Continue from here.

After boundary  $\Rightarrow$

$$w = f(z) + f(a^2/2) = -m \log \left( \frac{z - a/2}{z + a/2} \right) - m \log \left( \frac{\frac{a^2}{2} - \frac{a}{2}}{\frac{a^2}{2} + \frac{a}{2}} \right)$$

$$= -m \log(z - a/2) - m \log \left( \frac{a^2 - a}{2} \right) + m \log \left( \frac{a^2 + a^3/2 + az + a^2/4}{2z} \right)$$

$$= -m \log \left( a^2 - \frac{a^3}{2z} - \frac{az}{2} + \frac{a^2}{4} \right) + m \log \left( \frac{5a^2}{4} + \frac{a^3(x-iy)}{2(x^2+y^2)} + \frac{a}{2}(x+iy) \right)$$

$$= -m \tan^{-1} \left( \frac{-ay + a^3y}{\frac{5a^2}{2} - \frac{ax}{2} - \frac{a^3x}{2z(x^2+y^2)}} \right) + m \tan^{-1} \left( \frac{ay - \frac{a^3y}{2(x^2+y^2)}}{\frac{5a^2}{2} + \frac{ax}{2} + \frac{a^3x}{2(x^2+y^2)}} \right)$$

$$= -m \tan^{-1} \left( \frac{-ayx^2 - ay^3 + a^3y}{5a^2x^2 + 5a^2y^2 - 2ax^3 - 2axy^2 - 2a^3x} \right) + m \tan^{-1} \left( \frac{ayx^2 + ay^3 - a^3y}{5a^2x^2 + 5a^2y^2 + 2ax^3 + 2axy^2 + 2a^3x} \right)$$

$$= -m \tan^{-1} \left( \frac{ay(\sigma^2 - \gamma^2)}{5a^2\sigma^2 - 2ax\sigma^2 - 2a^3x} \right) + m \tan^{-1} \left( \frac{ayx^2 + ay^3 - a^3y}{5a^2x^2 + 5a^2y^2 + 2ax^3 + 2axy^2 + 2a^3x} \right)$$

## ⑦ Infinite vortex rows

$\omega$  = Due to an infinite row

vortices -  $k$  at  $z = na$

$$\omega = \frac{ik}{2\pi} \sum_{n=-\infty}^{\infty} \log(z-na)$$

$$= \frac{ik}{2\pi} \log [z(z^2-a^2)(z^2-2^2a^2) \dots]$$

$$= \frac{ik}{2\pi} \log \left[ \frac{\pi z}{a} \left( 1 - \frac{z^2}{a^2} \right) \left( 1 - \frac{z^2}{2^2 a^2} \right) \dots \right]$$

$$+ \frac{ik}{2\pi} \log \left[ (-1)^n \frac{a}{\pi} a^2 \cdot 2^2 a^2 \dots n^2 a^2 \right]$$

$$= \frac{ik}{2\pi} \log \left[ \frac{\pi}{a} z \left( 1 - \frac{z^2}{a^2} \right) \left( 1 - \frac{z^2}{2^2 a^2} \right) \dots \right] + C.$$

⑧ 2 infinite rows  $[x=xa, y=b, k], [x=xa, y=-b, -k]$ .

Prove the street moves as a whole along its length.

$$w = \frac{ik}{2\pi} \sum_{n=0}^{\infty} \log(z-na+ib) - \frac{ik}{2\pi} \sum \log(z-na+ib)$$

$$\cancel{\frac{ik}{2\pi} \log(z+iy-b)} + \log$$

For a pair, remember

$$w = \frac{ik}{2\pi} \log \left( \sin \frac{\pi(z-ib)}{a} \right) - \frac{ik}{2\pi} \log \left( \sin \frac{\pi(z+ib)}{a} \right)$$

# Vortex Motion

$i \hat{x}, j \hat{y}, k \hat{z}$

## ① Vorticity

$$\left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z}, \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x}, \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

$$② w = \frac{ik}{2\pi} \log(z-z_0) = \phi + i\psi$$

$k = \text{Strength of vortex}$

$\frac{k}{2\pi} \log z_0$

$-k \theta_0 \frac{i}{2\pi}$

$Z_0 = z_0 e^{i\alpha}$

$$u = -\frac{k}{2\pi} \frac{y-y_0}{r_0^2}$$

$$v = \frac{k}{2\pi} \frac{x-x_0}{r_0^2}$$

Mind

$$q = -\frac{k}{2\pi r_0} = \sqrt{u^2 + v^2}$$

IMP

$$-U + iV = + \frac{dw}{dz}$$

## ③ Combination of vortices

- Vortex pair [Equal and opposite]  $\Rightarrow v = \frac{k}{2\pi d}$
- Two of same direction (sign)  $\Rightarrow w_0 = \frac{k_1 + k_2}{2\pi d}$

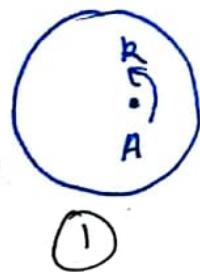
\* Finding  $w =$  superimpose velocity with E sign  
 Add  $\frac{k_2 z}{2\pi d}$  in  $w \Rightarrow \left( \frac{k}{2\pi d} dz = \frac{k z}{2\pi d} \right)$

$$\text{Laplace} = \nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{1}{\gamma} \frac{\partial}{\partial \gamma} + \frac{1}{\gamma^2} \frac{\partial^2}{\partial \theta^2}$$

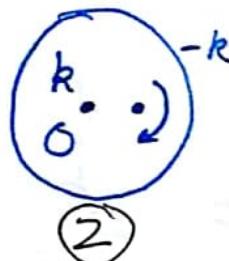
#### ④ Images

$\rightarrow$  Plane (opposite)  $\cdot k$       |       $\therefore -k.$

$\rightarrow$  Circular Cylinder (at A)



$A'$   
 $\cdot$   
 $-k$



$\cdot$   
 $k$   
 $A$

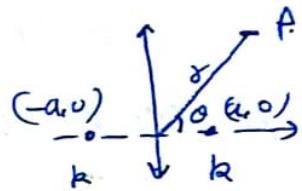
- To Show stationary (At points of vortex).
  - Find  $w'$  by considering others (excluding that vortex)
  - Show  $\left( \frac{dw'}{dz} \right) = 0$  at  $z=a$  (vortex point)

\* Show fluid motion is stationary  $\equiv$  Show  $q$  at all vertices  $= 0$

- Find stagnation points.  $\frac{dw}{dz} = 0$   
( $w$  = include all vortices)

① 2 vortices of same strength, find rel. stream lines

$$w = \frac{ik}{2\pi} \log(r e^{i\theta} - a) + \frac{ik}{2\pi} \log(r e^{i\theta} + a)$$



$$= \frac{ik}{2\pi} \log(r \cos \theta - a + i \sin \theta)$$

$$+ \frac{ik}{2\pi} \log(r \cos \theta + a + i \sin \theta)$$

$r, \theta, k, a$  given

$$\psi = \frac{k}{2\pi} \log \left[ (r^2 + a^2 - 2ar \cos \theta) (r^2 + a^2 + 2ar \cos \theta) \right]^{\frac{1}{2}}$$

Also, system rotates with  $\omega = \frac{2k}{2\pi(2a)^2} = \boxed{\frac{k}{4\pi a^2}}$

So, impose  $\Rightarrow -\omega r = \frac{-kr}{4\pi a^2}$  on system

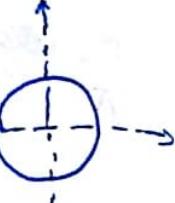
$$\frac{\partial \psi_1}{\partial r} = -\frac{kr}{4\pi a^2} \Rightarrow \boxed{\psi_1 = \frac{-kr^2}{8\pi a^2}}$$

IMP

$$\psi_{\text{fin}} = \psi + \psi_1$$

②  $n$  vortices ( $k$ ) are symmetrical in a cylinder  $|z|=a$ ,  
 find motion of vortices and fluid.  
 (velocity of fluid, time to move a vortex around  $|z|=a$ )

$\Rightarrow n$  vortices are at

$$(ae^{i\frac{2\pi}{n}}, ae^{i\frac{4\pi}{n}}, \dots, a^{i2\pi}) = a^{i\frac{2k\pi}{n}}$$


$$\omega = \frac{ik}{2\pi} \log(z-z_0) + \frac{ik}{2\pi} \log(z-z_1) + \dots + \frac{ik}{2\pi} \log(z-z_{n-1})$$

$$= \frac{ik}{2\pi} \log(z^n - a^n).$$

• Fluid velocity  $= |U-iV| = \left| \frac{dw}{dz} \right| = \left| \frac{ik}{2\pi} \frac{n z^{n-1}}{z^n - a^n} \right| = \frac{kn}{2\pi} \left| \frac{z^{n-1}}{z^n - a^n} \right|$

Velocity of vortex at  $z = ae^{i2\pi} = a$

$$w' = w - \frac{ik}{2\pi} \log(z-a)$$

$$= \frac{ik}{2\pi} \log\left(\frac{z^n - a^n}{z-a}\right) = \frac{ik}{2\pi} \log\left(z^{n-1} + az^{n-2} + \dots + a^{n-1}\right)$$

$$\left( \frac{dw}{dz} \right)_{z=a} = \frac{ik}{2\pi} \left( \frac{(n-1)z^{n-2} + a(n-2)z^{n-3} + \dots + a^{n-2}}{\cancel{na^{n-1}}} \right)_a$$

$$= \frac{ik}{2\pi} \left( \frac{1+2+\dots+n-1}{na} \right) = \frac{ik}{2\pi} \frac{n(n-1)}{2na} = \frac{ik(n-1)}{4\pi a}$$

$$q = \left| \frac{dw}{dz} \right|_{z=a} = \frac{k(n-1)}{4\pi a}$$

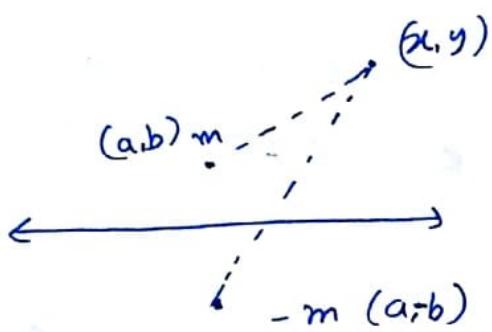
$$T = \frac{2\pi a \cdot 4\pi a}{k(n-1)} = \boxed{\frac{8\pi^2 a^2}{n-1/k}}$$

③ Rigid wall at  $y=0$  and a vortex ( $m$ ).

Find surfaces of equal pressure (need  $x, y$  expressions)  
coordinates of vortex ( $a, b$ )

→ Image system is as shown.

$$W = \frac{im}{2\pi} (- - - -)$$



⇒ For surfaces of equal pressure  $\Rightarrow q$  must be equal.

$$\boxed{\frac{P}{P} + \frac{q^2}{2} = \text{constant}}$$

$$q = \sum \frac{k}{2\pi r_0} = -\frac{k}{2\pi} \left[ \frac{1}{\sqrt{(x-a)^2 + (y-b)^2}} + \frac{1}{\sqrt{(x-a)^2 + (y+b)^2}} \right]$$

$$q^2 = \frac{k^2}{4\pi^2} \left[ \frac{1}{(x-a)^2 + (y-b)^2} + \frac{1}{(x-a)^2 + (y+b)^2} + \frac{2}{\sqrt{(x-a)^2 + (y-b)^2} \sqrt{(x-a)^2 + (y+b)^2}} \right]$$

$$q^2 = \text{constant} \Rightarrow \frac{(x-a)^2 + (y+b)^2 + (x-a)^2 + (y-b)^2 + 2\sqrt{(x-a)^2 + (y-b)^2} \sqrt{(x-a)^2 + (y+b)^2}}{(x-a)^2 + (y-b)^2 (x-a)^2 + (y+b)^2} = k$$

$$\frac{x_1^2 + x_2^2 + 2x_1 x_2}{x_1^2 x_2^2} = k \Rightarrow$$

$$U = -\frac{k}{2\pi} \frac{(y-y_0)}{x_1^2}; V = \frac{k}{2\pi} \frac{(x-x_0)}{x_1^2}$$

Instead

$$q^2 = U^2 + V^2$$

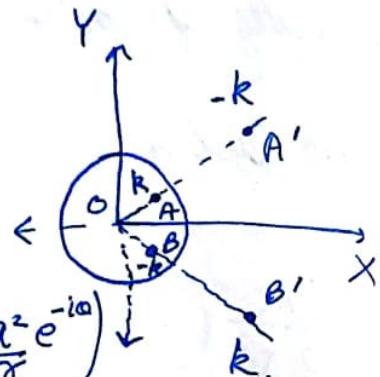
$$\frac{b}{P} + \frac{q^2}{2} = \text{constant}$$

④ A pair of equal and opposite recirculating vortices in  $|z|=a$ . Find path of each vortex.

$$A = r, \theta \quad A' = \frac{a^2}{r}, \theta \\ B = r, -\theta \quad B' = \frac{a^2}{r}, -\theta$$

$$w = \frac{ik}{2\pi} \log(z - re^{i\theta}) + \frac{ik}{2\pi} \log\left(z - \frac{a^2}{r} e^{-i\theta}\right)$$

$$- \frac{ik}{2\pi} \log(z - re^{-i\theta}) - \frac{ik}{2\pi} \log\left(z - \frac{a^2}{r} e^{i\theta}\right)$$



Velocity of vortex at P is due to w of other 3.

$$w_p = \frac{ik \log\left(re^{i\theta} - \frac{a^2}{r} e^{-i\theta}\right)}{2\pi} - \frac{ik \log\left(re^{i\theta} - re^{-i\theta}\right)}{2\pi} - \frac{ik \log\left(re^{i\theta} - \frac{a^2}{r}\right)}{2\pi}$$

$$\psi = \psi_p$$

$$\left(r \cos \theta - \frac{a^2}{r} \cos \theta\right) + i \left(r \sin \theta + \frac{a^2}{r} \sin \theta\right)$$

$$\frac{1}{2} \frac{k}{2\pi} \log \left[ \left(r - \frac{a^2}{r}\right)^2 \cos^2 \theta + \left(r - \frac{a^2}{r}\right)^2 \sin^2 \theta \right] - \frac{k}{2\pi} \log(2r \sin \theta) - \frac{k}{2\pi} \log\left(r - \frac{a^2}{r}\right)$$

$$\psi = - \frac{k}{4\pi} \log \frac{\left(\left(r - \frac{a^2}{r}\right)^2 \times (2r \sin \theta)\right)^2}{r^4 + a^4/r^2 - 2a^2 \cos 2\theta} = b^2$$

$$\text{Path} \Rightarrow b^2 \left(r^4 + a^4/r^2 - 2a^2 \cos 2\theta\right) = (r^2 - a^2)^2 (r^2 \sin^2 \theta)$$

$$b^2 \left(r^4 + a^4/r^2 - 2a^2 r^2 \cos 2\theta\right) = (r^2 - a^2)^2 (r^2 \sin^2 \theta)$$

$$b^2 \left((r^2 - a^2)^2 + 2a^2 r^2 (1 - \cos 2\theta)\right) = (r^2 - a^2)^2 r^2 \sin^2 \theta$$

$$\Rightarrow 4a^2 b^2 r^2 \sin^2 \theta = (r^2 - a^2)^2 (r^2 \sin^2 \theta - b^2) \quad \leftarrow \text{Read from}$$

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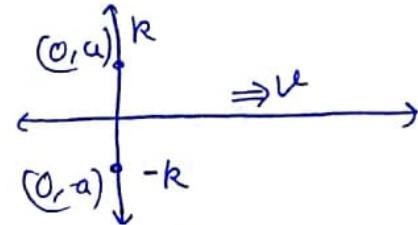
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③ Two parallel equal and opposite vortices at distance  $2a$ ,  
above streamlines relative to vortices

$$\log \frac{x^2 + (y-a)^2}{x^2 + (y+a)^2} + \frac{y}{a} = C$$

As a vortex pair (equal & opposite)

Move with constant  $V = \frac{k}{4\pi a}$



$q = \frac{dw}{dz} \Rightarrow w = \frac{kz}{4\pi a}$  (to be reversed added to given  $w$ )  
 (for relative streamlines)

$$w = \frac{i k}{2\pi} \log(z-ia) - \frac{i k}{2\pi} \log(z+ia) + \frac{kz}{4\pi a}$$

To find  $\psi \Rightarrow$  get imaginary part of  $w$

$$\psi = \frac{k}{2\pi} \log \sqrt{x^2 + (y-a)^2} - \frac{k}{2\pi} \log \sqrt{x^2 + (y+a)^2} + \frac{ky}{4\pi a}$$

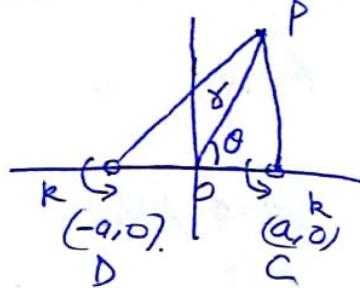
$$\psi = \frac{k}{4\pi} \left[ \log \frac{x^2 + (y-a)^2}{x^2 + (y+a)^2} + \frac{y}{a} \right] = C$$

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$$\log z = \log r + i\theta.$$

⑤ Find relative streamlines for two same vortices with same spin and strength.

(Given  $\theta$  is measured from origin,  $2a$ )



$$\omega = \frac{k_1 + k_2}{2\pi(A_1 A_2)^2} = \boxed{\frac{k}{4\pi a^2}}$$

$$CP = (\gamma^2 + a^2 - 2ar\cos\theta)^{1/2} \quad DP = (\gamma^2 + a^2 + 2ar\cos\theta)^{1/2}$$

$$\begin{aligned} \psi &= \frac{k}{2\pi} \log r_1 + \frac{k}{2\pi} \log r_2 = \frac{k}{2\pi} \log(r_1 r_2) \\ &= \frac{k}{4\pi} \log((\gamma^2 + a^2)^2 - 4a^2\gamma^2 \cos^2\theta). \end{aligned}$$

To find  $\psi$  relative we need to look at motion of vortices.

L Impose velocity  $(-w\gamma)$  on the system

$$\frac{\partial \psi_1}{\partial \gamma} = -w\gamma = -\frac{k\gamma}{4\pi a^2} \Rightarrow \boxed{\psi_1 = -\frac{k\gamma^2}{8\pi a^2}}$$

So, streamlines are (relative)

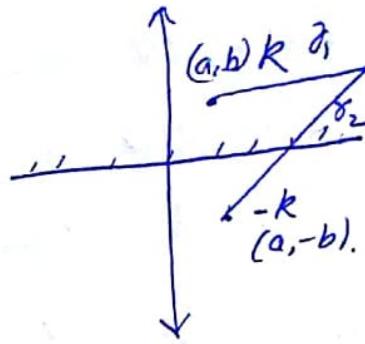
$$\psi + \psi_1 = \frac{k}{4\pi} \left( \log(\gamma^2 + a^2 - 2a^2\gamma^2 \cos 2\theta) - \frac{\gamma^2}{2a^2} \right) = \text{constant}$$

⑥ To find surfaces of equal pressure.

$$\frac{P}{P} + \frac{q^2}{2} = \text{constant}$$

$\Rightarrow q^2 = \text{constant}$  gives required surfaces.

$$q^2 = U^2 + V^2$$



$$U = \left( -\frac{k}{2\pi} \frac{(y-y_1)}{r_1^2} \right) = -\frac{k}{2\pi} \frac{(y-b)}{r_1^2} + \frac{k}{2\pi} \frac{(y+b)}{r_2^2}$$

$$V = \left( \frac{k}{2\pi} \frac{x-a}{r_1^2} \right) - \left( \frac{k}{2\pi} \frac{x+a}{r_2^2} \right)$$

• However, vortex pair is moving with velocity  $\frac{k}{2\pi \cdot 2b} = \left( \frac{k}{4\pi b} \right)$   
we need to superimpose this velocity.

$$U = -\frac{k}{2\pi} \frac{y-b}{r_1^2} + \frac{k}{2\pi} \frac{y+b}{r_2^2} - \frac{k}{4\pi b}$$

$$q^2 = \frac{k^2}{4\pi^2} \left[ \frac{(x-a)^2 + (y-b)^2}{r_1^4} + \frac{(x-a)^2 + (y+b)^2}{r_2^4} - \frac{2(x-a)^2}{r_1^2 r_2^2} + \frac{1}{4b^2} \right. \\ \left. - 2 \frac{(y^2 - b^2)}{r_1^2 r_2^2} - 2 \frac{(y+b)}{r_2^2} \cdot \frac{1}{2b} + 2 \cdot \frac{1}{2b} \frac{(y-b)}{r_1^2} \right]$$

$$= \frac{k^2}{4\pi^2} \left[ \frac{1}{r_1^2} + \frac{1}{r_2^2} - \frac{2((x-a)^2 + (y^2 - b^2))}{r_1^2 r_2^2} + \frac{1}{b} \left( \frac{y-b}{r_1^2} - \frac{(y+b)}{r_2^2} \right) + \frac{1}{4b^2} \right] \\ = \underline{\text{constant}}$$

$$r_1^2 + r_2^2 - 2((x-a)^2 + (y^2 - b^2)) + \frac{1}{b} ((y r_2^2 - b r_2^2) - y r_1^2 - b r_1^2) = k r_1^2 r_2^2$$

$$(y-b)^2 + (y+b)^2 - 2(y^2 - b^2) + \frac{1}{b} y(r_2^2 - r_1^2) - (r_2^2 + r_1^2) = k r_1^2 r_2^2$$

$$4b^2 + 4y^2$$

$$r_2^2 - r_1^2 = 4by$$

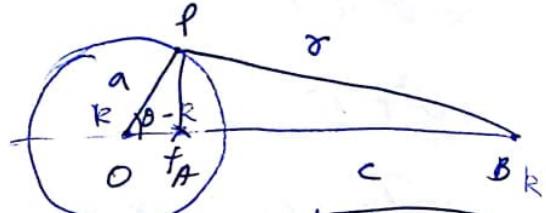
⑦ A vortex outside a cylinder (a) at distance  $c$  ( $\text{m}$ )

Show speed of fluid at cylinder surface is  $\frac{k}{2\pi a} \left( 1 - \frac{c^2 - a^2}{r^2} \right)$

$r$  is the distance of point from vortex.

Equivalent system is as shown.

$$w = \frac{ik}{2\pi} (\log z + \log(z-c) - \log(z-f))$$



$$\frac{a^2}{f} = c$$

$$\frac{dw}{dz} = \frac{ik}{2\pi} \left( \frac{1}{z} + \frac{1}{z-c} - \frac{1}{z-f} \right) = \frac{ik}{2\pi} \left( \frac{z^2 - 2fz + fc}{z(z-c)(z-f)} \right)$$

$$q = \left| \frac{dw}{dz} \right| = \frac{k}{2\pi} \frac{|z^2 - 2fz + fc|}{OP \cdot AP \cdot BP}$$

$$= \frac{k}{2\pi} \frac{2a(f - a \cos \theta)}{OP \cdot AP \cdot BP}$$

$\Delta OPA \sim \Delta OAP$ .

$$\frac{AP}{OP} = \frac{BP}{OB} \Rightarrow AP = \frac{OP \cdot r}{k}$$

$$q = \frac{k}{2\pi} \frac{2a(f - a \cos \theta)}{a \cdot \frac{ar}{c} \cdot r}$$

$$\cos \theta = \frac{a^2 + c^2 - r^2}{2ac}$$

$$= \frac{k}{2\pi} \frac{2c(f - \frac{a^2 + c^2 - r^2}{2ac} \cdot \frac{a^2 + c^2 - r^2}{2ac})}{ar^2}$$

$$= \frac{ka}{\pi a r^2} \left( \frac{2a^2 - a^2 - c^2 + r^2}{2f} \right)$$

$$= \frac{k}{2\pi a} \left( \frac{a^2 - c^2}{r^2} + 1 \right) \underline{\text{Ans}}$$

$$z = a e^{i\theta}$$

$$|z^2 - 2fz + fc| = |a^2 e^{2i\theta} - 2af e^{i\theta} + a^2|$$

$$= |a^2 + 2a^2 \cos 2\theta - 2af \cos \theta + i(a^2 \sin 2\theta - 2af \sin \theta)|$$

$$= (a^4 + a^4 \cos^2 2\theta + 4a^2 f^2 \cos^2 \theta + 2a^4 \cos 2\theta - 4a^3 f \cos \theta \cos 2\theta - 4a^3 f \sin \theta \cos \theta + a^4 \sin^2 2\theta + 4a^2 f^2 \sin^2 \theta - 4a^3 f \sin \theta \cos 2\theta - 6a^4 + a^4 + 4a^4 f^2 + 8a^3 f (-\cos \theta) + 2a^4 \cos 2\theta)^{1/2}$$

$$= 2a^4 + (4a^2 f^2 - 8a^3 f \cos \theta + 4a^4 (\cos^2 \theta))^{\frac{1}{2}}$$

$$= 2a \left( f^2 - 2af \cos \theta + a^2 \cos^2 \theta \right)^{\frac{1}{2}}$$

$$= \boxed{2a (f - a \cos \theta)}$$

③ Prove necessary and sufficient condition that vortex lines may be at right angles to streamlines are,

$$(U, V, W) = M \left( \frac{\partial \phi}{\partial x}, \frac{\partial \phi}{\partial y}, \frac{\partial \phi}{\partial z} \right)$$

$$\text{Vortex lines} \rightarrow \frac{dx}{\Omega_x} = \frac{dy}{\Omega_y} = \frac{dz}{\Omega_z} \Rightarrow \begin{aligned}\Omega_x &= \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \\ \Omega_y &= \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \\ \Omega_z &= \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\end{aligned}$$

$$\text{Streamlines} \rightarrow \frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w}$$

At right angles if

$$u(w_y - v_z) + v(u_z - w_x) + w(v_x - u_y) = 0$$

↓

This is the necessary and sufficient condition so that  $udx + vdy + wdz$  is a perfect differential.

$$udx + vdy + wdz = M d\phi = M \left( \frac{\partial \phi}{\partial x} dx + \frac{\partial \phi}{\partial y} dy + \frac{\partial \phi}{\partial z} dz \right)$$

$$\Rightarrow u = M \frac{\partial \phi}{\partial x}, \quad v = M \frac{\partial \phi}{\partial y}, \quad w = M \frac{\partial \phi}{\partial z}$$

- If vorticity is constant  $\Rightarrow$  Show  $\nabla^2 u = \nabla^2 v = \nabla^2 w = 0$   
(Use continuity equation)

Get 2  $\frac{\partial^2 u}{\partial y^2}, \frac{\partial^2 u}{\partial z^2}$  from constant vorticity

Then  $\frac{\partial^2 u}{\partial x^2}$  from continuity equation

$$(2) \rightarrow u dx + v dy + w dz = d\theta + \lambda d\chi$$

$\theta, \lambda, \chi$  are functions of  $x, y, z, t$

Prove vortex lines are lines of intersection of surfaces  $\lambda = \text{constant}$  and  $\chi = \text{constant}$ .

- $udx + vdy + wdz = \frac{\partial \theta}{\partial t} dt + \frac{\partial \theta}{\partial x} dx + \frac{\partial \theta}{\partial y} dy + \frac{\partial \theta}{\partial z} dz + \lambda \left[ \frac{\partial \chi}{\partial x} dx + \frac{\partial \chi}{\partial y} dy + \frac{\partial \chi}{\partial z} dz + \frac{\partial \chi}{\partial t} dt \right]$

- $u = \frac{\partial \theta}{\partial x} + \lambda \frac{\partial \chi}{\partial x}$

$$v = \vdots \quad \vdots$$

$$w = \vdots$$

- $0 = \frac{\partial \theta}{\partial t} + \lambda \frac{\partial \chi}{\partial t}$

Spin components

$$\begin{aligned} 2. \Omega_x &= \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} = \frac{\partial}{\partial y} \left( \frac{\partial \theta}{\partial z} + \lambda \frac{\partial \chi}{\partial z} \right) - \frac{\partial}{\partial z} \left( \frac{\partial \theta}{\partial y} + \lambda \frac{\partial \chi}{\partial y} \right) \\ &= \cancel{\frac{\partial^2 \theta}{\partial y \partial z}} + \frac{\partial \lambda}{\partial y} \cdot \frac{\partial \chi}{\partial z} + \lambda \cancel{\frac{\partial^2 \chi}{\partial y \partial z}} - \cancel{\frac{\partial^2 \theta}{\partial z \partial y}} - \lambda \cancel{\frac{\partial^2 \chi}{\partial z \partial y}} - \frac{\partial \lambda}{\partial z} \frac{\partial \chi}{\partial y} \\ &= \begin{vmatrix} \frac{\partial \chi}{\partial y} & \frac{\partial \chi}{\partial z} \\ \frac{\partial \chi}{\partial y} & \frac{\partial \chi}{\partial z} \end{vmatrix} \end{aligned}$$

$$\text{Hence } 2\Omega_y = \begin{vmatrix} \frac{\partial \chi}{\partial z} & \frac{\partial \chi}{\partial x} \\ \frac{\partial \chi}{\partial z} & \frac{\partial \chi}{\partial x} \end{vmatrix}, \quad 2\Omega_z.$$

$$\text{Hence } 2 \left( \Omega_x \frac{\partial \lambda}{\partial x} + \Omega_y \frac{\partial \lambda}{\partial y} + \Omega_z \frac{\partial \lambda}{\partial z} \right) = 0 \quad \left. \begin{array}{l} \text{Implies the} \\ \text{requirement} \end{array} \right\}$$

$$\Omega_x \frac{\partial \chi}{\partial z} + \Omega_y \frac{\partial \chi}{\partial x} + \Omega_z \frac{\partial \chi}{\partial y} = 0$$