2014-UPSC-ESEDELTA PORO

gq-5(e) eq. of motion of a compound pendulum hoose of hoo (T-kinetic energy of penduly)
let OG=h V= potential energy = - Mghcoso

(-ive because below fixed point o) 80, by 1=T-V= 1-M2+ Mghloso since Here O is the only generalised conordinate;

since Lis free of time so,

M-T+V=1mx262-mghcoso From (D), DL-po = Mk26 =) 6 = po mkz Now (2) becomes. H= I mk Po - mgh cos D

m²ky

DELTA PONO

so Hamiltonian equations are-

24=-Po 124-6 gweg

-po = Mgh sind , 0 = po Mgh sind , 0 = po Mk2

 $\Rightarrow \hat{\theta} = p\theta = -Mghsin\theta = -ghsin\theta$ $mkl \qquad jak 2 \qquad k^2$

is the equation of motion for Compaind perdulum

Problem 19. Show that velocity potential
$$= \frac{1}{2} \log \left[\frac{(x+a)^2 + y^2}{(x-a)^2 + y^2} \right]$$

gives a possible notion. Determine the form of stream lines and the curves of equal speed.

Solution. Given, $\phi = \frac{1}{2} \log ((x+a)^2 + y^2) - \frac{1}{2} \log ((x-a)^2 + y^2)$(1)

 $\frac{\partial}{\partial x} = \frac{x + a}{(x + a)^2 + y^2} = \frac{(x - a)}{(x - a)^2 + y^2}$

des, sinks & Doublets (Fluid Dynamics)

$$\frac{\partial^2 \phi}{\partial x^2} = \frac{[(x+\alpha)^2 + y^2] - 2(x+\alpha)^2}{[(x+\alpha)^2 + y^2]^2} - \frac{[(x-\alpha)^2 + y^2] - 2(x-\alpha)^2}{[(x-\alpha)^2 + y^2]^2}$$

$$\frac{\partial^2 \phi}{\partial x^2} = \frac{y^2 - (x+\alpha)^2}{[(x+\alpha)^2 + y^2]^2} - \frac{y^2 - (x-\alpha)^2}{[(x-\alpha)^2 + y^2]^2} \qquad \dots (2)$$

By (1),
$$\frac{\partial \phi}{\partial y} = \frac{y}{(x+a)^2 + y^2} - \frac{y}{(x-a)^2 + y^2}$$

$$\frac{\partial^2 \phi}{\partial y^2} = \frac{(x+a)^2 + y^2 - 2y^2}{[(x+a)^2 + y^2]^2} - \frac{(x-a)^2 + y^2 - 2y^2}{[(x-a)^2 + y^2]^2} \dots (3)$$

Adding (2) and (3), $\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial x^2} = 0$ or $\nabla^2 \phi = 0$.

Thus the equation of continuity is satisfied and so (1) gives a possible liquid

Second Part. To determine stream lines.

$$-\frac{\partial \phi}{\partial x} = u = -\frac{\partial \psi}{\partial y}, -\frac{\partial \phi}{\partial y} = v = \frac{\partial \psi}{\partial x}.$$

Hence
$$\frac{\partial x}{\partial x} = \frac{\partial y}{\partial y}$$
, $\frac{\partial y}{\partial y} = -\frac{\partial y}{\partial x}$

Now
$$\frac{\partial y}{\partial y} = \frac{x+\alpha}{(x+\alpha)^2 + y^2} - \frac{x-\alpha}{(x-\alpha)^2 + y^2}$$

Integrating w.r.t. y.

$$\Psi = \tan^{-1} \frac{y}{x+a} - \tan^{-1} \frac{y}{x-a} + F(x) \dots$$
 ... (4)

where F(x) is constant of integration. To determine F(x).

$$\frac{\partial y}{\partial x} = -\frac{\partial \delta}{\partial y} = \frac{y}{(x+a)^2 + y^2} + \frac{y}{(x-a)^2 + y^2} \qquad \dots (5)$$

By (4),
$$\frac{\partial y}{\partial x} = -\frac{y}{(x+a)^2 + y^2} + \frac{y}{(x-a)^2 + y^2} + F'(x)$$
 ... (6)

Equating (5) to (6), F'(x) = 0. Integrating this

F(x) = absolute const. and hence neglected.

Since it has no effect on the fluid motion.

Now (4) becomes

$$\forall = \tan^{-1} \frac{y}{x + \alpha} - \tan^{-1} \frac{y}{x - \alpha} \qquad ... (7)$$

$$= \tan^{-1} \frac{-2\alpha y}{x^2 - \alpha^2 + y^2}.$$

Stream lines are given by w = const., i.e.,

$$\tan^{-1}\left[\frac{-2ay}{x^2-a^2+y^2}\right] = \text{const. or } \frac{y}{x^2-a^2+y^2} = \text{const.}$$

If we take const. = 0, then we get y = 0, i.e., x-axis. If we take conts. = ∞ , then we get circle $x^2 - a^2 + y^2 = 0$.

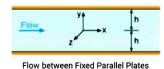
26 (8c)

20 (00)

Find Navier-Stokes equation for a steady laminar flow of a viscous incompressible fluid between two infinite parallel plates.

* * *

Flow between Fixed Parallel Plates



Consider steady, incompressible, laminar flow between two infinite parallel horizontal plates as shown in the figure. The flow is in the x- direction, hence there is no velocity component in either the y- or z- direction (i.e., v = 0 and w = 0). The steady-state continuity equation becomes

$$\frac{\partial u}{\partial x} = 0$$

[1]

From Eqn. 1, it can be concluded that the velocity u is a function of both y and z only. Since the plates are infinitely wide, it can be argued that the velocity u should not be a function of z, i.e., it must be a function of y only, u = u(y).

Applying the <u>Navier-Stokes</u> equations, along with the assumptions that v = 0, w = 0 and u = u(y), yields

$$\frac{\partial p}{\partial x} = \mu \frac{d^2 u}{dy^2}$$
[2]
$$\frac{\partial p}{\partial y} = -pg$$

$$\frac{\partial p}{\partial z} = 0$$
[4]

Eqn. 4 indicates that the pressure is a function of x and y. Integrate Eqn. 3 to yield

$$p = -\rho gy + g_1(x)$$

Hence it can be concluded that $\frac{\partial p}{\partial x}$ is a function of x only. Now, integrate Eqn. 2 twice with respect to y, and treat $\frac{\partial p}{\partial x}$ as a constant (with respect to y) to give:

$$u = \frac{1}{2\mu} \frac{\partial p}{\partial x} y^2 + c_1 y + c_2$$

Applying the no-slip conditions (i.e., the fluid is "stuck" to the

$$\frac{\partial p}{\partial z} = 0$$

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Applying the no-slip conditions (i.e., the fluid is "stuck" to the plates, or u = 0 at y = ±h) to determine the coefficients as follows:

$$c_1 = 0$$
 and $c_2 = -\frac{h^2}{2\mu} \frac{\partial p}{\partial x}$



Velocity Profile

The velocity profile becomes

$$u = \frac{1}{2\mu} \frac{\partial p}{\partial x} (y^2 - h^2)$$

which is a parabola. The total volumetric flow per linear depth can be obtained integrating the velocity to give

$$q = \int_{-h}^{h} u \, dy = -\frac{2h^3}{3\mu} \left(\frac{\partial p}{\partial x} \right)$$

Note, q is per linear depth, which is different than Q which is the total volumetric flow rate.

Also note that the flow is negative, i.e. to the left, for a positive pressure gradient, dp/dx. This is due to the gradient definition where decreasing pressure to the right is negative.