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Senior Honours Project Report

BSc GEOPHYSICS

Improving the prediction of magnetic disturbances from solar wind parameters

Student:

Kittiphon BOONMA

Supervisor:

Dr.Brian HAMILTON

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Abstract

Magnetic indices are used as a measure of the intensity of various geomagnetic activities. Dst index measures the strength of the equatorial Ring Current. A previous study by Newell *et al.* (2007) proposed a function, $d\Phi_{MP}/dt$, which is said to give the best correlation between the solar wind parameters and the geomagnetic indices. The coupling function E_{SR} was claimed to be the second best. This study aims to investigate these claims by applying corrections to indices and using new data (2000-2013) which the previous study did not have. Dst index measures both the ring current (Est index), and the induced current in the conducting Earth (Ist index). This gives rise to the idea of improving the prediction model of the geomagnetic disturbances by replacing the Dst index with Est index, and the Ist is eliminated because it is produced indirectly by the solar wind parameters. The analytical results from this study confirms that the correlation between the top coupling function with Dst can be improved by replacing it with Est index. However, this is not the case for the second best function, E_{SR} . An additional hypothesis is that the measure of auroral electrojet indices, AU and AL, contain contributions from the ring current. The results from the investigation confirmed the hypothesis since about 8% of the AU index and 6% of the AL index were the contribution from the x-component of the ring current. Removing the effect of ring current improves AU index but worsens AL index. Overall, since the claimed top function $d\Phi_{MP}/dt$ was not consistently better than E_{SR} when changes were introduced, $d\Phi_{MP}/dt$ function may not be as robust and correlate best with every indices as Newell *et al.* (2007) claimed it to be.

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Chapter 1

Introduction

The monitoring and accurate prediction of magnetic disturbances in the Earth's upper atmosphere, the ionosphere and the magnetosphere, has become increasingly important due to the expansion of human activity at high altitudes, in space, and on the ground where sensitive technologies, such as power grids and directional drilling for oil and gas, can be affected. Understanding the coupling mechanisms and how the solar wind affects the Earth's magnetic field is crucial. Therefore, it is necessary to quantitatively assess the underlying statistical and correlating models.

Geomagnetic indices represent a measure of geomagnetic activity, which is a key signature of how the Earth's ionosphere and magnetosphere respond to interactions with the solar wind. There are many types of indices, each describing different geomagnetic activity, and many try to isolate the activity from different sources (current systems) in the ionosphere/magnetosphere. For the past 15 years in the field of Solar Terrestrial studies, these indices have become a key parameter in detecting and describing space weather events.

The overall aim of this study is to investigate the improvements on geomagnetic indices (Dst, Est, AU, and AL) and their ability to be predicted by the solar wind parameters. There are three parts to the investigation:

- 1) The first hypothesis is that it may be more appropriate to predict the Est index, rather than the Dst index (as many scientists have done in the past). The correlation between the indices and a coupling function is used to test this hypothesis.
- 2) We investigate the possibility of improving the AU and AL indices by subtracting contaminating effects arise from the ring current. Again, using coupling functions, it can be tested whether the corrected AU and AL indices have a stronger correlation than the original indices.
- 3) A study previously done by Newell *et al.* (2007) has found a good correlation with the Dst index (and many others) and claimed that $d\Phi_{MP}/dt$ is the best coupling function. The analysis from this study will be used to test if the claim stands up to scrutiny, especially with the new data (up to 2013) which Newell *et al.* (2007) did not have. The corrected indices will also be used to test the claim, and the investigation will extend into looking at the second best coupling function E_{SR} (claimed by Newell *et al.* (2007)) as well.

This chapter introduces the reader to the project and the underlying physical background. The analytical methods are presented in Chapter 3 (p. 14). The results and corresponding

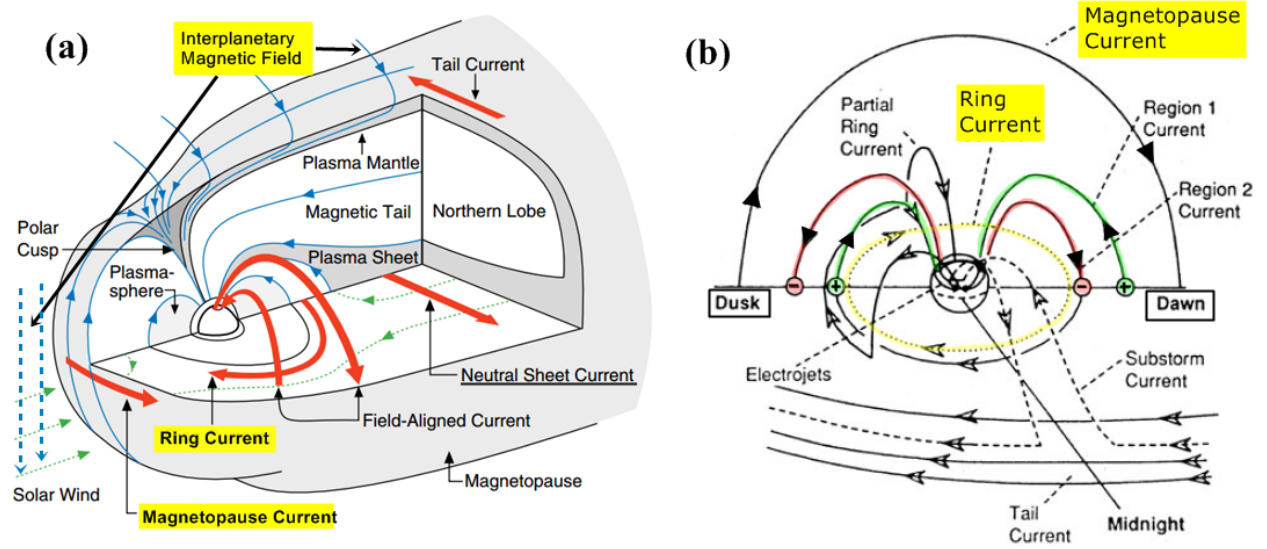


Figure 1.1: (a) A diagram of the magnetosphere showing the current systems and plasma regions (Source: Russell (2000)). (b) A diagram of a cross-section of the outer magnetosphere, looking towards the Sun. The diagram represents various current systems within the magnetosphere-ionosphere cavity, which give rise to magnetic activity (Source: Modified from McPherron (1995)).

discussions are presented in and Chapter 4(p. 17). The overall discussion and conclusion will be presented in Chapter 5(p. 24).

1.1 Solar wind and Magnetosphere

Other than solar radiation, solar wind is regarded as the second most crucial connection between the Earth and the Sun. The *solar wind* is defined as a stream of ionized particles ejected from the Sun. Within the solar wind stream, there is an embedded *Interplanetary Magnetic Field (IMF)*(Figure 1.1 (a)) which the solar wind carries through the solar system. The solar wind carries this weak magnetic field past the magnetized Earth, where it controls the strength of the coupling between the upper atmosphere (magnetosphere and ionosphere) and the solar wind.

As the Earth travels around the Sun in this dynamical solar-wind environment, the Earth's magnetic field interacts with the solar flux. The blast of solar wind envelops and deflects around the Earth's dipole geomagnetic field, compressing the magnetosphere to 11 Earth-radii (R_e) on the dayside, and stretches the magnetosphere outwards into a long tail, far past the Moon's orbit ($60R_e$) on the nightside (Campbell, 2003).

The energy, plasma, and momentum are transferred from the solar wind into the Earth's upper atmosphere through the merging of the Earth's magnetospheric field lines and the IMF. The magnetopause (Figure 1.1 (a)) is the boundary between the Earth's magnetic field and the solar wind. The merging process occurs within the magnetopause, however, the merging location depends on the orientation of IMF (Russell, 2000).

Figure 1.2 shows the *Geocentric Solar Magnetospheric (GSM)* coordinate system which is typically used when dealing with the orientation of IMF. B_x lies along the Earth-Sun plane and is positive towards the Sun. B_y is the cross product of the magnetic dipole

B_z and B_x , and B_z represents the cross product of B_x and B_y . It is considered the best coordinate system for the study of the effect of IMF components on the magnetospheric phenomena because by applying this system, the direction of the Earth's geomagnetic field near the 'nose' of the magnetosphere is well-ordered. This coordinate system is also applied to the components of the field used to calculate the IMF clock angle in section 1.4.1.

There are two scenarios in which the IMF can merge with the Earth's magnetic field, both of which are illustrated by the Dungey's Model in Figure 1.3:

1. Figure 1.3 (a) shows the case where IMF is directed southward ($B_z < 0$), and the dayside shrinks by flux ablation. According to the Dungey's Model, with this configuration, the reconnection occurs with closed magnetospheric field lines on the dayside magnetopause. This leads to the formation of the opened field lines, which are dragged by the magnetic tension away from the Sun. On the dayside, the energy flows from the magnetic field into the plasma, however, there is some electromagnetic energy flux at the tail, and this energy flows from the plasma into the magnetic field (Russell, 2000).

The energy transferred into the magnetic field is stored as magnetic energy in the magnetotail (Menvielle *et al.*, 2011). This stored energy can be released from the tail randomly during violent bursts of energy, called *substorms*, and causes the plasma to flow into the magnetosphere. The field and plasma continue to move towards the reconnection point on the dayside, as shown by the green dashed line in Figure 1.1 (a), and the cycle is repeated.

2. Figure 1.3 (b) illustrates the situation when IMF is directed solely northward ($B_z > 0$), and the dayside grows by flux accretion. In this case, the reconnection still occurs, however, the effect on the magnetosphere will be different from the previous case. The reconnection occurs with opened magnetospheric field lines tailward of the cusp (Figure 1.1 (a)).

However, the northward-IMF model is not consistent with the results from previous studies on electron data and on solar proton data. Therefore, it is important to note that it is highly doubtful that the same field line would connect to both the north and south neutral points at the tail (Russell, 1972a).

To put this into context, these merging events illustrate the physical processes of how the solar wind continually feeds energy into the Earth's atmosphere (Russell, 2000) and affects the magnetosphere and ionosphere. Ultimately, this produces the disturbances which are measured by the indices studied in this project. Therefore, it is justifiable and appropriate to consider solar wind based coupling functions (Section 1.4). The northward directed IMF usually causes less disturbance and, hence, smaller indices. This explains the importance of the clock angle term (Equation 1.7 in the coupling functions).

1.2 Current Systems

There are several current systems in the upper atmosphere that magnetic perturbations can originate from. This study is mainly concerned with two current systems, both of which are described in this section.

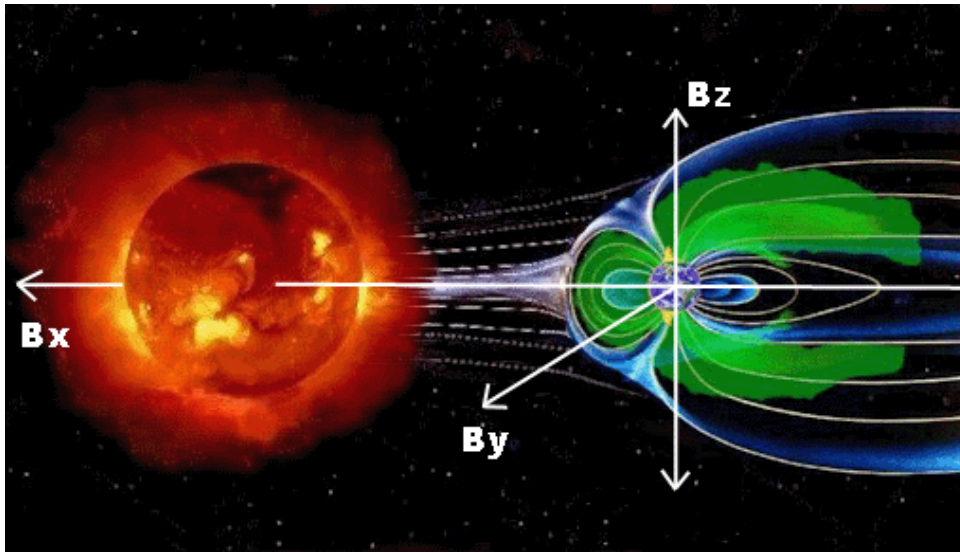


Figure 1.2: Diagram showing the *Geocentric Solar Magnetospheric (GSM)* coordinate system. x-axis is Earth-Sun line and is positive towards the Sun. y-axis is the cross product of the magnetic dipole axis and the x-axis. z-axis is the cross product of x- and y-axes. (Source: <http://poleshift.ning.com>)

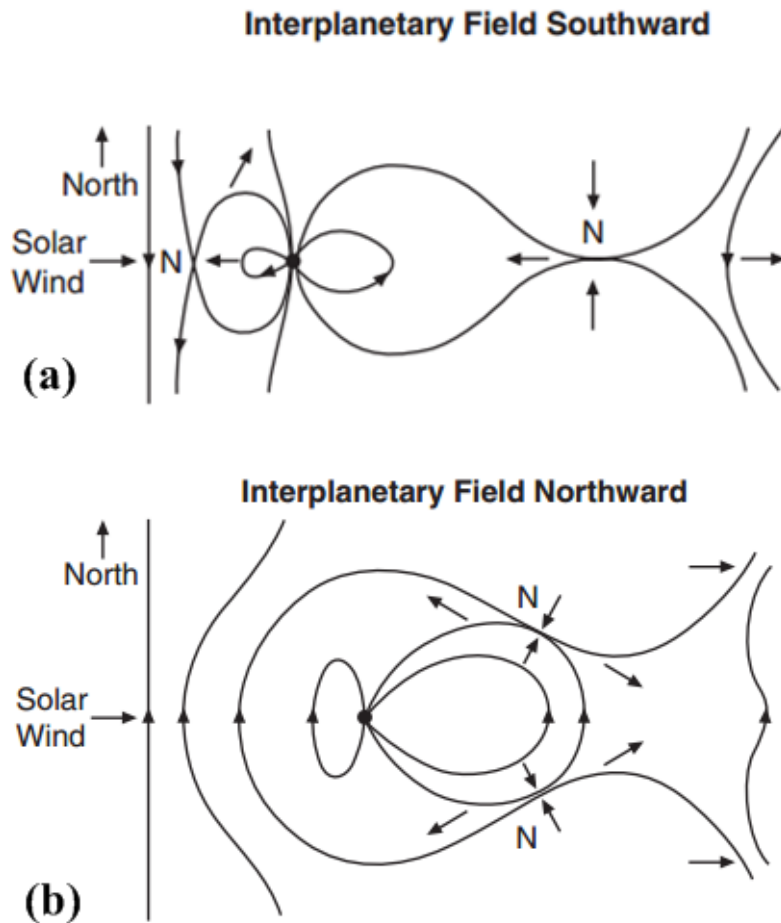


Figure 1.3: *The Dungey's Model of the Magnetosphere* showing the interaction between the solar wind Interplanetary Magnetic Field (IMF) and the Earth's magnetosphere. 'N' represents the reconnection area which is also known as Neutral point. (a) A schematic illustration of the open magnetosphere and the reconnection theory with IMF pointing southwards and the Earth's field pointing northward. (b) A schematic illustration of the closed magnetosphere with both IMF and Earth's magnetic field pointing northward. (Source: Russell (2000))

1.2.1 Magnetopause Current

There are several sources of electric currents co-existing within the magnetosphere-ionosphere system. They can be either directly or indirectly associated with the dynamical interaction between the Earth's environments and the Sun (Menvielle *et al.*, 2011). This study mainly focuses on the currents arising from the interaction between the Earth's magnetic field and the solar wind. One of the current systems, namely *Magnetopause Current*, will be described in this section.

The *Magnetopause Current*, or also known as '*Chapman-Ferraro Current*', is the basic nature of the interaction between the Earth's magnetic field and the solar wind. The existence of such a current system was first deduced in 1930's by Sydney Chapman and Vincenzo Ferraro (Chapman and Ferraro, 1931). Magnetopause represents a surface of current layer (Figure 1.1 (a)), and separates the solar wind (high density, low magnetic fields (\mathbf{B})) from the magnetosphere (low density, high magnetic fields (\mathbf{B})). In the first approximation, magnetopause forms at a distance of about 8-11 R_e where the magnetic pressure of the Earth's field equals the dynamic pressure of the solar wind (Menvielle *et al.*, 2011). As the solar wind particles begin to penetrate the magnetopause they experience the Lorentz force provided by the magnetopause current, which deflects the electrons and protons of the solar wind in opposite directions normal to the field, shielding the Earth's magnetic field from the solar wind. This magnetopause current, shown in Figure 1.1 (b), flows towards *dusk* on the dayside, and towards *dawn* on the nightside magnetopause where it is called the *tail current* (Figure 1.1 (b)). A type of current called the *Neutral Sheet Current* flows duskward in the magnetic equatorial plane and acts as a barrier confining the tail part of the magnetopause current inside the magnetosphere (Figure 1.1 (a)). From the right-hand rule, it can be speculated that the ground magnetic field will be affected by a southward perturbation. The magnetopause current is important to this study because it is used in the correction of the Dst index in Section 1.4.1.

1.2.2 Ring-current

Ring current is produced from the eastward drift of electrons and the westward drift of protons in a ring around the magnetic equator of the Earth at a distance of 3 – 5 R_e (Figure 1.1), where the magnetic field is approximately dipolar (Menvielle *et al.*, 2011). Burton *et al.* (1975) used what was understood, regarding the control of the reconnection process of IMF, to demonstrate that the ring current could be predicted from solar wind parameters. The amount of convected southward IMF is proportional to the injection rate of the energy that goes into the ring current (Le *et al.*, 2004). Similar to the tail current, the ring current flows duskward on the Earth's magnetic equatorial plane, and by applying the right-hand rule, its ground effect is the reduction in the main field strength at low latitudes due to the southward perturbation. This is similar to the effect of the magnetopause current.

For the purpose of this study, ring current is important because its strength directly varies in response to the solar wind. Since the ring current is associated with the solar wind-magnetosphere interaction, the solar wind parameters can be studied through Dst index, which measures the strength of the ring current.

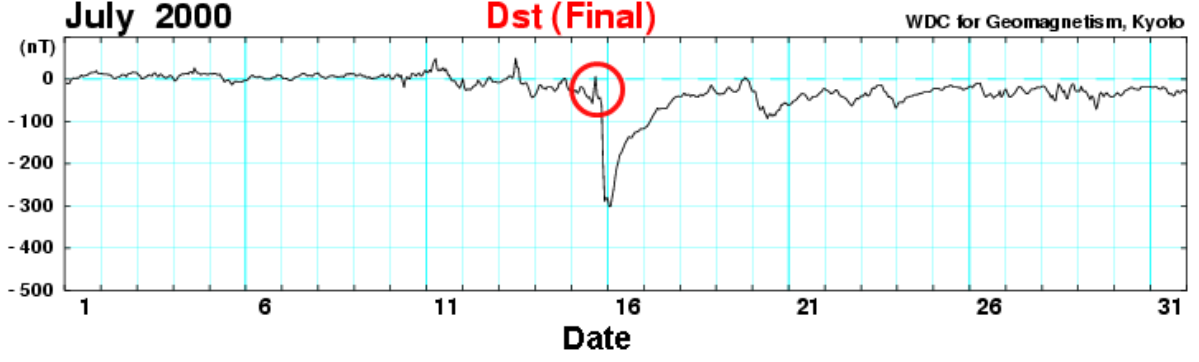


Figure 1.4: A graph showing Hourly Equatorial Dst Values of July 2000. The y-axis is the H-component of the field, while the x-axis represents each day of the month. The increase in Dst (circled) amplitude prior to a large drop is a result of magnetospheric compression. The large decrease in Dst amplitude signifies the main phase of the storm. (Source: World Data Centre for Geomagnetism, Kyoto)

1.3 Geomagnetic Indices

1.3.1 Disturbance Storm Time (Dst) Index

During a period of intense solar activity, the southward IMF that merge with the Earth’s magnetic field can be strong and have long duration. The reconnection occurs at the ‘nose’ of the magnetosphere (Figure 1.3 (a)) and causes compression of the magnetosphere on the day side and continues down the tail. This compression distorts the Earth’s dipole field and provides an entry for the solar wind particles, which cause the geomagnetic disturbances i.e. geomagnetic storms. These storms are the main source that provide energy for the ring current (Section 1.2.2) which has a westward-flowing current. This westward current causes a horizontal (H)-component field depression worldwide (Campbell, 2003).

Disturbance Storm Time index (Dst) is one of many geomagnetic indices. It monitors the magnetic storm level worldwide, as well as representing the axially symmetric magnetic disturbance at the dipole equator on the Earth’s surface (Menvielle *et al.*, 2011). Dst index is derived from averaging the H-component of the Earth’s magnetic field from well-spaced mid- or low-latitude observatories from all over the world.

The severity of the geomagnetic disturbance is represented by the magnitude of the decrease in H-component field, i.e. the disturbance field is southward. An example of Dst data is shown in Figure 1.4. The circled area emphasizes a sharp increase prior to a large drop in amplitude, which is caused by the compression of the magnetosphere, and the following drop after the upward spike is caused by the ring current.

The derivation of Dst described in this section is based on Sugiura and Kamei (1991). At a time ‘ t ’, the $Dst(t)$ is defined as the average of the disturbance variation of H-component of the field, $D_i(t)$, at ‘ n ’ observatories. In order to normalize the Dst index to the dipole equator, the expression is divided by the average of the cosines of magnetic dipole latitudes of the i -th observatory, λ_i . Thus, Dst can be expressed as:

$$Dst(t) = \frac{\sum_{i=1}^n D_i(t)}{\sum_{i=1}^n \cos(\lambda_i)} \quad (1.1)$$

For further technical details and in-depth derivation, the reader is referred to (Sugiura and Kamei, 1991). Many previous studies (Russell *et al.* (1974), Biktash (2012)) have been

done in attempt to correlate the *Dst* index, and therefore the ring current strength, with the solar wind parameters such as the bulk plasma speed, strength, and the orientation of the magnetic field that it carries. This project investigates if the correlation could be improved by replacing it with the *Est* Index.

1.3.2 Est and Ist Indices

Geomagnetic field disturbances are commonly thought to be caused solely by a field source outside of the Earth’s body. However, it has recently been found that a sole external field source may not always be the case. The disturbance magnetic field representations have some difficulties due to the *secondary field*, which arises from the external time-varying source field inducing electric currents in the Earth. The strength of this secondary field is approximately one third of the inducing field (Maus and Weidelt, 2004). Thus, the field disturbance observed on the Earth’s surface and measured by *Dst* index is, in fact, the sum of the external source field and its induced counterpart.

A physically correct representation of the symmetric disturbance field is provided by the *internal* and *external* contributions, which are derived by taking a Fourier transform of the *Dst* index as well as implementing a conductivity model (Utada et al, 2003). Therefore, the *Dst* index can be split into the **external dipole source field** (*Est Index*) and the **internal induced dipole field** (*Ist Index*). Thus, at any given time;

$$Dst(t) = Est(t) + Ist(t) \quad (1.2)$$

and the corresponding internal and external dipole fields are represented as

$$\mathbf{B}_{ext}(t) = -Est(t) \left(\sin \theta \hat{\theta} - \cos \theta \hat{r} \right) \quad (1.3)$$

$$\mathbf{B}_{int}(t) = -Ist(t) \left(\sin \theta \hat{\theta} + 2 \cos \theta \hat{r} \right) \quad (1.4)$$

where $\hat{\theta}$ is the a unit vector for the local southward, and \hat{r} is the local outward unit vector. The negative signs in Eqn. 1.3 and 1.4 are present because *Dst* represents the northward field, while $\hat{\theta}$ points southward. In an ideal conductor, the strength of the induced field would have its radial component cancelling out the radial component of the inducing field on its surface. However, in the case of the real Earth, the amplitude and phase lag relation between the induced internal and induced external field depends on the *frequency content* of the external source field (Maus and Weidelt, 2004). For more mathematical details on separating external and internal parts of the symmetric disturbance field, the reader is referred to Maus and Weidelt (2004). B_{ext} and B_{int} will be used later in this project (Section 3.2) for the correction of the AU and AL indices, which will be described next.

1.3.3 Auroral-Electrojet (AE) Indices

The Auroral-Electrojet (AE) indices were introduced as a measure of global disturbances resulting from the auroral electrojet activity, and thus specifying the state of the Earth’s upper atmosphere. It should be noted that the term “AE Indices” includes the *AU* and *AL* indices, which this study will also focus on, whereas the term “AE Index” is just the AE index itself. AE indices data are maintained at the World Data Center for

Geomagnetism in Kyoto, Japan, and has been exploited both qualitatively and quantitatively as a correlation index in the fields of radio propagation, substorm morphology, the behaviour of communication satellites, and the magnetosphere-IMF coupling.

The AE indices are derived from the geomagnetic variations in the horizontal component (H) of the main field observed at the selected observatories along the auroral zone in the northern hemisphere. The variations are measured from a baseline, which is determined for each individual observatory, and then go through a normalization process. For more detail of the normalization process, please refer to Menvielle *et al.* (2011).

In a superimposed plot of one-minute H values from all the AE observatories at each given universal time (UT) the AU index represents the upper envelope of the plotted points, while the AL index represents the lower envelope. The AE value gives the separation of these envelopes:

$$AE = AU - AL \quad (1.5)$$

The AU and AL indices are defined by the largest and the smallest values respectively, selected from the data from all the stations (Menvielle *et al.*, 2011). While the AE index represents a measure of global auroral electrojet activity, the AU index is intended to represent a measure of the strongest current intensity of the eastward auroral electrojet, and the AL index represents a measure of a westward auroral electrojet. Some caution should be taken when using the indices for substorm studies for a number of reasons. One in particular, which is relevant to our study, is that although the indices are generated from ionospheric electrojet current, magnetospheric current such as the equatorial ring current can also affect the indices. Davis and Sugiura (1966) carried out a study and suggested that the negative values of the AU index may have occurred due to the contribution from the ring current. This project will also investigate this speculation to see whether the corrected AU and AL indices (Section 3.2) would correlate with the coupling functions better.

1.4 Coupling Functions

1.4.1 $d\Phi_{MP}/dt$

The velocity, density, or pressure of the solar wind themselves are unlikely to have much predictive power for the magnetospheric indices. However, it was realized by Dungey (1961) that the IMF is continual and continuously varying, thus, suggesting the importance of B_z and the merging between the IMF and the Earth's fields (Section 1.1). There are many different coupling functions representing a variety of magnetospheric activities due to the interaction between the solar wind and the magnetosphere. The one particular coupling function this project focuses on was proposed by Newell *et al.* (2007). This function represents the rate of which magnetic flux is open to the magnetosphere, and is presented as;

$$d\Phi_{MP}/dt = v^{4/3} B_T^{2/3} \sin^{8/3}(\theta_c/2) \quad (1.6)$$

where v is the approximated rate at which IMF field lines approach the magnetopause, B_T is the magnitude of the interplanetary field, and $\sin^{8/3}(\theta_c/2)$ represents the percentage of the merging IMF lines. The IMF clock angle (θ_c), here, is defined by;

$$\theta_c = \arctan(B_y/B_z) \quad (1.7)$$

where B_y and B_z represent the y-component and the z-component of the field in GSM coordinates (Figure 1.2). The reader is referred to Newell *et al.* (2007) for an in-depth derivation.

The study by Newell *et al.* (2007) shows that the AU and AL indices both correlate best with this coupling function. However, the Dst index does not correlate best with just $d\Phi_{MP}/dt$ alone but with an additional factor of $p^{1/2}$, which gives $p^{1/2}d\Phi_{MP}/dt$. The term $p^{1/2}$ was used to correct for the magnetopause currents. This effect was taken into consideration when applying the coupling function to our data set.

The coupling function $d\Phi_{MP}/dt$ greatly depends on the term involved the IMF clock angle, $\sin^{8/3}(\theta_c/2)$. Despite the awkward 8/3 exponent, which was a result based on ‘trial and error search’ through all the possible values, the dependence of $d\Phi_{MP}/dt$ on the clock angle is rather straightforward.

The physical meaning of $\sin^{8/3}(\theta_c/2)$ is the fraction of the field lines impacting the magnetosphere which merge and is a function of the magnetic shear. Newell *et al.* (2007) also carried out a statistical comparison for the relationship between the IMF clock angle and various coupling functions. All the coupling functions, including the $d\Phi_{MP}/dt$, predicted that the greatest amount of field lines merging is due to the southward IMF (Figure 1.3), and the functions go to zero for a northward IMF. This can also be used to support the statement that says the southward IMF model is preferable over the northward IMF, as previously discussed in section 1.1.

All in all, the coupling function $d\Phi_{MP}/dt$ is chosen by this study because, according to a study done by Newell *et al.* (2007), it gives a reasonable prediction of a variety of magnetospheric phenomena, and is claimed to have the best correlation with every indices ($r > 0.80$ in some cases). This study will, therefore, investigate the robustness of this function and how it performs when the corrected indices are used. The second best coupling function E_{SR} , as claimed by Newell *et al.* (2007), will also be used to test the robustness

1.4.2 E_{SR}

This coupling function was proposed by Scurry and Russell (1991) in the study of the transferring of energy into the magnetosphere. The function is represented as;

$$E_{SR} = vB_T \sin^4(\theta_c/2)p^{1/2} \quad (1.8)$$

It examines the relation between the reconnection mechanism and the orientation of the IMF as well as the upstream solar wind parameter. The solar wind parameters required for this coupling function are the same as those used in $d\Phi_{MP}/dt$. In the study done by Newell *et al.* (2007), E_{SR} gave the second best correlation coefficient for Dst (1995-2002) and the third best for Dst (1984-1994). However, E_{SR} does not correlate well with the AU and AL indices at all.

In this study, the coupling function E_{SR} is used to test whether the function $d\Phi_{MP}/dt$ really is the best coupling function, as claimed by Newell *et al.* (2007). This investigation will be also be comparing the correlation coefficient between E_{SR} and OMNI 2 data with the results from Newell *et al.* (2007). Only the basic description of this coupling function is given here because it is not the main focus of this study, however, for more details on the E_{SR} function the reader is referred to Scurry and Russell (1991).

Chapter 2

Data

2.1 Solar Wind Parameters, Dst, AU, and AL Data

The solar wind dataset used in this study is called “OMNI 2” data and is obtained from NASA’s OMNIWeb Data and Service site (<http://omniweb.gsfc.nasa.gov>) provided by the Space Physics Data Facility (SPDF) at NASA/Goddard Space Flight Center. The OMNI 2 data are a multi-source dataset which contains continuous time-series of solar wind data contributed by several different spacecraft. Most of these spacecrafts are in geocentric orbit, while some are orbiting the Sun at the Lagrange point. Within the data set, factors such as time-delays between satellites at different locations have already been dealt with before the data were published on the site.

The particular data set that was used in this study is low-resolution with hourly averages from the year 1963 to 2013. The data set has 55 columns, and the information selected for the purpose of this investigation is:

Parameters	Column
Year	1
Decimal Day	2
Decimal Hour	3
IMF Field Strength, B_T (nT)	10
B_y (nT)	16
B_z (nT)	17
Solar wind speed, v (km/s)	25
Solar wind pressure, p (nPa)	29
Dst Index (nT)	41
AL Index (nT)	53
AU Index (nT)	54

Table 2.1: The column number in which the solar wind parameters are extracted from OMNI 2 data. Note: B_x and B_y are field components in GSM coordinates (Figure 1.2)

This selected information is then used to calculate the coupling functions, $d\Phi_{MP}/dt$ (Equation 1.6) and E_{SR} (Equation 1.8). Some modifications have to be made on some sets of data, such as Dst, AU, and AL indices, as well as the solar wind pressure. These modifications will be discussed later in Section 2.3 and 2.4.

2.2 Est and Ist Data

In this study, the Est index data would be used to replace Dst in the hope of improving the correlation with the coupling function, and therefore improving the prediction of the magnetic disturbances from solar wind parameters.

The Est and Ist data were extracted from a set of data obtained from the File Transfer Protocol (FTP) Directory of the National Geophysical Data Center (NGDC) which is a part of the USA's National Oceanic and Atmospheric Administration (NOAA). The dataset contains the data from 1957 - 2013, and has 4 columns with the first column being the Julian date, the second is Dst index values, the third column is the Est index value, and the fourth being the Ist index value.

2.3 Correction on Dst Index data

As mentioned previously in Section 1.4.1, there is an additional factor of $p^{1/2}$ when correlating Dst index with the coupling function.

On the sunward side, the magnetopause is *compressed* and drawn out tailward until the magnetic pressure from the magnetosphere balances the dynamic pressure of the incoming solar wind (Russell *et al.*, 1974). This leads to the distortion of the magnetosphere, which is associated with the magnetopause currents. Therefore, the Dst index will need to correct for this effect by removing an amount proportional to a value of solar wind pressure (p).

Due to this correction, the correlation of the coupling function is not with “Dst” but actually with $\mathbf{Dst} - 18.9\mathbf{p}^{1/2}$ (Newell *et al.*, 2007).

2.4 Time Integration of Solar Wind for Each Index

In their study, Newell *et al.* (2007) worked with 10 indices, each representing a different magnetospheric phenomena. Therefore, it was expected that the response of these indices to the solar wind would show a variety of timescales and decay with differing hysteresis. So, it became apparent that time-integration is needed, and that the integration time for a given index is somewhat invariant of the coupling function used (Newell *et al.*, 2007).

Although this investigation is focusing on the coupling function $d\Phi_{MP}/dt$, time integration process can be applied to other coupling functions as well, such as E_{SR} . The coupling function, $d\Phi_{MP}/dt$, was calculated first at an hourly time resolution before the time integration was done.

The process involved separately altering the number of hours of IMF data integrated over, and the weight, w ($0 < w < 1$), used in the integration.

The *weighted mean* is applied to each hour. The weight takes the form of w^n (n = the hour before the present), and the number of hour is altered for different indices (Table 2.2). Since the present hour influences the indices prediction the most, it is therefore the most weighted ($w^0 = 1$), and the hours following that incrementally increase in n (weighting factor decreases) (Newell *et al.*, 2007). The solar wind integration time optimizing each index is shown in Table 2.2.

Index	No. of Hours of IMF	w
Dst	72	0.95
AU	3	0.77
AL	3	0.69

Table 2.2: The Solar Wind Integration Time for Each Index

The Dst index, a measure of the ring current, requires the longest solar wind time integration (72 hours) because large ring currents can occur continuously for several days even after dayside merging ceases. This is due to substorm, from the crossing neutral point at the tail of the magnetosphere (Figure 1.3 (a)), injecting energy into the ring current system (refer to section 1.1).

Only 3 hours of solar wind time integration is needed for the AU and AL indices because that is roughly the amount of time it takes for the auroral electrojets (refer to section 1.3.3) to return to their normal state after dayside merging ceases (Newell *et al.*, 2007).

Chapter 3

Method

The data was processed mainly using MATLAB[®]. Some of the codes used throughout this study were written by the student, while some help and Function codes were provided by the project supervisor, Dr. Brian Hamilton (BGS). Several different codes were written since this study investigates the correlation between the coupling functions, $d\Phi_{MP}/dt$ and E_{SR} , with different geomagnetic indices. This section will describe the processing steps implemented within different codes.

3.1 Main Code

The main code was written by the student to investigate the correlation between the coupling functions and the indices. The steps implemented in the main code (Appendix A) are used to process different indices and coupling functions, therefore the description of each step is kept as general as possible. The implemented steps are:

1. The code started by importing the OMNI 2 data and the solar wind parameters needed were extracted. In the case of Est, the code imports the data from the NOAA data set and the Est index values were then extracted.
2. The clock angle (θ_c) was then calculated using the Equation 1.7. This is the same for all coupling functions used in this study.
3. The coupling function was calculated using the Equation 1.6 for $d\Phi_{MP}/dt$, and Equation 1.8 for E_{SR} .
4. The OMNI 2 data contains a lot of missing data. The missing data appeared to have very large value such as 99.99 nPa for the solar wind pressure (p), 9999 km/s for the solar wind speed (v), and 999.9 nT for the IMF field strength (B_T), the y-component (B_y), and the z-component (B_z). Any row of data with at least one missing datum would be marked as a faulty row.
5. The proceeding 72 or 3 hours (Table 2.2) were checked. The selected blocks should have 72 rows for Dst and Est indices, and 3 rows for AU and AL indices. By doing this, data points with the required number of good and complete hours would be selected. Any data points that did not have the required number of good hours preceding them were flagged as bad.
6. The time integration of solar wind for the index was then implemented. The Dst and Est indices are best predicted by integrating over 72 hours of the solar wind data and 3 hours for the AU and AL indices, with the hour n -hours previous to the present receiving the relative weight of w^n . For example, the current hour influences

the index prediction the most so it is weighted with w^0 which is 1, the previous hour is weighted w^1 , two hours earlier is weighted w^2 , and etc. The expression was also, then, divided by the sum of all the weighting factor terms i.e. $\sum_0^n w^n$ (weighted mean).

As mentioned in Section 2.3, the Dst index were corrected with a term proportional to $p^{1/2}$. Therefore, the solar wind pressure term, $p^{1/2}$ was also integrated over previous 72 hours.

7. The data selection process in step 5 picked out blocks of either 72 or 3 complete data, this meant leaving an empty array of 71 or 2 rows at the beginning. The resulting time-series data would, therefore, need to be padded in order to build the data array up to the correct dimension required for future process. Thus, both the coupling function and solar wind pressure data were padded.
8. The faulty rows marked earlier in step 4 could now be removed from the data set.
9. The correlation coefficient between the coupling function and the index was calculated using the built-in ‘corrcoef’ function in MATLAB®.
10. Lastly, the results were plotted.

These ten steps are the core of all the codes used in this study. Equations and parameters were altered in order to investigate other different correlations.

In order to investigate this correlation between Est and $p^{1/2}d\Phi_{MP}/dt$, Dst index values (from OMNI 2 data) used in the processes described in Section 3.1 were simply replaced by Est index values (NOAA). The rest of the code remained unchanged.

For the investigation of the correlation between $d\Phi_{MP}/dt$ and AU and AL, again, the Dst values were simply replaced by either AU or AL values from the same data set (OMNI 2). Due to both AU and AL having different solar wind time integration from that of Dst (refer to Table 2.2), this would need to be adjusted accordingly.

Lastly, in order to investigate the correlation between Dst index and its corresponding second best coupling function (Newell *et al.*, 2007), E_{SR} , the $d\Phi_{MP}/dt$ was replaced by E_{SR} (Equation 1.8). The solar wind parameters and the time-integration stayed unaltered (from Section 3.1).

3.2 Correction of the AU and AL indices

There are two spherical harmonic expansions in the general solution to Laplace’s equation in spherical coordinates, the field associated with the external and internal source. The spherical harmonic expansions take the form;

$$V = a \sum_{l=1}^{\infty} \sum_{m=0}^l \left[\left(\frac{a}{r} \right)^{l+1} (g_l^m \cos m\phi + h_l^m \sin m\phi) P_l^m(\theta) + \left(\frac{r}{a} \right)^l (q_l^m \cos m\phi + s_l^m \sin m\phi) P_l^m(\theta) \right] \quad (3.1)$$

where V is the potential function and a is the Earth’s radius. The integers, l and m , are called *degree* and *order*, respectively. The degree, l , has a value of 1 or greater, while m is always less than or equal to l . The weighting coefficients g_l^m , h_l^m , q_l^m , and s_l^m are known as Gauss Coefficient (Campbell, 2003) and are measured in nT. The term $P_l^m(\theta)$ is the Legendre polynomials, and ϕ and θ are longitude and colatitude, respectively. The part of Equation 3.1 with the coefficients g_l^m and h_l^m represents the internal sources, while the part with q_l^m and s_l^m represents external sources.

As previously mentioned in Section 1.3.2 that the Est and Ist indices are, respectively, the external and internal parts of the Dst index. The Est and Ist values are the external and internal spherical harmonic coefficients of a simple degree 1 representation of the magnetospheric field and the induced field. Since the Est and Ist indices are strictly of degree 1 ($l=1, m=0$) representation, it is assumed that all other coefficients are zero; all $l > 1$ are zero and all non-zero orders ($m > 0$) are also zero. This transforms Equation 3.1 into;

$$V = a \cos(\theta) \left[\left(\frac{a}{r} \right)^2 g_1^0 + \left(\frac{r}{a} \right) q_1^0 \right] \quad (3.2)$$

with the standard definition of $P_1^0 = \cos(\theta)$. Equation 3.2 indicates that this is a zonal model composing of the coefficients g_1^0 and q_1^0 , which are Ist and Est, respectively. It should also be noted that harmonics representation is in geomagnetic coordinates because the Dst index is a component of the field along the direction of Earth's dipole field, which means the Ist and Est (which arise from Dst) are effectively and respectively g_1^0 and q_1^0 coefficients in geomagnetic coordinates.

The Est and Ist indices are actually a simple spherical harmonic representation of the ring current (Section 1.2.2) fields, so they can be used to generate the magnetic field produced by the ring current and the currents it induces in the Earth. Equation 3.2 shows that the field potential is colatitude dependent and not longitude dependent. This makes it easier to calculate a single correction to AU and AL i.e. the longitude that the observatories used to produce the indices can be neglected.

All of the above is used to build up the correction of the AU and AL indices. The time-series of Est and Ist can be used to generate a time-series of the corresponding magnetic field values. These field values can then be subtracted from the AU and AL indices, and the resulting effect are investigated in this study.

The MATLAB® code for this correction was provided by Dr. Hamilton. The code took in the Est and Ist values and calculated magnetic field values at fixed location i.e. the magnetic latitudes of the observatories that were used to produce AU and AL indices (Menvielle *et al.*, 2011). The observatories have a narrow range of latitudes (refer to Table 8.2 in (Menvielle *et al.*, 2011)), so as a first approximation the latitudes were averaged and estimated to be about 65° .

The code produced three components of the field, B_X , B_Y , and B_Z . The AU and AL are calculated from the field horizontal to the ground, so in order to subtract from the indices, B_X and B_Y components are used. However, the Est and Ist were only spherical harmonic of order 0 ($m = 0$) terms, therefore the output B_Y values from the code were all zero. So, the only horizontal field component which would be subtracted from AU and AL hourly values was the B_X component. This is an assumption since B_Y would not be completely zero in reality. It can be seen in the circled area in Figure 4.3 that the negative AU values coincide with the most negative B_X values, indicating that B_X is the most dominant component of the field from the ring current. From the geometry of the ring current (around the Earth's geomagnetic equator), the field is expected to be perpendicular to this ring i.e. in the geomagnetic B_X -direction. Therefore, neglecting B_Y component would not cause significant error.

Chapter 4

Results

This section presents the results *and* the discussion of the results. All of the numerical results are shown in Table 4.1. This study presents results from a recent 2000-2013 data, which covers a complete solar cycle and therefore, provide a better overall analysis of the correlation. For each correlation, we first start off by trying to reproduce the results by Newell *et al.* (2007) before moving on to look at the results from using corrected indices.

4.1 Correlation between Dst and $p^{1/2}d\Phi_{MP}/dt$

The study by Newell *et al.* (2007) shows only the correlation of the data from 1984-1994 and 1995-2002. The first obvious difference between the two studies (Table 4.1) is that the valid number of events in this study is only 112 events from a total of 96432 possible events, while the previous study has 21418 events. The lower number of valid events in this study, therefore, gives a higher correlation coefficient ($r = 0.929$) than that of the previous study ($r = 0.857$).

This large difference in the number of valid events could be due to a number of reasons. Dr. Hamilton also wrote a code for the same purpose and ran it alongside my code. This acted as a confirmation, and if our results were different then we knew there were some errors. It could, therefore, be said with confidence that the code really did produced 112 valid events during the 1984-1994 interval with the time integration method suggested by Newell *et al.* (2007) (Section 2.4).

One possible cause of such a large difference could be that, in the earlier years, the data were just not as continuous as the more current set. We were very strict with the selection of the valid events, so the solar wind time integration implemented in the MATLAB[®] code only picked out blocks of completely continuous 72 time-series data and rejected the whole data block even if there was a single missing datum. It is also possible that Newell *et al.* (2007) were more relaxed about their selection criteria, since their number of valid events is quite high. Perhaps, they threw away bad values but, nevertheless, carried out the calculation, which could explain the slight differences in the correlation coefficient values. Also, the number of valid events in the later years was much closer to the previous study, and there were fewer rejected data, this supports the idea that strict data selection was the problem. In either case, it should be noted here

Index	Year	No. of Event (n)		Magnitude of r		Gradient	Y-intercept
		Newell 's	This study	Newell's	This study		
Coupling Function = $p^{1/2}d\Phi_{MP}/dt$							
Dst	1984 - 1994	21418	112	0.857	0.929	-0.0037	-23.50
Dst	1995 - 2002	59666	63589	0.866	0.846	-0.0041	-14.76
Dst	2000 - 2013		115139		0.856	-0.0037	-14.14
Est	1984 - 1994		112		0.939	-0.0031	-28.06
Est	1995 - 2002		63598		0.847	-0.0030	-20.83
Est	2000 - 2013		115139		0.864	-0.0030	-18.92
Coupling Function = $d\Phi_{MP}/dt$							
AU	1983 - 1987	10352	11668	0.765	0.729	-0.0134	-21.34
AU(c)	1983 - 1987		11668		0.752	-0.0146	-21.67
AU	2000 - 2013		118434		0.683	-0.0123	-15.83
AU(c)	2000 - 2013		118434		0.717	-0.0137	-15.02
AL	1983 - 1987	10352	11668	0.528	0.762	-0.0326	4.21
AL(c)	1983 - 1987		11668		0.754	-0.0314	4.63
AL	2000 - 2013		118434		0.768	-0.0277	7.99
AL(c)	2000 - 2013		118434		0.755	-0.0263	7.25
Coupling Function = E_{SR}							
Dst	1984 - 1994	21418	112	0.838	0.934	-0.0159	-27.47
Dst	1995 - 2002	59666	63589	0.860	0.852	-0.0159	-21.21
Dst	2000 - 2013		115139		0.853	-0.0145	-19.80
Est	1984 - 1994		112		0.942	-0.0126	-26.05
Est	1995 - 2002		63598		0.849	-0.0117	-23.71
Est	2000 - 2013		115139		0.851	-0.0117	-23.71

Table 4.1: The analytical results from the investigation of the relationship between different coupling functions and geomagnetic indices. ‘ r ’ is the correlation coefficient. The ‘(c)’ after AU and AL indices signifies the indices after correction (Section 3.2). The analytical results of the more recent data, which are not present in the study done by Newell *et al.* (2007), are also shown here. The Gradient and Y-intercept columns are the corresponding values from the line of best-fit as shown as red lines in Figure 4.1.

that due to the small amount of valid events, any interpretations based on the 1984-1994 data are avoided in this study.

In the case of the data from 1995-2002, the previous study has 59666 valid events while this study has 63589 valid events. The corresponding correlation coefficients for the previous study and this study are 0.866 and 0.846, respectively. Again, the data set with more valid events appears to have a lower correlation coefficient value.

As for more up-to-date data set, from 2000-2013, the previous study does not have the analytical results presented. However, this study showed that from the data set of 2000-2013, 115139 valid events were used, which produced $r = 0.856$. The scatter plot of the correlation during 2000-2013 interval is shown in Figure 4.1(bottom right). Note that the correlation with Dst should be negative according to how the $p^{1/2}d\Phi_{MP}/dt$ is defined. However, the negative signs are neglected because we are interested in the magnitude of the correlation coefficients and also to better compare the results.

4.2 Correlation between Est and $p^{1/2}d\Phi_{MP}/dt$

The previous section attempted to reproduce the results from the previous study by Newell *et al.* (2007). We will now look at the effect on the correlation when the Dst index is replaced by the Est index.

After the Dst index is replaced with the Est index, the resulting correlation coefficients appear to have improved (Table 4.1). The minimum improvement is in the 1995-2002 data set, with r going from 0.846 (Dst) to 0.847 (Est). The maximum improvement is in the 1984-1994 data set, with $r = 0.929$ (Dst) increased to $r = 0.939$ (Est), but again, due to the a small number of valid events, the results from 1984-1994 data are not robust enough for any conclusions to be drawn from.

For the 2000-2013 data, the analytical results from this study shows that the correlation coefficient increased from $r = 0.856$ (Dst) to $r = 0.864$ (Est). The scatter plot of the data from the 2000-2013 interval is shown in Figure 4.1 (top).

The changes in the correlation are very small, but this was to be expected. The scatter plot, Figure 4.2, shows the correlation between the Est and Dst indices. The high correlation coefficient ($r = 0.99$) indicates that the two indices are very similar. This, therefore, explains the very small change in the correlation with both coupling functions when the Dst index is replaced with the Est index.

4.3 Correlation between Dst and E_{SR}

So far, it had been speculated that the correlation with the top coupling function $p^{1/2}d\Phi_{MP}/dt$, as claimed by Newell *et al.* (2007), could be improved by replacing the Dst index with the Est index. We then moved on to look at the second best coupling function E_{SR} to see whether the claim made by Newell *et al.* (2007) about the top function would still hold after the corrections have been made and using new data (2000-2013). If the $p^{1/2}d\Phi_{MP}/dt$ function is knocked off its top spot, this would mean that function is not as robust and may not correlate best with every indices as claimed by Newell *et al.* (2007).

For the 1984-1994 data, the previous study (21418 valid events) produced $r = 0.838$, while our study (112 valid events) produced $r = 0.934$. For the 1995-2002 data, Newell *et al.* (2007)'s study (59666 valid events) produced $r = 0.860$, while our study (63589 valid events) produced $r = 0.852$.

The 2000-2013 data has 115139 valid events and produced $r = 0.853$, (corresponding scatter plot is shown in Figure 4.1, bottom left).

Now, let us compare the correlation between the Dst index and the $p^{1/2}d\Phi_{MP}/dt$ and the E_{SR} functions from this study. For the data of 1984-1994 and 1995-2002 the coupling function E_{SR} appears to correlate with Dst better than $p^{1/2}d\Phi_{MP}/dt$ by 0.006 on average. However, $p^{1/2}d\Phi_{MP}/dt$ function correlated with Dst better than E_{SR} function for the 2000-2013 data by about 0.003 (Table 4.1).

Although the differences between the correlations from the two coupling functions are very small, it is enough to cast doubt upon the claim by Newell *et al.* (2007) because the results show that the $p^{1/2}d\Phi_{MP}/dt$ function is not constantly better than the E_{SR} function. We then moved on to look at the correlation between $d\Phi_{MP}/dt$ function and AU, and AL indices.

4.4 Correlation between AU and $d\Phi_{MP}/dt$

Similar to previous sections, we first attempted to reproduce the AU results from the previous study, and then moved on to look at the effect of replacing the AU index values with the corrected values on the correlation.

As for the correlation between the coupling function, $d\Phi_{MP}/dt$, and the AU index, the 1983-1987 OMNI 2 data set produced $r = 0.729$ (1668 valid events), while $r = 0.765$ (10352 valid events) for the previous study. This is not a bad attempt at reproducing the results from the study by Newell *et al.* (2007).

The more current data from 2000 to 2013 is also presented in Table 4.1; the number of valid event is 118434, which is roughly 10 times the data from earlier years. This high number of valid events only produced $r = 0.683$. Again, these comparisons between the two studies only act as a guideline and are not so significant for the purpose of this study since we are focusing on ways of improving the correlation of the data from the same source.

Let us now compare the results, from OMNI 2 data, between the original AU index and the AU index that has been corrected using the process described in Section 3.2. For the 1983-1987 data, the correlation coefficient has increase from 0.729 to 0.752, and from 0.683 to 0.717 for the 2000-2013 data. It is clear that the correction did improve the correlation between the AU index and the coupling function, but why?

The correction, involved subtracting the northward (X) component of the field, calculated from the Est and Ist indices values at the geomagnetic latitude of $65^\circ N$, from the AU and AL indices. This process was done based on the idea that Est and Ist indices are a simple harmonic representation of the ring current field.

The top panel in Figure 4.3 shows a time-series plot of the AL index values (red), the AU index values (green), and the x-component values of the ring current (blue). The circled area on the plot contains some negative AU values, which seem to correspond to the period where the x-component of the ring current is most negative. This indicates a possible correlation between the two data sets. From a numerical analysis, it was found that, on average, approximately 8% of the AU index and 6% of the AL index were due to the contribution from B_X .

The bottom panel in Figure 4.3 shows a time-series plot, from the 1995-2002 data, of the AU index values (blue) and the *corrected* AU index (AU-BX) values (red). It can be seen that most of the negative AU index values before the correction (blue) are now shifted to the less negative values. This is somewhat consistent with the theory that the negative and uncorrected AU index values are, in part, caused by the ring current. Thus, applying the correction to the AU index helped improve its correlation with the coupling function.

4.5 Correlation between AL and $d\Phi_{MP}/dt$

Similar to the previous section, for the 1983-1987 data, the study by Newell *et al.* (2007) has 10352 valid events while this study has 11668 valid events. However, the correlation coefficient between the coupling function and the AL index from this study ($r = 0.762$) is higher than that of the previous study ($r = 0.528$). There is no real explanation for this except that Newell *et al.* (2007) may have made a mistake. For the

2000-2013 interval, OMNI 2 data produced $r = 0.768$.

Let us now compare the AL index with corrected AL index (Section 3.2) from OMNI 2 data set. For the 1983-1987 data, the correlation coefficient decreased from 0.762 to 0.754, and also decreased from 0.768 to 0.755 for the 2000-2013 data. The correction on the AL values appears to have a contrasting effect, comparing to the corrected AU values.

There is no certain explanation why our AL results are better than Newell *et al.* (2007)'s, and worsened after the correction is applied. However, there is a possible explanation for this. Newell *et al.* (2007) proposed the solar wind integration time (Table 2.2) optimizing each index. The optimal valued of the weighting factor (w) and the number of hours of IMF (n) were chosen after a 'trial and error search' which was done based on the results of AU and AL indices. This means that the w and n used are not optimal for the corrected version of the AU and AL indices. Newell *et al.* (2007) has fine-tuned the values used in the calculation of the coupling function $d\Phi_{MP}/dt$ assuming uncorrected AU and AL indices, so the correlation with the corrected AU and AL might have been improved if Newell *et al.* (2007)'s analysis was to be repeated using the corrected AU and AL indices.

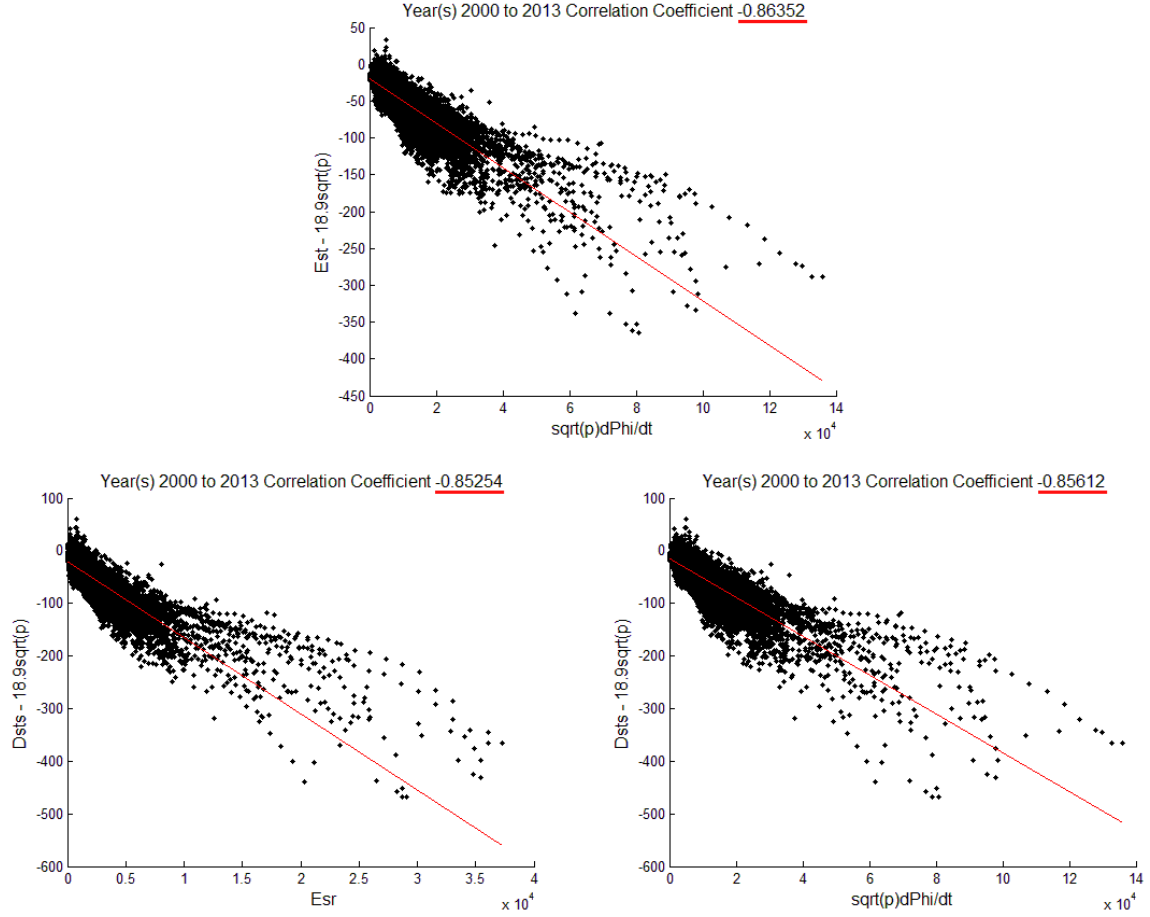


Figure 4.1: Examples of some output scatter plots. Scatter plots of Est versus $p^{1/2}d\Phi_{MP}/dt$ (top), Dst versus E_{SR} (bottom left), and Dst versus $p^{1/2}d\Phi_{MP}/dt$ (bottom right).

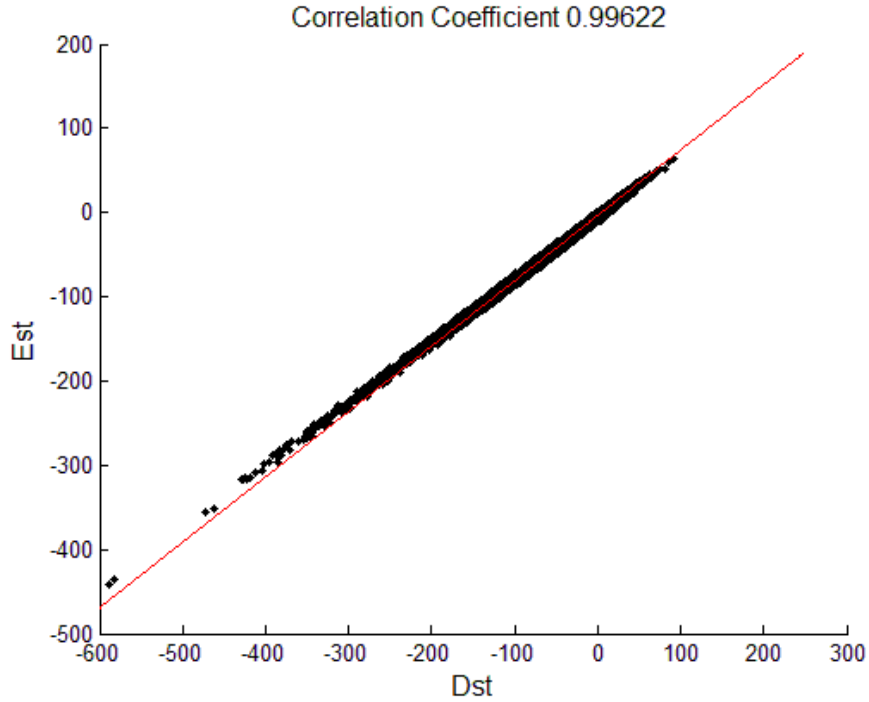


Figure 4.2: A scatter plot illustrating that the Dst index correlates very well with the Est index ($r = 0.99$). Dst and Est are so well correlated that comparing $p^{1/2}d\Phi/dt$ with Est will never be that much different from comparing it with Dst.

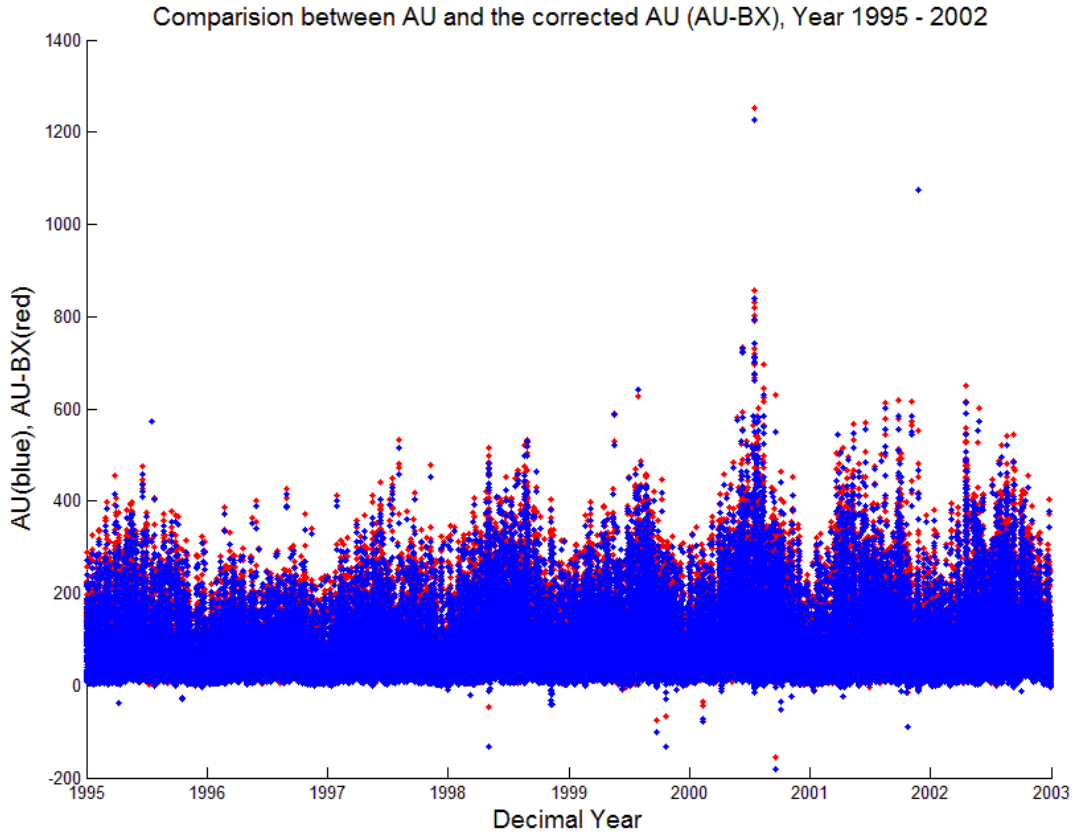
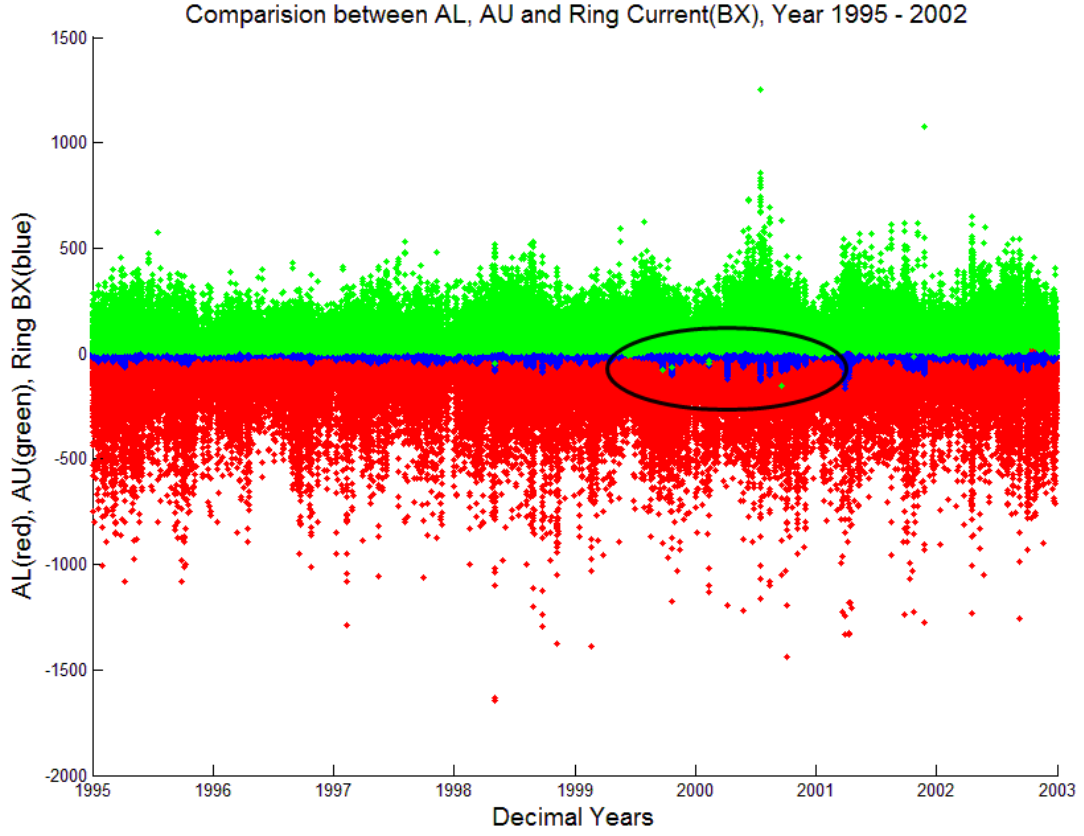


Figure 4.3: Top panel: A time-series scatter plot comparing the AL values (red), AU values (green), and x-component of the Ring Current (blue). The circled area shows that there is a possible correlation between the negative AU values and the x-component of ring current. Bottom panel: A times-series scatter plot comparing the AU index values(blue) with the corrected AU index values (red).

Chapter 5

Discussion and Conclusion

Newell *et al.* (2007) was only interested in the correlation coefficients from the plots, such as those shown in Figure 4.1, but the (red) best-fit lines on those correlation plots could be used to predict the Dst/Est/AU/AL from solar wind parameters. The ‘Gradient’ and the ‘Y-intercept’ are presented in Table 4.1. The indices prediction could be done by using the gradient to scale down the x-axis (coupling functions) and apply an offset (y-intercept) to determine the index value on the y-axis.

Regardless of the exact values, the difference between $p^{1/2}d\Phi_{MP}/dt$ and E_{SR} for the 1995-2002 Dst data in the study by Newell *et al.* (2007) is very small. It is possible that what Newell *et al.* (2007) is arguing is that their coupling function, $d\Phi_{MP}/dt$, is the best coupling function across all the indices they studied. However, the analysis from this investigation does not seem to support this claim since the correlations between $d\Phi_{MP}/dt$ and Dst, Est, AU, and AL appear to be very inconsistent.

This project proposed and investigated the hypothesis that magnetic disturbances can be better predicted by replacing the Dst index with the Est index. The hypothesis is true for the coupling function $p^{1/2}d\Phi_{MP}/dt$, however, this is not the case for the coupling function E_{SR} . Once the Dst index is replaced with the Est index, the correlation coefficients produced from the 1995-2002 and 2000-2013 data are lower than before the replacement. The first conjecture was the possibility of the E_{SR} function having different solar wind integration time (Table 2.2) from the $p^{1/2}d\Phi_{MP}/dt$ function. However, the time integration was done after the calculation of the coupling functions, which means that the optimization of the solar wind parameters and the optimization of the coupling functions are two independent operations. Therefore, the integration time should not be specific to a particular coupling function.

However, after all the corrections and using more recent data from 2000-2013, the coupling function $p^{1/2}d\Phi_{MP}/dt$ produced higher correlation coefficients than the coupling function E_{SR} , overall. The difference is very small which is due to the fact that Est and Dst are extremely well correlated ($r = 0.99$), even if they can differ slightly in amplitude. So, when correlating with Solar Wind parameters, any improvement will be small compared to the failings of the coupling function ($r \approx 0.85$). It is, nevertheless, an improvement on the prediction model. Therefore, it can be used to support our hypothesis that better prediction of the magnetic disturbances could be done by replacing the Dst index with the Est index, which correlates better with solar wind parameters, through the coupling function, $p^{1/2}d\Phi_{MP}/dt$.

Regarding the results from the correlation between the AU and AL indices with the coupling function $d\Phi_{MP}/dt$ (Table 4.1), it was found that roughly 8% of the AU index and 6% of the AL index were the contribution from the x-component of the ring current. However, it can be said that the correction improved the correlation for the AU index, but was not so for the AL index. Overall, the improvement in the AU index is quantitatively better than the worsening of the AL index.

All of the factors that went into calculating the coupling functions proposed by Newell *et al.* (2007) were fine-tuned, some through trial and error, to be optimal for the indices values they used. However, in this study, many magnetic indices were corrected, therefore, the factors which had been optimized for uncorrected values may not have worked as well.

If more time was available I would like to extend this investigation across more indices, in order to really test Newell *et al.* (2007)'s claims. Also, I would like to focus on just one complete solar cycle and investigate the seasonal effect on the AU index. From the bottom panel of Figure 4.3, it can be seen that the AU values are at maximum around the middle of the year while at minimum around the end and the beginning of each year. It would be interesting to see whether these trends occur recursively and can be predicted. If the trends are there, then how do they affect the prediction model. For example, the ability to predict might work better at some points in the solar cycle than other.

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Appendix A

Main MATLAB[®] code

```
1 %% A code loading a DAT file and apply the dPhi/dt fuction to
   the data
2 clear all;
3 clc
4 %% Read in required data data from OMNI hourly set:
5 FID=fopen('omni2-all-years.dat','r');
6 FormatStr = ['%d %d %d %s %s %s %s %s %s %f ' ...
7             '%s %s %s %s %s %f %f %s %s %s ' ...
8             '%s %s %s %s %f %s %s %s %f %s ' ...
9             '%s %s %s %s %s %s %s %s %s %s ' ...
10            '%f %s %s %s %s %s %s %s %s %s ' ...
11            '%s %s %s %s %s '];
12
13 Data=textscan(FID,FormatStr);
14 fclose(FID);
15 % Extract data to vs_Rawvariables:
16 Years_All = double(Data{1});
17 Days_All = double(Data{2});
18 Hours_All = double(Data{3});
19 BTs_All = Data{4};
20 Bys_All = Data{5};
21 Bzs_All = Data{6};
22 vs_All = Data{7};
23 ps_All = Data{8};
24 Dsts_All = Data{9};
25 % Clear the Data vs_Rawvariable to free memory:
26 clearvars Data;
27 % Generate mask to select desired date range. As it's written
   below,it'll only work for wholeyears, not fractions of a year
   :
28 StartYear = 1984;
29 EndYear = 1994;
30 YearRangeMask = Years_All >= StartYear & ...
31                Years_All <= EndYear;
32
```

```

33 % Extract Years_Raw range using the mask:
34 Years_Raw = Years_All(YearRangeMask);
35 Days_Raw = Days_All(YearRangeMask);
36 Hours_Raw = Hours_All(YearRangeMask);
37 BTs_Raw = BTs_All(YearRangeMask);
38 Bys_Raw = Bys_All(YearRangeMask);
39 Bzs_Raw = Bzs_All(YearRangeMask);
40 vs_Raw = vs_All(YearRangeMask);
41 ps_Raw = ps_All(YearRangeMask);
42 Dsts_Raw = Dsts_All(YearRangeMask);
43
44 % Clear unneeded vs_Raw variables to free memory:
45 clearvars Years_All Days_All Hours_All BTs_All Bys_All ...
46           Bzs_All ps_All vs_All Dsts_All...
47
48 theta_c = atan2(abs(Bys_Raw),Bzs_Raw); %the clock angle(radians)
49
50 %% Convs_Rawert to Julian day
51 JD = YearDayOfYearHourToMJD2000o0(Years_Raw, Days_Raw, Hours_Raw
   );
52
53 %% Calculate Coupling Function
54 dPhi_dt = vs_Raw.^(4/3) .* BTs_Raw.^(2/3) .* (sin(theta_c./2))
   .^(8/3);
55 dPhi_mod_BeforeIntegration = sqrt(ps_Raw).*dPhi_dt;
56
57 %% Masking — marking the fault elements
58 MaskForGoodData = BTs_Raw<999.9 & ...
59                  Bys_Raw<999.9 & ...
60                  Bzs_Raw<999.9 & ...
61                  vs_Raw<9999 & ...
62                  ps_Raw<99.99;
63
64 ConvolvedMasks=conv(double(MaskForGoodData),ones(72,1),'valid');
65 MaskForGoodData72HrsBeforePadding = ConvolvedMasks > 71.5;
66 MaskForGoodData72Hrs = [false(71,1);
   MaskForGoodData72HrsBeforePadding];
67
68 %% 72-hour time integration
69 n = 0:71;
70 weight_function = (0.95.^n);
71
72 dPhi_dt_AfterIntegrationBeforePadding = conv(
   dPhi_mod_BeforeIntegration, weight_function,'valid')/sum(
   weight_function);
73
74 dPhi_dt_AfterIntegration = [false(71,1);
   dPhi_dt_AfterIntegrationBeforePadding];

```

```

75
76 ps_AfterIntegrationBeforePadding = conv(sqrt(ps_Raw),
      weight_function, 'valid')/sum(weight_function);
77
78 ps_AfterIntegration = [false(71,1);
      ps_AfterIntegrationBeforePadding];
79
80 Dsts_Raw_mod = Dsts_Raw - 18.9*(ps_AfterIntegration);
81
82 %% Pick out the FALSE values
83 dPhi_dt_AfterIntegration(MaskForGoodData72Hrs==0,:)=[];
84 Dsts_Raw_mod(MaskForGoodData72Hrs==0,:)=[];
85 JD(MaskForGoodData72Hrs==0,:)=[];
86
87 %% Correlation Coefficient
88 coef = corrcoef(dPhi_dt_AfterIntegration, Dsts_Raw_mod);
89
90 %% Straight line fit to Corrected Dst versus the coupling
      function
91 [Gradient, YIntercept] = BestFitLineFromPerpendicularOffsets(
      dPhi_dt_AfterIntegration, Dsts_Raw_mod);
92
93 %% Create plots
94 figure(1)
95 scatter(dPhi_dt_AfterIntegration, Dsts_Raw_mod, 'k. ');
96 xlabel({'sqrt(p)dPhi/dt'}, 'FontSize', 12);
97 ylabel({'Dsts - 18.9sqrt(p)'}, 'FontSize', 12);
98 title(['\fontsize{12} Year(s) ' num2str(Years_Raw(1,1)) ' to '
      num2str(Years_Raw(end,1)) ' Correlation Coefficient ' num2str
      (coef(2,1))])
99 hold on
100 x=0:1:max(dPhi_dt_AfterIntegration);
101 plot(x, Gradient*x+YIntercept, 'r');
102 hold off
103 %%%%%%%%%% END OF CODE %%%%%%%%%%

```