

# Physics 137B Homework 8

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April 10, 2024

## Question 1: Fermi's Golden Rule for Three Body Decays

- (a) **Calculating the total rate:** Using the set up laid out above, we can integrate to get an expression for the total decay.
- (b) **The lifetime of the neutron:** Assuming that  $\mathcal{M} \sim 1$ , calculate the lifetime of the neutron. Comment on the order of magnitude compared to the experimental value, which is  $\tau \sim 887.7\text{s}$ .
- (c) **The lifetime of the muon:** Do the same with the muon, whose experimentally measured lifetime is  $\tau \sim 2.197 \times 10^{-6}\text{s}$ .

### Solution:

- (a) The initial particle (neutron) is at rest, so by conservation of momentum we have

$$0 = \mathbf{P}_f + \mathbf{p}_1 + \mathbf{p}_2$$

where  $P_f, p_1, p_2$  denote the momenta of the proton, electron, and anti-neutrino after decay.

The proton is effectively static and we ignore its momentum, whereas we treat the electron and anti-neutrino relativistically. For simplicity, we assume that  $m_{\bar{\nu}_e} = 0$  since the mass of the anti-neutrino is much smaller than even that of the electron.

Now, to get the Decay Rate, we can integrate over the momenta of the electron and anti-neutrino:

$$W = \frac{2\pi}{\hbar} \left| \frac{G_F \mathcal{M}}{V} \right|^2 \int \frac{V d^3 \mathbf{p}_1}{(2\pi\hbar)^3} \frac{V d^3 \mathbf{p}_2}{(2\pi\hbar)^3} \delta(E_0 - E_1 - E_2)$$

Note that  $d^3 \mathbf{p}_1 \rightarrow 4\pi p_1^2 dp_1$  and similarly  $d^3 \mathbf{p}_2 \rightarrow 4\pi p_2^2 dp_2$ . And since we're assuming  $m_{\bar{\nu}} = 0$ ,

$$\begin{aligned} E_2^2 &= (p_2 c)^2 + (\underbrace{m}_{=0} c^2)^2 \\ \implies E_2^2 &= c^2 p_2^2 \\ \implies 2E_2 dE_2 &= 2c^2 p_2 dp_2 \\ \implies E_2 dE_2 &= c^2 p_2 dp_2 \\ \implies p_2 dp_2 &= \frac{E_2}{c^2} dE_2 \\ \implies p_2^2 dp_2 &= \frac{E_2 p_2}{c^2} dE_2 = \frac{E_2 \cdot \left(\frac{E_2}{c}\right)}{c^2} dE_2 \\ \implies p_2^2 dp_2 &= \frac{(E_2)^2}{c^3} dE_2 \end{aligned}$$

Now,

$$\begin{aligned}
W &= \frac{2\pi}{\hbar} \left| \frac{G_F \mathcal{M}}{V} \right|^2 \iint \frac{V d^3 \mathbf{p}_1}{(2\pi\hbar)^3} \frac{V d^3 \mathbf{p}_2}{(2\pi\hbar)^3} \delta(E_0 - E_1 - E_2) \\
&= \frac{2\pi}{\hbar} \frac{|G_F \mathcal{M}|^2}{(2\pi\hbar)^6} \iint d^3 \mathbf{p}_1 d^3 \mathbf{p}_2 \delta(E_0 - E_1 - E_2) \\
&= \frac{2\pi}{\hbar} \frac{|G_F \mathcal{M}|^2}{(2\pi\hbar)^6} \iint 4\pi p_1^2 dp_1 \cdot 4\pi p_2^2 dp_2 \delta(E_0 - E_1 - E_2) \\
&= \frac{2\pi \cdot (4\pi)^2}{\hbar} \frac{|G_F \mathcal{M}|^2}{(2\pi\hbar)^6} \iint p_1^2 dp_1 \cdot p_2^2 dp_2 \delta(E_0 - E_1 - E_2) \\
&= \frac{|G_F \mathcal{M}|^2}{2\pi^3 \hbar^7} \iint p_1^2 dp_1 \cdot \frac{(E_2)^2}{c^3} dE_2 \delta(E_0 - E_1 - E_2)
\end{aligned}$$

The  $\delta(E_0 - E_1 - E_2)$  factor selects only the contribution in which  $E_0 - E_1 - E_2 = 0 \iff E_0 = E_1 + E_2$  so we get

$$\begin{aligned}
W &= \frac{|G_F \mathcal{M}|^2}{2\pi^3 \hbar^7} \int p_1^2 \frac{(E_2)^2}{c^3} dp_1 \\
&= \frac{|G_F \mathcal{M}|^2}{2\pi^3 \hbar^7 c^3} \int p_1^2 (E_0 - E_1)^2 dp_1
\end{aligned}$$

In the relativistic limit ( $E_1 \approx p_1 c$ ), we have

$$\begin{aligned}
\int_0^{p_1^{max}} p_1^2 (E_0 - E_1)^2 dp_1 &= \int_0^{p_1^{max}} p_1^2 (E_0 - p_1 c)^2 dp_1 \\
&= \int_0^{p_1^{max}} p_1^2 (E_0^2 + p_1^2 c^2 - 2E_0 p_1 c)^2 dp_1 \\
&= \int_0^{p_1^{max}} (E_0^2 p_1^2 + p_1^4 c^2 - 2E_0 p_1^3 c) dp_1 \\
&= \left[ \frac{E_0^2 (p_1^{max})^3}{3} + \frac{(p_1^{max})^5 c^2}{5} - 2E_0 \frac{(p_1^{max})^4 c}{4} \right]
\end{aligned}$$

And recall that  $p_1^{max} = \frac{E_0}{c}$ . Thus,

$$\int_0^{p_1^{max}} p_1^2 (E_0 - E_1)^2 dp_1 = \frac{E_0^5}{30c^3}$$

so,

$$\boxed{W = \frac{|G_F \mathcal{M}|^2}{60\pi^3 \hbar^7 c^6} \cdot E_0^5}$$

or, in terms of the neutron and proton masses,

$$\begin{aligned}
W &= \frac{|G_F \mathcal{M}|^2}{60\pi^3 \hbar^7 c^6} \cdot [c^2 \cdot (m_n - m_p)]^5 \\
\implies &\boxed{W = \frac{|G_F \mathcal{M}|^2}{60\pi^3 \hbar^7 c^4} \cdot (m_n - m_p)^5}
\end{aligned}$$

(b) Assuming  $\mathcal{M} \sim 1$ , the lifetime of a neutron,  $\tau$ , is given by

$$\tau = \frac{1}{W} = \frac{60\pi^3 \hbar^7 c^6}{|G_F \mathcal{M}|^2} \cdot \frac{1}{(E_0)^5} = \frac{60\pi^3 \hbar^7 c^4}{|G_F \mathcal{M}|^2} \cdot \frac{1}{(m_n - m_p)^5} \sim 2500s = 2.5 \times 10^3 s$$

The experimentally measured value is  $\tau \sim 887.7s$ , so our estimate is nearly on the same order of magnitude. It's still a very crude approximation because we're off by a factor of about 3.

(c) One way in which the muon decays is

$$\mu \rightarrow e + \bar{\nu}_e + \nu_\mu$$

This decay follows mechanics similar to the beta decay we studied earlier, except in this case the light particles are  $\bar{\nu}_e, \nu_\mu$  while the heavy particles are  $\mu, e$ . Thus in this case  $E_0 = c^2(m_\mu - m_e)$ . Then, using the formula found earlier, we calculate the lifetime of the muon to be

$$\begin{aligned} \tau &= \frac{60\pi^3 \hbar^7 c^4}{|G_F \mathcal{M}|^2} \cdot \frac{1}{(m_\mu - m_e)^5} \\ &\sim 6.5 \times 10^{-7} s \end{aligned}$$

This is one order of magnitude off from the experimental value of  $\tau \approx 2.197 \times 10^{-6} s$ .

## Question 2: Sudden Approximation

A particle of mass  $m$  is in the ground state of a harmonic oscillator with the standard Hamiltonian  $\hat{H}_{HO} = \frac{\hat{p}^2}{2m} + \frac{k\hat{x}^2}{2}$ . At time  $t = 0$ , the value of  $k$  changes suddenly to  $k' = 4k$ . Find the probability that the oscillator remains in its ground state.

**Solution:**

$$\begin{aligned} \text{Before time } t = 0 : \hat{H}|\psi_n\rangle &= \hbar\omega \left(n + \frac{1}{2}\right) |\psi_n\rangle \text{ with } \hat{H} = \frac{\hat{p}^2}{2m} + \frac{k\hat{x}^2}{2} \\ \text{After time } t = 0 : \hat{H}'|\phi_{n'}\rangle &= \hbar\omega' \left(n' + \frac{1}{2}\right) |\phi_{n'}\rangle \text{ with } \hat{H}' = \frac{\hat{p}^2}{2m} + \frac{4k\hat{x}^2}{2} \end{aligned}$$

where  $\phi_n$  is the  $n^{\text{th}}$  state of the new harmonic oscillator with spring constant  $4k$ .

The particle starts off in the ground state, and the probability that we find it *still* in the ground state after the sudden change is

$$\begin{aligned} P_{(n'=0)} &= |\langle\phi_n(t)|\psi(t)\rangle|^2 \\ &= |d_{(n'=0)}|^2 \end{aligned}$$

In the position basis,

$$\begin{aligned} \psi_0(x) &= \langle x|\psi_0\rangle \\ &= \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar} \\ &= \left(\frac{m}{\pi\hbar}\right)^{1/4} \left(\frac{k}{\hbar}\right)^{1/8} e^{-x^2\sqrt{km}/2\hbar} \end{aligned}$$

and

$$\begin{aligned}
\phi_0(x) &= \langle x | \phi_0 \rangle \\
&= \left( \frac{m\omega'}{\pi\hbar} \right)^{1/4} e^{-m\omega' x^2 / 2\hbar} \\
&= \left( \frac{m}{\pi\hbar} \right)^{1/4} \left( \frac{4k}{m} \right)^{1/8} e^{-x^2 \sqrt{4km} / 2\hbar}
\end{aligned}$$

$$\begin{aligned}
\langle \phi_0 | \psi_0 \rangle &= \int_{-\infty}^{\infty} dx \langle \phi_0 | x \rangle \langle x | \psi_0 \rangle \\
&= \int_{-\infty}^{\infty} dx \phi_0^*(x) \psi_0(x) \\
&= \left( \frac{m}{\pi\hbar} \right)^{1/4} \left( \frac{m}{\pi\hbar} \right)^{1/4} \left( \frac{k}{m} \right)^{1/8} \left( \frac{4k}{m} \right)^{1/8} \int_{-\infty}^{\infty} dx \exp \left( -\frac{x^2 \sqrt{km}}{2\hbar} \right) \exp \left( -\frac{x^2 \sqrt{4km}}{2\hbar} \right) \\
&= \left( \frac{m}{\pi\hbar} \right)^{1/2} \cdot \underbrace{\left( \frac{k}{m} \right)^{1/4}}_{\sqrt{\omega}} \cdot (4)^{1/8} \int_{-\infty}^{\infty} dx e^{-\frac{3x^2 \sqrt{km}}{2\hbar}} \\
&= 2^{1/4} \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \sqrt{\frac{2}{3}} \sqrt{\frac{\pi\hbar}{2m\omega}} \\
&= 2^{1/4} \left( \frac{2}{3} \right)^{1/2} \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \cdot \left( \frac{m\omega}{\pi\hbar} \right)^{-1/2} \\
&= \frac{2^{3/4}}{3^{1/2}}
\end{aligned}$$

So,

$$\begin{aligned}
d_{(n'=0)} &= |\langle \phi_0 | \psi_0 \rangle|^2 \\
&= \frac{2^{3/2}}{3} \\
&= 2 \frac{\sqrt{2}}{3}
\end{aligned}$$

### Question 3: More Sudden Approximation

If the atom is initially in the ground state, what is the probability that the  ${}^3\text{He}^+$  ion remains in the ground state after the transition?

**Solution:**

The effect of the nuclear decay is to change the nuclear charge at  $t = 0$  without affecting the orbital electrons. We're interested in the probability that the ion remains in the ground state after the transition.

$$\begin{aligned}
\text{Before time } t = 0 : \hat{H} | \psi_{nlm} \rangle &= -\frac{E_1}{n^2} | \psi_{nlm} \rangle \\
\text{After time } t = 0 : \hat{H}' | \phi_{nlm} \rangle &= -\frac{2^2 E_1}{n^2} | \phi_{nlm} \rangle
\end{aligned}$$

where  $\hat{H}$  is the hydrogen atom hamiltonian wherein the nuclear charge is  $Z = 1$  and  $\hat{H}'$  is the same hamiltonian modified with  $Z = 2$  in this case.

The ground state solution for the a Hydrogenic Hamiltonian is

$$\psi_{100}(r) = \left( \frac{2}{\sqrt{4\pi a_0^3}} \right) e^{-r/a_0}$$

with the Bohr radius being

$$a_0 = \frac{\hbar}{mZe^2}$$

So, the Bohr radii for Tritium is

$$a_0^t = \frac{\hbar}{me^2}$$

and since the Helium ion has  $Z = 2$ , its bohr radius is

$$a_0^+ = \frac{\hbar}{2me^2} = \frac{a_0^t}{2}$$

the probability that the ion remains in the ground state after the sudden change is given by  $|\langle \phi_{100} | \psi_{100} \rangle|^2$ .

$$\begin{aligned} \langle \phi_{100} | \psi_{100} \rangle &= \iiint d^3\vec{r} \phi_{100}^*(r) \psi_{100}(r) \\ &= 4\pi \int_0^\infty r^2 dr \phi_{100}^*(r) \psi_{100}(r) \\ &= 4\pi \int_0^\infty r^2 dr \left( \frac{2}{\sqrt{4\pi (a_0^t/2)^3}} e^{-2r/a_0^t} \right)^* \left( \frac{2}{\sqrt{4\pi (a_0^t)^3}} e^{-r/a_0^t} \right) \\ &= 4\pi \int_0^\infty r^2 dr \left( \frac{4}{4\pi} \cdot \frac{2^{3/2}}{(a_0^t)^3} \right) e^{-3r/a_0^t} \\ &= 4\pi \frac{\sqrt{8}}{(a_0^t)^3 \pi} \int_0^\infty r^2 dr e^{-3r/a_0^t} \end{aligned}$$

Let  $y = r/a_0^t$ . Then,

$$\begin{aligned} \langle \phi_{100} | \psi_{100} \rangle &= \frac{4\sqrt{8}}{(a_0^t)^3} \int_0^\infty (a_0^t y)^2 (a_0^t dy) e^{-3y} \\ &= 4\sqrt{8} \int_0^\infty dy y^2 e^{-3y} \\ &= 4\sqrt{8} \cdot \left( \frac{2}{27} \right) \text{ Using Integral calculator} \\ &= \frac{8\sqrt{8}}{27} \\ &= \frac{16\sqrt{2}}{27} \\ &= 0.838 \end{aligned}$$

Thus the probability that we find the ion still in the ground state is  $|\langle \phi_{100} | \psi_{100} \rangle|^2 = (0.838)^2 \approx 0.703$