

Physics 137B Lecture (Not sure)

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These are notes taken from lectures on Quantum Mechanics delivered by Professor Raúl A. Briceño for UC Berkeley's Physics 137B class in the Spring 2024 semester.

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1 March 1 -

Recap

- Last time we discussed two particle states:

$$|a, b\rangle_B = \frac{1}{\sqrt{2}} (|a\rangle \otimes |b\rangle + |b\rangle \otimes |a\rangle)$$

$$|a, b\rangle_F = \frac{1}{\sqrt{2}} (|a\rangle \otimes |b\rangle - |b\rangle \otimes |a\rangle)$$

- Since the Fermionic wavefunction has to be antisymmetric, it must be the case that

$$|a, a\rangle_F = 0$$

1.1 Slater Determinant

If, instead of 2, we have N particles a useful tool for constructing antisymmetric states $\langle \phi_i | \phi_j \rangle = S_{ij}$ is the **Slater Determinant**.

$$|\psi\rangle_{A,N} = \frac{1}{\sqrt{N!}} \begin{bmatrix} |\phi_1(\vec{r}_1)\rangle & |\phi_2(\vec{r}_1)\rangle & \cdots & |\phi_N(\vec{r}_1)\rangle \\ |\phi_1(\vec{r}_2)\rangle & |\phi_2(\vec{r}_2)\rangle & \cdots & |\phi_N(\vec{r}_2)\rangle \\ \vdots & \vdots & \ddots & \vdots \\ |\phi_1(\vec{r}_N)\rangle & |\phi_2(\vec{r}_N)\rangle & \cdots & |\phi_N(\vec{r}_N)\rangle \end{bmatrix}$$

1.2 Non-interacting N particles

[Fill in some stuff from recording]

Let's consider : distinguishable
 spinless bosons
 spin-1/2 fermions

1.2.1 Distinguishable

[Write later]

1.2.2 Spinless Bosons

[Fill later]

Note: This is the starting point for Bose-Einstein condensates (at low temperatures, all of the bosons enter the same energy state.)

1.2.3 Spin-1/2 fermions

This time, it depends on whether N is odd or even.

Next, let's take the large N limit of a system with N fermions. This gives us the *Fermi Gas Model*.

1.3 Fermi Gas Model

- Large N limit of N non-interacting fermions
- Metals
- Heavy Nuclei
- White Dwarfs and Neutron Stars

Let's assume our fermions live in a cube of side-length L which is so large that our boundary conditions *don't quite matter*. i.e. it doesn't matter much to the behavior in the interior whether we apply periodic boundary conditions or something else.

Let's impose periodic boundary conditions. Then,

$$\begin{aligned}\phi(x+L, y, z) &= \phi(x, y, z) \\ \phi(x, y, z) &\sim e^{i(k_x x + k_y y + k_z z)} \text{ where } k_x L = 2\pi n_x\end{aligned}$$

It's useful to define the vector

$$\vec{k} = \frac{2\pi}{L} (n_x, n_y, n_z)$$

Let's try to count the number of modes which can fit in this box, with the goal of finding the total number of spin-1/2 particles inside the box.

We have

$$\begin{aligned}dn &= 2 \times dn_x dn_y dn_z \\ &= 2 \times \left(\frac{L}{2\pi}\right)^2 dk_x dk_y dk_z\end{aligned}$$

We're going to engage in a slight *oopsie* by assuming that even though our momenta are discrete, the steps between consecutive momenta are infinitesimally small. We're going to take the large N limit, where this assumption does hurt, so it should work out.

[Write motivation for Fermi Energy and Fermi Momentum]

Fermi Energy and Fermi Momentum

We define the ***Fermi Energy*** to be the largest energy that any particle in the gas takes.

[Write the rest from recording]