

**Phys 3810, Spring 2010**  
**Problem Set #4, Hint-o-licious Hints**

1. *Griffiths, 4.2* Show that the energies are given by

$$E_{(n_x, n_y, n_z)} = \frac{\pi^2 \hbar^2}{2ma^2} (n_x^2 + n_y^2 + n_z^2)$$

The latter part of the problem is a math–puzzle sort of thing where you figure out how many ways one can get the same  $E$  with different sets of integers.

You can find the values and the degeneracies from  $E_1$  up to  $E_{14}$ , but if you run out of patience, I'll tell you that

$$E_{14} = \frac{\pi^2 \hbar^2}{2ma^2} (27)$$

So what's different about this value from the first few?

2. *Griffiths, 4.38* The separated solution is

$$\psi(\mathbf{r}) = X(x)Y(y)Z(z)$$

Note that when you write out the Schrödinger equation in cartesian coordinates for this potential with the separated solution (substitute and then divide by  $\psi = XYZ$ ) it becomes a sum of three terms each of which depends only on  $x$ ,  $y$  or  $z$ . This means that you can write three separate Schrödinger equations, i.e.

$$-\frac{\hbar^2}{2m} \frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{2} m \omega^2 x^2 = E_x \quad , \text{etc.}$$

which we know how to solve. (Be as clear as you can about this part.) The total energy is just sum of energies from each separate Schrödinger equation:

$$E = E_x + E_y + E_z = \hbar \omega (n_x + n_y + n_z + \frac{3}{2})$$

The second part involves more thought. The energy of the oscillator state just depends on

$$n \equiv n_x + n_y + n_z \quad \text{for} \quad n_x = 0, 1, 2, \dots \quad \text{etc.}$$

so the question is how many ways can you get  $n$  from the three separate indices. That's a sort of puzzle–math problem. First, see if you can spot the pattern for the lowest  $n$ 's. (To give you one:  $n = 3$  has 10 possible states.)

3. *Griffiths, 4.8* Test that the  $u(r)$  radial solution

$$u_1(r) = A r j_1(kr)$$

really does solve Eq. (4.41).

For the infinite spherical well, the boundary condition  $u(a) = 0$  (continuity of the wave function) leads to  $ka = \tan(ka)$ . (Show this, of course.) Use a graph to demonstrate that for the higher solutions, we have

$$ka \approx \frac{(2n+1)}{2}\pi \quad \text{for } n = 1, 2, 3, \dots$$

4. *Griffiths, 4.10* Start off with  $c_0$  arbitrary and then use (4.76) to get the succeeding coefficients for the polynomial  $v(\rho)$ . They will truncate after very few terms. The math isn't hard it's just to get familiar with the notation and see how the math works out. The correct value for  $c_0$  would come from normalizing the radial function but that's not necessary here.

5. *Griffiths, 4.45* Obviously you want to do the integral  $\int_V |\psi_{100}|^2 dV$  for a sphere of radius  $b$  centered at the origin,  $b$  being the radius of the nucleus. Normally, I work these out by hand, but when my tables just gave a recursion formula for  $\int x^2 e^{-ax} dx$ , I just turned to Maple!

6. *Griffiths, 4.11* The normalization condition on  $R(r)$  is

$$\int_0^\infty R(r)^2 r^2 dr = 1$$

Find the  $c_0$  (that is, the overall constant) which makes this true and then write out  $R(r)$ . You can compare with Table 4.7.

7. *Griffiths, 4.14* The trick here is that the probability to find the electron at a certain radius  $r$  (within  $dr$ ) is *not*  $R(r)^2 dr$ . A hint comes from how  $R(r)$  is normalized (as used in Prob 6, Griff 4.11); the integral of any probability distribution must give 1.

When you do get the right probability function, find its maximum to get the most probable value.

8. *Griffiths, 4.16* Note that *only* place in the whole derivation where the electric charges showed up was in the potential

$$V(r) = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

Obviously if the central charge is replaced by  $+Ze$ , then the  $e^2$  in the potential gets replaced by  $Ze^2$ , but then  $e^2$  should be replaced *everywhere* by this. See how the result for the energy levels changes. (Is it proportional to  $Z$ ? To  $Z^2$ ?)