## Phys 3810, Spring 2009 Problem Set #4, Hint-o-licious Hints

1. Griffiths, 4.2 Show that the energies are given by

$$E_{(n_x.n_y,n_z)} = \frac{\pi^2 \hbar^2}{2ma^2} (n_x^2 + n_y^2 + n_z^2)$$

The latter part of the problem is a math–puzzle sort of thing where you figure out how many ways one can get the same E with different sets of integers.

You can find the values and the degeneracies from  $E_1$  up to  $E_{14}$ , but if you run out of patience, I'll tell you that

$$E_{14} = \frac{\pi^2 \hbar^2}{2ma^2} (27)$$

So what's different about this value from the first few?

2. Griffiths, 4.38 The separated solution is

$$\psi(\mathbf{r}) = X(x)Y(y)Z(z)$$

Note that when you write out the Schrödinger equation in cartesian coordinates for this potential with the separated solution (substitute and then divide by  $\psi = XYZ$ ) it becomes a sum of three terms each of which depends only on x y or z. This means that you can write three separate Schrödinger equations, i.e.

$$-\frac{\hbar^2}{2m}\frac{1}{X}\frac{d^2X}{dx^2} + \frac{1}{2}m\omega^2 = E_x$$
, etc.

which we know how to solve. (Be as clear as you can about this part.) The total energy is just sum of energies from each separate Schrödinger equation:

$$E = E_x + E_y + E_z = \hbar\omega(n_x + n_y + n_z + \frac{3}{2})$$

The second part involves more thought. The energy of the oscillator state just depends on

$$n \equiv n_x + n_y + n_z$$
 for  $n_x = 0, 1, 2, \dots$  etc.

so the question is how many ways can you get n from the three separate indices. That's a sort of puzzle—math problem. First, see if you can spot the pattern for the lowest n's. (To give you one: n = 3 has 10 possible states.)

**3.** Griffiths, **4.8** Test that the u(r) radial solution

$$u_1(r) = Arj_1(kr)$$

really does solve Eq. (4.41).

For the infinite spherical well, the boundary condition u(a) = 0 (continuity of the wave function) leads to  $ka = \tan(ka)$ . (Show this, of course.) Use a graph to demonstrate that for the higher solutions, we have

$$ka \approx \frac{(2n+1)}{2}\pi$$
 for  $n = 1, 2, 3, ...$ 

- 4. Griffiths, 4.10 Start off with  $c_0$  arbitrary and then use (4.76) to get the succeeding coefficients for the polynomial  $v(\rho)$ . They will truncate after very few terms. The math isn't hard it's just to get familiar with the notation and see how the math works out. The correct value for  $c_0$  would come from normalizing the radial function but that' snot necessary here.
- **5.** Griffiths, **4.45** Obviously you want to do the integral  $\int_V |\psi_{100}|^2 dV$  for a sphere of radius b centered at the origin, b being the radius of the nucleus. Normally, I work these out by hand, but when my tables just gave a recursion formula for  $\int x^2 e^{-ax} dx$ , I just turned to Maple!
- **6.** Griffiths, **4.11** The normalization condition on R(r) is

$$\int_0^\infty R(r)^2 r^2 \, dr = 1$$

Find the  $c_0$  (that is, the overall constant) which makes this true and then write out R(r). You can compare with Table 4.7.

7. Griffiths, 4.14 The trick here is that the probability to find the electron at a certain radius r (within dr) is not  $R(r)^2 dr$ . A hint comes from how R(r) is normalized (as used in Prob 6, Griff 4.11); the integral of any probability distribution must give 1.

When you do get the right probability function, find its maximum to get the most probable value.

8. Griffiths, 4.16 Note that only place in the whole derivation where the electric charges showed up was in the potential

$$V(r) = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

Obviously if the central charge is replaced by +Ze, then the  $e^2$  in the potential gets replaced by  $Ze^2$ , but then  $e^2$  should be replaced *everywhere* by this. See how the result for the energy levels changes. (Is it proportional to Z? To  $Z^2$ ?)