

**Phys 4610, Fall 2004**  
**Exam #1, Answers**

1. Find

a)

$$\nabla \cdot \mathbf{v} , \quad \text{where} \quad \mathbf{v} = 3xy\hat{\mathbf{x}} + 4yz\hat{\mathbf{y}} - zx\hat{\mathbf{z}}$$

$$\nabla \cdot \mathbf{v} = \frac{\partial}{\partial x}(3xy) + \frac{\partial}{\partial y}(4yz) + \frac{\partial}{\partial z}(-zx) = 3y + 4z - x = -x + 3y + 4z$$

b)

$$\nabla V , \quad \text{where} \quad V = s^2 z \cos \phi$$

$$\nabla V = \frac{\partial V}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial V}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial V}{\partial z} \hat{\mathbf{z}} = 2sz \cos \phi \hat{\mathbf{s}} - sz \sin \phi \hat{\boldsymbol{\phi}} + s^2 \cos \phi \hat{\mathbf{z}}$$

c)

$$\nabla \times \mathbf{v} , \quad \text{where} \quad \mathbf{v} = r \sin \theta \hat{\mathbf{r}} + r \cos \theta \hat{\boldsymbol{\phi}}$$

$$\begin{aligned} \nabla \times \mathbf{v} &= \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (r \sin \theta \cos \theta) \hat{\mathbf{r}} + \frac{1}{r} \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \phi} (r \sin \theta) - \frac{\partial}{\partial r} (r^2 \cos \theta) \right] \hat{\boldsymbol{\theta}} - \frac{1}{r} \frac{\partial}{\partial \theta} (r \sin \theta) \hat{\boldsymbol{\phi}} \\ &= \frac{(\cos^2 \theta - \sin^2 \theta)}{\sin \theta} \hat{\mathbf{r}} - 2 \cos \theta \hat{\boldsymbol{\theta}} - \cos \theta \hat{\boldsymbol{\phi}} \end{aligned}$$

d)

$$\nabla^2 V , \quad \text{where} \quad V = e^{-r/a} \cos \theta \sin \phi$$

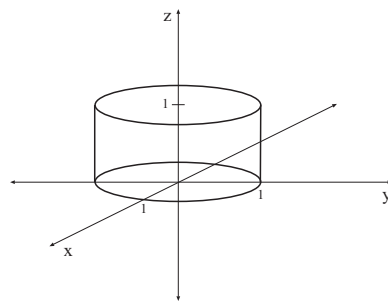
$$\begin{aligned} \nabla^2 V &= \frac{1}{r^2} \frac{\partial}{\partial r} \left( \frac{-r^2}{a} e^{-r/a} \cos \theta \sin \phi \right) - \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} e^{-r/a} \sin^2 \theta \sin \phi - \frac{1}{r^2 \sin^2 \theta} e^{-r/a} \cos \theta \sin \phi \\ &= \left[ -\frac{2}{ra} + \frac{1}{a^2} \right] e^{-r/a} \cos \theta \sin \phi - \frac{2}{r^2} e^{-r/a} \cos \theta \sin \phi - \frac{1}{r^2} e^{-r/a} \frac{\cos \theta}{\sin^2 \theta} \sin \phi \\ &= \left[ -\frac{2}{ra} + \frac{1}{a^2} - \frac{1}{r^2} \left( 2 + \frac{1}{\sin^2 \theta} \right) \right] e^{-r/a} \cos \theta \sin \phi \end{aligned}$$

2. The vector field  $\mathbf{v}$  is given by

$$\mathbf{v} = s \cos^2 \phi \hat{\mathbf{s}} + s \cos \phi \hat{\boldsymbol{\phi}} + z^2 \hat{\mathbf{z}}$$

a) Find the divergence of  $\mathbf{v}$ .

$$\begin{aligned}\nabla \cdot \mathbf{v} &= \frac{1}{s} \frac{\partial}{\partial s} (s^2 \cos^2 \phi) + \frac{1}{s} \frac{\partial}{\partial \phi} (s \cos \phi) + \frac{\partial}{\partial z} (z^2) \\ &= 2 \cos^2 \phi - \sin \phi + 2z\end{aligned}$$



b) Show that the divergence theorem is satisfied using, as the volume, a cylinder of radius 1 coaxial with the  $z$  axis and extending from  $z = 0$  to  $z = 1$ .

Test  $\int_V \nabla \cdot \mathbf{v} d\tau = \oint_S \mathbf{v} \cdot d\mathbf{a}$  over the given volume.

The lhs is

$$\int_{z=0}^1 \int_{\phi=0}^{2\pi} \int_{s=0}^1 (2 \cos^2 \phi - \sin \phi + 2z) s ds d\phi dz$$

The  $s$  integral can be done first:  $\int_0^1 s ds = \frac{1}{2}$ . For the second term, the  $\phi$  integral gives zero. This leaves:

$$\int_V \nabla \cdot \mathbf{v} d\tau = \frac{1}{2} \int_0^1 \int_0^{2\pi} (2 \cos^2 \phi + 2z) d\phi dz$$

use

$$\int_0^{2\pi} \cos^2 \phi d\phi = \left[ \frac{1}{2} \phi + \frac{1}{4} \sin 2\phi \right]_0^{2\pi} = \pi \quad \text{and} \quad \int_0^1 2z dz = 1$$

Then we have:

$$\rightarrow \pi + \frac{1}{2} \cdot 1 \cdot (2\pi) = 2\pi$$

Now do the rhs. There are 3 parts to the surface:

(i) Bottom:  $d\mathbf{a} = -s ds d\phi \hat{\mathbf{z}}$  so

$$\int_S \mathbf{v} \cdot d\mathbf{a} = \int_0^{2\pi} \int_0^1 s ds d\phi (-z^2) \Big|_{z=0} = 0$$

(ii) Sides:  $d\mathbf{a} = s d\phi dz \hat{\mathbf{s}}$ , so with  $s = 1$  on the sides,

$$\int_S \mathbf{v} \cdot d\mathbf{a} = \int_0^1 \int_0^{2\pi} 1 \cdot d\phi dz (1 \cdot \cos^2 \phi) = 1 \cdot \pi = \pi$$

where we used  $\int_0^{2\pi} \cos^2 \phi d\phi = \pi$  again.

(iii) Top:  $d\mathbf{a} = +s ds d\phi \hat{\mathbf{z}}$  with  $z = 1$ :

$$\int_S \mathbf{v} \cdot d\mathbf{a} = \int_0^{2\pi} \int_0^1 (s ds d\phi) (z^2) \Big|_{z=1} = \frac{1}{2} (2\pi) = \pi$$

Add them up

$$\oint_S \mathbf{v} \cdot d\mathbf{a} = \pi + \pi = 2\pi$$

The two sides of the divergence theorem agree.

3. Evaluate

$$\int_{-3}^3 (x^3 + 7x + 5)\delta(x - 1) dx$$

Since the integration interval does contain  $x = 1$ , the  $\delta$ -function “picks out” the value of the function under the integral sign at  $x = 1$  to give

$$\rightarrow = (1^3 + 7 \cdot 1 + 5) = 13$$

4. Consider a rectangle with a uniform surface charge density  $\sigma$ . The observation point  $P$  is in the plane of the rectangle on the bisector of the side of length  $L$ , a distance  $a$  from the nearest side. The other side of the rectangle has length  $b$ . See the figure.

Give the direction and magnitude of the  $E$  field at  $P$ . It will be sufficient for you to *clearly* set up any necessary integrals if they are at all difficult to work out!

Consider a strip of surface charge which lies at a distance  $x$  from  $P$  and has length  $L$ . The result from the text (given on the exam, at the end) gives the  $E$  field  $dE_z$  due to this strip. It is:

$$dE_z = \frac{1}{4\pi\epsilon_0} \frac{\lambda L}{x \sqrt{x^2 + (L/2)^2}}$$

where  $\lambda$  is the linear charge density of the strip,

$$\lambda = \frac{\sigma(Ldx)}{L} = \sigma dx$$

Then

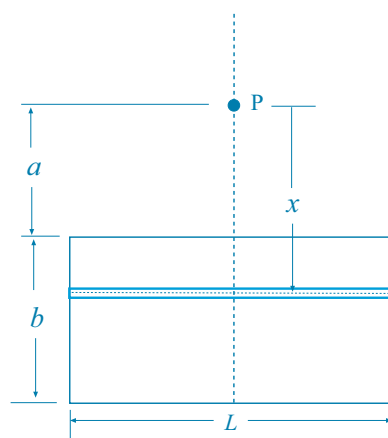
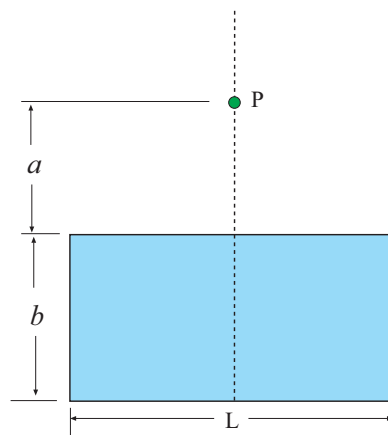
$$dE_z = \frac{1}{4\pi\epsilon_0} \frac{\sigma L dx}{x \sqrt{x^2 + (L/2)^2}}$$

Add up all the contributions as  $x$  ranges from  $a$  to  $a + b$ . Get:

$$E_z = \frac{\sigma L}{4\pi\epsilon_0} \int_a^{a+b} \frac{dx}{x \sqrt{x^2 + (L/2)^2}}$$

The integral is listed in tables and we get:

$$\begin{aligned} E_z &= \frac{\sigma L}{4\pi\epsilon_0} \left[ -\frac{2}{L} \log \left( \frac{L/2 + \sqrt{x^2 + (L/2)^2}}{x} \right) \right] \Big|_a^{a+b} \\ &= \frac{\sigma}{2\pi\epsilon_0} \log \left[ \frac{(a+b)}{a} \frac{(L/2 + \sqrt{a^2 + (L/2)^2})}{(L/2 + \sqrt{(a+b)^2 + (L/2)^2})} \right] \end{aligned}$$



Maybe one can simplify this...

5. If the electric potential in a certain region of space is given by

$$V(\mathbf{r}) = V_0 e^{-r^2/a^2}$$

a) What is the electric field in that region?

Use  $\mathbf{E} = -\nabla V$  and the expression for the gradient in spherical coordinates:

$$\mathbf{E} = -\frac{\partial V}{\partial r} \hat{\mathbf{r}} = +\frac{2rV_0}{a^2} e^{-r^2/a^2} \hat{\mathbf{r}} = \frac{2V_0}{a^2} r e^{-r^2/a^2} \hat{\mathbf{r}}$$

b) What is the charge density  $\rho(\mathbf{r})$ ?

Use  $\rho = \epsilon_0 \nabla \cdot \mathbf{E}$ . Then

$$\begin{aligned} \rho &= \epsilon \nabla \cdot \mathbf{E} = \frac{2V_0\epsilon_0}{a^2} \frac{1}{r^2} \frac{\partial}{\partial r} [r^3 e^{-r^2/a^2}] \\ &= \frac{2V_0\epsilon_0}{a^2} \left[ 3 - \frac{2r^2}{a^2} \right] e^{-r^2/a^2} \end{aligned}$$

c) How much charge is contained within a sphere of radius  $R$  centered at the origin?

Use Gauss's law,  $\oint \mathbf{E} \cdot d\mathbf{a} = Q_{\text{enc}}/\epsilon_0$ . The  $E$  field is radial so that  $\mathbf{E}$  is parallel to the area vector  $d\mathbf{a}$  everywhere and the magnitude of  $\mathbf{E}$  is the same everywhere on the sphere,

$$E_r(R) = \frac{2V_0}{a^2} R e^{-R^2/a^2}$$

. This gives

$$\oint \mathbf{E} \cdot d\mathbf{a} = (4\pi R^2) E_r(R) = \frac{8\pi R^3 V_0}{a^2} e^{-R^2/a^2} = Q_{\text{enc}}/\epsilon_0$$

Then  $Q_{\text{enc}}$  is

$$Q_{\text{enc}} = \frac{8\pi\epsilon_0 R^3 V_0}{a^2} e^{-R^2/a^2}$$

6. What is the electric potential at a point on the axis of a uniformly charged disk of radius  $R$  and surface charge density  $\sigma$ , a distance  $z$  from the center of the disk?

Assume  $V = 0$  at infinity.

The  $E$ -field at point  $P$  is

$$E_z = \frac{\sigma}{2\epsilon_0} \left[ 1 - \frac{z}{\sqrt{z^2 + R^2}} \right]$$

If we consider the path from  $\infty$  to  $P$  along the  $z$  axis, the change in potential is

$$V = - \int_{\infty}^z \mathbf{E}(z') \cdot d\mathbf{l}$$

where  $d\mathbf{l} = dz\hat{\mathbf{z}}$ , so

$$V = - \int_{\infty}^z \frac{\sigma}{2\epsilon_0} \left[ 1 - \frac{z'}{\sqrt{z'^2 + R^2}} \right] dz' = - \frac{\sigma}{2\epsilon_0} \left( z' - \sqrt{z'^2 + R^2} \right) \Big|_{\infty}^z$$

Now while we can substitute  $z' = z$  easily enough, substituting  $\infty$  gives “ $\infty - \infty$ ”. We need to take this limit carefully:

$$\lim_{z' \rightarrow \infty} (z' - \sqrt{z'^2 + R^2}) = \lim_{z' \rightarrow \infty} \left( \frac{1 - \sqrt{1 + R^2/z'^2}}{1/z'} \right)$$

which is now of the form  $0/0$ , so use L'Hopital's rule: Derivatives of numerator and denominator give

$$\lim_{z' \rightarrow \infty} \left( \frac{(R^2/z'^3)(1 + R^2/z'^2)^{-1/2}}{-1/z'^2} \right) = \lim_{z' \rightarrow \infty} -\frac{R^2}{z'} (1 + R^2/z'^2)^{-1/2}$$

which is now clearly zero. Using this,  $V$  is:

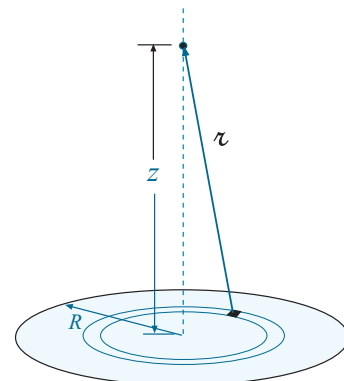
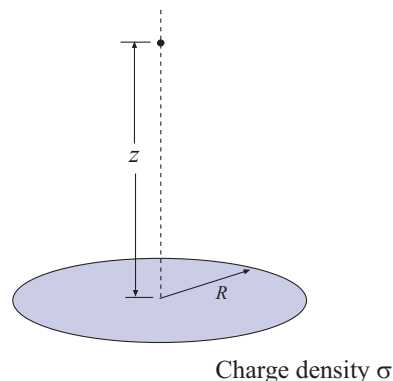
$$V(z) = -\frac{\sigma}{2\epsilon_0} (z - \sqrt{z^2 + R^2}) = \frac{\sigma}{2\epsilon_0} (\sqrt{z^2 + R^2} - z).$$

We could also have found the potential from scratch with

$$V(P) = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma da'}{r} = \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{s ds d\phi}{\sqrt{s^2 + z^2}}$$

This integral is easy, giving

$$V(P) = \frac{\sigma}{4\pi\epsilon_0} (2\pi) \sqrt{s^2 + z^2} \Big|_0^R = \frac{\sigma}{2\epsilon_0} (\sqrt{R^2 + z^2} - z)$$



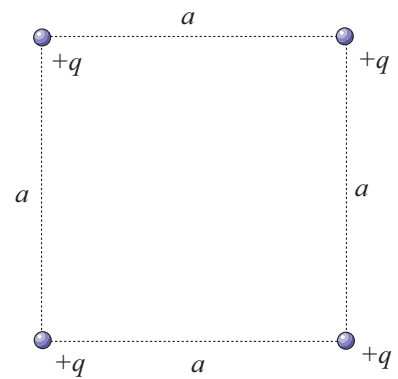
7. Find the work required to assemble four point charges  $+q$  in a square with side  $a$ .

We can use

$$W = \frac{1}{2} \frac{1}{4\pi\epsilon_0} \sum_{\substack{i,j \\ i \neq j}} \frac{q_i q_j}{r_{ij}} ,$$

(a sum over distinct pairs) then with 4  $q \cdot q$  pairs separated by  $a$  and 2  $q \cdot q$  pairs separated by  $\sqrt{2}a$ , the required work is

$$W = \frac{1}{4\pi\epsilon_0} \left[ 4 \frac{q^2}{a} + 2 \frac{q^2}{\sqrt{2}a} \right] = \frac{q^2}{4\pi\epsilon_0 a} [4 + \sqrt{2}]$$



## Useful Equations

$$\int_{\mathbf{a}}^{\mathbf{b}} (\nabla T) \cdot d\mathbf{l} = T(\mathbf{b}) - T(\mathbf{a}) \quad \int_V (\nabla \cdot \mathbf{v}) d\tau = \oint_S \mathbf{v} \cdot d\mathbf{a} \quad \int_S (\nabla \times \mathbf{v}) \cdot d\mathbf{a} = \oint_{\mathcal{P}} \mathbf{v} \cdot d\mathbf{l}$$

**Spherical:**

$$d\mathbf{l} = dl_r \hat{\mathbf{r}} + dl_\theta \hat{\boldsymbol{\theta}} + dl_\phi \hat{\boldsymbol{\phi}} \quad d\tau = r^2 \sin \theta dr d\theta d\phi$$

Gradient:

$$\nabla T = \frac{\partial T}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial T}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi} \hat{\boldsymbol{\phi}} \quad (1)$$

Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} \quad (2)$$

Curl:

$$\nabla \times \mathbf{v} = \frac{1}{r \sin \theta} \left[ \frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}} + \frac{1}{r} \left[ \frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} + \frac{1}{r} \left[ \frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}} \quad (3)$$

Laplacian:

$$\nabla^2 T = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2} \quad (4)$$

**Cylindrical:**

$$d\mathbf{l} = dl_s \hat{\mathbf{s}} + dl_\phi \hat{\boldsymbol{\phi}} + dl_z \hat{\mathbf{z}} \quad d\tau = s ds d\phi dz$$

Gradient:

$$\nabla T = \frac{\partial T}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial T}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial T}{\partial z} \hat{\mathbf{z}} \quad (5)$$

Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z} \quad (6)$$

Curl:

$$\nabla \times \mathbf{v} = \left( \frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right) \hat{\mathbf{s}} + \left( \frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right) \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[ \frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}} \quad (7)$$

Laplacian:

$$\nabla^2 T = \frac{1}{s} \frac{\partial}{\partial s} \left( s \frac{\partial T}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 T}{\partial \phi^2} + \frac{\partial^2 T}{\partial z^2} \quad (8)$$

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## More Math

$$\int \sin^2 x \, dx = \frac{1}{2}x - \frac{1}{4}\sin 2x \qquad \int \cos^2 x \, dx = \frac{1}{2}x + \frac{1}{4}\sin 2x$$

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## Physics:

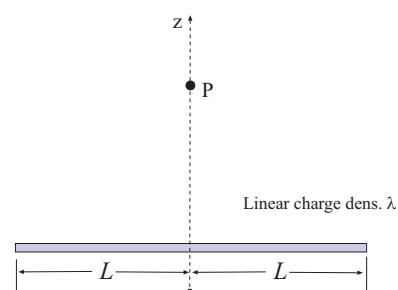
$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q \hat{\mathbf{r}}}{r^2} \qquad V(\mathbf{r}) = - \int_{\mathcal{O}}^{\mathbf{r}} \mathbf{E} \cdot d\mathbf{l}$$

$$\mathbf{E} = -\nabla V \qquad -\nabla^2 V = \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

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## Specific Results:

$$E_z = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z\sqrt{z^2 + L^2}}$$



$$E_z = \frac{\sigma}{2\epsilon_0} \left[ 1 - \frac{z}{\sqrt{z^2 + R^2}} \right]$$

