

**Phys 3810, Spring 2011**  
**Problem Set #3, Hint-o-licious Hints**

**1. Griffiths, 2.26** Following the hint (and assuming that Plancherel's theorem is good for *anything* we want to stick into it), show that the Fourier transform of  $\delta(x)$  is

$$F(k) = \frac{1}{\sqrt{2\pi}}$$

that is, a constant function. If we now merrily transform *back* to  $x$  space with the other transform formula, we get the desired (highly flaky) result.

**2. Griffiths, 2.29** Actually, this one isn't so bad. You just need to go through the finite-well derivation in the book and make the appropriate changes for the asymmetric state.

The main results you should get to are the results of applying the boundary conditions:

$$-\kappa = \ell \cot(\ell a)$$

and, using the same definitions of  $z$  and  $z_0$ , the new condition for finding the energies (graphically, perhaps)

$$-\cot(z) = \sqrt{(z_0/z)^2 - 1}$$

which is *not* guaranteed to have a root since the right side is positive and the left side starts off being negative for small  $z$ .

**3. Griffiths, 2.45** Follow the given hints and integrate to show that for any  $x_1$  and  $x_2$  we have

$$\int_{x_1}^{x_2} \left( \psi_2 \frac{\partial \psi_1}{\partial x} - \psi_1 \frac{\partial \psi_2}{\partial x} \right) dx = 0$$

Integration by parts of both terms leads to a cancellation and the result that

$$\psi_2 \frac{\partial \psi_1}{\partial x} - \psi_1 \frac{\partial \psi_2}{\partial x}$$

is constant (the *same* for any  $x$ . Since it is clearly zero as  $a \rightarrow \infty$  it is zero everywhere.

But we're still not done! Considering  $\frac{d}{dx}(\psi_1/\psi_2)$  will give the result that  $\psi_1$  and  $\psi_2$  are not independent.

**4. Griffiths, 3.7** (a) is pretty easy; (b) is also easy; just try the simplest linear combinations of  $f(x)$  and  $g(x)$  you can think of. Big hint, they're common functions accessed on calculators with a key marked "hyp".

**5. Griffiths, 3.12** Some steps: Start with

$$\langle x \rangle = \int_{-\infty}^{\infty} \Psi(x, t)^* x \Psi(x, t) dx$$

and substitute (3.55),

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{ipx/\hbar} \Phi(p, t) dp$$

Watch for complex conjugates and be sure that the dummy variable of  $p$  integration is *different* in the two integrals. When you do this you have a big integral over three variables:  $x$ ,  $p$  and  $p'$ .

Do what Griffiths says, use

$$xe^{ipx/\hbar} = \frac{\hbar}{i} \frac{d}{dp} e^{ipx/\hbar}$$

and then do an integration by parts. Be careful with variable names to know what the derivatives are acting on.

Use (and review) the result of problem 2.26,

$$\int_{-\infty}^{\infty} dx e^{iqx} = 2\pi\delta(q)$$

Use this for one of the integrations, but as usual be careful with the variables.

The next integration collapses the delta function and gives the desired result.

**6. Griffiths, 3.23** As discussed in class, from considering the action of this  $H$  on the state  $|1\rangle$  and  $|2\rangle$ , one can show that

$$H = \begin{pmatrix} E & E \\ E & -E \end{pmatrix}$$

With all of the  $\sqrt{2}$ 's flying around, the normalization is a little funky.