

For given λ : wavelength

$I(\lambda) \equiv$ Photon Flux $I_{\infty}(\lambda) \equiv$ Flux @ infinity (above atmosphere)

$n_a(\lambda) \equiv$ Number density of species a (O , N_2 and O_2)

$\sigma_a(\lambda) \equiv$ absorption cross section for species a as a function of λ
 $s \equiv$ path length

this is "a" (sigma - x sect)

$ds = -ds \sec \chi$
 χ : solar zenith angle
 χ : Altitude

Note: χ , s , σ_a are independent of I or λ

$$dI(\lambda) = -I(\lambda)(\lambda) \sigma_a^- ds$$

$$\frac{dI}{I} = -\sigma_a^- \sec \chi ds$$

$$\ln\left(\frac{I}{I_{\infty}}\right) = -\sigma_a^- \sec \chi \int_{\infty}^s n_a(\lambda) ds$$

$$= -\sigma_a^- \sec \chi \int_{\infty}^s n_a(\lambda) ds$$

For isothermal at atmosphere, this is $n_a(\lambda) \cdot H$

$$I_a = I_{\infty} e^{-\sigma_a^- \sec \chi n_a(\lambda) H}$$

Approximate! better to use \int

A bit more specific:

$$I(\lambda, \chi, z) = I_{\infty}(\lambda) e^{-\int_{\infty}^z \sigma_a^-(\lambda) n_a(\lambda) dz}$$

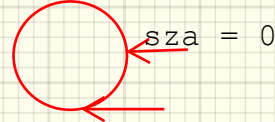
$$\text{Define } \mathcal{I} = \sigma_a^-(\lambda) n_a(\lambda) \Rightarrow \mathcal{I}(\lambda) = \sigma_a^-(\lambda) n_a(\lambda) \sum_a n_a(\lambda) dz$$

remember that $\sum_a n_a(\lambda) dz = n_a(\lambda) H_a(\lambda)$

Can use this for above the model domain

Assumption of $ds = -ds \sec \chi$ works to about $2^\circ - 7^\circ \rightarrow$ Smith & Smith 1972 for better result.

Assumes no horizontal variations in atmosphere! Yikes!



Now what?

Technically Absorption drives ionization
 \rightarrow creation of ions + electrons

Exothermic recombination
 \rightarrow Heating of atmosphere

$\sim 30\%$ of energy absorbed is turned to heat

Q EuV = energy added due to EuV

$$Q_{EuV} = \int_{\lambda=0}^{\infty} \int_{\chi=0}^{\pi/2} \int_{s=0}^{\infty} \sigma_a^-(\lambda) I(\lambda, \chi, s) n_a(\lambda) ds d\chi d\lambda$$

$$= \int_{\lambda=0}^{\infty} \sigma_a^-(\lambda) \int_{\chi=0}^{\pi/2} I(\lambda, \chi) n_a(\lambda) d\chi d\lambda$$

$$= \int_{\lambda=0}^{\infty} \sigma_a^-(\lambda) I_{\infty}(\lambda) e^{-\int_{\infty}^z \sigma_a^-(\lambda) n_a(\lambda) dz} d\lambda$$

See $\int_{\infty}^z n_a(\lambda) dz$ is a Chapman integral

$$\mathcal{I} = \sigma_a^-(\lambda) n_a(\lambda) H_a(\lambda)$$

number this is $\int_{\infty}^z n_a(\lambda) dz$ for exponentially decreasing, isothermal at atmosphere, and technically n_a is constant

$$\frac{dI}{dt} = \frac{Q_{EuV}}{P_{eq}}$$

q_p specific heat at constant pressure (or c_p @ constant volume)

- change with composition

- Can approximate as $1500 \frac{J}{kg \cdot K}$

$$\frac{k}{\lambda} = \frac{3}{2} \frac{R}{\lambda}$$

$$\frac{dT}{dt} = \frac{E \sum_a n_a(\lambda) \sigma_a^-(\lambda) I_{\infty}(\lambda) e^{-\int_{\infty}^z \sigma_a^-(\lambda) n_a(\lambda) dz}}{c_p \sum_a n_a(\lambda) H_a(\lambda)}$$

$$= \frac{E \sum_a n_a(\lambda) \sigma_a^-(\lambda) I_{\infty}(\lambda) e^{-\int_{\infty}^z \sigma_a^-(\lambda) n_a(\lambda) dz}}{c_p \sum_a n_a(\lambda) H_a(\lambda)}$$

$$\Rightarrow \text{Then } \frac{T^{n+1} - T^n}{\Delta t} = R$$

mean mass, could approximate as $\sim 16 \text{ AMU}$ (O), or could compute

Where to get $I_{\infty}(\lambda)$ & $\sigma_a^-(\lambda)$?
 \rightarrow Schwab & Noyes book

$$\rightarrow \text{EUVAC} \rightarrow I_{\infty}(\lambda) = F_{7413} [1 + A_{\lambda} (P_{800})]$$

- I have uploaded CSV file containing: $P = (F_{7417} + \langle F_{717} \rangle) / 2$

- Wave length λ (nm, E (J))
- F_{7413}
- A_{λ}
- $\sigma_a^-(\lambda)$ for O , M_2 , O_2

\rightarrow 81 day avg (3 solar rotation)

$$\text{Local time} = \left[\text{Lan} / 15 + \text{UT (hours)} \right] \% 24$$

$$\text{Equator } \chi = \arcsin \left(\cos \left(\left(\text{Local Time} - 12 \right) / 15 \right) \right)$$

(noon $\Rightarrow \chi = 0$; midnight $\Rightarrow \chi = 180^\circ$)

Steps:

1. Download CSV file

2. Get λ , F_{7413} , A_{λ} , $\sigma_a^-(\lambda)$, $\sigma_a^-(\lambda)$, $\sigma_a^-(\lambda)$

3. Make an initial function of T ; $T = 800 - 500 \cdot \lambda$ (compare) - create function initialize Temp

4. Start with one species: O calculate $H_a(\lambda)$ (compare) - create function that takes T, s, m , returns H

5. Calculate $n_a(\lambda)$ from 100-500 given $n_a(\text{use}) = 5 \times 10^{18} / \text{m}^3$ [for nit $n_a(\text{use}) = 1.7 \times 10^{18} / \text{m}^3$; $n_a(100) = 4 \times 10^{18} / \text{m}^3$] (compare) - create function

6. Assume $\chi = 0$; calculate \mathcal{I} from 100-500 km for λ_{use} (compare) - create function that takes $\lambda, n_a(\lambda), H_a(\lambda), \sigma_a^-(\lambda)$ (50-100)

7. Calculate $I_{\infty}(\lambda)$; then $I_{\infty}(\lambda)$ from 100-500 km (compare) - create EUVAC function

8. Calculate $E(50-100)$; calculate $\frac{dT}{dt}$ (for O only; $\lambda = 50-100$; $E = 100-500 \text{ km}$) (compare) - create function that takes $I_{\infty}(\lambda), \sigma_a^-(\lambda), \lambda$ returns $\frac{dT}{dt}$

9. Rescale functions to allow multiple wavelengths (plot $\frac{dT}{dt}$ for all wavelengths on same plot)

10. Rescale functions ($\mathcal{I}, \sigma_a^-(\lambda)$) to allow multiple species

11. Allow different χ

12. Update T , then $n_a, H_a, \sigma_a^-(\lambda)$

13. choose Δt ; t_{end} ; update $T, n_a, H_a, \sigma_a^-(\lambda)$ over 24 hrs