Heat transfer

Equation of heat conduction

In Cartesian co-ordinates,

$$\nabla^2 T + \frac{\dot{q}}{k} = \frac{1}{\alpha} \frac{dT}{dt}$$

where,

 ∇ = Laplace operator

 \dot{q} = Heat generation per unit volume

 α = Thermal diffusivity

- ullet Steady state condition: $abla^2 T + rac{\dot{q}}{k} = 0$ Poisson equation
- No heat sources: $\dot{q}=0$ Diffusion equation
- No heat source and steady state: $\dot{q}=0$ and $\frac{dT}{dt}=0$ Laplace equation

In cylinderical co-ordinates,

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial T}{\partial r}\right) + \frac{\dot{q}}{k} = \frac{1}{\alpha}\frac{dT}{dt}$$

In spherical co-ordinators,

$$\frac{\partial^2 T}{\partial r^2} + \frac{2}{r} \frac{\partial T}{\partial r} + \frac{1}{r^2 \sin \psi} \frac{\partial}{\partial \psi} \left(\sin \psi \frac{\partial T}{\partial \psi} \right) + \frac{1}{r^2 \sin \psi} \frac{\partial}{\partial \psi} \left(\frac{\partial^2 T}{\partial \phi} \right) + \frac{\dot{q}}{k} = \frac{1}{\alpha} \frac{dT}{dt}$$

If temperature variation is only in radial direction, then for steady state conditions with no heat generations, conduction equation becomes,

$$\frac{\partial^2 T}{\partial r^2} + \frac{2}{r} \frac{\partial T}{\partial r} = 0 \Rightarrow \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) = 0$$

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Boiling

Nucleate boiling: $q_s = \mu_l h_{fg} \left(g\sigma^{-1}(\rho_l - \rho_v)\right)^{\frac{1}{2}} \left(\left(\frac{C_{pl}}{C_{sg}}\right) \frac{\Delta T_e}{h_{fg}Pr^{1.7}}\right)^3$ Critical heat flux / Burnout point / Boiling crisis: $q_s = 0.18 h_{fg} \left(g\sigma(\rho_l - \rho_v)\rho_v^2\right)^{\frac{1}{4}}$