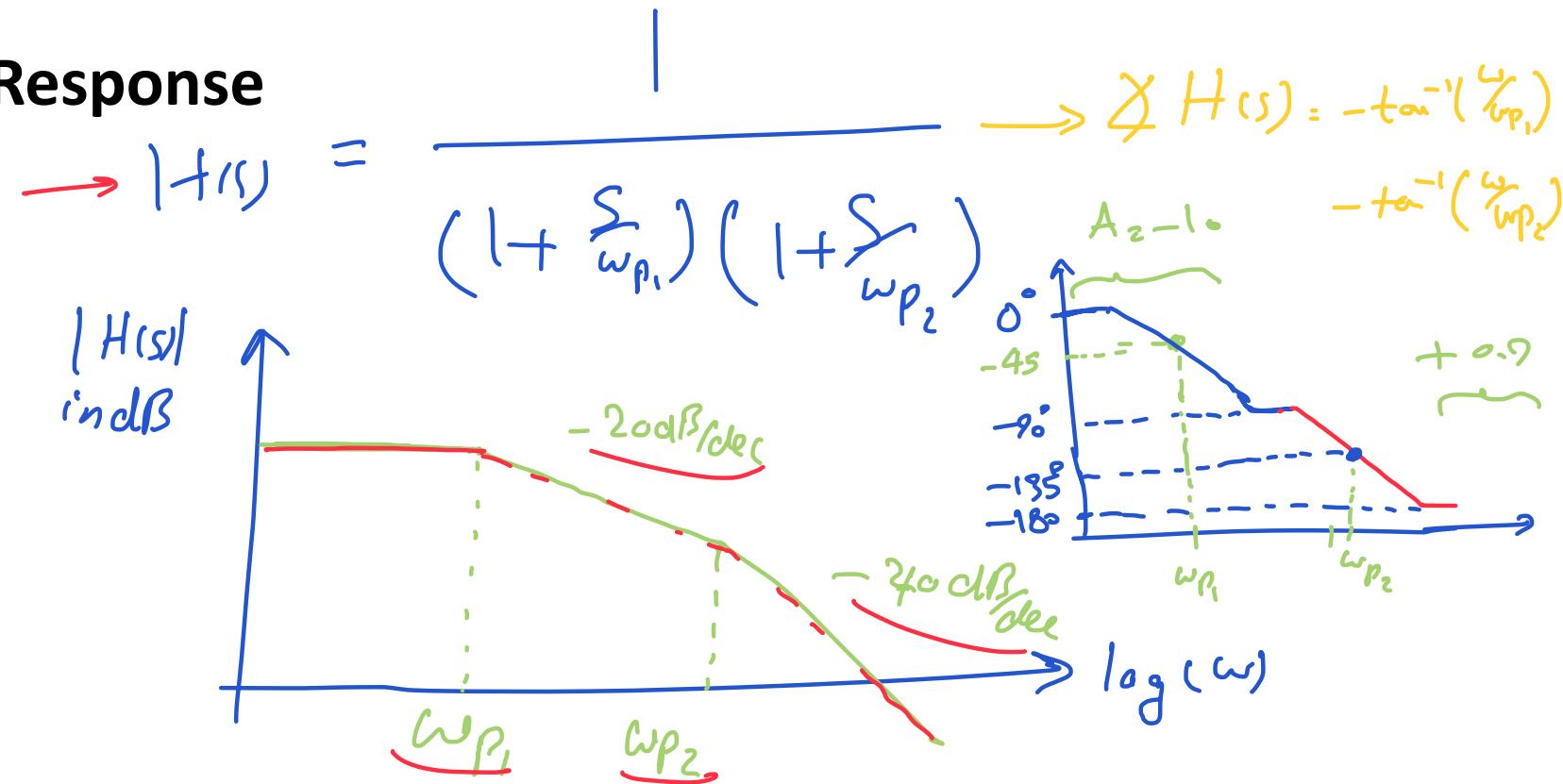
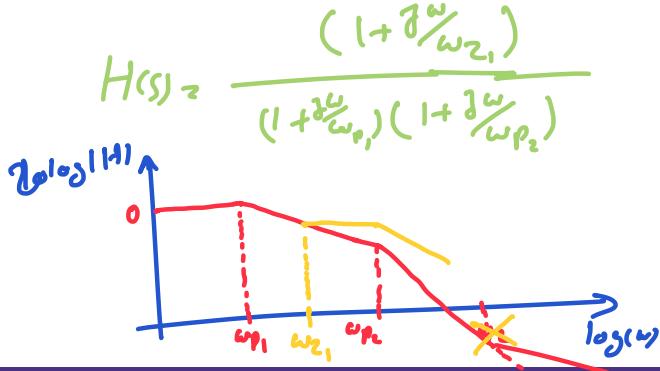


EE 332: Devices and Circuits II

Lecture 6: Frequency Response

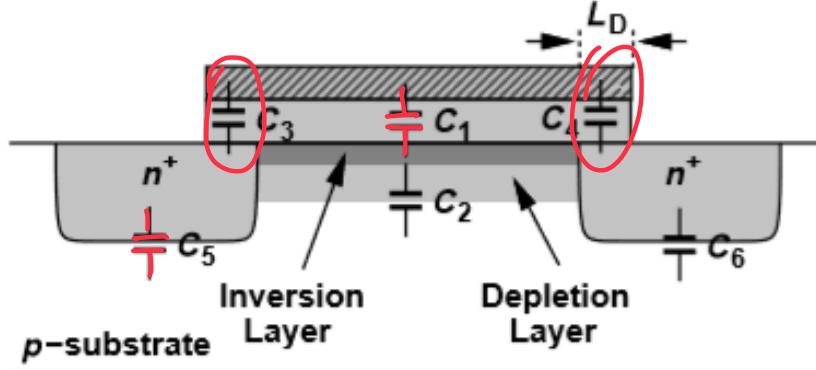
Prof. Sajjad Moazeni
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Autumn 2022



MOS Device Capacitances

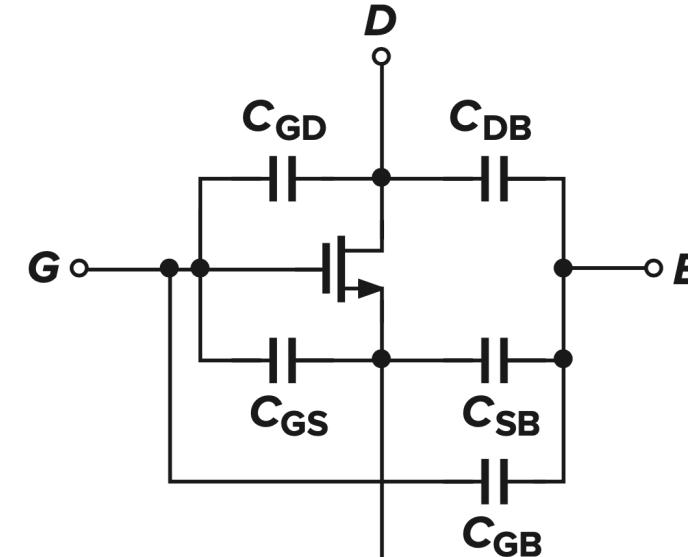
$$Q_L \rightarrow Z_L = SL \rightarrow j\omega L$$



$$C \propto WL$$

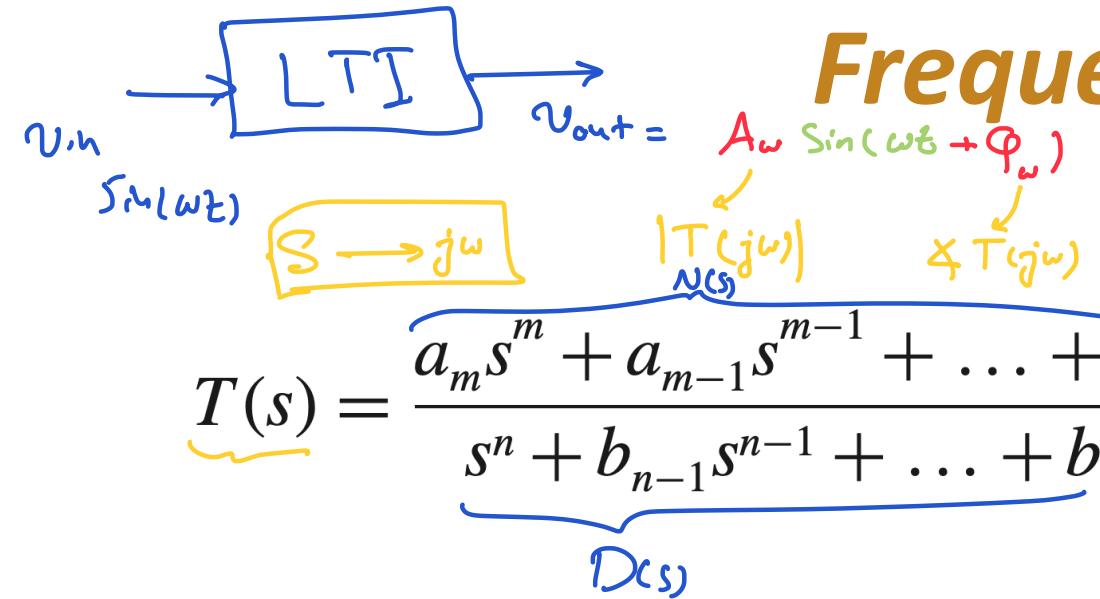
$$\propto W, L, E$$

$$C \frac{1}{T} \rightarrow \text{Inductor symbol} \quad Z_c = \frac{1}{SC} = \frac{1}{j\omega C}$$



- Device capacitances cause frequency dependent Av, Z_{in} , Z_{out} , etc.
- Capacitance exists almost between every two of the four terminals
- Larger device \rightarrow Larger caps!
except D-S

Frequency Response



$$T(s) = \frac{a_m s^m + a_{m-1} s^{m-1} + \dots + a_0}{s^n + b_{n-1} s^{n-1} + \dots + b_0} \Rightarrow T(s) = a_m \frac{(s - Z_1)(s - Z_2) \dots (s - Z_m)}{(s - P_1)(s - P_2) \dots (s - P_n)}$$

$N(s) = 0 \rightarrow s_z$ zero
 $D(s) = 0 \rightarrow s_p$ pole
in Zeros
n poles

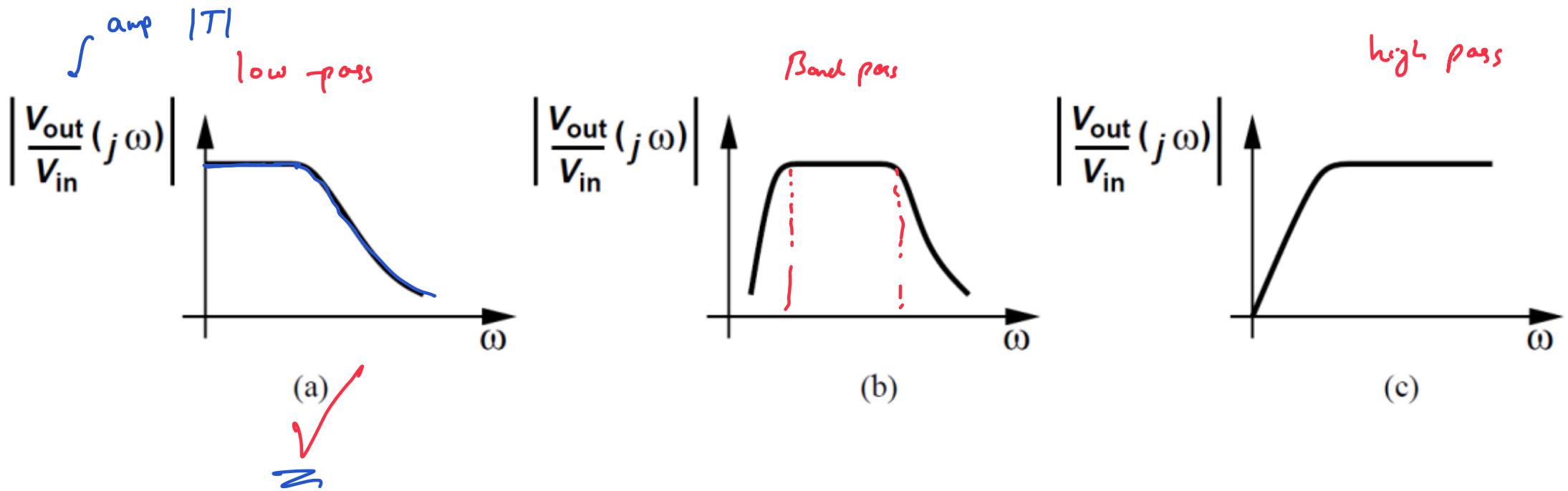
$$Z = -\omega_{z1}, -\omega_{z2}, \dots, -\omega_{zn}$$

$$P = -\omega_{p1}, -\omega_{p2}, \dots, -\omega_{pn}$$

$$H(s) = \frac{b_n}{a_m} \cdot \frac{\left(1 + \frac{s}{\omega_{z1}}\right) \left(1 + \frac{s}{\omega_{z2}}\right) \dots \left(1 + \frac{s}{\omega_{zn}}\right)}{\left(1 + \frac{s}{\omega_{p1}}\right) \left(1 + \frac{s}{\omega_{p2}}\right) \dots \left(1 + \frac{s}{\omega_{pm}}\right)}$$

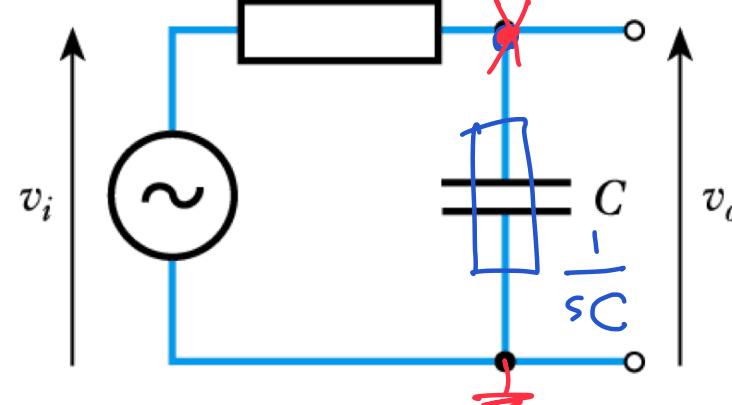
DC gain

General Considerations



- Zeros and poles are respectively defined as the roots of the numerator and denominator of the transfer function.
- In this chapter, we are primarily interested in the magnitude of the transfer function.
- The magnitude of a complex number $a + jb$ is given by $\sqrt{a^2 + b^2}$

$$\frac{1}{j\omega} = \frac{-j}{\omega C}$$



$$T(s) = \frac{Z_c}{Z_c + Z_R} = \frac{\frac{1}{sC}}{\frac{1}{sC} + R} = \frac{1}{1 + sRC}$$

(pole @ $-j\frac{1}{RC}$) = $\omega_{p_1} = \frac{1}{RC}$

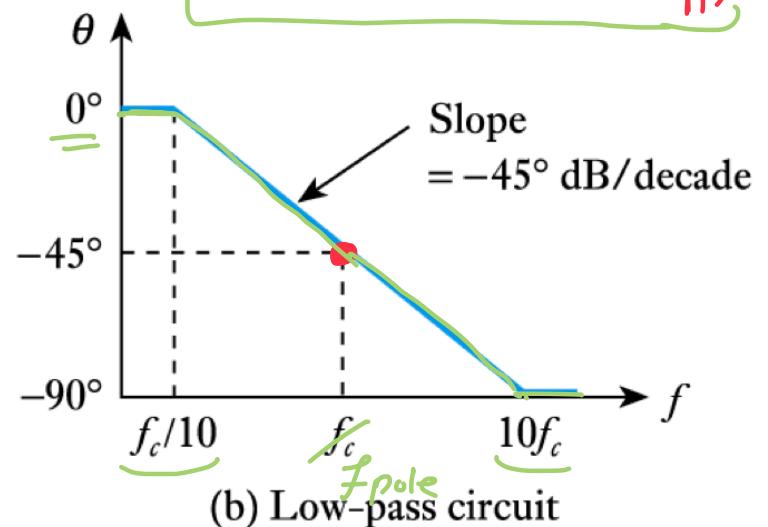
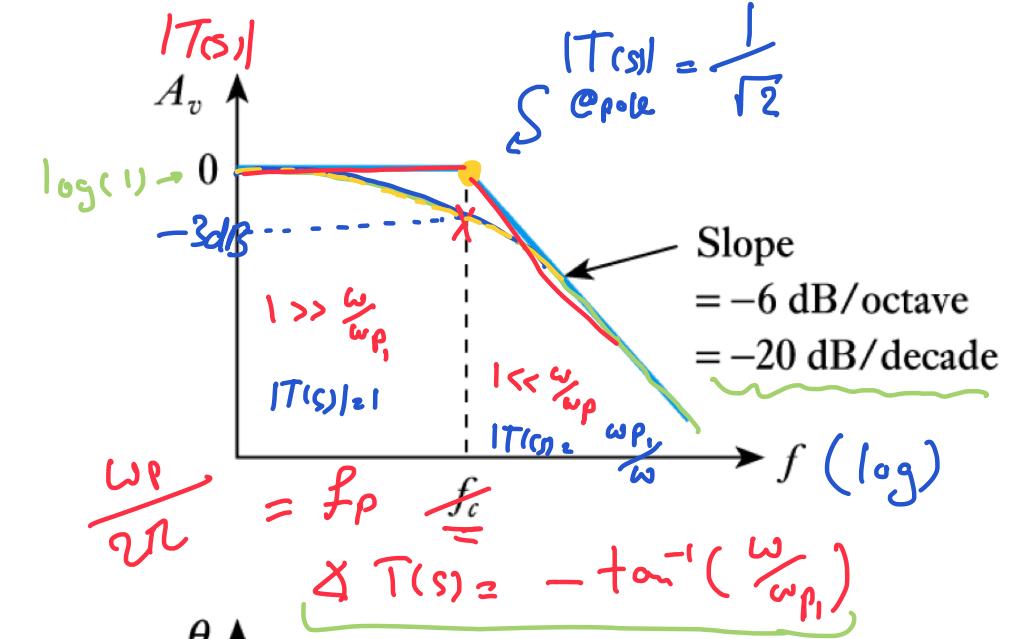
Bode Plot

$$T = \frac{v_o}{v_i} = ?$$

$* \omega = 2\pi f$

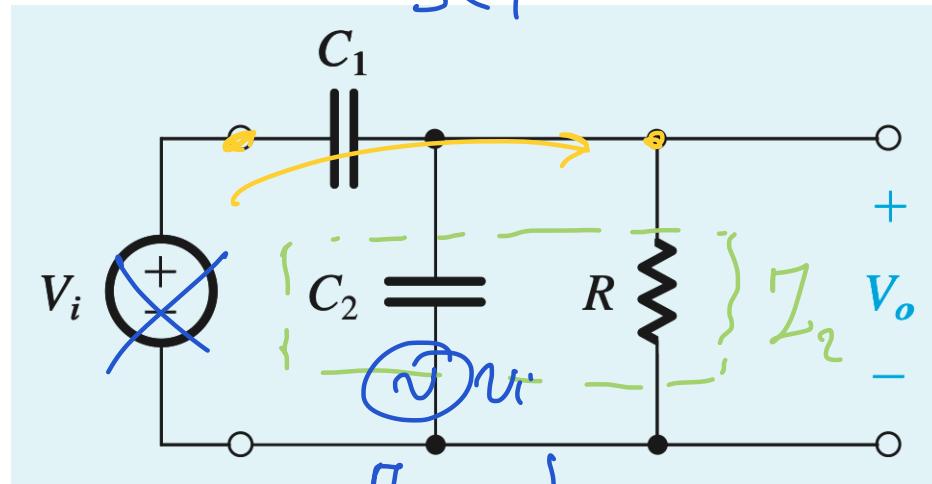
$$A_v = 20 \log (|T(s)|)$$

$$|T(s)| = \frac{1}{\sqrt{1 + (\omega/\omega_{p_1})^2}} \rightarrow 20 \log |T(s)| = \dots$$



Example 1

$$Z_{C_1} = \frac{1}{SC_1}$$



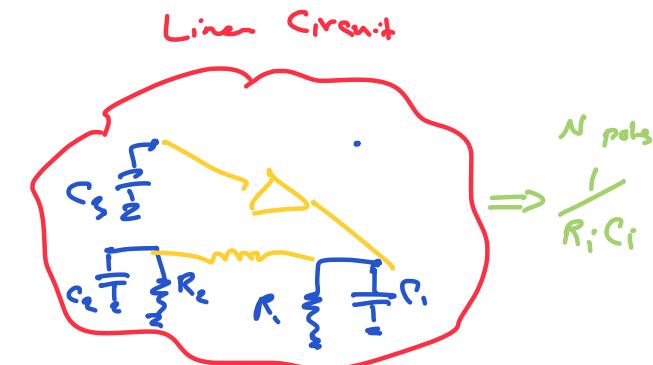
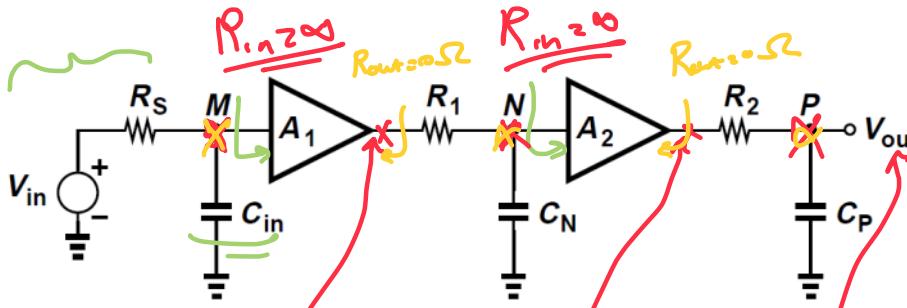
$$\frac{V_o}{V_i}(s) = \frac{Z_2}{Z_2 + Z_{C_1}} = \frac{sRC_1}{1 + sR(C_1 + C_2)}$$

$$I_2(s) = ? \quad \frac{1}{\frac{1}{R} + \frac{1}{Z_{C_2}}} = \frac{R}{1 + SRC_2}$$

$$\frac{RSC_1}{1 + SRC_2} = \frac{\cancel{RSC_1}}{1 + \cancel{SRC_2}} + \frac{(1 + sRC_2)}{\cancel{SC_1}}$$

\rightarrow zero @ $\omega_2 \rightarrow 0$
 \rightarrow pole @ $\omega_p = \frac{1}{R(C_1 + C_2)}$

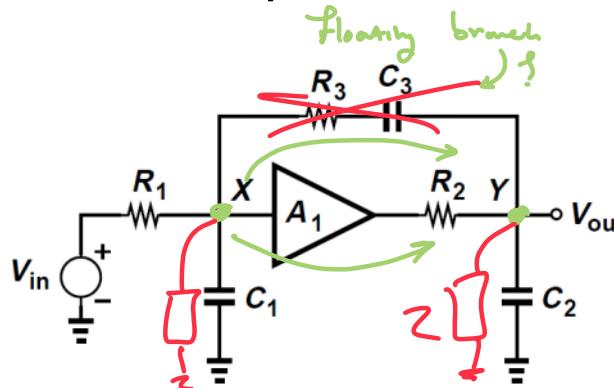
Association of Poles with Nodes



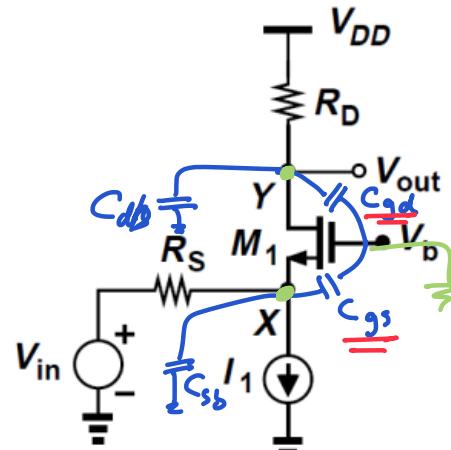
- The overall transfer function can be written as

$$\frac{V_{in}}{V_{in}} \times \frac{V_N}{V_M} \times \frac{V_{out}}{V_N} = \left\{ \frac{V_{out}}{V_{in}}(s) = \frac{A_1}{1 + R_S C_{in} s} \cdot \frac{A_2}{1 + R_1 C_N s} \cdot \frac{1}{1 + R_2 C_P s} \right.$$

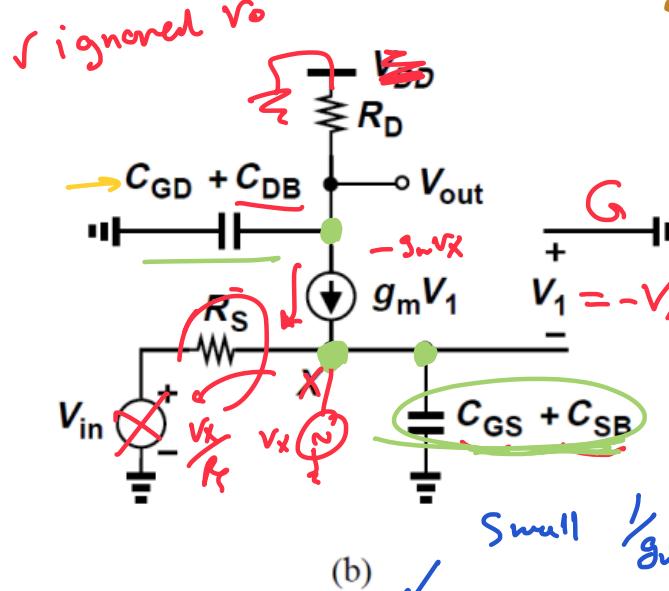
- Each node in the circuit contributes one pole to the transfer function.
- Not valid in general. Example:



Example: Common-Gate Stage ✓



(a)



(b)

- At node X:
- At node Y:
- The overall transfer function:

$$\omega_{in} = \left[(C_{GS} + C_{SB}) \left(R_S \parallel \frac{1}{g_m + C_{GS}} \right) \right]^{-1}$$

$\omega_{out} \ll \omega_{in}$

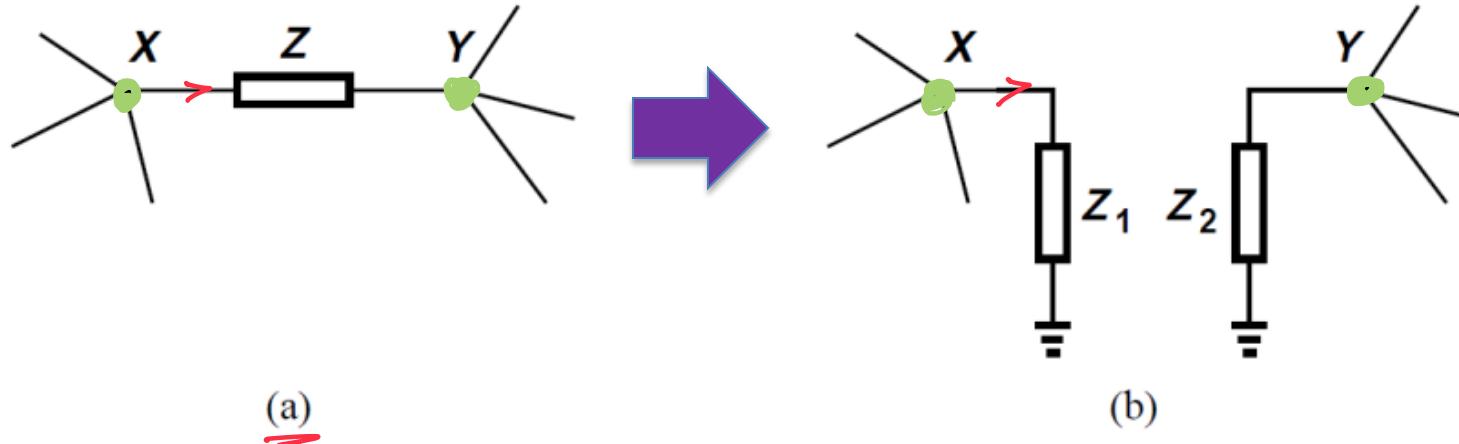
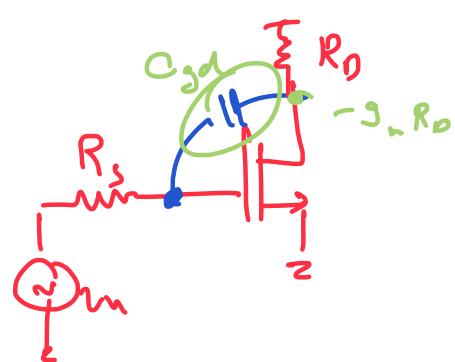
$$\omega_{out} = [(C_{DG} + C_{DB})R_D]^{-1}$$

$$\frac{1}{g_m} \ll R_D$$

$$\frac{V_{out}(s)}{V_{in}(s)} = \frac{(g_m + \cancel{sC_{GS}})R_D}{1 + (g_m + \cancel{sC_{GS}})R_S} \cdot \frac{1}{\left(1 + \frac{s}{\omega_{in}}\right)\left(1 + \frac{s}{\omega_{out}}\right)}$$

DC gain
 2 pols

Miller effect



$$I_Z = \frac{V_X - V_Y}{Z} = \frac{V_X}{Z_1}$$

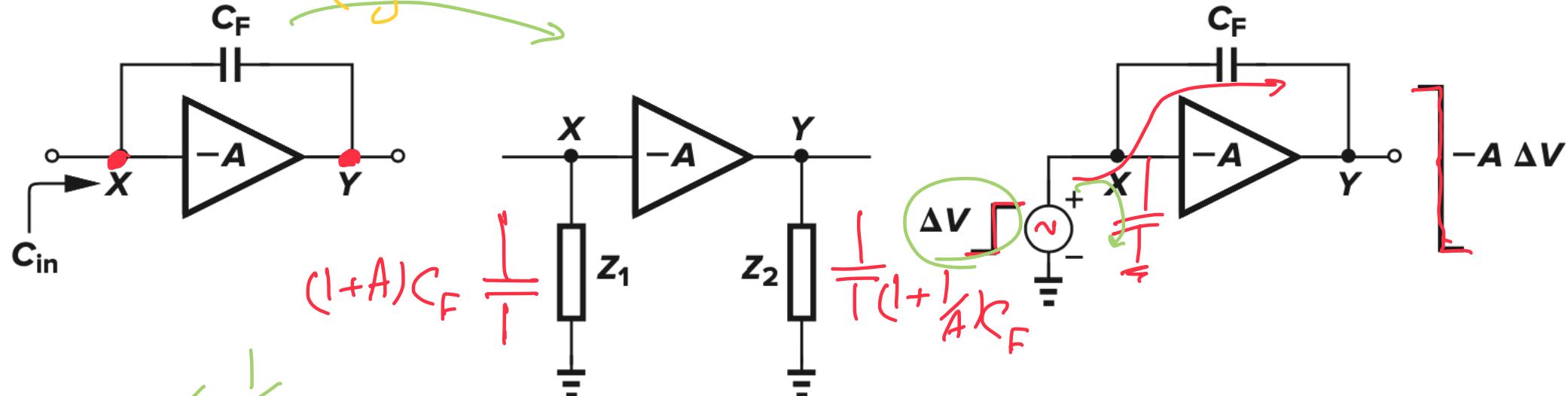
$$Z_1 = \frac{Z}{1 - \frac{V_Y}{V_X}}$$

$$Z_2 = \frac{Z}{1 - \frac{V_X}{V_Y}}$$

$$\frac{V_y}{V_x} = -A \rightarrow \text{low freq gain}$$

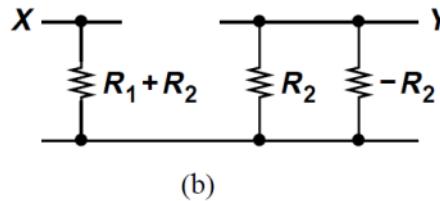
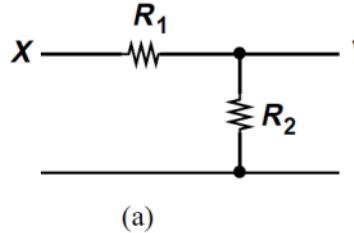
(ignore caps ✓)

Example

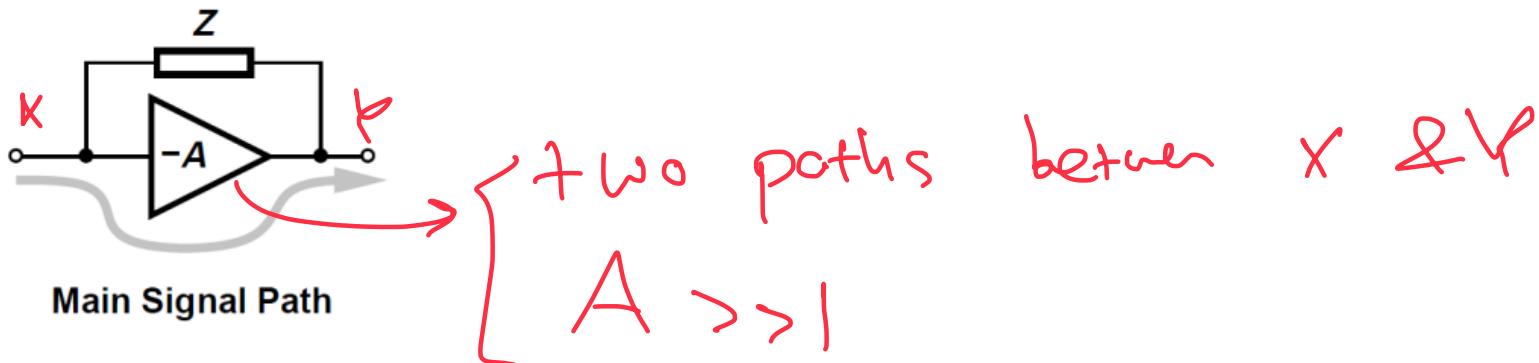


$$\left\{ \begin{array}{l} Z_{1e} = \frac{Z_F s C_F}{1+A} \\ Z_{2e} = \frac{Z_F}{1+\frac{1}{A}} \end{array} \right. \Rightarrow \left\{ \begin{array}{l} Z_{1e} = \frac{1}{s(1+A)C_F} \\ Z_{2e} = \frac{1}{s(1+\frac{1}{A})C_F} \end{array} \right. \begin{array}{l} A \gg 1 \\ C_1 = (1+A)C_F \\ C_2 \approx C_F \end{array}$$

Validity of Miller's Theorem



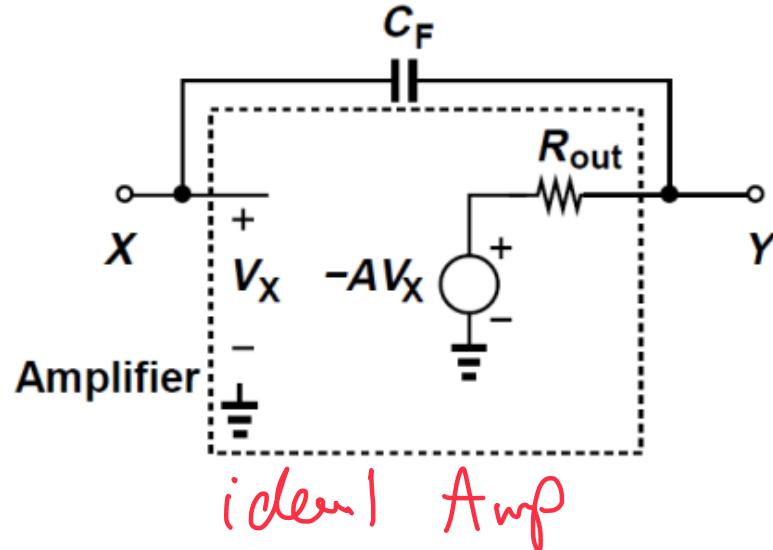
Improper application of Miller's theorem



Typical case for valid application of Miller's theorem.

- Miller's theorem does not stipulate the conditions under which this conversion is valid.
- If the impedance Z forms the only signal path between X and Y , then the conversion is often invalid.

Example



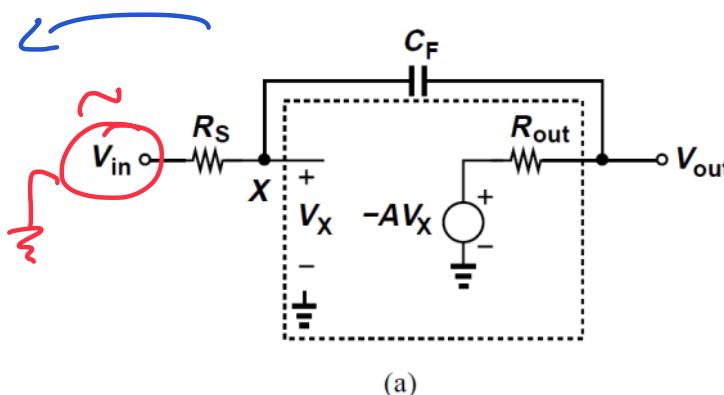
- The value of $A_v = V_Y / V_X$ must be calculated at the frequency of interest.
- In the figure, the equivalent circuit reveals that $V_Y \neq -A V_X$ at high frequencies.
- In many cases we use the low-frequency value of V_Y / V_X to gain insight.
- We call this approach “Miller’s approximation.”

Example

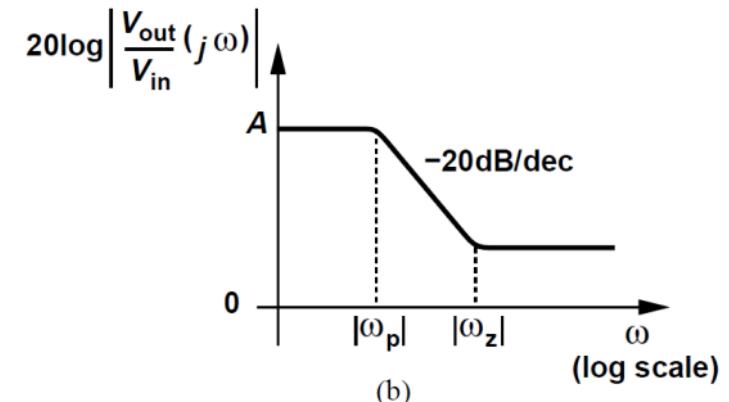
- Direct Calculation:

$$\frac{V_{out}}{V_{in}}(s) = \frac{R_{out}C_F s - A}{[(A+1)R_S + R_{out}]C_F s + 1}$$

1 Zero & 1 pole



(a)

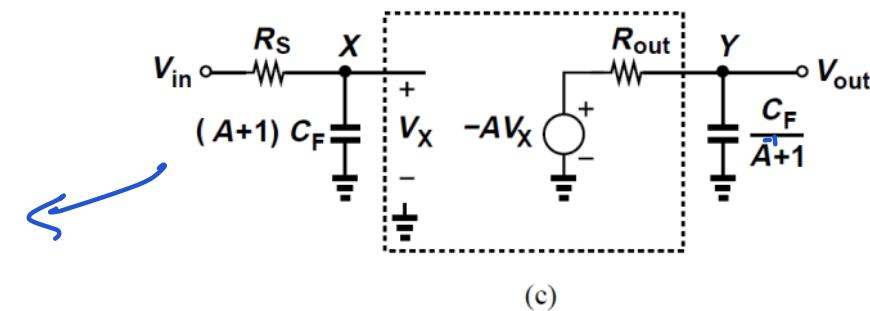


(b)

- Miller Approximation:

$$\frac{V_{out}}{V_{in}}(s) = \frac{-A}{[(1+A)R_S C_F s + 1] \left(\frac{1}{1+A^{-1}} C_F R_{out} s + 1 \right)}$$

No zeros & 2 poles



(c)

Miller Est.

- Miller's approximation has eliminated the zero and predicted two poles for the circuit!

✓ Dominant pole

Miller's approximation

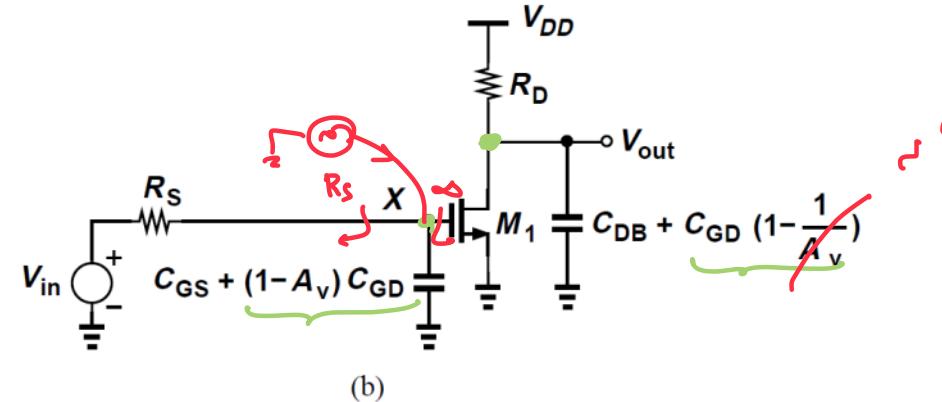
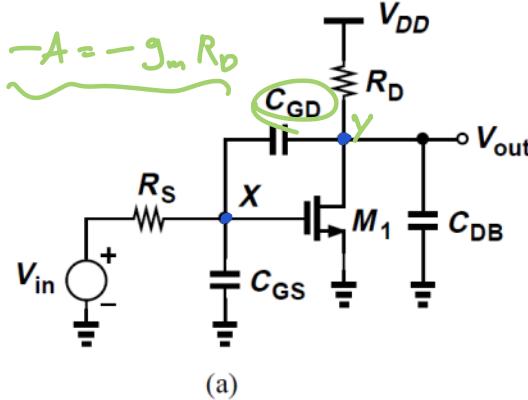
- Miller's approximation:

- (1) it may eliminate zeros
- (2) it may predict additional poles
- (3) it does not correctly compute the “output” impedance

✓ Predicts the dominant pole
 $\sim 2\text{dB}$ BW of Amp.

Miller effect ✓

Common-Source Stage



- The magnitude of the “input” pole (using Miller approximation):

$$A_{le} = -g_m R_D \Rightarrow$$

$$\omega_{in} = \frac{1}{R_S [C_{GS} + (1 + g_m R_D) C_{GD}]} \quad \checkmark$$

- At the output node

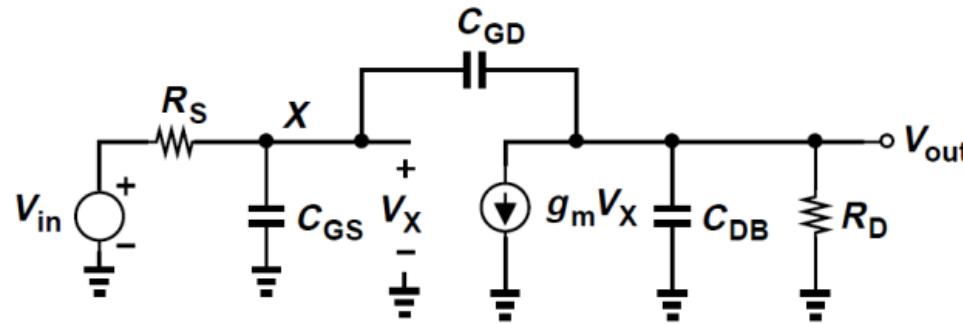
$$\omega_{out} = \frac{1}{R_D (C_{DB} + C_{GD})} \quad \checkmark$$

$$\frac{V_{out}}{V_{in}}(s) = \frac{-g_m R_D}{\left(1 + \frac{s}{\omega_{in}}\right) \left(1 + \frac{s}{\omega_{out}}\right)}$$

$-\omega_m, -\omega_{out}$

Direct Analysis

1 zero
2 poles



$$\frac{V_{out}}{V_{in}}(s) = \frac{(C_{GDS} - g_m)R_D}{R_S R_D \xi s^2 + [R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})]s + 1}$$

- While the denominator appears rather complicated, it can yield intuitive expressions for the two poles. $|\omega_{p1}| \ll |\omega_{p2}|$
- “Dominant pole” approximation.

$$\underline{\omega_{p1}} = \frac{1}{\underbrace{R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})}_{\text{approximate dominant pole term}}} \quad \text{green bracket and arrow}$$

- The intuitive approach provides a rough estimate with much less effort.

“Dominant Pole” Approximation

- Assume a two pole transfer function:

$$\omega_{p_1} \ll \omega_{p_2}$$

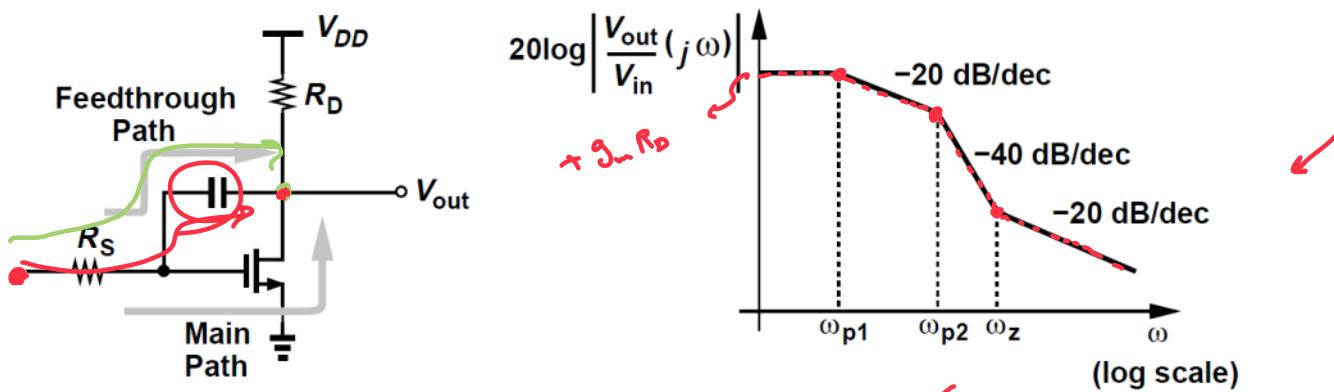
$$D = \left(\frac{s}{\omega_{p_1}} + 1 \right) \left(\frac{s}{\omega_{p_2}} + 1 \right)$$
$$= \frac{s^2}{\omega_{p_1}\omega_{p_2}} + \left(\frac{1}{\omega_{p_1}} + \frac{1}{\omega_{p_2}} \right) s + 1$$
$$A s^2 + \cancel{\beta s} + G$$

- If one pole will be dominant ($|\omega_{p_1}| \ll |\omega_{p_2}|$), then:



$$\left\{ \begin{array}{l} \omega_{p_1} \approx \frac{1}{\beta} \\ \omega_{p_2} \approx \beta \end{array} \right.$$

Zero in Transfer Function



$$\frac{V_{out}(s)}{V_{in}} = \frac{(C_{GD}s - g_m)R_D}{R_S R_D \xi s^2 + [R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})]s + 1}$$

- The transfer function exhibits a zero given by

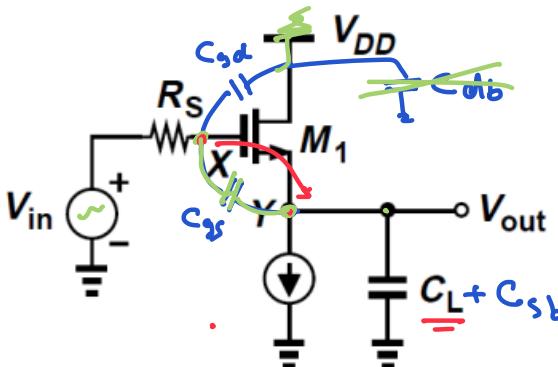
$$\omega_z = \underline{+g_m/C_{GD}} \quad \gg \quad \omega_{p_1}, \omega_{p_2}$$

- CGD provides a feedthrough path that conducts the input signal to the output at very high frequencies.

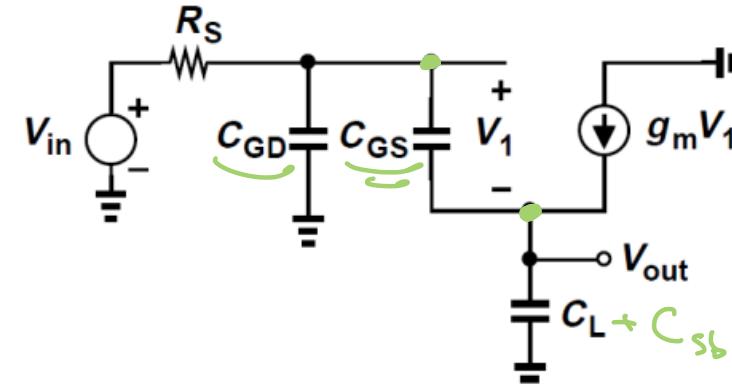
$A_V < 1$

don't use Miller X

Source Followers



(a)



(b)

- The strong interaction between nodes X and Y through C_{GS} makes it difficult to associate a pole with each node.

$$\frac{V_{out}(s)}{V_{in}} = \frac{g_m + C_{GSS}}{R_s(C_{GS}C_L + C_{GS}C_{GD} + C_{GD}C_L)s^2 + (g_m R_s C_{GD} + C_L + C_{GS})s + g_m}$$

$\omega_{Zn} - \frac{C_{GS}}{g_m}$

} 1 zero
 } 2 poles
 Roots

- Contains a zero in the left half plane. Why?

↓ No need
for Miller

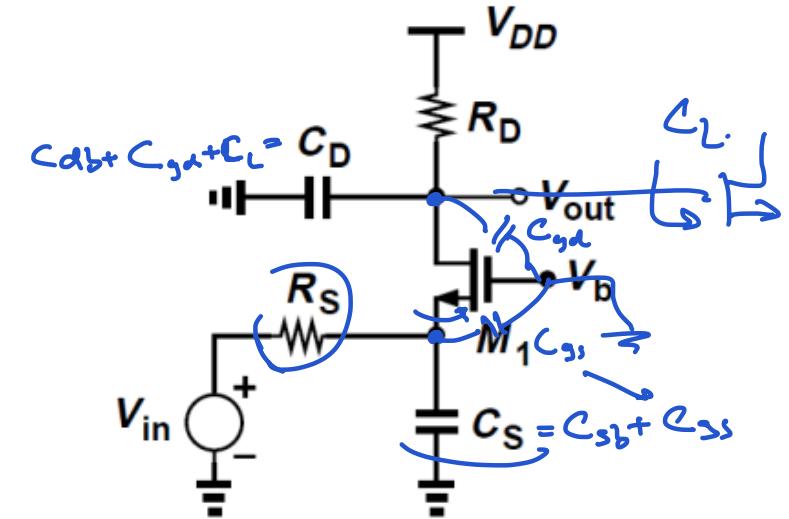
Common-Gate Stage

- A transfer function (w/o channel length modulation):

$$\frac{V_{out}}{V_{in}}(s) = \frac{(g_m + g_{mb})R_D}{1 + (g_m + g_{mb})R_S} \frac{1}{\left(1 + \frac{C_S}{g_m + g_{mb} + R_S^{-1}s}\right)(1 + R_D C_{DS}s)}$$

R_S || g_m

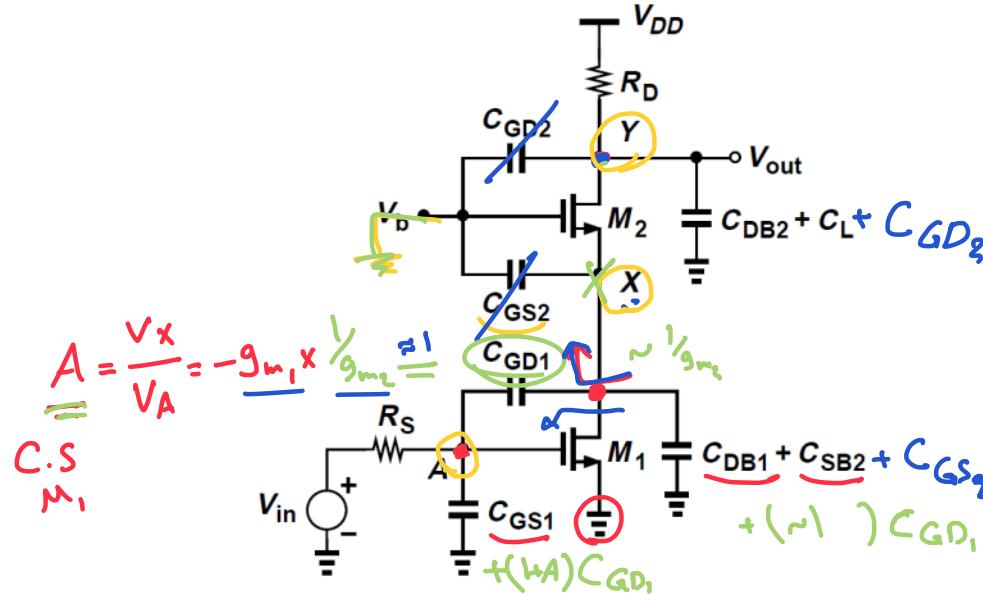
- No Miller multiplication of capacitances.
- R_D is typically maximized, so the dc level of the input signal must be quite low.
- As an amplifier in cases where a low input impedance is required in cascode stages.



$$r_o \text{ & } R_{o2} \gg \frac{1}{g_{m1,2}}$$

Cascode Stage

✓ Miller effect
is acceptable
to use



High-frequency model of a cascode stage.

input pole

$$\underline{\omega_{p,A}} = \frac{1}{R_S \left[C_{GS1} + \left(1 + \frac{g_{m1}}{g_{m2} + g_{mb2}} \right) C_{GD1} \right]}$$

$$R_X = r_o \parallel \frac{1}{g_{m2}} \approx \frac{1}{g_{m2}}$$

$$+ \frac{1}{A} \approx 2$$

Miller Cap of C_{GD1}

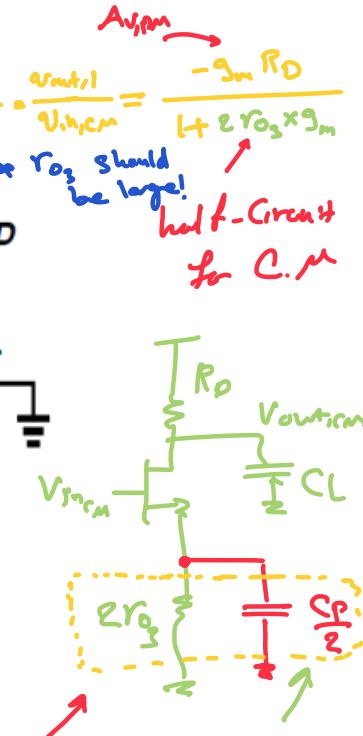
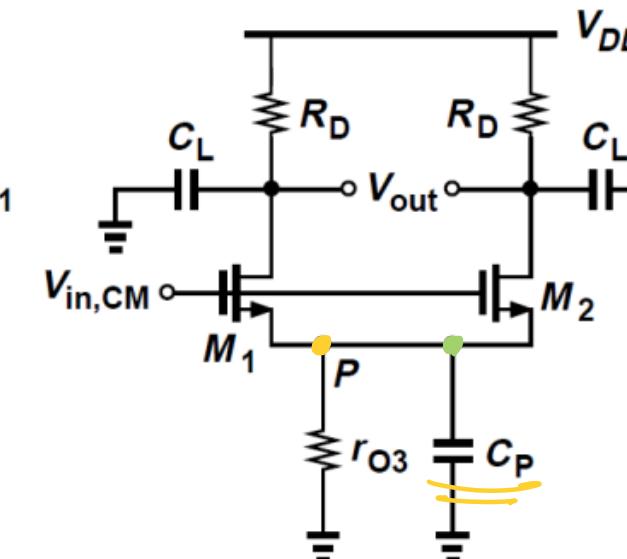
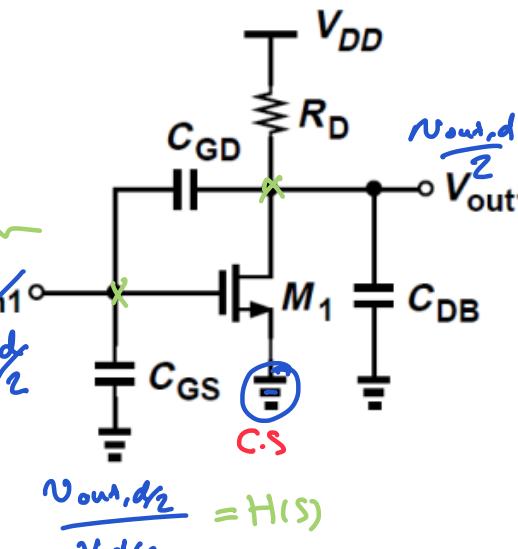
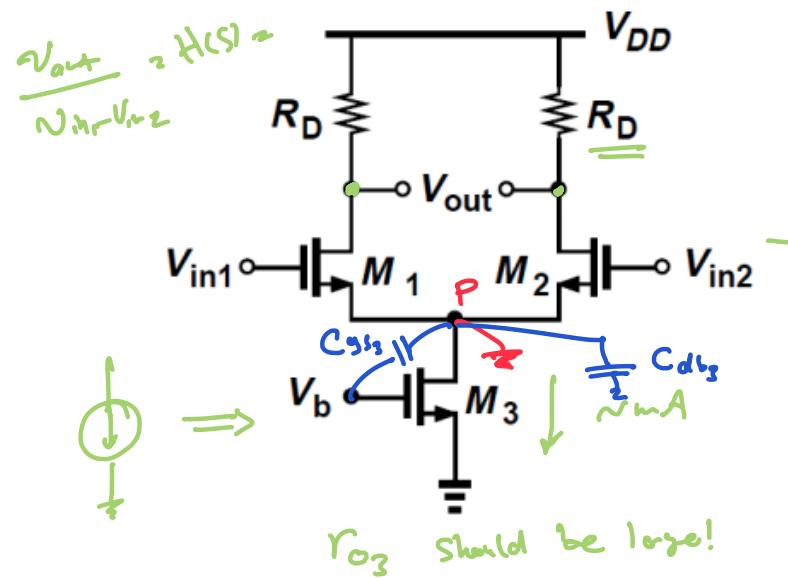
$$\underline{\omega_{p,X}} = \frac{1}{2C_{GD1} + C_{DB1} + C_{SB2} + C_{GS2}}$$

output pole

$$\underline{\omega_{pX}} \gg \underline{\omega_{pY}}, \underline{\omega_{pA}}$$

- Miller effect is less significant in cascode amplifiers than in common-source stages.
- But $\omega_{p,X}$ is typically quite higher than the other two.
- What if R_D is replaced by a current source?
 - Pole at node X may be quite lower, but transfer function will not affect much by this. See example.

Differential Pair



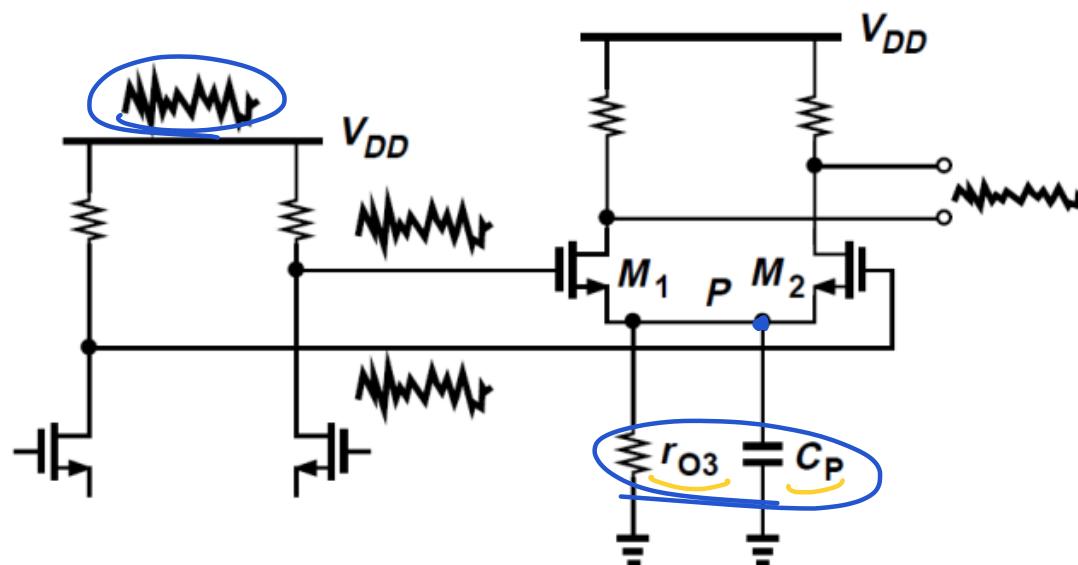
- For differential signals, the response is identical to that of a common-source stage.
- The common-mode rejection of the circuit degrades considerably at high frequencies.

$$A_{v,CM} = -\frac{\Delta g_m \left[R_D \parallel \left(\frac{1}{C_L s} \right) \right]}{(g_m + g_m) \left[r_{O3} \parallel \left(\frac{1}{C_{PS}} \right) \right] + 1}$$

Mismatch
 M_1 & M_2
 C_{Zss}

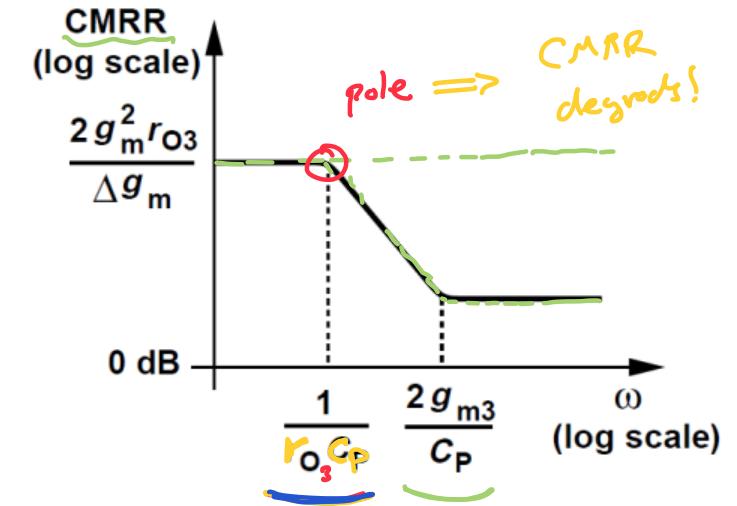
- Channel-length modulation, body effect, and other capacitances are neglected.

Differential Pair



$$\left| \frac{A_{om}}{A_{cm}} \right| = \text{CMRR} \approx \frac{g_m}{\Delta g_m} \left[1 + 2g_m \left(r_{O3} || \frac{1}{C_{PS}} \right) \right]$$

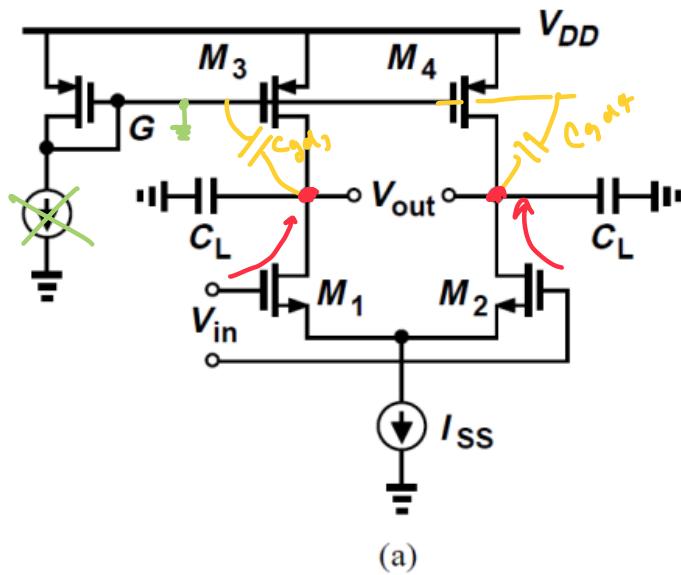
$$\approx \frac{g_m}{\Delta g_m} \frac{r_{O3}C_{PS} + 1 + 2g_m r_{O3}}{r_{O3}C_{PS} + 1}$$



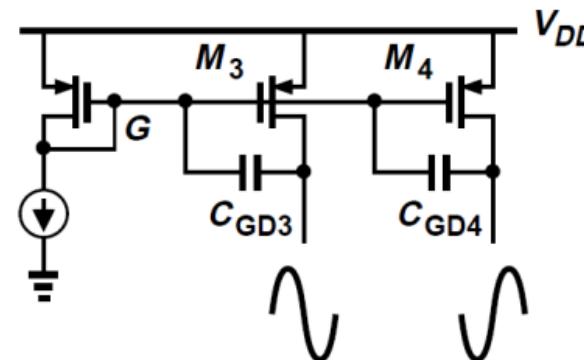
- This transfer function contains a zero and a pole.
- The magnitude of the zero is much greater than the pole.

↓
how large C_m3
be μ_3

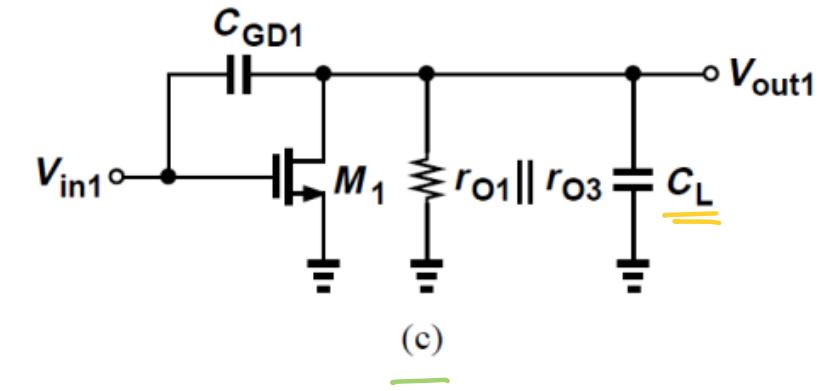
Differential Pair



(a)



(b)



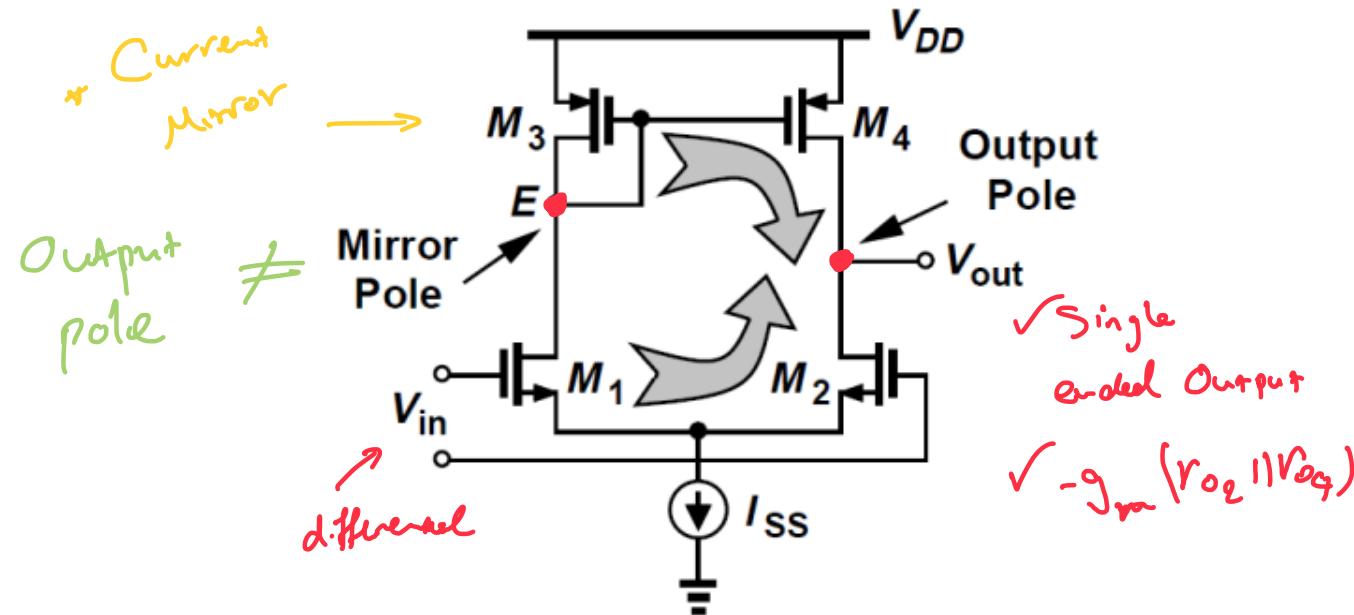
(c)

- Frequency response of differential pairs with high-impedance loads.
- Fig (b) C_{GD3} and C_{GD4} conduct equal and opposite currents to node G , making this node an ac ground.
- The differential half circuit is depicted in Fig. (c).
- More in chapter 10 ...

$$CMRR_{book} = \frac{A_{DM}}{A_{CM-DM}}$$

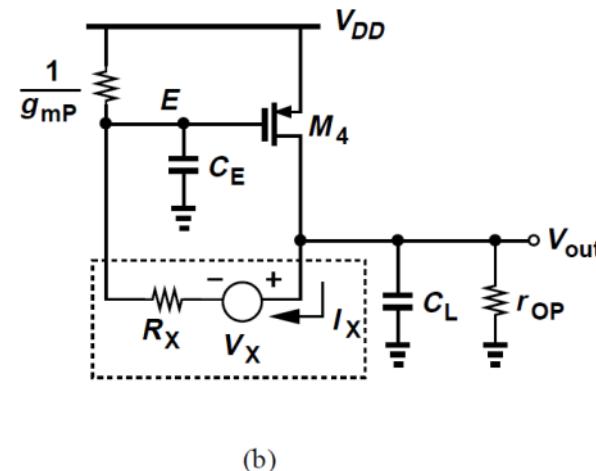
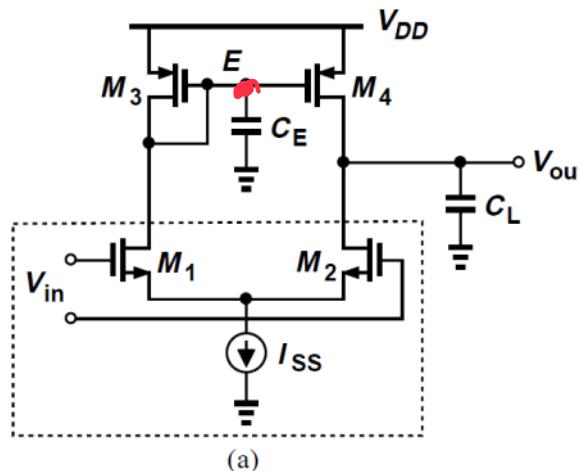
$$CMRR_2 = \frac{A_{DM}}{A_{CM}}$$

Single ended output



- How many poles does this circuit have?
- The severe trade-off between g_m and C_{GS} of PMOS devices results in a pole that impacts the performance of the circuit.
- The pole associated with node E is called a “mirror pole.”

Active Load



- Replacing Vin, M1, and M2 by a Thevenin equivalent (Direct calculation method)

$$V_X = g_{mN} r_{ON} V_{in} \quad R_X = 2r_{ON}$$

1 Zero

$$\frac{V_{out}}{V_{in}} = \frac{g_{mN} r_{ON} (2g_{mP} + C_E s) r_{OP}}{2r_{OP} r_{ON} C_E C_L s^2 + [(2r_{ON} + r_{OP}) C_E + r_{OP} (1 + 2g_{mP} r_{ON}) C_L] s + 2g_{mP} (r_{ON} + r_{OP})}$$

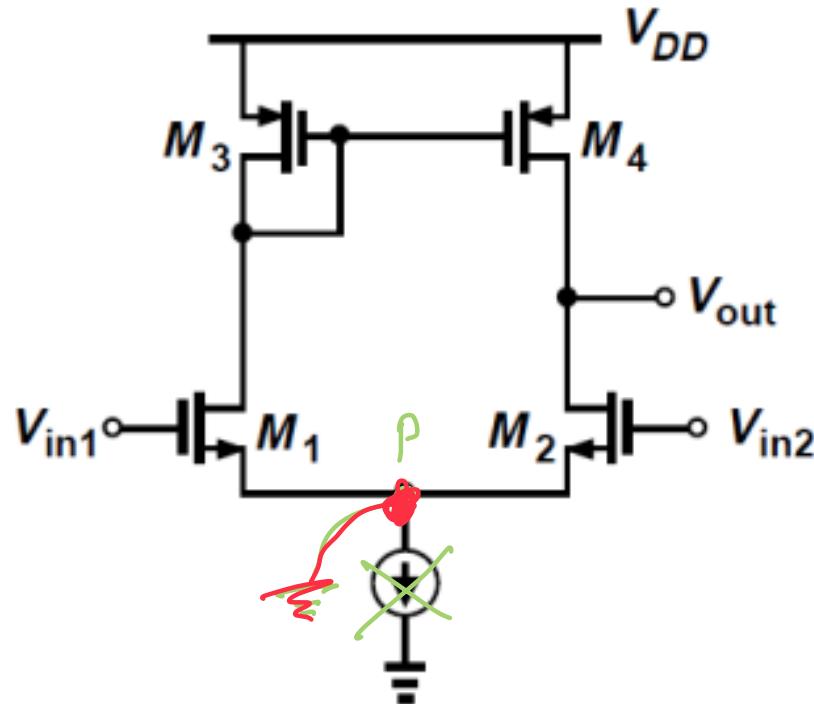
$$\omega_L = \frac{-2g_{mP}}{C_E} = 2\omega_{p2}$$

$$\omega_{p1} \approx \frac{1}{(r_{ON} \parallel r_{OP}) C_L}$$

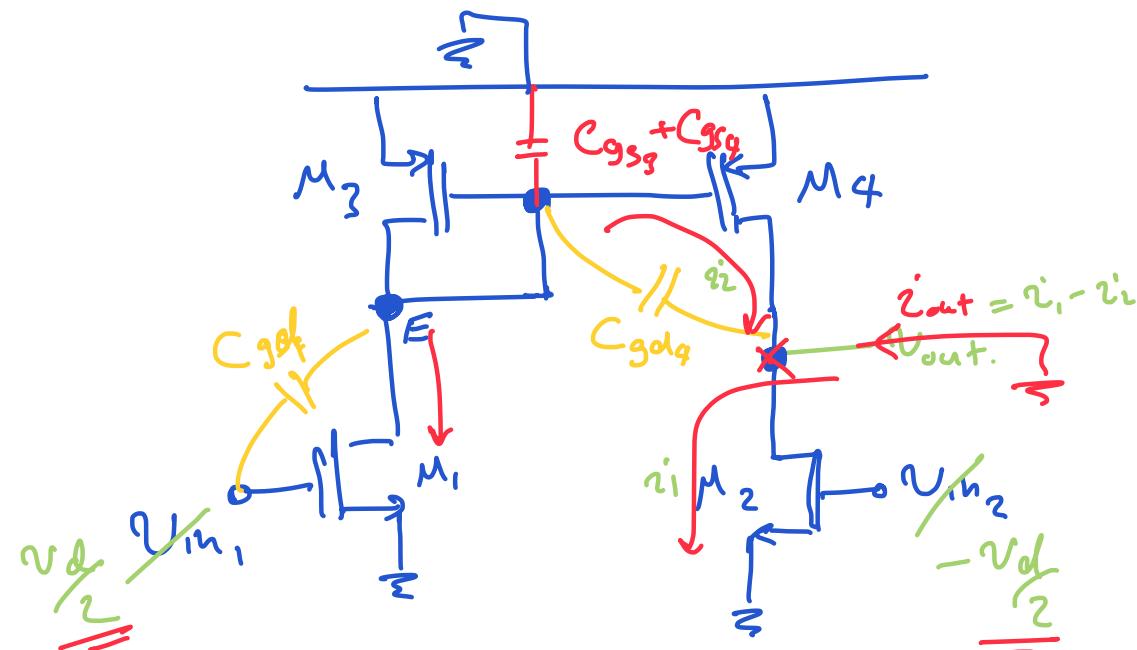
$$\omega_{p2} \approx \frac{g_{mP}}{C_E}$$

$\leftarrow R_E \approx \frac{1}{g_{mP}} \parallel r_{O1} \approx \frac{1}{g_{m3}}$

5T-OTA



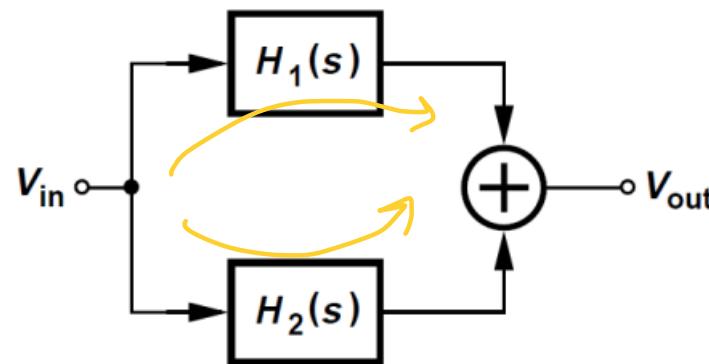
$$\left\{ \begin{array}{l} C_E = C_{GS_3} + C_{GS_4} + C_{dbs_1} + C_{dbs_3} + (1 + \frac{\omega}{\omega_0}) C_{gd4} + C_{gd_1} \\ C_{out} = C_{dbs_2} + C_{dbs_4} + (1 + \omega_1) C_{gd4} + C_{gd_2} \end{array} \right.$$



$$\frac{V_{out}}{V_E} = -g_{m4} \times (r_o 2 || r_o 4)$$

$$G_m(s) = \frac{-g_{m4}}{(1 + j\frac{\omega}{\omega_{pout}})} + \frac{-g_{m1}}{(1 + j\frac{\omega}{\omega_{pout}})(1 + j\frac{\omega}{\omega_E})}$$

Example



$$H(s) = H_1(s) + H_2(s)$$

- Can this (the zero) occur if $H_1(s)$ and $H_2(s)$ are first-order low-pass circuits?
- $H_1(s) = A_1/(1 + s/\omega_{p1})$ and $H_2(s) = A_2/(1 + s/\omega_{p2})$
- The overall transfer function contains a zero.

$$\frac{V_{out}}{V_{in}}(s) = \frac{\left(\frac{A_1}{\omega_{p2}} + \frac{A_2}{\omega_{p1}}\right)s + A_1 + A_2}{\left(1 + \frac{s}{\omega_{p1}}\right)\left(1 + \frac{s}{\omega_{p2}}\right)}$$

Active Load

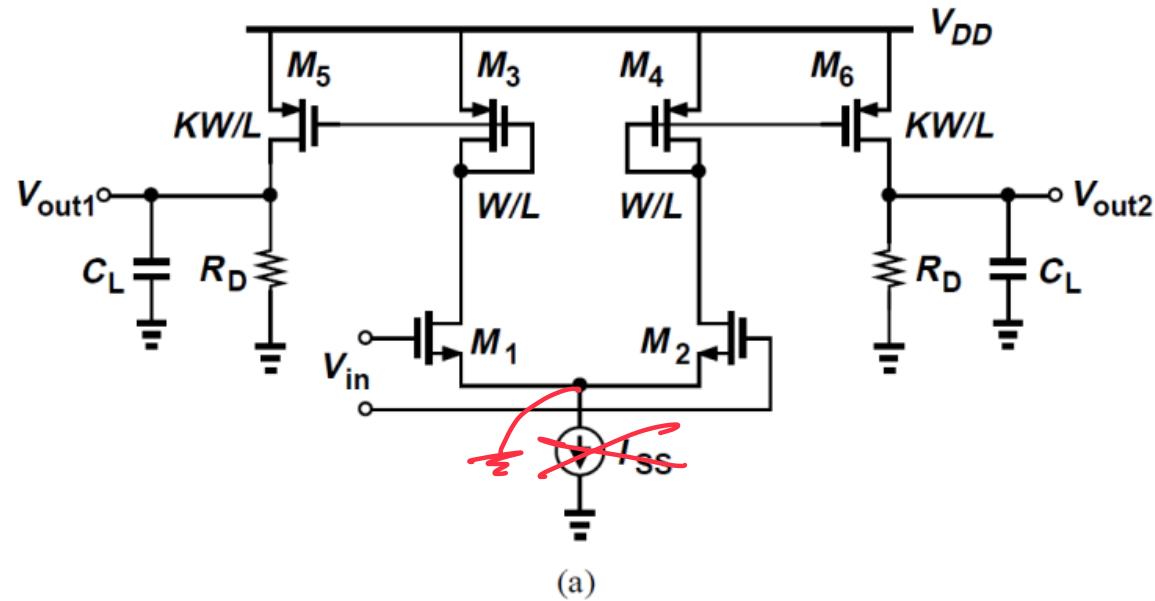
$$\begin{aligned}
 & \frac{1}{(1 + \frac{s}{\omega_{p_1}})} + \frac{1}{(1 + \frac{s}{\omega_{p_1}})(1 + \frac{s}{\omega_{p_2}})} \\
 = & \frac{2 + \frac{s}{\omega_{p_2}}}{(1 + \frac{s}{\omega_{p_1}})(1 + \frac{s}{\omega_{p_2}})} \quad \cancel{+} \quad 2 \left(1 + \frac{\cancel{s}}{2\omega_{p_2}} \right) \text{ mirror pole}
 \end{aligned}$$

$$\begin{aligned}
 \frac{V_{out}}{V_{in}} &= \frac{A_0}{1 + s/\omega_{p1}} \left(\frac{1}{1 + s/\omega_{p2}} + 1 \right) \\
 &= \frac{A_0(2 + s/\omega_{p2})}{(1 + s/\omega_{p1})(1 + s/\omega_{p2})}
 \end{aligned}$$

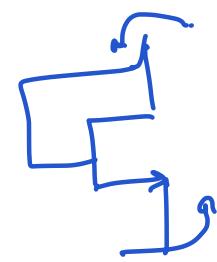
- A zero with a magnitude of $2g_m P / C_E$ in the left half plane.
- The appearance of this zero can be understood by noting that the circuit consists of:
- a “slow path” (M1,M3 and M4) $A_0/[(1 + s/\omega_{p1})(1 + s/\omega_{p2})]$ in parallel with
- a “fast path” (M1 and M2) by $A_0/(1 + s/\omega_{p1})$

Example

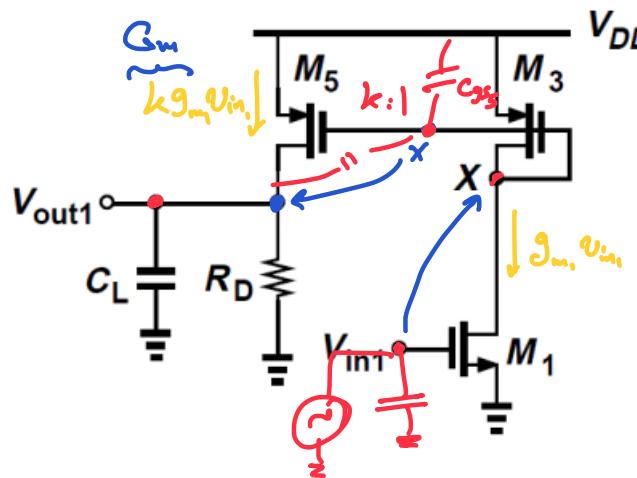
- Estimate the low-frequency gain and the transfer function of this circuit.



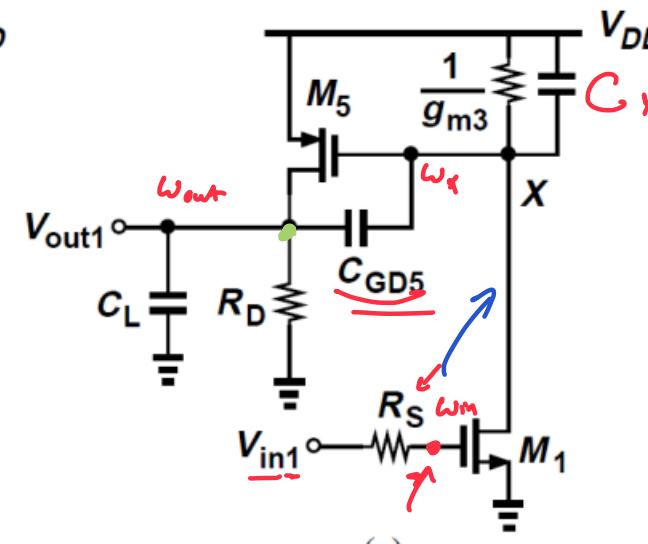
Example



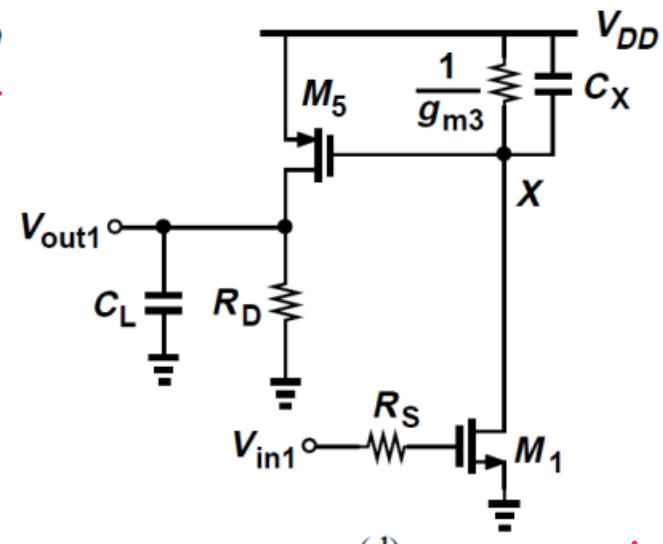
D.M half circuit



(b)



(c)



(d)

$$\omega = \frac{1}{RC} \quad \checkmark$$

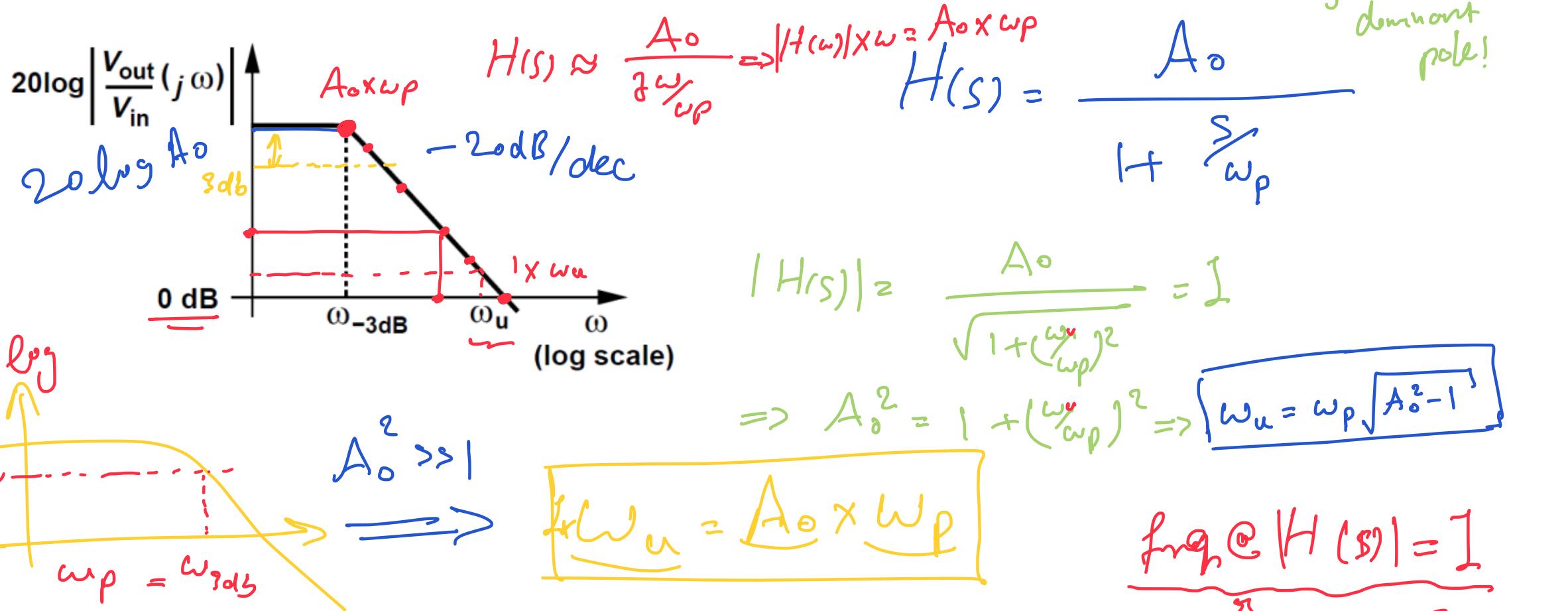
$$A_{\theta} = \frac{V_{out1}}{V_{in1}} = k \cdot g_{m1} (R_D || \omega_5)$$

$$A_{\theta} = \frac{V_X}{V_{in1}} \times \frac{V_{out1}}{V_X} = \dots$$

↑ CS

Poles	R_{eq}	C_{eq}
ω_m	R_S	$C_{gss_1 + C_{gd_1}} (1 + \frac{g_{m1}}{g_{m3}})$
ω_X	$\frac{1}{g_{m3} (R_D \omega_5)}$	$C_{gss_1 + C_{ss_5} + C_{db_1} + C_{db_3} + C_{gds_5} (1 + g_{m3} R_S)}$
ω_{out}	$R_D \omega_5$	$C_L + C_{db_5} + C_{gds_5} (1 + n)$

Gain-Bandwidth Trade-Offs

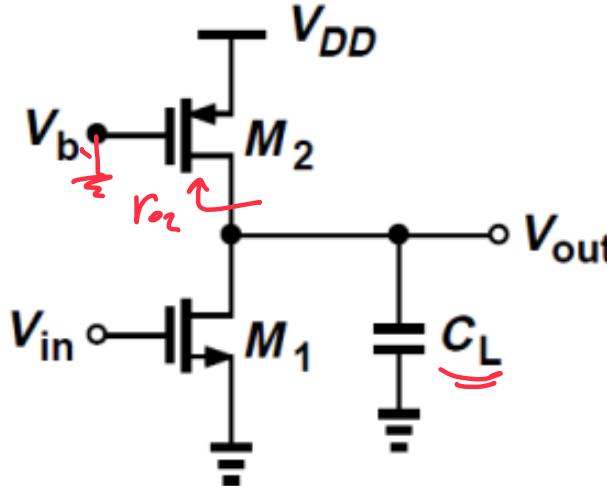


- We wish to maximize both the gain and the bandwidth of amplifiers.
- We are interested in both the 3-dB bandwidth, ω_{-3dB} , and the “unity-gain” bandwidth, ω_u

$$A_{v2} = G_m R_{out}$$

$$\omega_p = \frac{1}{R_{out} C_L}$$

One pole circuit



* ω_u is one of

first metrics to find!

↳ $G_m \rightarrow R_{out}, \dots$

$$\frac{2I}{V_{DD}}$$

$$\underline{\text{GBW}} = A_0 \omega_p$$

$$= g_{m1} (r_{o1} \| r_{o2}) \frac{1}{2\pi (r_{o1} \| r_{o2}) C_L}$$

$$G_m = \frac{g_{m1}}{2\pi C_L}$$

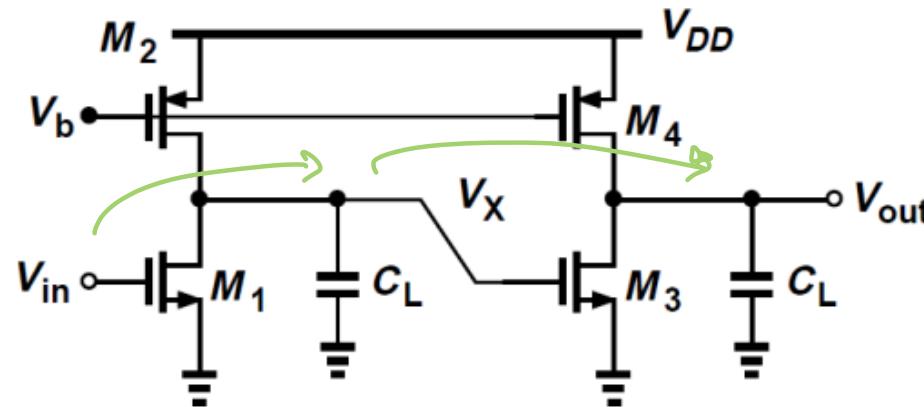
$$\omega_u = \sqrt{A_0^2 - 1} \omega_p$$

$$\approx A_0 \omega_p = \underline{\text{GBW}}$$

Multi-Pole Circuits

$$H_1(s) = H_2(s)$$

$$= \frac{A_0}{1 + \frac{s}{\omega_p}}$$



- It is possible to increase the GBW product by cascading two or more gain stages.
- Assume the two stages are identical and neglect other capacitances.

$$\frac{V_{out}}{V_{in}} = \frac{-A_0^2}{(1 + \frac{s}{\omega_p})^2}$$

$H_1 \times H_2 \rightarrow$

$$\begin{aligned} \omega_{-3dB} &= \sqrt{\sqrt{2} - 1} \omega_p \\ &\approx 0.64 \omega_p \end{aligned}$$

$$\text{GBW} = \sqrt{\sqrt{2} - 1} A_0^2 \omega_p$$

- While raising the GBW product, cascading reduces the bandwidth.