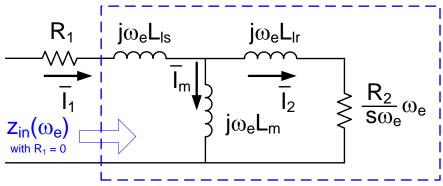
Constant Volts/Hz operation:



Induction Motor Per Phase Equivalent Circuit

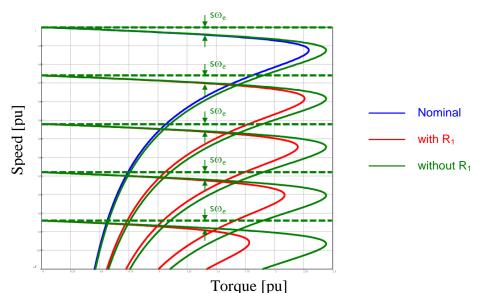
Neglecting R₁:

$$Z_{in}(\omega_e) = j\omega_e L_{ls} + j\omega_e L_m / \left(\frac{R_2}{s\omega_e} \omega_e + j\omega_e L_{lr} \right)$$

$$T_{em} = \frac{3P}{2} \frac{I_2^2 R_2}{S\omega_e}$$

Maintain constant V/f:

- Maintains circuit (machine) impedance
- Maintains stator current
- Maintains rotor current (and thus torque)
- Maintains magnetizing current (and thus flux)



Induction Motor Speed Torque Curves with Constant V/Hz

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V/f control is often referred to as 'scalar' control.

Scalar Control:

- Low speed torque limitations
- Low dynamic response (low acceleration capability)
- Load torque determines speed transiently → does not directly control torque

Vector Control (Field Oriented Control):

- Full torque at zero speed
- High dynamic response (high acceleration capability)

Pioneers in FOC:

K. Hasse (1969) Zur Dynamik drehzahlgeregelter Antriebe mit Stromrichtergespeisten Asynchronkurzschlufermotoren, Ph.D. dissertation, T.U. Darmstadt

F. Blaschke (1972) The principle of field orientation as applied to the new transvector closed loop control system for rotating field machines, Siemens Rev

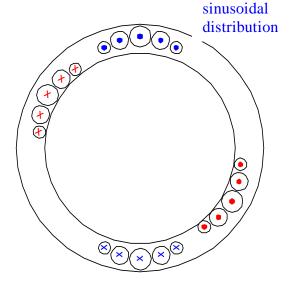
In order to understand this, we need a deeper understanding of the induction motor.

Coupled Circuit Modeling of AC Machines: Begin with uniform airgap idealized model Assumptions: 1) uniform airgap 2) ideal iron a. infinite permeability b. no core loss 3) no end effects 4) no slotting effects 4) no slotting effects To g (air gap) Note: assumptions 2,3, and 4 are handled by adding "correction terms" stator and rotor structure of induction machine

motgen/b7097 2.htm

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Winding Inductances:



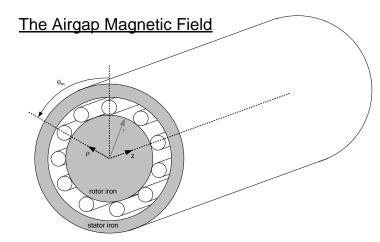
Winding Inductances:

$$\lambda_a \; = \; L_{aa} \, i_a + L_{ab} \, i_b + L_{ac} \, i_c + \dots$$

$$L_{aa} \; = \; \frac{\lambda_a}{i_a} \bigg|_{i_b \, = \, i_c \, = \, \ldots \, = \, 0} \label{eq:Laa}$$

$$L_{ab} \; = \; \frac{\lambda_a}{i_b} \bigg|_{i_a \, = \, i_c \, = \, \ldots \, = \, 0} \label{eq:Lab}$$

Sinusoidally Distributed Stator and Rotor Windings



Cylindrical Coordinate System chosen for simplicity

Coordinates: ρ , θ_m , z

z-axis chosen along axis of machine

with $\rho = r$ (rotor radius), θ_m defines a point in the air gap

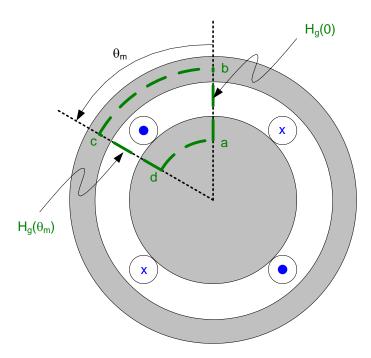
In general, $H = H(\rho, \theta_m, z)$ - Magnetic Field Intensity in air gap will be a function of, or in other words, have components in each of the three directions. Since we're concerned with finding the flux linking a winding placed upon the surface of the rotor or stator, only the component normal to the surface (radial component) is significant. We need to find the radial component of $H(\rho, \theta_m, z)$.

With assumption #3 (no end effects), magnetic field intensity is independent of the z coordinate. $H = H(\rho, \theta_m)$.

If r >> g (rotor radius much greater than air gap), variation with ρ will be negligibly small. Therefore, $H = H(\theta_m)$.

Ampere's Law:

$$\oint \overline{H} \cdot d\overline{\ell} = \iint \overline{J} \cdot d\overline{A} = \text{current enclosed}$$
any closed current surface density area



Integrate along abcda

Note: H_{bc} & H_{ad} must equal zero with assumption #2a (infinite permeability). Along air gap portions of path, MMF is simply air gap length (g) times air gap field intensity (H_g). Choose reference to be positive for H_g directed inward.

$$g[H_g(\theta)-H_g(0)] = \text{current enclosed} = \eta(\theta) \, i$$
 net number of conductors carrying positive current
$$\text{Radial component of air gap field intensity at } \theta_m$$

Note: $\eta(\theta)$ is dependent upon where you start

Guass' Law:

$$\oint _s \overline{B} \bullet d\overline{A} \; = \; \mu_o \oint _s \overline{H} \bullet d\overline{A} \; = \; 0$$

(total flux crossing air gap, or leaving rotor, is zero)

dA =

$$\oint_s \overline{B} \bullet d\overline{A} \ = \ \int_Z \int_{\theta_m} \, B_g(\theta_m) \, r \, d\theta_m \, d_z \ = \ r \, \ell \int_{\theta_m} \, B_g(\theta_m) \, d\theta_m \ = \ \mu_o \, r \, \ell \int_0^{2\pi} \, H_g(\theta_m) \, d\theta_m$$

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from above

$$H_g(\theta_m) \, = \, \frac{i}{g} \, \eta(\theta) + H_g(0)$$

$$\int_{0}^{2\pi} \left[\frac{i}{g} \eta(\theta_{m}) + H_{g}(0) \right] \! d\theta_{m} \; = \; 0 \; \; \boldsymbol{\rightarrow} \; \frac{i}{g} \int_{0}^{2\pi} \eta(\theta_{m}) d\theta_{m} + 2\pi H_{g}(0) \; = \; 0$$

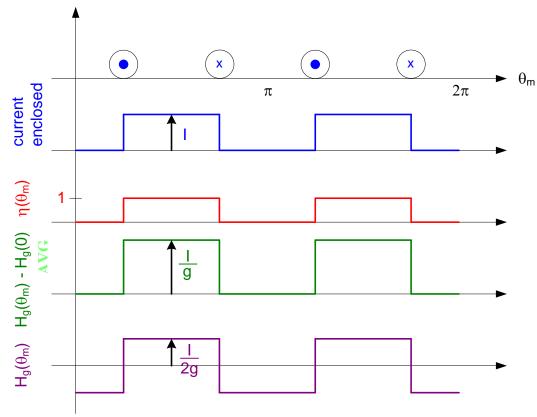
$$H_{g}(0) = -\frac{i}{g} \left\{ \frac{1}{2\pi} \int_{0}^{2\pi} \eta(\theta_{m}) d\theta_{m} \right\} = -\frac{i}{g} AVG \left\{ \eta(\theta_{m}) \right\}$$

therefore $H_g(\theta_m) =$

define
$$N(\theta_m) = \left[\eta(\theta_m) - AVG\left\{ \eta(\theta_m) \right\} \right]$$
 (called the winding function)

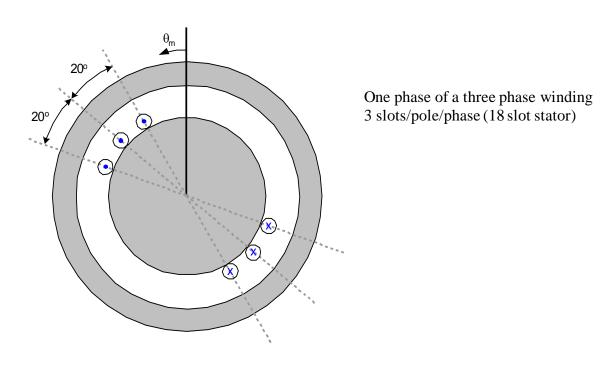
then

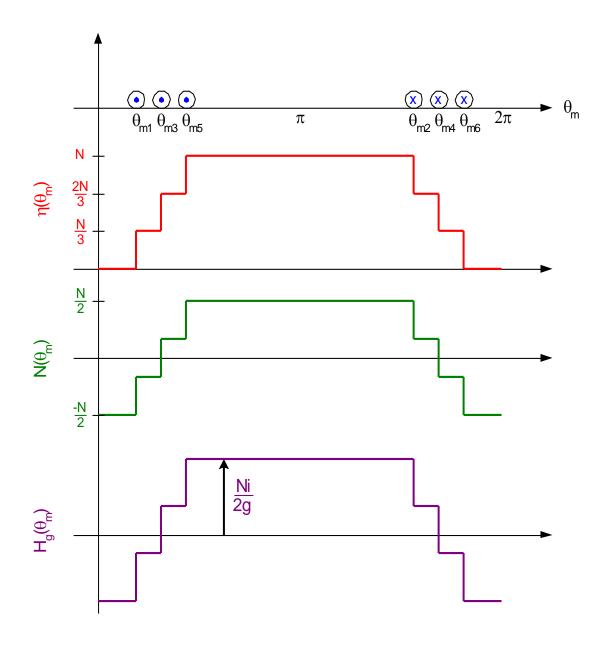
$$H_{g}(\theta_{m}) \stackrel{\Delta}{=} \frac{i}{g}N(\theta_{m})$$



Winding Function from Winding Distribution Above

Example:

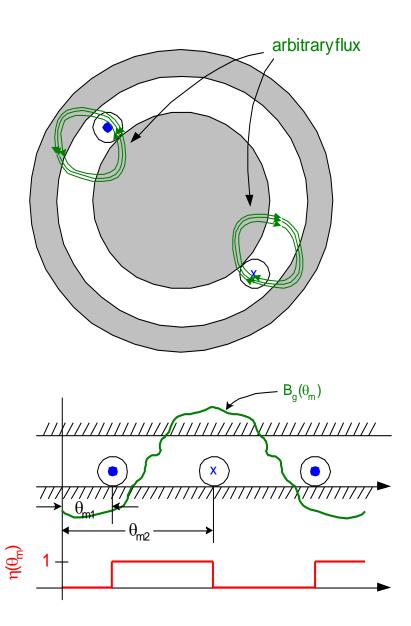




Peak value of $N(\theta_m)$:

$$N_{pk}(\theta_m) \, = \, \frac{total \, \# \, of \, turns}{\# \, of \, poles}$$

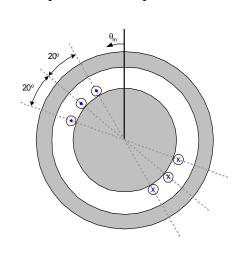
Flux Linkage and Inductance:

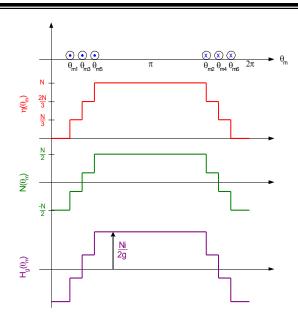


$$\lambda \; = \; r \, \mathit{l} \int_{\theta_1}^{\theta_2} B_g(\theta_m) \, d\theta_m \; = \;$$

$$\lambda \; = \; r \, \ell \int_0^{2\pi} \, \eta \! \left(\theta_m \right) B_g \! \left(\theta_m \right) d\theta_m$$
 Counting Function

from previous example:





 $\lambda =$

$$= r \, \mathcal{l} \, \left[\int_{\theta_1}^{\theta_3} B_g(\theta_m) \, d\theta_m + 2 \int_{\theta_3}^{\theta_5} B_g(\theta_m) \, d\theta_m + 3 \int_{\theta_5}^{\theta_2} B_g(\theta_m) \, d\theta_m + 2 \int_{\theta_2}^{\theta_4} B_g(\theta_m) \, d\theta_m + \int_{\theta_4}^{\theta_6} B_g(\theta_m) \, d\theta_m \right]$$

$$= r \, \mathcal{l} \int_{0}^{2\pi} \, \eta(\theta_m) \, B_g(\theta_m) \, d\theta_m$$

but this is dependent upon where we choose our reference

$$N(\theta_m) \, = \, \left[\, \eta(\theta_m) \, - \, AVG \left\{ \, \eta(\theta_m) \, \right\} \, \right] \, \, \boldsymbol{\rightarrow} \, \, \eta(\theta_m) \, = \, \left[\, N(\theta_m) \, + \, AVG \left\{ \, \eta(\theta_m) \, \right\} \, \right]$$

So

$$\begin{split} \lambda &= \, r \, \ell \int_0^{2\pi} \, \eta(\theta_m) \, B_g(\theta_m) \, d\theta_m \, = \, r \, \ell \, \int_0^{2\pi} \left[N(\theta_m) + AVG \left\{ \eta(\theta_m) \right\} \right] \, B_g(\theta_m) \, d\theta_m \\ &= \, r \, \ell \int_0^{2\pi} \, N(\theta_m) \, B_g(\theta_m) \, d\theta_m + r \, \ell \, AVG \left\{ \eta(\theta_m) \right\} \, \int_0^{2\pi} \, B_g(\theta_m) \, d\theta_m \end{split}$$

$$\boxed{ \lambda = r \, \ell \int_0^{2\pi} \, N(\theta_m) \, B_g(\theta_m) \, d\theta_m }$$

for the case of self inductance, $B_{\rm g}(\theta_{m})$ is the field produced by the winding itself:

$$B_{ga}(\theta_m) \; = \; \mu_o \, H_{ga}(\theta_m) \; = \; \frac{\mu_o \, i_a}{g} N_a(\theta_m) \label{eq:bga}$$

The self-inductance is therefore:

$$\left| L_a \right| = \frac{\lambda_{aa}}{i_a} = \frac{\mu_0 r \ell}{g} \int_0^{2\pi} \left[N_a(\theta_m) \right]^2 d\theta_m$$

for Mutual Inductance:

$$B_{ga}(\theta_m) \; = \; \frac{\mu_o \, i_b}{g} N_b(\theta_m)$$

$$\left| L_{ab} \right| = \frac{\lambda_{ab}}{i_b} = \frac{\mu_o \, r \, \ell}{g} \int_0^{2\pi} N_a(\theta_m) N_b(\theta_m) \, d\theta_m$$

Recall for machines with more than 2 poles to express quantities in electrical (rather than mechanical) radians, we have:

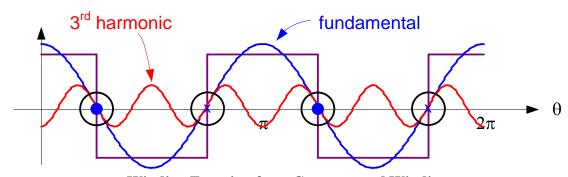
$$\theta \ = \ \frac{P}{2} \theta_m$$

 θ - electrical angle

 $\theta_{m}-mechanical \ angle$

P – number of motor poles

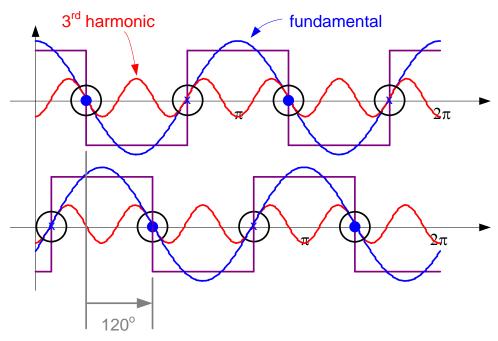
Concentrated Windings:



Winding Function for a Concentrated Winding

$$N(\theta) = \frac{4}{\pi} Np \left[\cos(\theta) - \frac{1}{3} \cos(3\theta) + \frac{1}{5} \cos(5\theta) \dots \right]$$

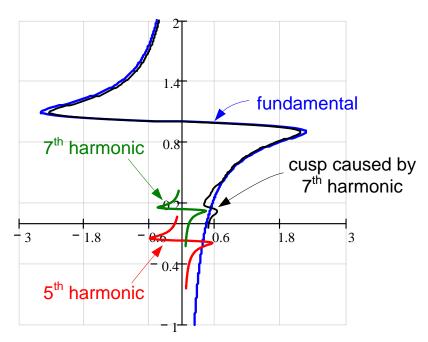
Elimination of 3rd Harmonic:



Winding Function for a Concentrated Winding of a 3 Phase Machine Showing 2 Phases

Displace turns by 120 (electrical) degrees to eliminate 3rd spatial harmonic

Side note on space harmonics with induction machines:



Speed Torque Curve of Induction Machine with Space Harmonics

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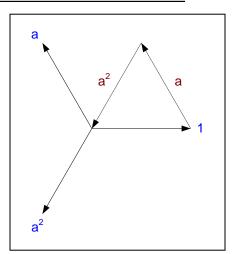
Typically

- a) No even harmonics (north pole is symmetric with south pole)
- b) Only sine or only cosine waves (1/4 wave symmetry)

$$N(\theta) \; = \; \sum_{h=1}^{\infty} N_h \, cos \big(h \, \theta + \phi_h \big)$$

For sinusoidal windings:

$$\begin{split} N(\theta) &= N_h \cos(h\,\theta + \phi_h) & \qquad h \frac{P}{2} \theta_m \\ L_h &= \frac{\mu_0 \, r \, \ell}{g} \int_0^{2\pi} \, N_h^2 \cos^2(h\,\theta + \phi_h) d\theta_m \end{split}$$



$$\boxed{L_h = \frac{\mu_o r \ell \pi}{g} \, N_h^2} \quad \textit{Self-Inductance}$$

$$\begin{split} L_{hk} &= \frac{\mu_o \, r \, \ell}{g} \int_0^{2\pi} N_h cos(h \, \theta + \phi_h) \, N_k cos(k \, \theta + \phi_k) d\theta_m \\ &= \frac{\mu_o \, r \, \ell \, N_h \, N_k}{g} \int_0^{2\pi} cos\! \left(h \frac{P}{2} \theta + \phi_h\right) \! cos\! \left(k \frac{P}{2} \theta + \phi_k\right) \! d\theta_m \\ &= 0 \, \, for \, h \neq k \end{split}$$

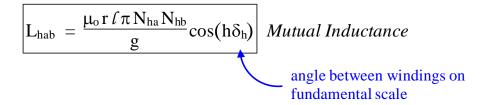
There is no magnetic interaction between sinusoidal windings with different number of poles.

Mutual Inductances

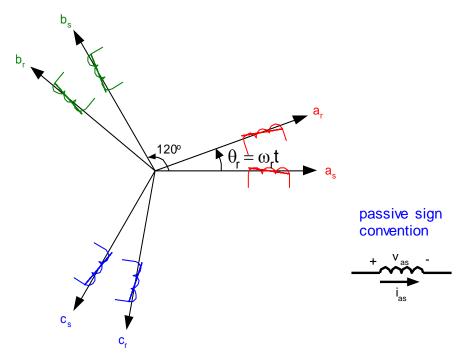
$$N_a(\theta) = N_{ha} \cos(h\theta + \phi_{ha})$$

$$N_b(\theta) = N_{hb} \cos(h\theta + \phi_{hb})$$

$$\begin{split} L_{hab} &= \frac{\mu_{o} \, r \, \ell}{g} \int_{0}^{2\pi} \, N_{ha} \cos(h \, \theta + \phi_{ha}) \, N_{hb} \cos(h \, \theta + \phi_{hb}) \, d\theta_{m} \\ &= \frac{\mu_{o} \, r \, \ell \, \pi \, N_{ha} \, N_{hb}}{g} \cos(\phi_{ha} - \phi_{hb}) \end{split}$$



3¢Machine Model:



Space Vectors of a Three Phase Machine

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- Assume sinusoidal winding distribution

$$N_a(\theta) = N_s \cos(\theta)$$

$$N_b(\theta) = N_s \cos(\theta - 120^\circ)$$

$$N_c(\theta) = N_s \cos(\theta + 120^\circ)$$

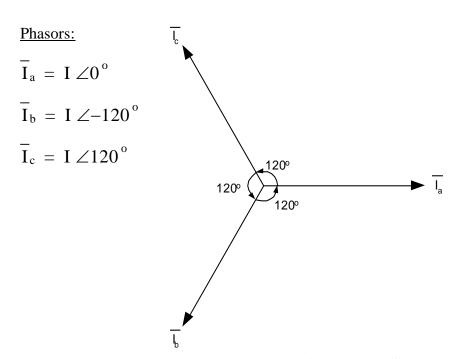
Note: Do not confuse phasors with space vectors (space vectors are spatial location of winding axis, phasors are temporal (time varying) vectors)

Example: for a time based system:

$$i_a(t) = \sqrt{2I}\cos(\omega t)$$

$$i_b(t) = \sqrt{2}I\cos(\omega t - 120^{\circ})$$

$$i_c(t) = \sqrt{2I}\cos(\omega t + 120^{\circ})$$



Phasor Representation of Temporal Signals

<u>Total flux linkage for a 3φ machine in 1φ:</u>

$$\lambda_{as} =$$

$$v_{as} = r_s i_{as} + p \lambda_{as} \rightarrow voltage equation$$

with sinusoidal windings,

Stator Self Inductance:

$$L_{as} \,=\, \frac{\mu_o\,r\ell\pi}{g}\,N_s^2 + L_{ls}$$

$$\frac{\mu_o \, r \ell \pi}{g} \, N_s^2 \, \stackrel{\Delta}{=} \, \, \text{Magnetizing Inductance} \, = \, L_{ms}$$

 $L_{ls} \stackrel{\Delta}{=} Leakage Inductance (accounts for end turn, slot, etc.)$

$$\left|L_{as}\right| = L_{ms} + L_{ls} = L_{bs} = L_{cs}$$

Stator Mutual Inductance (with other Stator Windings):

$$\left|L_{abs}\right| = \frac{-1}{2}L_{ms} = L_{acs} = L_{bcs}$$

Note: Negative sign denotes flux from other winding is opposite in direction of that produced by current in its own coil.

<u>Stator Mutual Inductance (with Rotor Windings):</u>

$$L_{asar} \, = \, \frac{\mu_o \, r \, \ell \, \pi}{g} \, N_s \, N_r cos(\theta_r) \, \boldsymbol{\rightarrow} \quad \text{from} \quad L_{hab} \, = \, \frac{\mu_o \, r \, \ell \, \pi \, N_{ha} \, N_{hb}}{g} cos(h \delta_h) \quad \text{with $h = 1$, $\delta = \theta_r$}$$

$$L_{asar} = \frac{N_r}{N_s} L_{ms} \cos(\theta_r)$$

$$\left| L_{asbr} = \frac{N_r}{N_s} L_{ms} \cos(\theta_r + 120^{\circ}) \right|$$

$$L_{ascr} = \frac{N_r}{N_s} L_{ms} \cos(\theta_r - 120^\circ)$$

Stator Voltage equation:

$$\begin{split} \lambda_{as} &= \ L_{as} \, i_{as} + L_{abs} \, i_{bs} + L_{acs} \, i_{cs} + L_{asar} \, i_{ar} + L_{asbr} \, i_{br} + L_{ascr} \, i_{cr} \\ v_{as} &= \ r_s \, i_{as} + p \, \lambda_{as} \\ &= \ r_s \, i_{as} + L_{ls} \, p \, i_{as} + L_{ms} \, p \bigg(i_{as} - \frac{1}{2} i_{bs} - \frac{1}{2} i_{cs} \bigg) + \\ &\qquad \qquad \frac{N_r}{N_s} \, L_{ms} \, p \bigg[i_{ar} cos(\theta_r) + i_{br} cos(\theta_r + 120^\circ) + i_{cr} cos(\theta_r - 120^\circ) \bigg] \end{split}$$

$$p(i_{ar}cos(\theta_r)) =$$

Observations:

- These equations describe a set of nonlinear (because of product of current, position, and speed), coupled set of differential equations with time varying coefficients!
- Will show this is an 8th order (with mechanical system) nonlinear, cross coupled set of differential equations with time varying coefficients!
- Special case: $\frac{d\theta_r}{dt} = constant = \omega_r$

$$\theta_r = \omega_r t + \theta_{ro}$$

equations are linear but still have time varying (periodic) coefficients