Exercise sheet 3

Kai Keller

January 8, 2017

Contents

| 1 | \mathbf{pro} | problem 1 - Weiss-Curie Ferromagnetism | | | | | | | | | | | | | | | 1 | | | | | | | | | |
|---|----------------|--|-----|--|--|--|--|--|--|--|--|--|--|--|--|--|---|--|--|--|--|--|--|--|--|---|
| | 1.1 | (a, 2P) |) . | | | | | | | | | | | | | | | | | | | | | | | 1 |
| | 1.2 | (a, 2P) |) . | | | | | | | | | | | | | | | | | | | | | | | 2 |

1 problem 1 - Weiss-Curie Ferromagnetism

Consider a System with the following Hamiltonian:

$$H(s) = -\frac{J}{2N} \sum_{i} \sum_{j} s_i s_j - h \sum_{i} s_i \tag{1}$$

1.1 (a, 2P)

show: That its Partition function can be written as:

$$Z(\beta, J, h) = \sqrt{\frac{N\beta J}{2\pi}} \int_{-\infty}^{+\infty} dx \, e^{-N\beta J x^2/2} \, C(x) \tag{2}$$

With:

$$C(x) = \prod_{i} \sum_{\{s_i\}} e^{\beta s_i (Jx+h)} \tag{3}$$

Let's consider first the general case of the summation done in eq. 3:

$$\Omega = \prod_{i} \sum_{s_{i[,j]}} f(s_{i[,j]}) \tag{4}$$

Where the index j = 1, ..., M runs over the members of the configuration s_i , i = 1, ..., N. We may expand the expression to:

$$\Omega = \prod_{i} (f(s_{i[,1]}) + \dots + f(s_{i[,M]}))$$
(5)

$$= (f(s_{1[,1]}) + \dots + f(s_{1[,M]})) \times \dots \times (f(s_{N[,1]}) + \dots + f(s_{N[,M]}))$$
 (6)

Which can be reordered to:

$$\Omega = (f(s_{1[,1]}) \dots f(s_{N[,1]})) + \dots + (f(s_{1[,M]}) \dots f(s_{N[,M]}))$$
(7)

$$= \prod_{i} f(s_{i[,1]}) + \dots + \prod_{i} f(s_{i[,M]})$$
 (8)

Where the last expression corresponds to a sum over all possible configurations. Hence we may write:

$$\Omega = \prod_{i} \sum_{s_{i[,j]}} f(s_{i[,j]}) = \sum_{\{\vec{s}\}} \prod_{i} f(s_{i[,j]})$$
(9)

Using this, we may write eq. 2 as:

$$Z(\beta, J, h) = \sum_{\{\vec{s}\}} \sqrt{\frac{N\beta J}{2\pi}} \int_{-\infty}^{+\infty} dx \, e^{-N\beta J x^2/2} \prod_{i} e^{\beta s_i (Jx + h)}$$
 (10)

$$= \sum_{I \in I} \sqrt{\frac{N\beta J}{2\pi}} \int_{-\infty}^{+\infty} dx \, e^{-N\beta J x^2/2} \, e^{\beta(\sum_i s_i J x + \sum_i s_i h)}$$
 (11)

$$= \sum_{\{\vec{s}\}} \sqrt{\frac{N\beta J}{2\pi}} e^{\beta \sum_{i} s_{i} h} \int_{-\infty}^{+\infty} dx \, e^{-N\beta J x^{2}/2} \, e^{\beta J \sum_{i} s_{i} x}$$
 (12)

We perform a variable ransform $y=\sqrt{N\beta J}x \implies dx=dy/\sqrt{N\beta J}$. Now we may write:

$$Z(\beta, J, h) = \sum_{\{\vec{s}\}} \sqrt{\frac{1}{2\pi}} e^{\beta \sum_{i} s_{i} h} \int_{-\infty}^{+\infty} dx \, e^{-y^{2}/2 - by}$$
 (13)

with:

$$b = -\frac{\beta J}{\sqrt{N\beta J}} \sum_{i} s_{i} \tag{14}$$

$$b^2 = \frac{\beta J}{N} \sum_{i} \sum_{j} s_i s_j \tag{15}$$

This is a general form of a gaussian integral with solution:

$$\int_{-\infty}^{+\infty} dx \, e^{-y^2/2 - by} = \sqrt{2\pi} \, e^{b^2/2} \tag{16}$$

This leads to:

$$Z(\beta, J, h) = \sum_{\{\vec{s}\}} e^{\beta \sum_{i} s_{i} h + \frac{\beta J}{N} \sum_{i} \sum_{j} s_{i} s_{j}}$$

$$\tag{17}$$

$$=\sum_{\{\vec{s}\}} e^{-\beta H(s)} \tag{18}$$

which is precisely the definition of the partition function.

1.2 (a, 2P)

Compute: C(x) explicitely.

We have only 2 possible values for each s_i . We hence have:

$$C(x) = \prod_{i} \left(e^{\beta(Jx+h)} + e^{-\beta(Jx+h)} \right) \tag{19}$$

$$= \prod_{i} 2\sin(\beta(Jx+h)) \tag{20}$$

$$= \sinh \left(\beta (Jx+h)\right)^N \tag{21}$$