

A Linear Response Framework for Simulating Bosonic and Fermionic Correlation Functions Illustrated on Quantum Computers

Alexander (Lex) Kemper



Department of Physics
North Carolina State University
<https://go.ncsu.edu/kemper-lab>

ORNL QC Users Forum
07/19/2023



What do you do with a quantum state once you've prepared one?

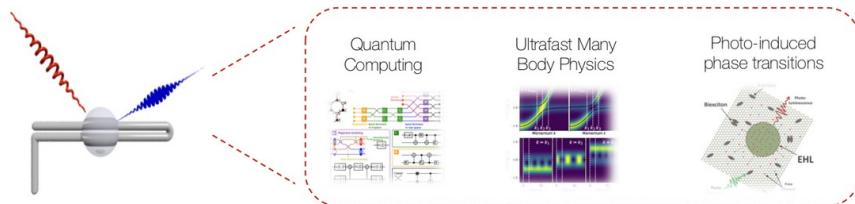
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Kemper Lab

Quantum materials in and out of equilibrium.

Collaborations with:

- Jim Freericks (Georgetown)
- Bert de Jong, Katie Klymko, Daan Camps, Roel van Beeumen, Akhil Francis (LBNL)
- Thomas Steckmann (UMD)

Current members



Alexander (Lex)
Kemper
Principal investigator



Efekan Kökcü
Graduate Researcher



Anjali Agrawal
Graduate Researcher



Heba Labib
Graduate Researcher



Jack Howard
Undergraduate
Researcher



Natalia Wilson
Undergraduate
Researcher



Daniel Brandon
Undergraduate
Researcher



Sarah Klas
Undergraduate
Researcher



Norman Hogan
Graduate Researcher



Ethan Blair
Undergraduate
Researcher

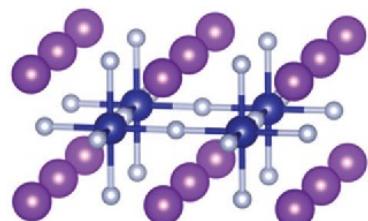


Your Name
New lab member

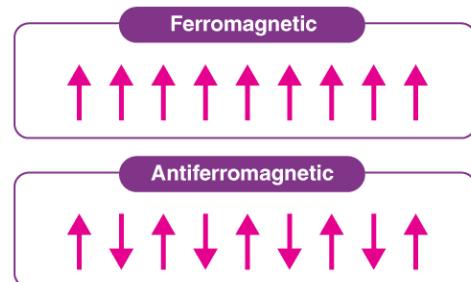
A Tale of Two Transitions

Ising Magnet

Rb_2CoF_4

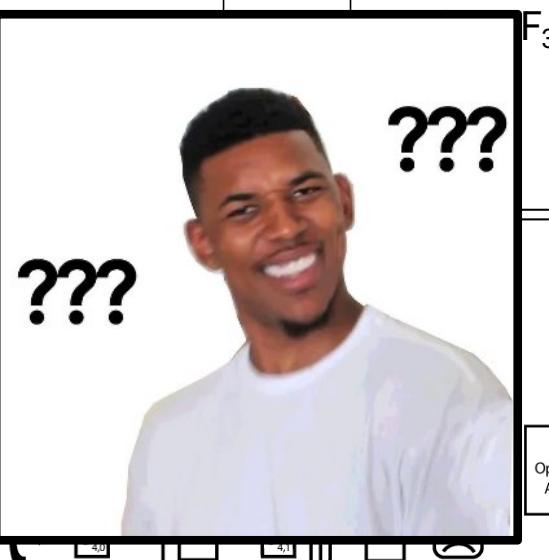
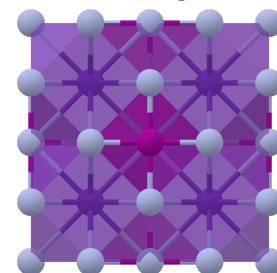


$$\mathcal{H} = -J \sum_i \sigma_i^z \sigma_{i+1}^z + h_x \sum_i$$



[10.1039/c6cp02362b](https://doi.org/10.1039/c6cp02362b)

Heisenberg Magnet



[Optimization of the Variational Quantum Eigensolver for Quantum Chemistry Applications](#)

Materials project

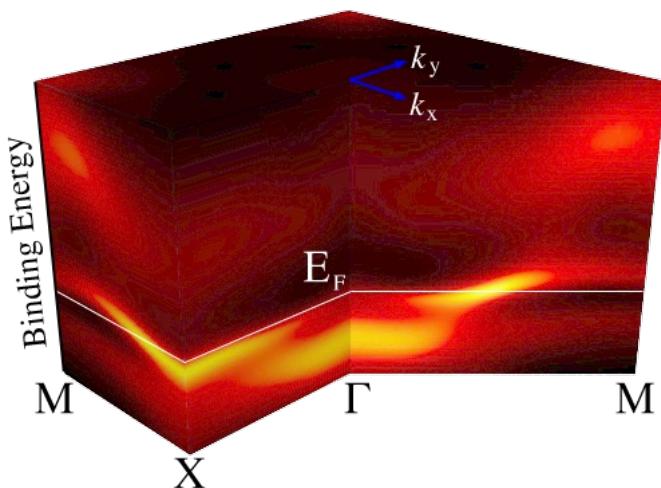
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Q: What do you do with a quantum state once you've prepared one?

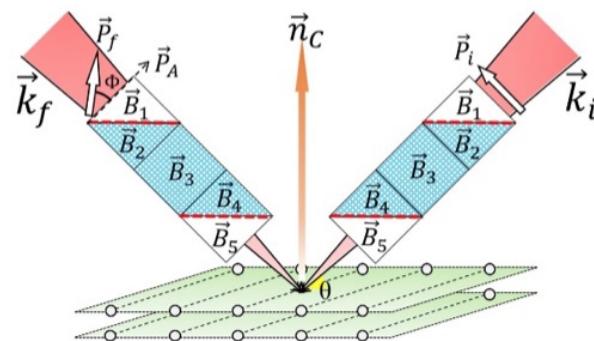
A: You measure its excitations.

Measuring Excitations

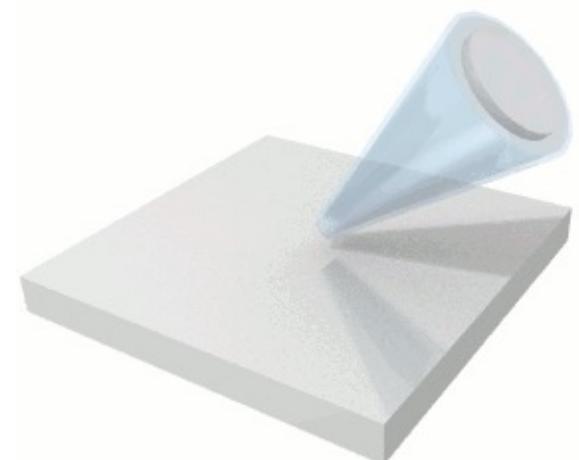
Figures courtesy of
Devereaux/Shen group
and ORNL



Angle-resolved Photoemission
(ARPES)



Neutron Scattering

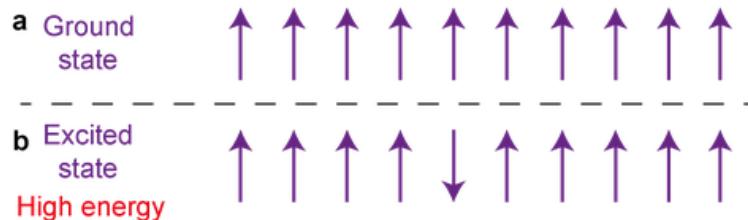


Time-resolved ARPES

Measuring Excitations

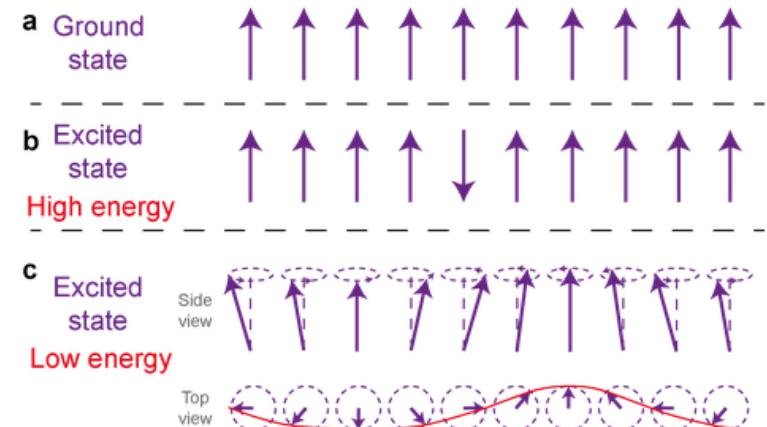
Ising Model

$$\mathcal{H} = -J \sum_i \sigma_i^z \sigma_{i+1}^z + h_x \sum_i \sigma_i^x$$

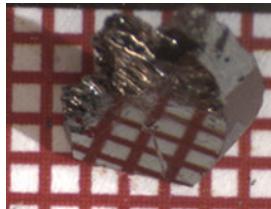


Heisenberg model

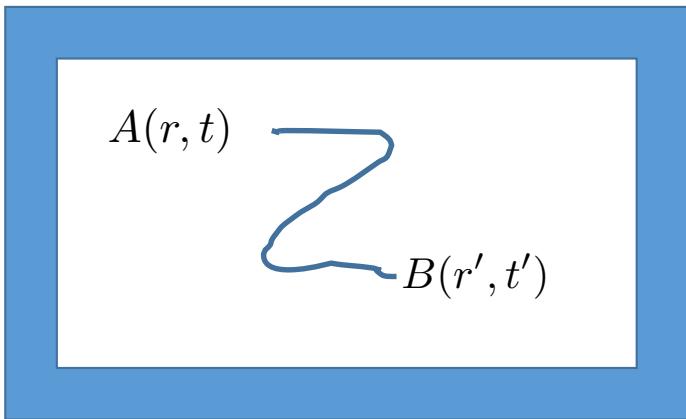
$$\mathcal{H} = -J \sum_i \vec{\sigma}_i \cdot \vec{\sigma}_{i+1} + h_x \sum_i \sigma_i^x$$



Correlation functions



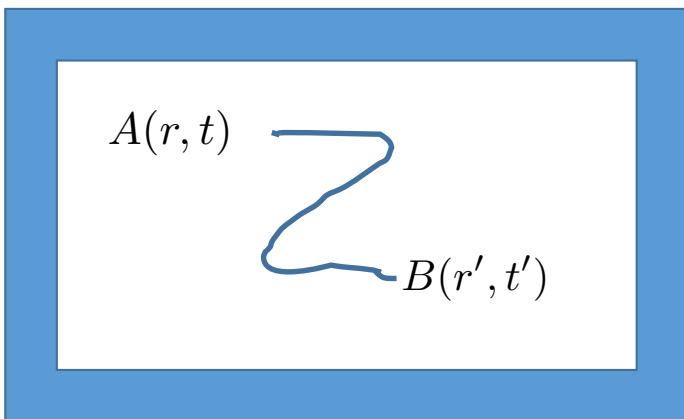
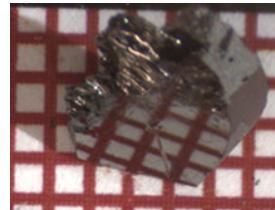
$$\langle A(r, t)B(r', t') \rangle$$



Given some (observable) operator B at (r', t') , what is the likelihood of some (observable) operator A at (r, t) ?

Optical conductivity, γ /X-ray scattering, photoemission, neutron scattering, Raman, IR absorption, etc.

Correlation functions



$e^{iE_0 t} \langle \phi_0 | B e^{-i\mathcal{H}t} A | \phi_0 \rangle$

Interfere with ground state

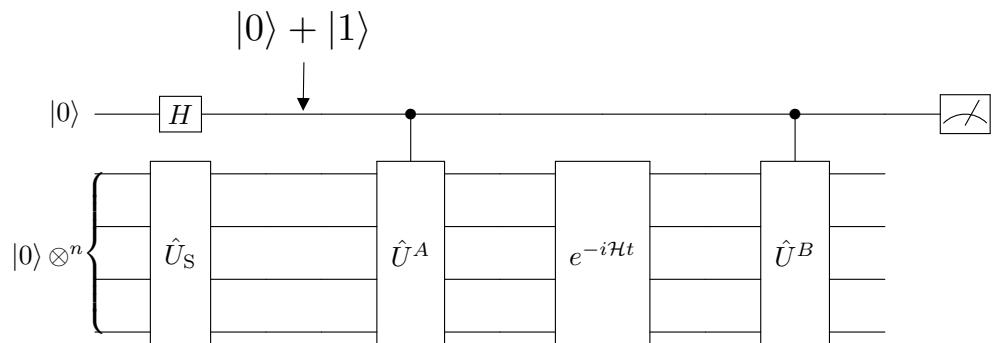
Complete expectation value

Time evolve

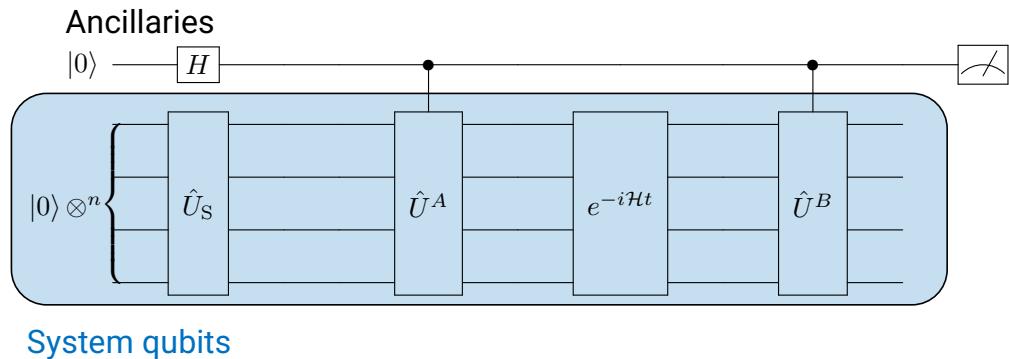
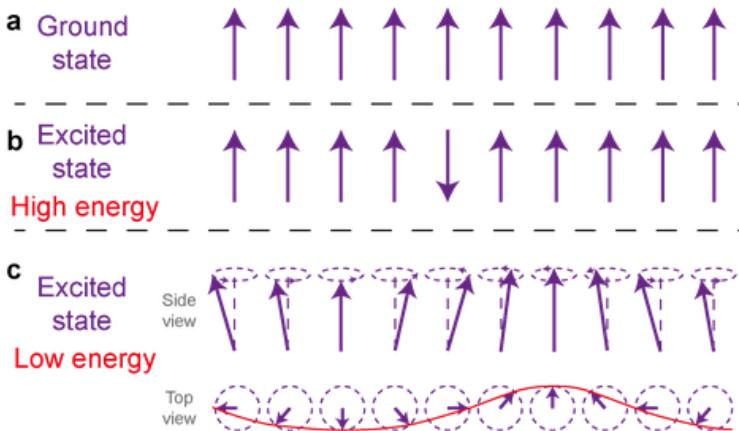
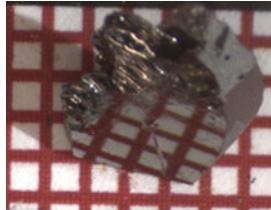
Apply excitation B

Apply excitation A

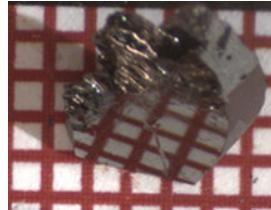
Prepare state of interest



Correlation functions



Correlation functions

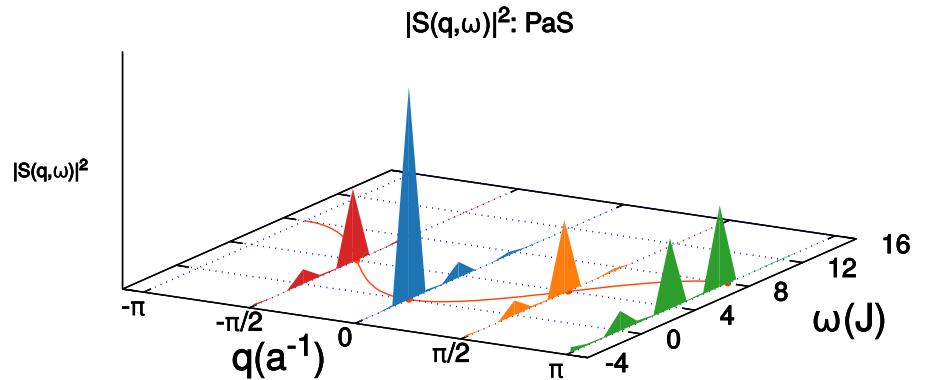
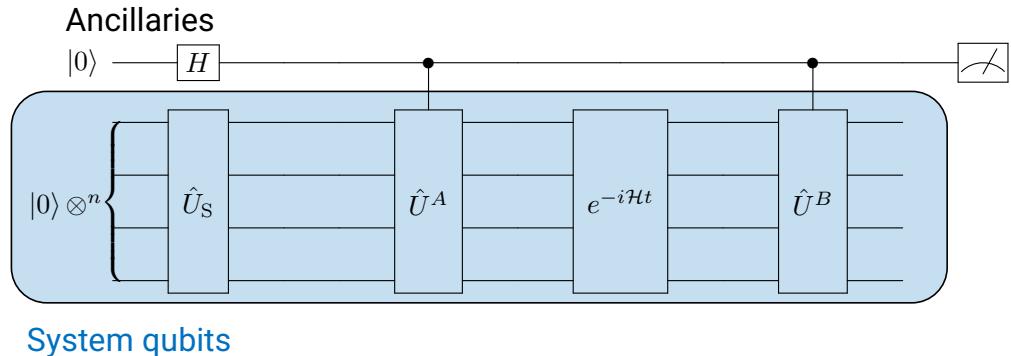


This works!

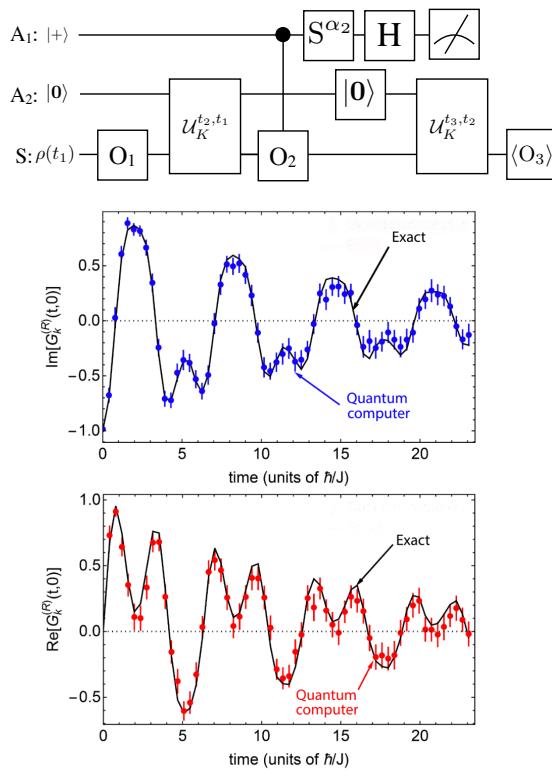
But:

- Need an ancilla with long coherence
- A and B need to be unitary & controlled
- More complex A,B need post-processing

$$\langle A(r, t)B(r', t') \rangle$$



(A few) Quantum Algorithm(s) for correlation functions



(Anti-)Commutators, dissipative

L. Del Re, B. Rost, M. Foss-Feig, AFK, J.K. Freericks
2204.12400

PHYSICAL REVIEW A 96, 022127 (2017)

Noninvasive measurement of dynamic correlation functions

Philipp Uhrich,^{1,2} Salvatore Castrignano,³ Hermann Uys,^{2,4} and Michael Kastner^{1,2,*}

¹National Institute for Theoretical Physics (NITheP), Stellenbosch 7600, South Africa

²Institute of Theoretical Physics, Department of Physics, University of Stellenbosch, Stellenbosch 7600, South Africa

³Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany

⁴Council for Scientific and Industrial Research, National Laser Centre, Pretoria, Brummeria, 0184, South Africa

(Received 24 November 2016; revised manuscript received 16 January 2017; published 21 August 2017)

1. Initial state preparation
2. Time evolution until time t_1
3. Weak coupling of ancilla and system site i .
4. Measuring the ancilla
5. Time evolution until time t_2
6. Projective measurement at site j
7. Correlating the measured outcomes

10.1103/PhysRevA.96.022127

Anti-commutators

PRL 111, 147205 (2013)

PHYSICAL REVIEW LETTERS

week ending
4 OCTOBER 2013

Probing Real-Space and Time-Resolved Correlation Functions with Many-Body Ramsey Interferometry

Michael Knap,^{1,2,*} Adrian Kantian,³ Thierry Giamarchi,³ Immanuel Bloch,^{4,5} Mikhail D. Lukin,¹ and Eugene Demler¹

¹Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

²JTAMP, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138, USA

³DPMC-MaNEP, University of Geneva, 24 Quai Ernest-Ansermet CH-1211 Geneva, Switzerland

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(Received 2 July 2013; revised manuscript received 18 September 2013; published 4 October 2013)

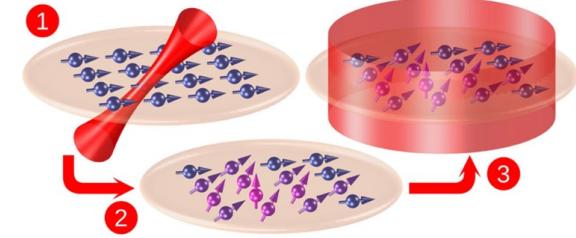
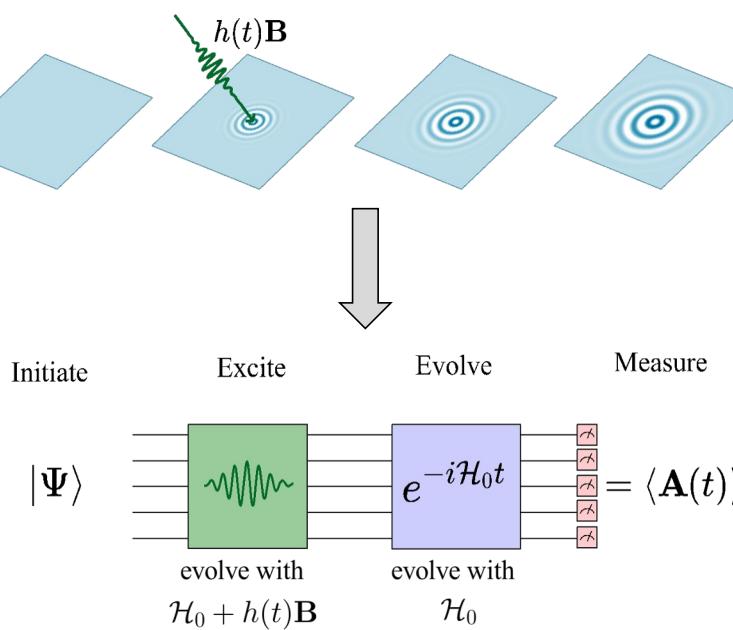
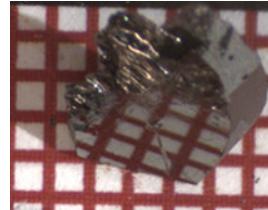


FIG. 1 (color online). Many-body Ramsey interferometry consists of the following steps: (1) A spin system prepared in its ground state is locally excited by $\pi/2$ rotation; (2) the system evolves in time; (3) a global $\pi/2$ rotation is applied, followed by the measurement of the spin state. This protocol provides the dynamic many-body Green's function.

Commutators

10.1103/PhysRevLett.111.147205

Linear Response



A linear response framework for simulating bosonic and fermionic correlation functions illustrated on quantum computers

Efekan Kökü¹, Heba A. Labib¹, J. K. Freericks² and A. F. Kemper^{1,*}

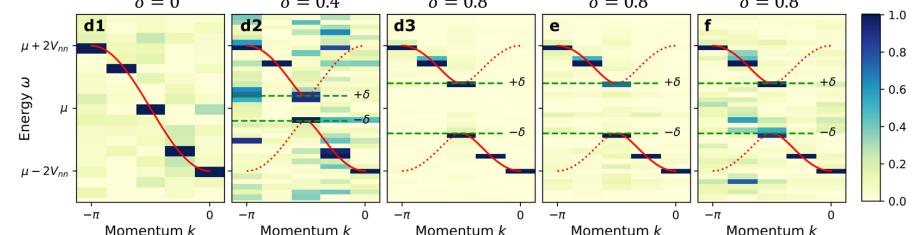
¹Department of Physics, North Carolina State University, Raleigh, North Carolina 27695, USA

²Department of Physics, Georgetown University, 37th and O Sts. NW, Washington, DC 20057 USA

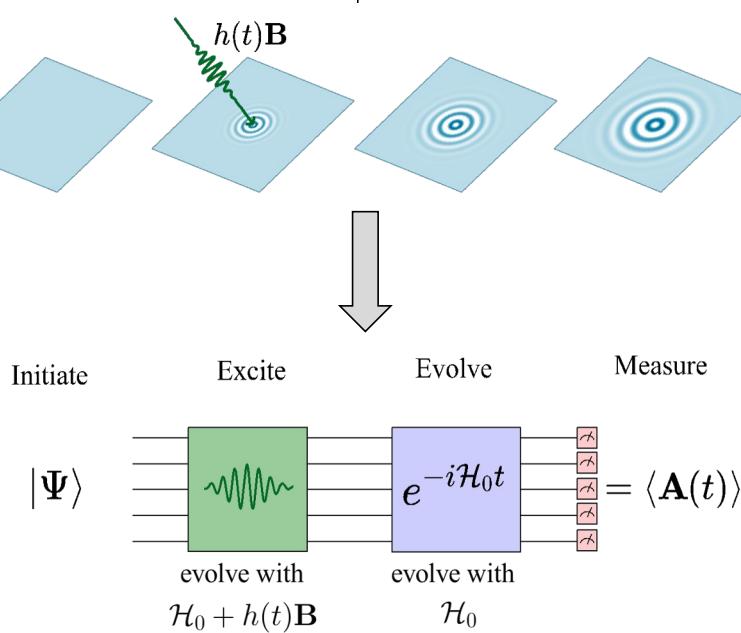
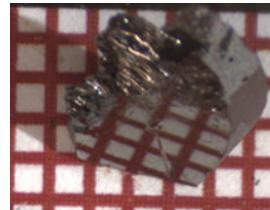
(Dated: February 22, 2023)

1. Make the excitation part of the quantum simulation
2. Post-process the data to get the response functions

$$\left. \frac{\delta A(t)}{\delta h(t')} \right|_{h=0} = -i\theta(t-t') \langle \psi_0 | [\mathbf{A}(t), \mathbf{B}(t')] | \psi_0 \rangle$$



Linear Response



A linear response framework for simulating bosonic and fermionic correlation functions illustrated on quantum computers

Efekan Kökcü ,¹ Heba A. Labib ,¹ J. K. Freericks ,² and A. F. Kemper ,^{1,*}

¹Department of Physics, North Carolina State University, Raleigh, North Carolina 27695, USA

²Department of Physics, Georgetown University, 37th and O Sts. NW, Washington, DC 20057 USA

(Dated: February 22, 2023)

Benefits

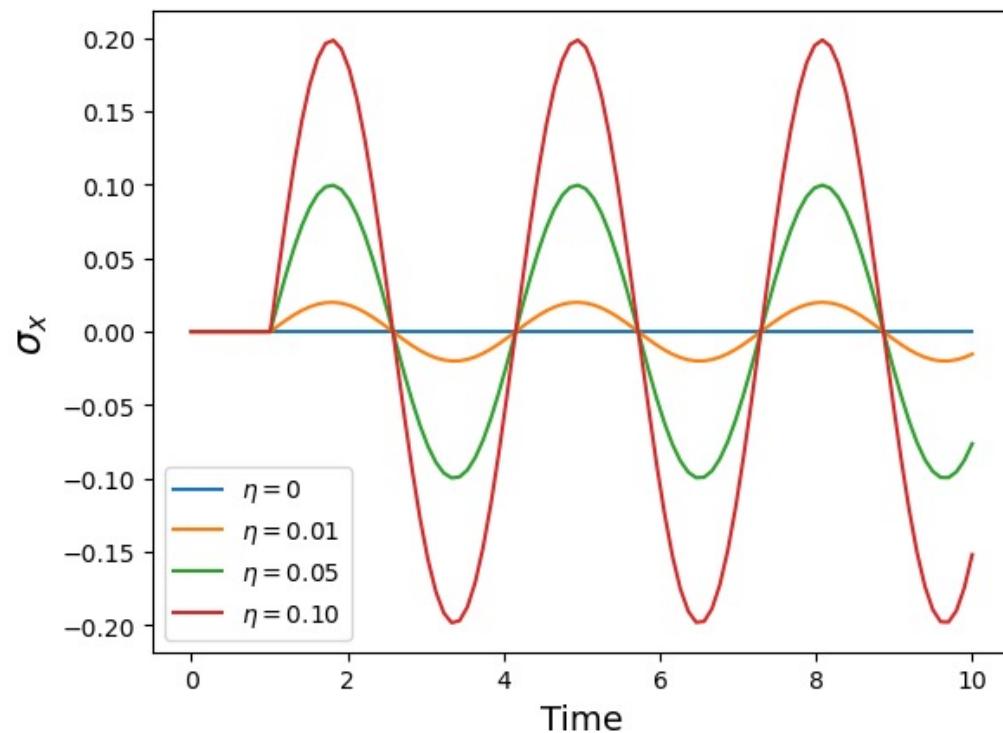
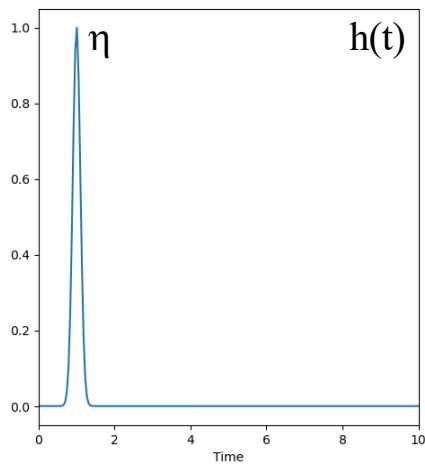
- Any operator A,B you desire (as long as it is Hermitian*)
- No ancillas/controlled operations needed
- Many correlation functions at the same time
- Less post-processing (less noise)
- Frequency/momentum selective

Linear Response

A simple example: single spin with energy level difference = 2

$$\mathbf{H}_0 = \sigma^z$$

$$\mathbf{A} = \mathbf{B} = \sigma^x$$

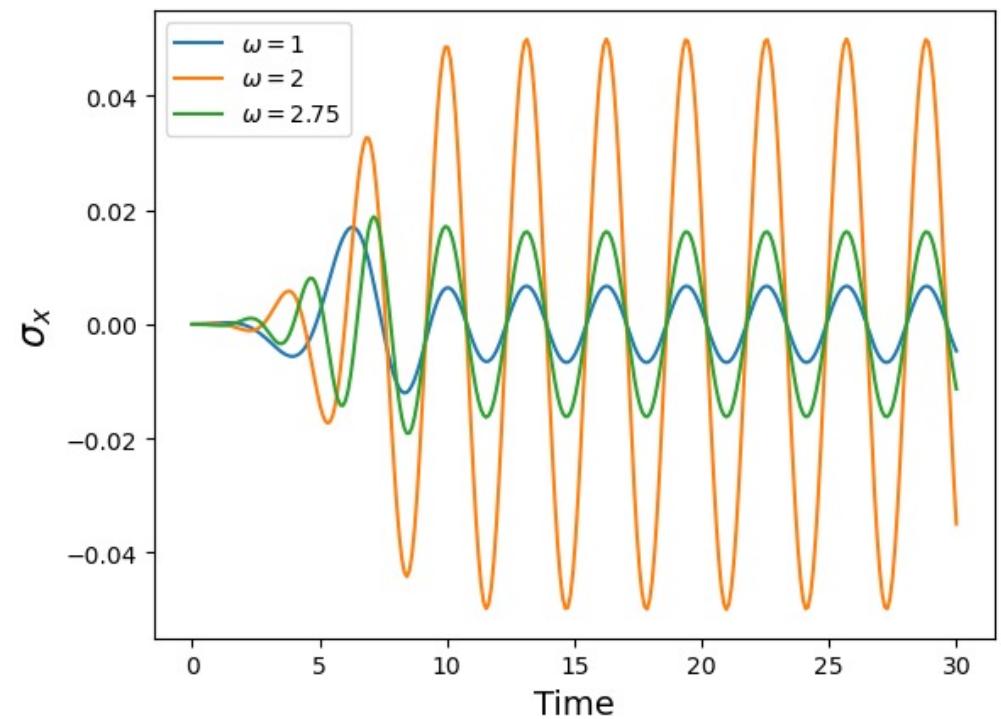
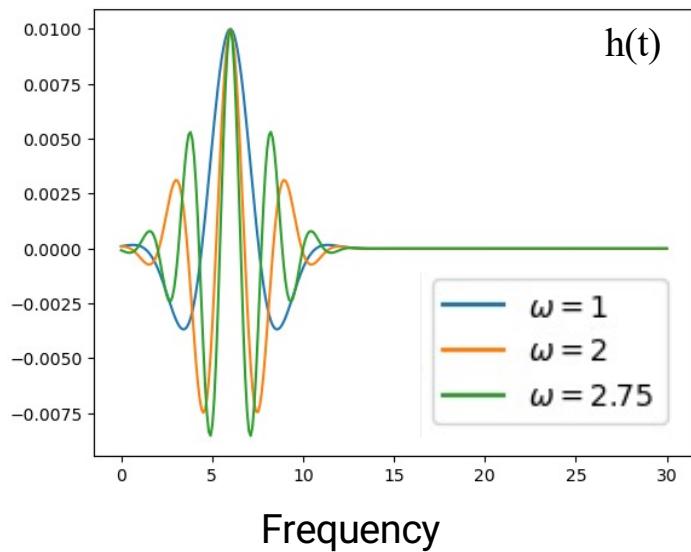


Linear Response

A simple example: single spin with energy level difference = 2

$$\mathbf{H}_0 = \sigma^z$$

$$\mathbf{A} = \mathbf{B} = \sigma^x$$

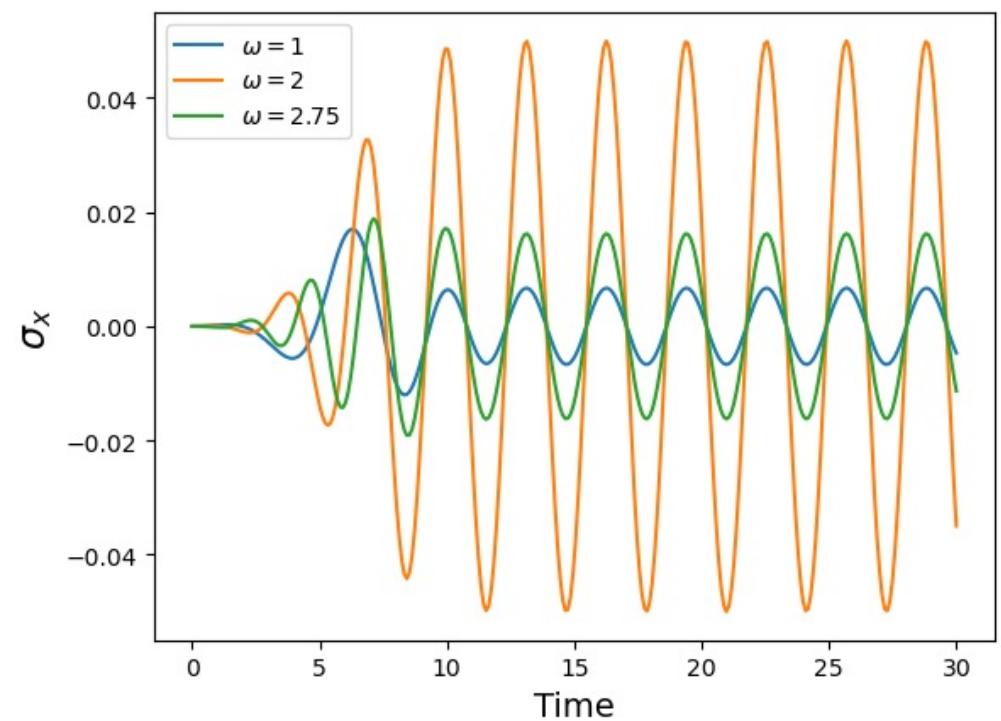
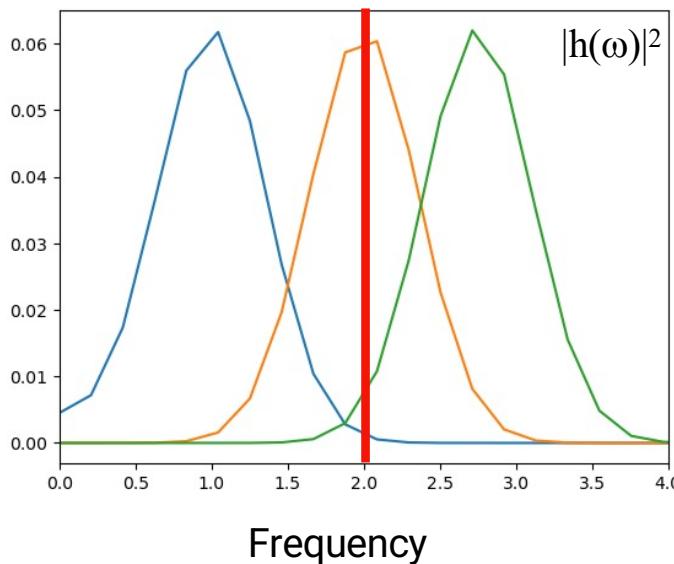


Linear Response

A simple example: single spin with energy level difference = 2

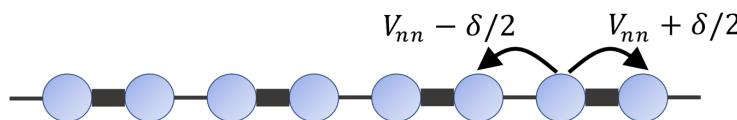
$$\mathbf{H}_0 = \sigma^z$$

$$\mathbf{A} = \mathbf{B} = \sigma^x$$

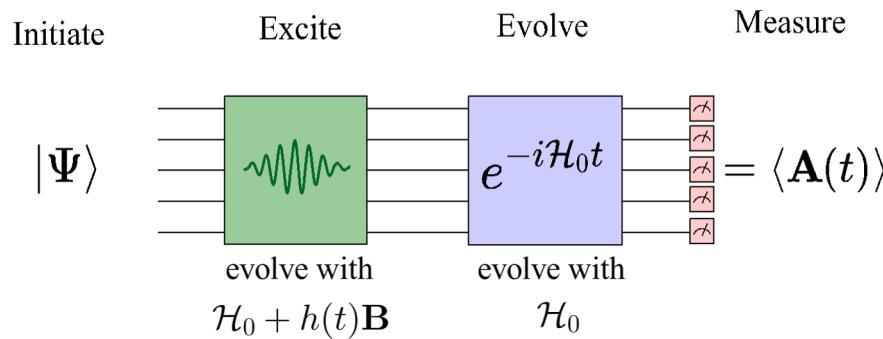


A Bosonic Correlation function: Polarizability

Su-Schrieffer-Heeger model for polyacetylene



$$\mathcal{H}_0 = - \sum_{\langle i,j \rangle} \left[V_{nn} + (-1)^i \delta/2 \right] c_i^\dagger c_j - \mu \sum_i c_i^\dagger c_i$$

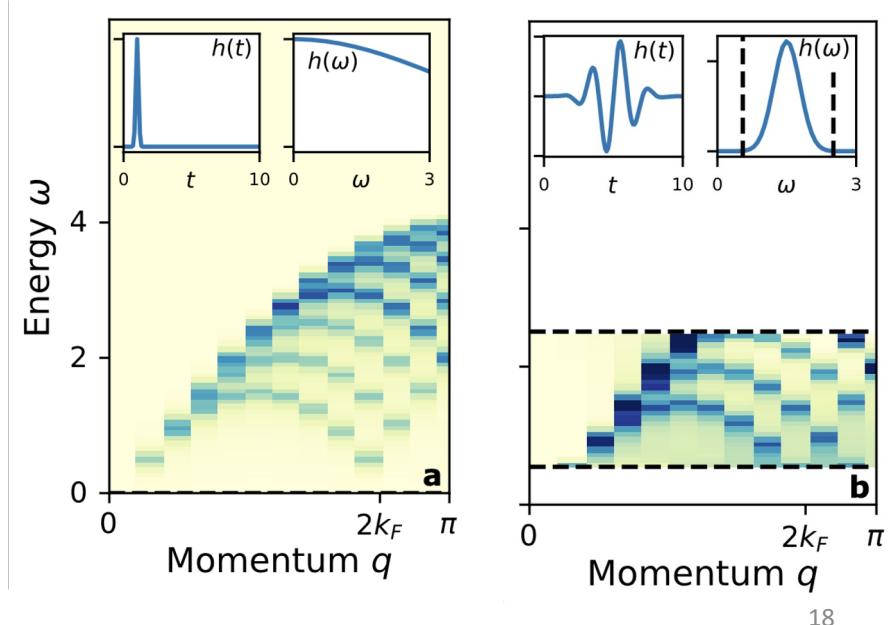


$$A(t) = \int \omega dt' \chi^R(t'(t)) h(t') + \mathcal{O}(h^{3/2})$$

$$\chi(r, t) = -i \langle \psi_0 | \delta n(r, t) \delta n(r=0, t=0) | \psi_0 \rangle$$

Measure density
on all sites ($\mathbf{A}=n_i$)

Wiggle potential
on site 0 ($\mathbf{B}=n_0 V_0$)



Fermionic Linear Response

$$\frac{\delta A(t)}{\delta h(t')} \Big|_{h=0} = -i\theta(t-t') \langle \psi_0 | [\mathbf{A}(t), \mathbf{B}(t')] | \psi_0 \rangle$$

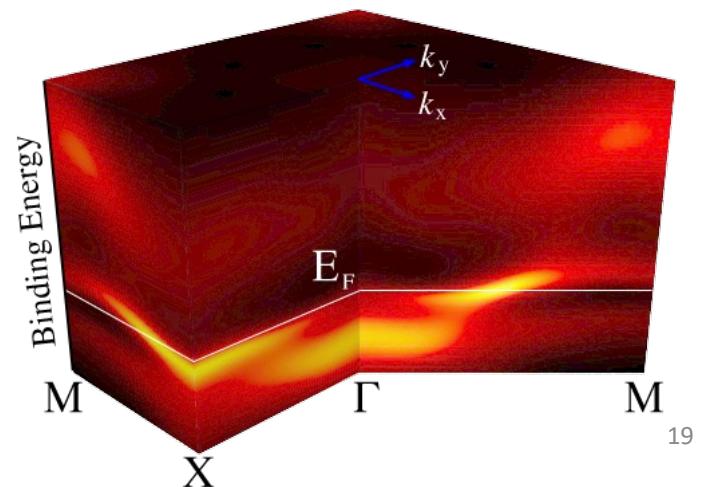
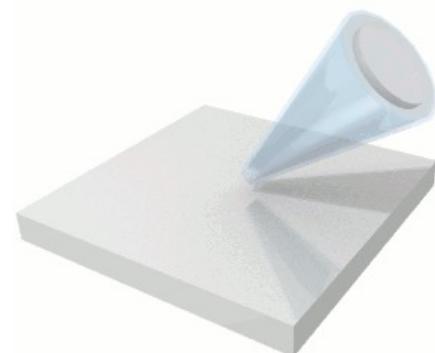
Notice this is a commutator...
... we might also want to have an anti-commuter

$$G(t, t') = -i\theta(t-t') \langle \psi_0 | \{ \mathbf{A}(t), \mathbf{B}(t') \} | \psi_0 \rangle$$

Why?

$$G^R(r_i, t; r_j, t') = -i\theta(t-t') \langle \psi_0 | \{ c_i(t), c_j^\dagger(t') \} | \psi_0 \rangle$$

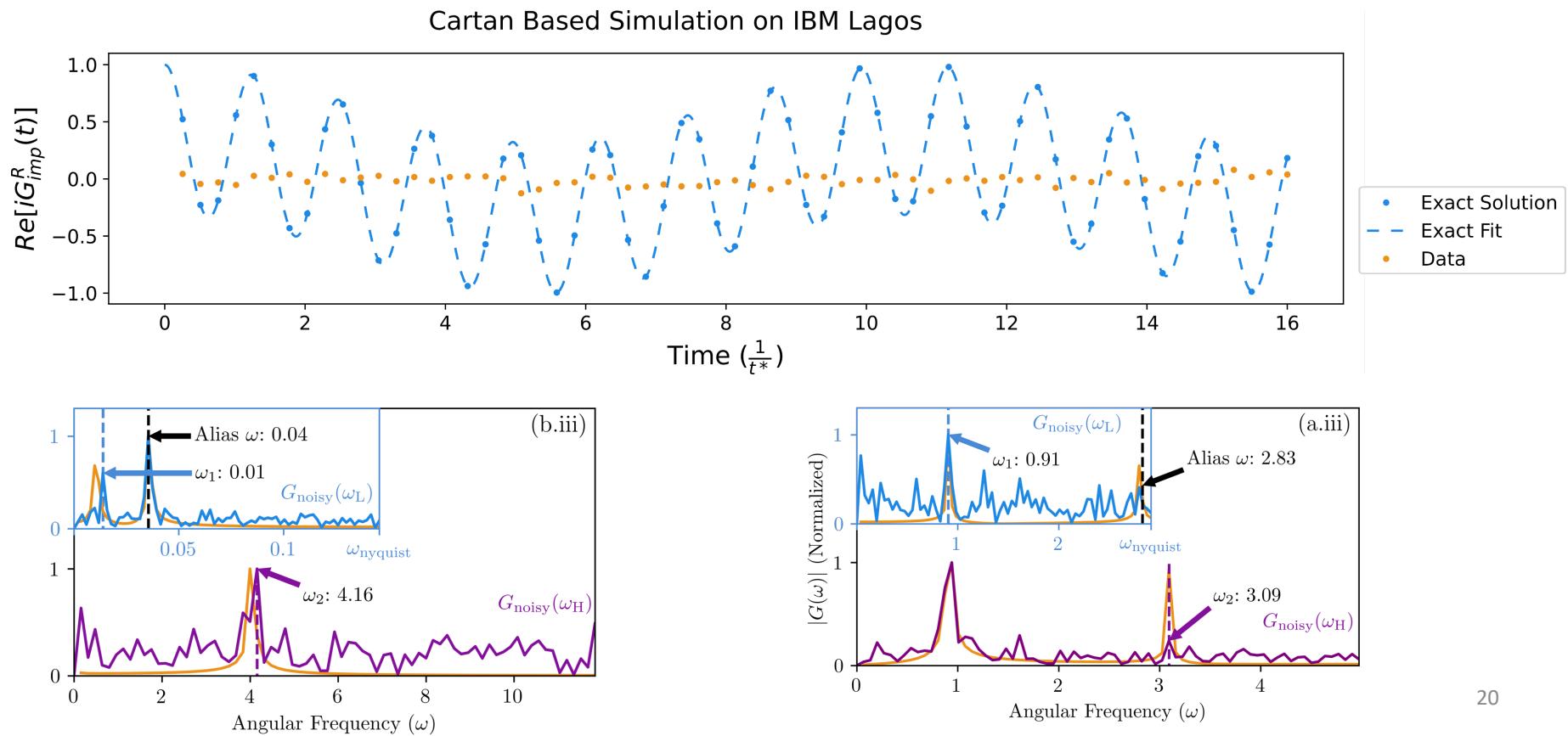
Fermionic creation/
annihilation operators



Application of Green's functions: DMFT

T. Steckmann et al.,
Phys. Rev. Research 5, 023198 (2023)

2-site Hubbard DMFT (5 qubits)

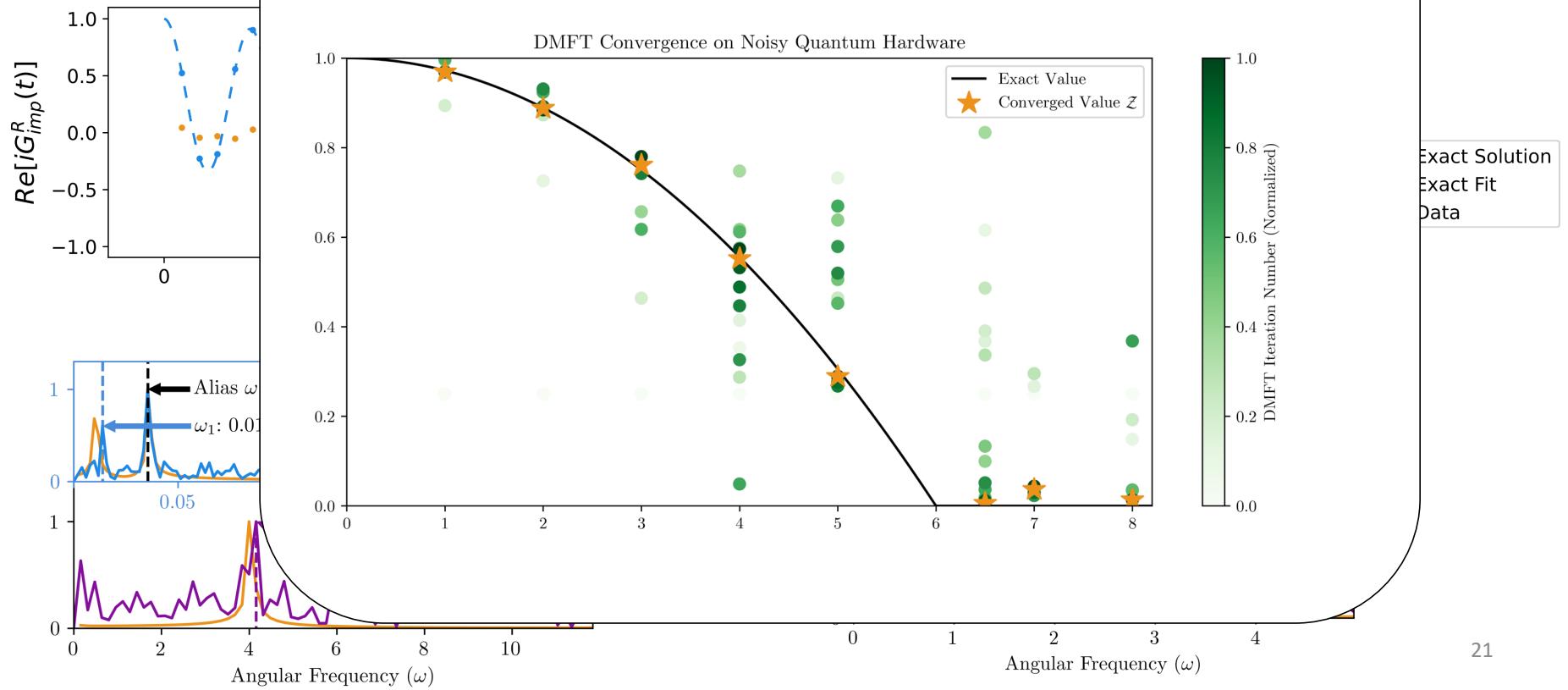


2-site Hubbard DMFT

T. Steckmann et al.,

Phys. Rev. Research 5, 023198 (2023)

Self-consistent DMFT phase diagram showing the metal-insulator transition for 2-site Hubbard model



Option 1: Auxiliary operator

$$\frac{\delta A(t)}{\delta h(t')} \Big|_{h=0} = -i\theta(t-t')\langle\psi_0|[\mathbf{A}(t), \mathbf{B}(t')]\rangle|\psi_0\rangle$$

Find an operator \mathbf{P} such that:

$$\{\mathbf{B}(t), \mathbf{P}\} = 0$$

$$[\mathcal{H}_0, \mathbf{P}] = 0$$

$$\mathbf{P}|\psi_0\rangle = s|\psi_0\rangle$$

Then:

$$\begin{aligned} G(t, t') &= -i\theta(t-t')\langle\psi_0|\{\mathbf{A}(t), \mathbf{B}(t')\}|\psi_0\rangle \\ &= \frac{i}{s}\theta(t-t')\langle\psi_0|[\mathbf{A}(t)\mathbf{P}(t), \mathbf{B}(t')]\rangle|\psi_0\rangle \end{aligned}$$

Example: parity

$$\mathbf{P} = Z_1 Z_2 \dots Z_n$$

Option 2: Post-selection

Option 1: Auxiliary operator

$$\frac{\delta A(t)}{\delta h(t')} \Big|_{h=0} = -i\theta(t-t')\langle\psi_0|[\mathbf{A}(t), \mathbf{B}(t')]\rangle|\psi_0\rangle$$

Find an operator \mathbf{P} such that:

$$\{\mathbf{B}(t), \mathbf{P}\} = 0$$

$$[\mathcal{H}_0, \mathbf{P}] = 0$$

$$\mathbf{P}|\psi_0\rangle = s|\psi_0\rangle$$

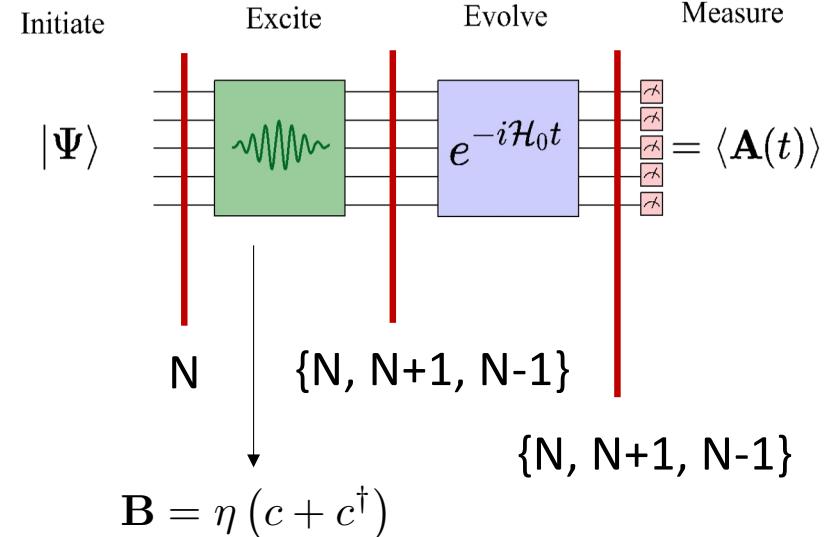
Then:

$$\begin{aligned} G(t, t') &= -i\theta(t-t')\langle\psi_0|\{\mathbf{A}(t), \mathbf{B}(t')\}|\psi_0\rangle \\ &= \frac{i}{s}\theta(t-t')\langle\psi_0|[\mathbf{A}(t)\mathbf{P}(t), \mathbf{B}(t')]\rangle|\psi_0\rangle \end{aligned}$$

Example: parity

$$\mathbf{P} = Z_1 Z_2 \dots Z_n$$

Option 2: Post-selection



Post-selection on particle number gives us

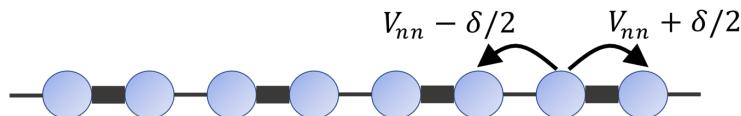
$$G_{ij}^<(t) = i \langle \psi_0 | c_j^\dagger(0) c_i(t) | \psi_0 \rangle$$

$$G_{ij}^>(t) = -i \langle \psi_0 | c_i(t) c_j^\dagger(0) | \psi_0 \rangle$$

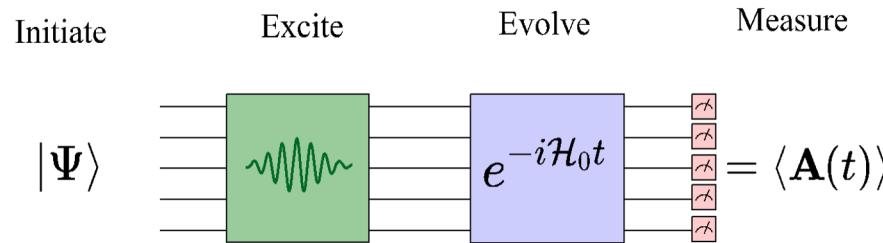
Linear Response -> Green's function

2302.10219

Su-Schrieffer-Heeger model for polyacetylene



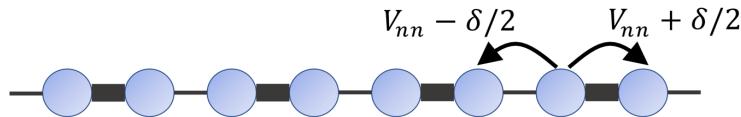
$$\mathcal{H}_0 = - \sum_{\langle i,j \rangle} \left[V_{nn} + (-1)^i \delta/2 \right] c_i^\dagger c_j - \mu \sum_i c_i^\dagger c_i$$



$$G^R(r_i, t; r_j, t') = -i\theta(t - t') \langle \psi_0 | \{c_i(t), c_j^\dagger(t')\} | \psi_0 \rangle$$

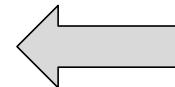
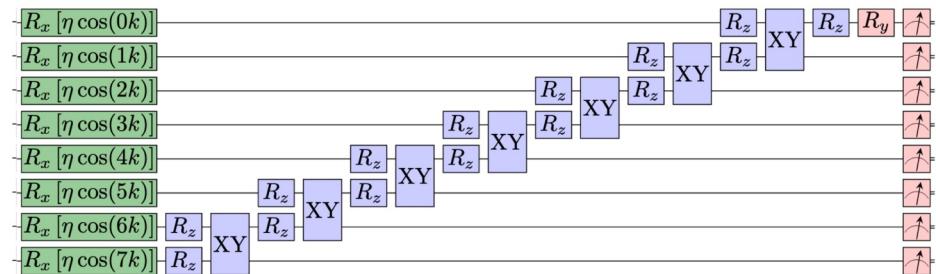
Linear Response -> Green's function

Su-Schrieffer-Heeger model for polyacetylene



$$\mathcal{H}_0 = - \sum_{\langle i,j \rangle} \left[V_{nn} + (-1)^i \delta/2 \right] c_i^\dagger c_j - \mu \sum_i c_i^\dagger c_i$$

Compressed circuit run on *ibm_auckland*



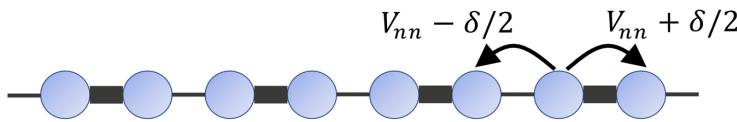
$$\mathbf{B} = \sum_i 2 \cos(kr_i) \left[c_i + c_i^\dagger \right]$$

Choose \mathbf{B} to create a momentum eigenstate

$$G_k^R(t) = -i\theta(t) \langle \psi_0 | \{c_k(t), c_k^\dagger(0)\} | \psi_0 \rangle$$

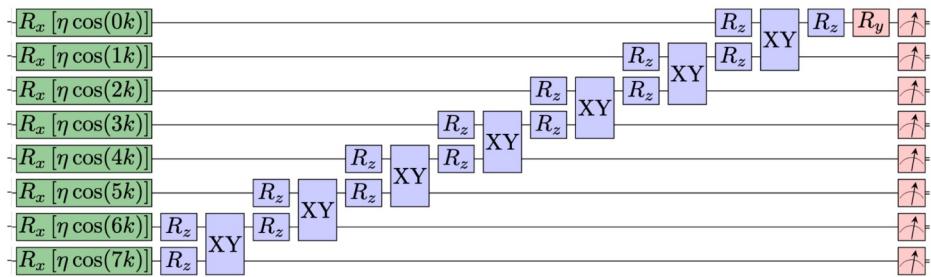
Linear Response -> Green's function

Su-Schrieffer-Heeger model for polyacetylene



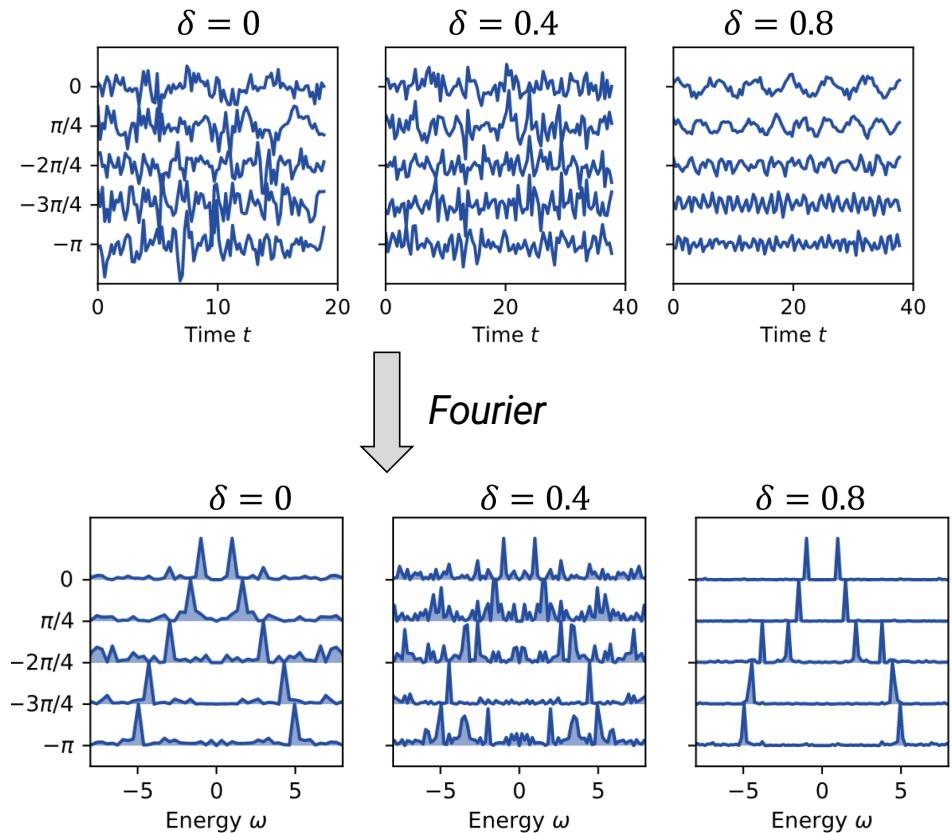
$$\mathcal{H}_0 = - \sum_{\langle i,j \rangle} \left[V_{nn} + (-1)^i \delta/2 \right] c_i^\dagger c_j - \mu \sum_i c_i^\dagger c_i$$

Compressed circuit run on *ibm_auckland*



Choose **B** to create a momentum eigenstate

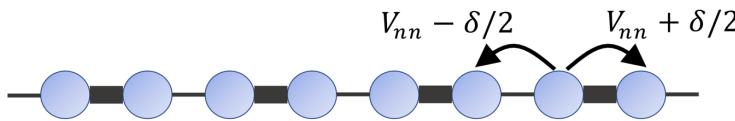
$$G_k^R(t) = -i\theta(t)\langle\psi_0|\{c_k(t), c_k^\dagger(0)\}|\psi_0\rangle$$



Linear Response -> Green's function

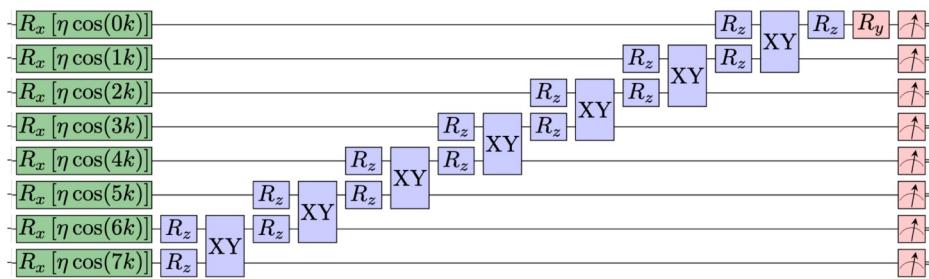
2302.10219

Su-Schrieffer-Heeger model for polyacetylene



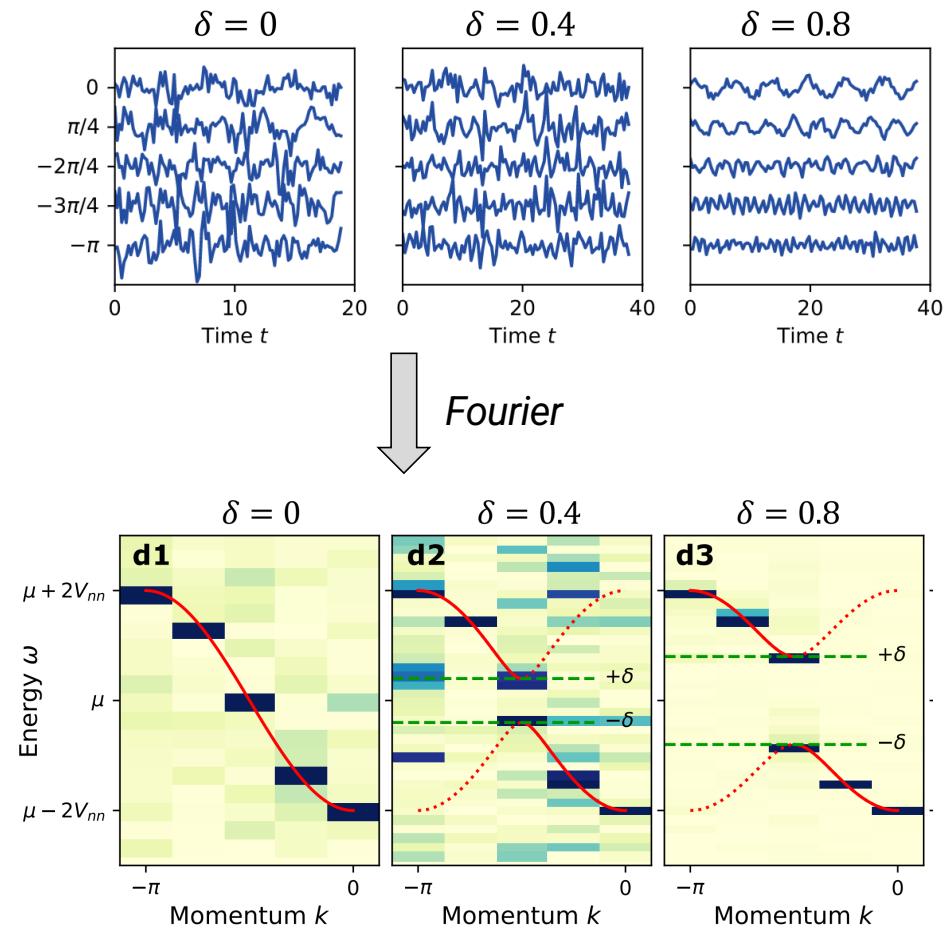
$$\mathcal{H}_0 = - \sum_{\langle i,j \rangle} \left[V_{nn} + (-1)^i \delta/2 \right] c_i^\dagger c_j - \mu \sum_i c_i^\dagger c_i$$

Compressed circuit run on *ibm_auckland*

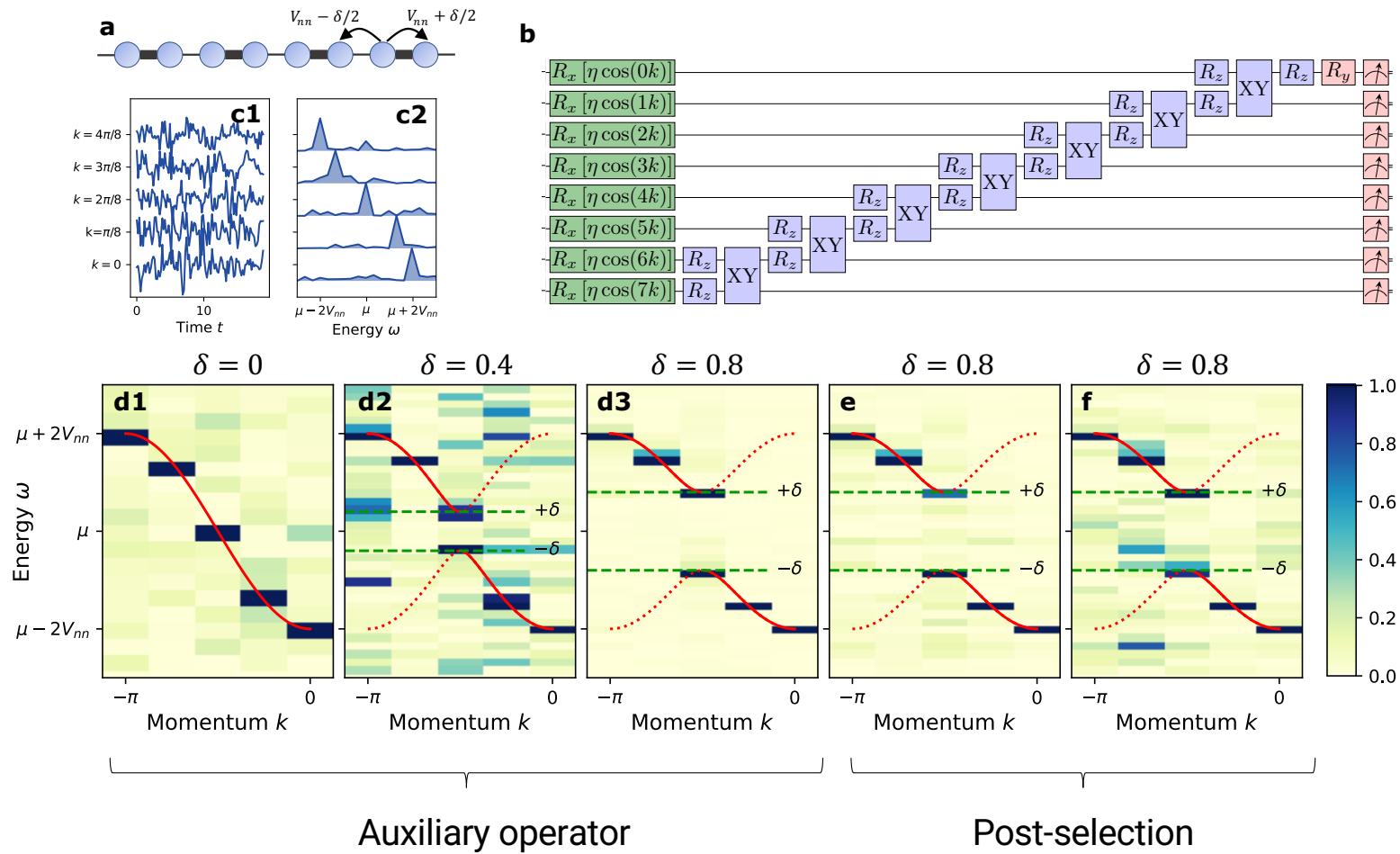


Choose **B** to create a momentum eigenstate

$$G_k^R(t) = -i\theta(t)\langle\psi_0|\{c_k(t), c_k^\dagger(0)\}|\psi_0\rangle$$



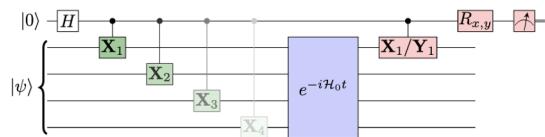
Linear Response -> Green's function



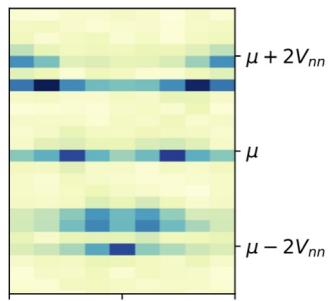
Linear Response -> Green's function

Why does this work so well?

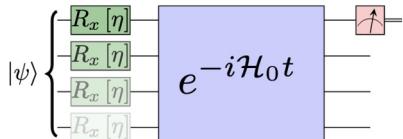
Hadamard test method



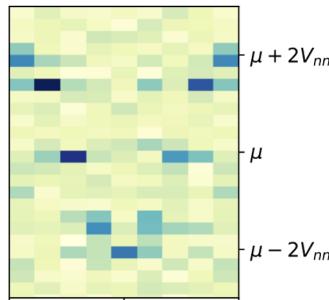
$\xrightarrow{\text{FT}}$
 $t \rightarrow \omega$
 $r \rightarrow k$



Position-selective linear response

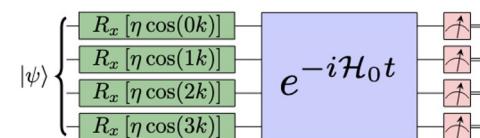


$\xrightarrow{\text{FT}}$
 $t \rightarrow \omega$
 $r \rightarrow k$

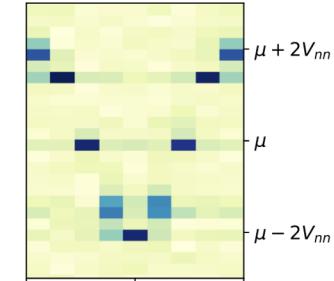


$$\mathbf{B} = \sum_i 2 \cos(kr_i) [c_i + c_i^\dagger]$$

Momentum-selective linear response

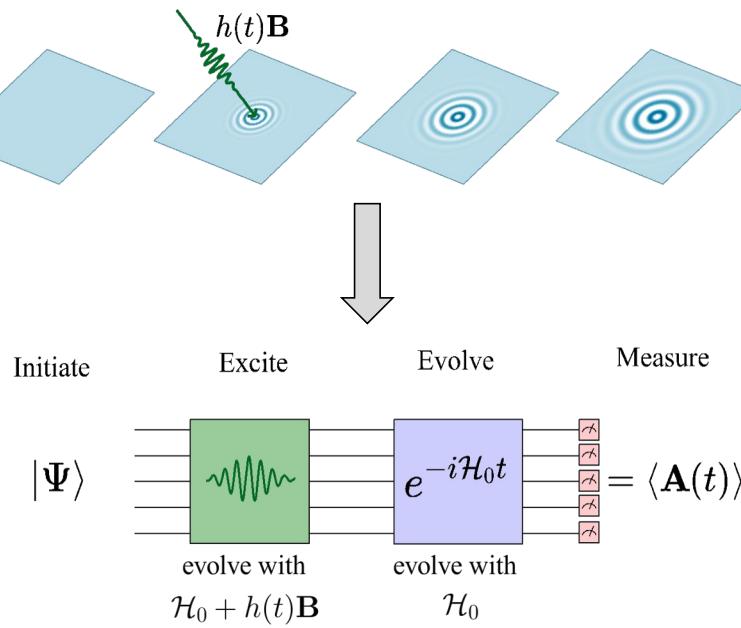
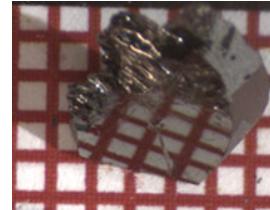


$\xrightarrow{\text{FT}}$
 ω

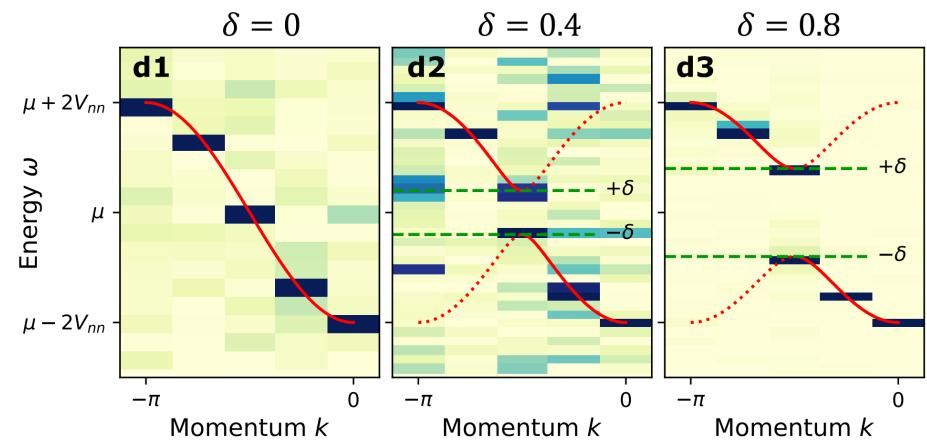


Data from noisy simulator with one/two qubit noise of 1% and 10%

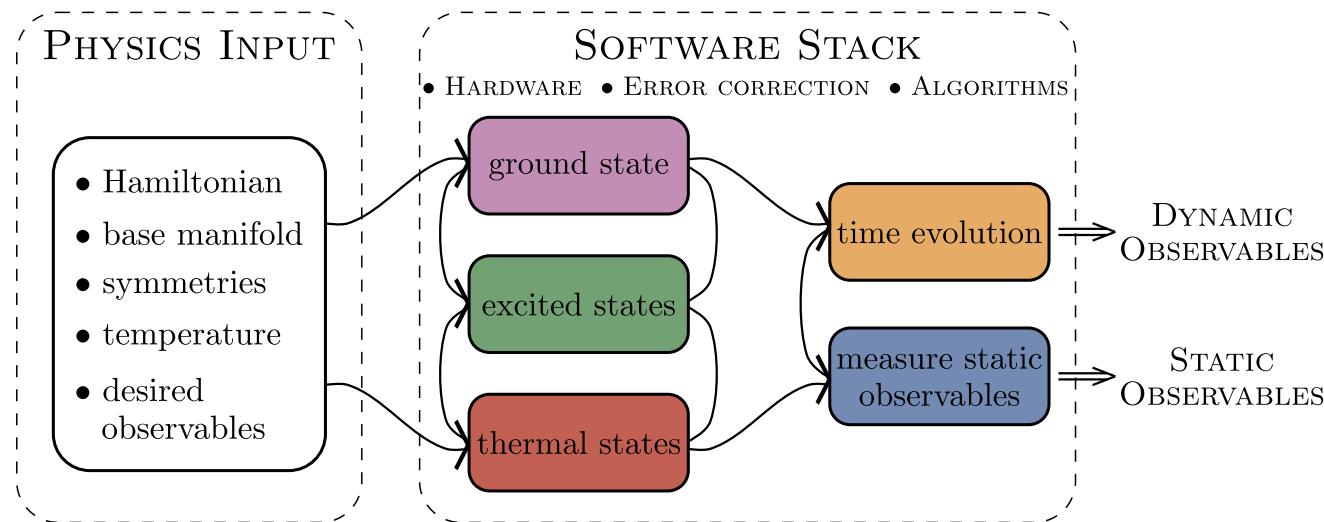
Linear Response



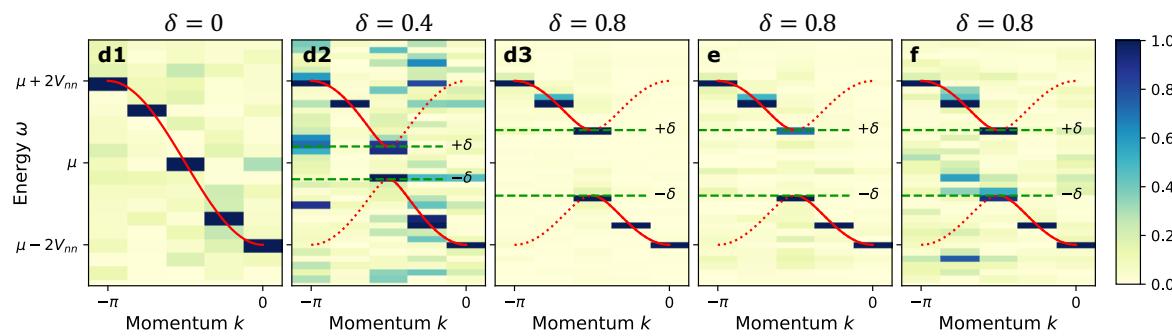
- Ancilla free
- Momentum and frequency selectivity
- Both bosonic and fermionic correlators
- More noise robust compared to existing methods



E. Kökcü, H. Labib, J.K. Freericks, AFK., arXiv:2302.10219



<https://go.ncsu.edu/kemper-lab>



- Experimental relevance:
Measuring correlation functions
- Measuring exact integer Chern numbers for topological states
- Driven/dissipative systems and fixed points (1000 Trotter steps)
- Time evolution via Lie algebraic decomposition and compression
- Thermodynamics via Lee-Yang Zeros
- Physics-Informed Subspace Expansions