

Letter

Structure near the $K^- + p + p$ threshold in the in-flight ${}^3\text{He}(K^-, \Lambda p)n$ reaction

J-PARC E15 Collaboration

Y. Sada^{1,*}, S. Ajimura¹, M. Bazzi², G. Beer³, H. Bhang⁴, M. Bragadireanu⁵, P. Buehler⁶, L. Busso^{7,9}, M. Cargnelli⁶, S. Choi⁴, C. Curceanu², S. Enomoto⁸, D. Faso^{7,9}, H. Fujioka¹⁰, Y. Fujiwara¹¹, T. Fukuda¹², C. Guaraldo², T. Hashimoto¹³, R. S. Hayano¹¹, T. Hiraiwa¹, M. Iio⁸, M. Iliescu², K. Inoue¹, Y. Ishiguro¹⁰, T. Ishikawa¹¹, S. Ishimoto⁸, T. Ishiwatari⁶, K. Itahashi¹³, M. Iwai⁸, M. Iwasaki^{13,14}, Y. Kato¹³, S. Kawasaki¹⁵, P. Kienle^{†,16}, H. Kou¹⁴, Y. Ma¹³, J. Marton⁶, Y. Matsuda¹⁷, Y. Mizoi¹², O. Morra⁷, T. Nagae¹⁰, H. Noumi¹, H. Ohnishi^{13,1}, S. Okada¹³, H. Outa¹³, K. Piscicchia², A. Romero Vidal², A. Sakaguchi¹⁵, F. Sakuma¹³, M. Sato¹³, A. Scordo², M. Sekimoto⁸, H. Shi², D. Sirghi^{2,5}, F. Sirghi^{2,5}, K. Suzuki⁶, S. Suzuki⁸, T. Suzuki¹¹, K. Tanida¹⁸, H. Tatsuno¹⁹, M. Tokuda¹⁴, D. Tomono¹, A. Toyoda⁸, K. Tsukada²⁰, O. Vazquez Doce^{2,21}, E. Widmann⁶, B. K. Wuenschek⁶, T. Yamaga¹⁵, T. Yamazaki^{11,13}, H. Yim²², Q. Zhang¹³, and J. Zmeskal⁶

¹Research Center for Nuclear Physics (RCNP), Osaka University, Osaka, 567-0047, Japan

²Laboratori Nazionali di Frascati dell' INFN, I-00044 Frascati, Italy

³Department of Physics and Astronomy, University of Victoria, Victoria BC V8W 3P6, Canada

⁴Department of Physics, Seoul National University, Seoul, 151-742, South Korea

⁵National Institute of Physics and Nuclear Engineering – IFIN HH, Romania

⁶Stefan-Meyer-Institut für subatomare Physik, A-1090 Vienna, Austria

⁷INFN Sezione di Torino, Torino, Italy

⁸High Energy Accelerator Research Organization (KEK), Tsukuba, 305-0801, Japan

⁹Dipartimento di Fisica Generale, Università di Torino, Torino, Italy

¹⁰Department of Physics, Kyoto University, Kyoto, 606-8502, Japan

¹¹Department of Physics, The University of Tokyo, Tokyo, 113-0033, Japan

¹²Laboratory of Physics, Osaka Electro-Communication University, Osaka, 572-8530, Japan

¹³RIKEN Nishina Center, RIKEN, Wako, 351-0198, Japan

¹⁴Department of Physics, Tokyo Institute of Technology, Tokyo, 152-8551, Japan

¹⁵Department of Physics, Osaka University, Osaka, 560-0043, Japan

¹⁶Technische Universität München, D-85748, Garching, Germany

¹⁷Graduate School of Arts and Sciences, The University of Tokyo, Tokyo, 153-8902, Japan

¹⁸ASRC, Japan Atomic Energy Agency, Ibaraki 319-1195, Japan

¹⁹Department of Chemical Physics, Lund University, Lund, 221 00, Sweden

²⁰Department of Physics, Tohoku University, Sendai, 980-8578, Japan

²¹Excellence Cluster Universe, Technische Universität München, D-85748, Garching, Germany

²²Korea Institute of Radiological and Medical Sciences (KIRAMS), Seoul, 139-706, South Korea

*E-mail: sada@rcnp.osaka-u.ac.jp

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[†]Deceased.

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 To search for an $S = -1$ di-baryonic state which decays to Λp , the ${}^3\text{He}(K^-, \Lambda p)n_{\text{missing}}$ reaction was studied at 1.0 GeV/c. Unobserved neutrons were kinematically identified from the missing mass M_X of the ${}^3\text{He}(K^-, \Lambda p)X$ reaction in order to have a large acceptance for the Λpn final state. The observed Λpn events, distributed widely over the kinematically allowed region of the Dalitz plot, establish that the major component comes from a three-nucleon absorption process. A concentration of events at a specific neutron kinetic energy was observed in a region of low momentum transfer to the Λp . To account for the observed peak structure, the simplest S-wave pole was assumed to exist in the reaction channel, having a Breit–Wigner form in energy and with a Gaussian form factor. A minimum χ^2 method was applied to deduce its mass, $M_X = 2355^{+6}_{-8}$ (stat.) ± 12 (syst.) MeV/ c^2 , and decay width, $\Gamma_X = 110^{+19}_{-17}$ (stat.) ± 27 (syst.) MeV/ c^2 , respectively. The form factor parameter $Q_X \sim 400$ MeV/ c implies that the range of the interaction is about 0.5 fm.

Subject Index D01, D33

1. *Introduction* The $\bar{K}N$ interaction is known to be strongly attractive from low-energy scattering data [1] and X-ray spectroscopy of kaonic atoms [2]. By assuming that the $\Lambda(1405)$ is a $K^- p$ bound state, the existence of a kaonic nuclear bound state has been predicted [3,4]. Observation of a kaonic nuclear bound state would provide definitive information on the $\bar{K}N$ interaction in the sub-threshold region, as well as the nature of $\Lambda(1405)$.

Both theoretical and experimental advances have been made in the last decade. In particular, careful attention has been paid to the simplest kaonic nuclear $\bar{K}NN$ state. Theoretically, all calculations predict the existence of a bound state. However, the predicted $\bar{K}NN$ pole positions, depending on $\bar{K}N$ interaction models, are scattered. For the energy-independent model (static calculation), the binding energy is 47–95 MeV [4–11]. On the other hand, it becomes 9–32 MeV [11–13] for the energy-dependent case. The widths are also widely scattered over 34–110 MeV/ c^2 .

Experimentally, there are many reports on observed peak structure ~ 100 MeV below the $\bar{K}NN$ production threshold. The FINUDA group reported a peak structure in the back-to-back Λp invariant mass spectra via the stopped kaon reaction on ${}^6\text{Li}$, ${}^7\text{Li}$, and ${}^{12}\text{C}$ targets [14] with binding energy (B.E.) 115^{+6}_{-5} MeV and width (Γ) 67^{+14}_{-11} (stat.) $^{+2}_{-3}$ (syst.) MeV/ c^2 . The DISTO group observed $\bar{K}NN$ decaying to Λp in pp collisions at 2.85 GeV, with a B.E. of 103 ± 3 (stat.) ± 5 (syst.) MeV and Γ of 118 ± 8 (stat.) ± 10 (syst.) MeV/ c^2 [15]. In the pion induced reaction, $d(\pi^+, K^+)$ at 1.69 GeV/ c , the E27 group observed “ $K^- pp$ ”-like structure in the $\Sigma^0 p$ decay mode at B.E. 95^{+18}_{-17} (stat.) $^{+30}_{-21}$ (syst.) MeV, $\Gamma = 162^{+87}_{-45}$ (stat.) $^{+66}_{-78}$ (syst.) MeV/ c^2 [16]. Conversely, no significant structure was observed in a SPring-8/LEPS γ induced inclusive experiment [17] or in a proton–proton interaction at HADES/GSI [18]. Also, for the kaon stopped reaction, the other interpretations (i.e. two-nucleon absorption of kaons which have the final state Λp or $\Sigma^0 p$) are widely discussed [19,20]. Thus, the evidence for kaonic nuclei remains controversial.

In an attempt to clarify this situation, the E15 experiment on the $K^- + {}^3\text{He}$ reaction at $p_{K^-} = 1$ GeV/ c is under way at J-PARC. The first physics data were accumulated in May 2013. By using an in-flight reaction and a light nuclear target, backgrounds from multi-nucleon absorption processes and hyperon decays are expected to be kinematically discriminated from the $\bar{K}NN$ signal. The semi-inclusive forward neutron spectrum in the E15 experiment has a long sub-threshold tail reaching ~ 100 MeV below the $\bar{K}NN$ threshold, but no significant structure was seen in the deeply bound region [21].

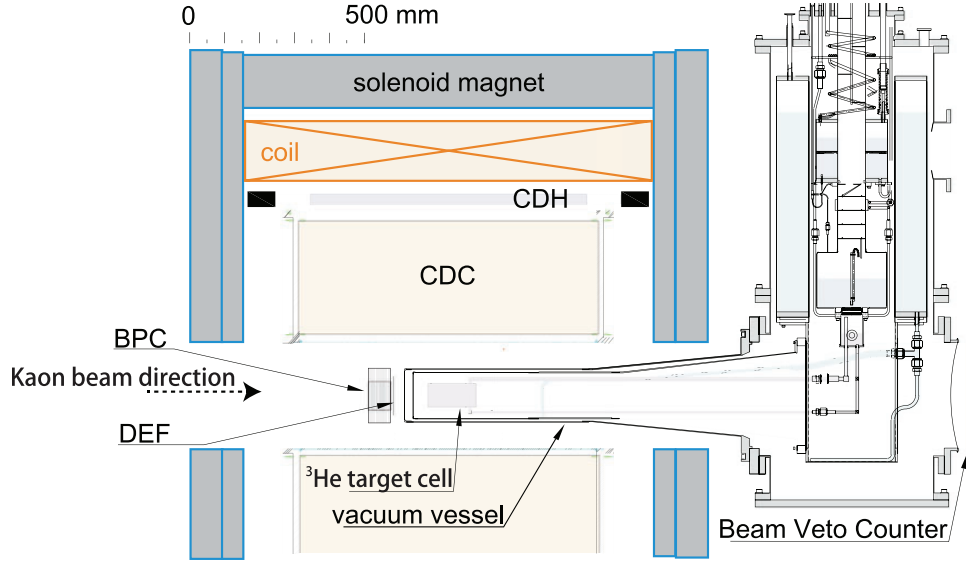


Fig. 1. Schematic diagram of detectors in the CDS and of the target system [22].

In this paper, we present a study of the $K^- + {}^3\text{He}$ reaction at $p_{K^-} = 1 \text{ GeV}/c$ focusing on the Λp invariant mass spectrum, in particular on the $\Lambda p n$ final state for the May 2013 data.

2. The experimental apparatus Here, the experimental apparatus relevant to the $\Lambda p n$ study is briefly described. The ongoing experiment is being conducted at the J-PARC 30 GeV synchrotron (MR). In May 2013, typical beam intensity of the primary proton beam in MR was 30×10^{12} per spill, where the spill length was 2 seconds with a 6 second repetition cycle. A high intensity proton beam from MR impinges on a gold target, and $1 \text{ GeV}/c$ K^- s are selected by the K1.8BR beamline spectrometer and purified with an electrostatic separator. The typical K^-/π ratio was 0.45. At the online trigger level, K^- s are selected with an aerogel Cherenkov counter. The purity of the K^- at this level, determined by time-of-flight (ToF) analysis, was 99%, with an intensity of 1.5×10^5 per spill. The momentum resolution of the beamline spectrometer is $(2.0 \pm 0.5) \times 10^{-3}$ with an absolute precision of $2 \text{ MeV}/c$ at $1 \text{ GeV}/c$. A more detailed description can be found in [22].

Figure 1 shows a schematic diagram of the setup around the ${}^3\text{He}$ target. A cylindrical target cell, 137 mm long and 68 mm in diameter, filled with liquid ${}^3\text{He}$, was placed at the final focus of the beam line. The density of the target was 0.081 g cm^{-3} at a temperature of 1.4 K. Details of the target system are given in [23]. To detect particles incident on the target, a small scintillator (DEF) was placed in front of the target cell. The reconstructed beam track from the K1.8BR beam line spectrometer was finally connected to the hit point on a drift chamber (BPC), installed just upstream of the DEF, to improve reconstructed vertex resolution. To measure the K^- reaction products, the ${}^3\text{He}$ target was surrounded by a cylindrical detector system (CDS), consisting of a cylindrical drift chamber (CDC) and a cylindrical detector hodoscope (CDH) operating in a uniform $\sim 0.7 \text{ T}$ magnetic field. The secondary particle tracks are reconstructed by the CDC using a helix fitting method, and their momenta are measured. Energy loss of each track, where major loss is mainly in the target region, was evaluated and corrected by the Runge–Kutta method. The transverse momentum resolution (σ_{P_t}) is found to be $5.3\% P_t \oplus 0.5\%/\beta$, where P_t and β are the transverse momentum in GeV/c and the velocity of the charged particle, respectively. Particle identification (PID) was performed based on

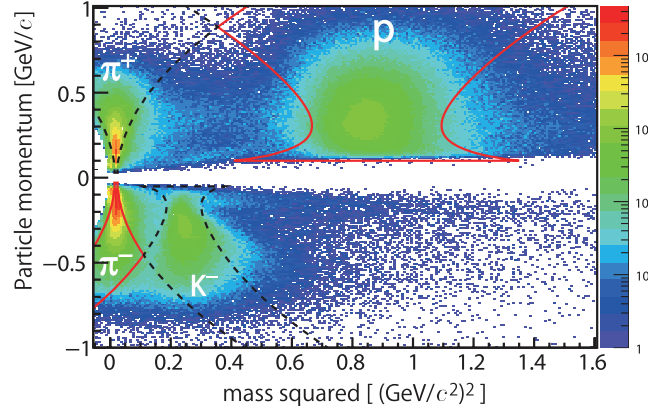


Fig. 2. PID distribution of the CDS. The cut region for each particle is defined as 2.5σ . Overlapped regions are excluded.

the ToF between incoming kaon timing and the CDH. Details of the CDS detectors are also described in [22].

The data acquisition (DAQ) trigger signal was generated by a kaon in the beamline spectrometer, a hit on the DEF counter, and two or more charged-particle hits in the CDH. The trigger rate was typically $\sim 10^3$ per spill with a DAQ live time of $\sim 80\%$. During the May 2013 run period of ~ 90 hours, data corresponding to 3.4×10^9 effective kaons on the ^3He target were accumulated and used in the following analysis.

3. Analysis The Λpn final state was identified by establishing that the $\pi^- p$ pair came from the Λ decay following detection of two protons and one negative pion in the CDS. The missing neutron was identified kinematically. To reconstruct the Λp tracks in the CDS, precise PID was performed based on a mass calculation using the momentum of the reconstructed track along with the CDH-based ToF information. Energy loss corrections based on properties of the inner CDC materials yielded improved PID functions and permitted more accurate momentum reconstructions at the reaction vertex. The calculated mass-squared is shown as a function of momentum in Fig. 2. The mass-squared distribution was sliced into momentum regions and fitted with Gauss functions for each particle species. The standard deviation was used for the PID function of each particle. The cuts defining protons, kaons, and pions, shown as lines on Fig. 2, were set to 2.5 times the sigma of the respective mass-squared distributions. To avoid misidentification, regions of overlap were excluded, resulting in an estimated purity of 99.5% for protons and pions.

To identify the $\pi^- p$ pair associated with a specific Λ decay in a $\pi^- pp$ event, a log-likelihood method was used on the product of five probability density functions: (1) distance of closest approach (DCA) between a $\pi^- p$ pair candidate for Λ decay ($f_{\text{DCA}(\pi p)}$), (2) DCA between the kaon beam and an unpaired proton ($f_{\text{DCA}(K-p)}$), (3) DCA between the kaon beam and a reconstructed Λ track ($f_{\text{DCA}(K-\Lambda)}$), (4) DCA between the unpaired proton and the reconstructed Λ track on the kaon beam ($f_{\text{DCA}(\Lambda p)}$), and (5) invariant mass distribution of the $\pi^- p$ pair candidate ($f_{M_{\text{inv},\pi^- p}}$). Probability density functions were calculated based on the distributions of these five quantities for a simulated Λpn final state. In this simulation, we assume a flat distribution (S-wave) over the three-body phase space based on the phase volume, namely:

$$\frac{d^2\sigma_{3\text{NA}(\Lambda pn)}}{dT_n^{\text{CM}} d\cos\theta_n^{\text{CM}}} \propto \rho_3(\Lambda pn), \quad (1)$$

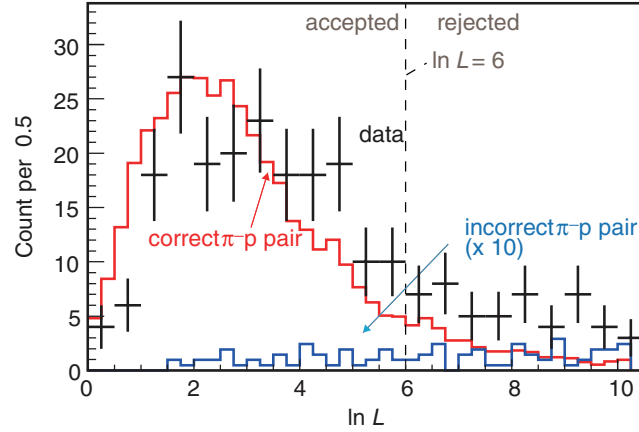


Fig. 3. Likelihood function ($\ln L$) distribution. Black crosses are $\ln L$ of data. Since there are two possible pairs, the pair having smaller $\ln L$ is plotted. The red histogram shows the simulated $\ln L$ of the correct $p\pi^-$ pair from Λ decay in a $pp\pi^-$ event. The blue histogram is $\ln L$ of the incorrect pair of $p\pi^-$ in a Λ decay. The blue histogram is vertically scaled 10 times.

where $\sigma_{3\text{NA}(\Lambda pn)}$ is the simulated event distribution for the Λpn final state, T_n^{CM} is the kinetic energy of the neutron in the CM frame, and $\cos \theta_n^{\text{CM}}$ is the neutron emission angle in the CM frame. GEANT4 [24] is utilized to take into account geometrical information for all detectors and their resolution for all simulations.

Using the distribution functions given above, the log-likelihood function ($\ln L$) is defined as:

$$\ln L = -\ln \left(f_{\text{DCA}(\pi p)} \times f_{\text{DCA}(K^- p)} \times f_{\text{DCA}(K^- \Lambda)} \times f_{\text{DCA}(\Lambda p)} \times f_{M_{\text{inv.}\pi^- p}} \right). \quad (2)$$

$\ln L$ distributions for MC simulations as well as for data are shown in Fig. 3. The pairs having smaller $\ln L$ were chosen as the correct $\pi^- p$ pair from Λ decay; events with $\ln L > 6$ were rejected. Based on our simulation, the incorrect $\pi^- p$ pair selection probability is estimated to be less than 0.5% of the total simulated events. Although utilization of the $\ln L$ function is valid for the Λpn three-body state, it would be somewhat less effective for other reaction channels.

The reaction vertex was defined to be at the center of vertex($K\Lambda$) and vertex(Kp), where vertex($K\Lambda$) is the closest approach point of the kaon and Λ tracks on the kaon track, and vertex(Kp) that of kaon and proton tracks on the kaon track. The vertex was constrained to lie within the fiducial volume of the target, 30 mm in radius and 100 mm long. The contamination from reactions on the material in the target cell is estimated to be less than 2% of events in the fiducial volume from the empty-target data. Furthermore, by focusing on the final state Λpn , it becomes negligible because of kinematical conservation.

4. Results Previous reports of an $S = -1$ di-baryonic state (the so-called “ $\bar{K}NN$ ” state) were based on Λp invariant mass spectra [14,15]. The present Λp invariant mass spectrum, based on the above analysis, is shown in Fig. 4(a). There are many channels having the Λp in the final state. To focus on the Λpn final state, the missing mass of the ${}^3\text{He}(K^-, \Lambda p)X_{\text{miss.}}$ reaction was calculated kinematically. A neutron peak is clearly seen as shown in Fig. 4(b). The resolution of the Λp invariant mass $M_{\text{inv.}\Lambda p}$ and the missing mass $M(X_{\text{miss.}})$ were ~ 10 MeV at $M_{\text{inv.}\Lambda p} \sim 2.37 \text{ GeV}/c^2$ and ~ 40 MeV at $M(X_{\text{miss.}}) \sim 0.94 \text{ GeV}/c^2$, respectively. To estimate the purity of our Λpn event selection, we applied a “multi-channel global fit” simultaneously to the Λp invariant mass and the

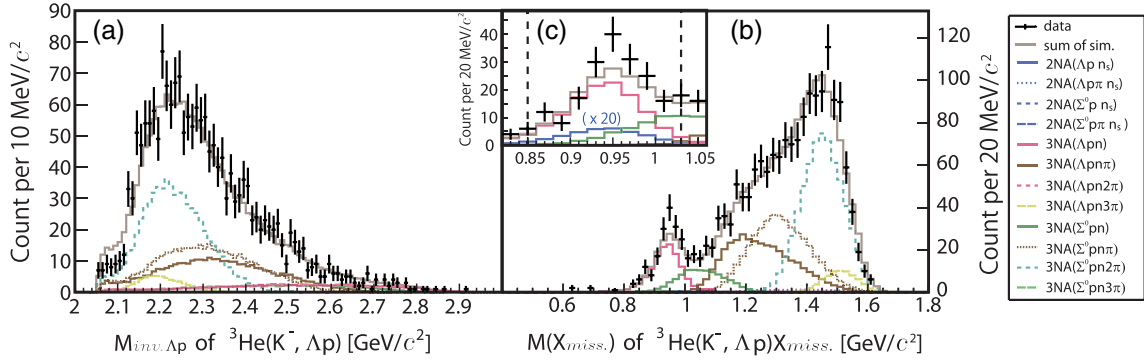


Fig. 4. Inclusive spectra of the ${}^3\text{He}(K^-, \Lambda p)$ reaction and the global fit result of simulation with multi-nucleon absorption processes. (a) Λp invariant mass distribution. (b) Missing mass $M(X_{\text{miss.}})$ spectra of ${}^3\text{He}(K^-, \Lambda p)X_{\text{miss.}}$. (c) Close-up view of (b) around the missing neutron region. In panel (c), $2\text{NA}(\Lambda p n_s)$ is vertically scaled 20 times. The dashed vertical lines in (c) show n -window selection.

missing mass spectra, with simulated physics processes which might contribute to these two spectra, namely multi-nucleon kaon absorption with multi-pion emission. We denoted each process as $2\text{NA}(YNN_s + \#n\pi)$ and $3\text{NA}(YNN + \#n\pi)$, where Y is a Λ or Σ^0 , N is a nucleon, N_s is a spectator nucleon, and $\#n$ is the number of pions. For simplicity, we haven't taken into account the final state interaction. In this simulation, we assumed that particles in the final state are distributed proportionally to the phase space volume in the same way as given in Eq. 1.

$$\frac{d^2\sigma(YNN + \#n\pi)}{dT_n^{\text{CM}} d\cos\theta_n^{\text{CM}}} \propto \rho_{3+\#n}(YNN + \#n\pi). \quad (3)$$

If there is a spectator nucleon, the Fermi momentum distribution is considered to reproduce the ${}^3\text{He}(e, e'p)$ result [25]. The generated events were converted to the event data format, and examined with the same analysis routine applied for the real data. A list of physics processes taken into account for the global fit is shown in Table 1, together with the relative yields obtained for each process to the number of obtained events. The fit results are given as histograms in Fig. 4, and the two spectra are described well by the processes listed in Table 1. The χ^2 and DOF of the fit are 122 and 147, respectively. Figure 4(c) shows a close-up view in the missing neutron region. As shown in the figure, the Λpn final state events can be selected by setting the neutron window to be $0.85 < M(X_{\text{miss.}}) < 1.03$ [GeV/ c^2]. For simplicity, we denote those events in the neutron window as “ Λpn events.” This close-up view also indicates that the fit result of the relative Λpn yield is weaker than the data (about 83% compared to the data). This implies that there could be an unidentified channel which contributes to the yield of the Λpn final state, other than $3\text{NA}(\Lambda pn)$ as discussed later. From the global fit, it is expected that three channels, $3\text{NA}(\Lambda pn)$, $3\text{NA}(\Sigma^0 pn)$, and $2\text{NA}(\Lambda pn_s)$, will remain in the neutron window (n -window) at the ratio of 0.62 : 0.20 : 0.01 compared with the data in the n -window.

The Λpn event distribution over the phase space in the CM frame is shown in Fig. 5(a) as a Dalitz plot with kinetic energies of Λ , proton and neutron, normalized by the Q -value of the reaction. And for reference, the detection efficiency of the CDS is plotted in Fig. 5(b). It shows that our detector system has a fairly flat acceptance over the Λpn phase space, except for the upper-left and the upper-right corners. Acceptance reduction of these corners comes from a lower kinetic energy of the p and Λ (below the detection threshold). There is, however, sufficient acceptance at the bottom, where two-nucleon kaon absorption with a spectator neutron n_s , $K^- + {}^3\text{He} \rightarrow \Lambda p + n_s$, is expected.

Table 1. Relative yield of each component to the number of obtained events in the global fit, normalized to that of data (the same for each component in the n -window normalized by the yield of data in the n -window). Note that the spectral shapes of some reaction channels are quite similar, especially for the channels given in a row.

Process	Relative yield		Process	Relative yield	
	all	n -window		all	n -window
$2\text{NA}(\Lambda p n_s)$	0.001	0.01	$2\text{NA}(\Sigma^0 p n_s)$	$<10^{-4}$	$<10^{-2}$
$2\text{NA}(\Lambda p n_s + \pi)$	$<10^{-4}$	$<10^{-2}$	$2\text{NA}(\Sigma^0 p n_s + \pi)$	0.010	$<10^{-2}$
$3\text{NA}(\Lambda p n)$	0.072	0.62	$3\text{NA}(\Sigma^0 p n)$	0.058	0.20
$3\text{NA}(\Lambda p n + \pi)$	0.199	$<10^{-2}$	$3\text{NA}(\Sigma^0 p n + \pi)$	0.239	$<10^{-2}$
$3\text{NA}(\Lambda p n + 2\pi)$	$<10^{-4}$	$<10^{-2}$	$3\text{NA}(\Sigma^0 p n + 2\pi)$	0.354	$<10^{-2}$
$3\text{NA}(\Lambda p n + 3\pi)$	0.039	$<10^{-2}$	$3\text{NA}(\Sigma^0 p n + 3\pi)$	$<10^{-4}$	$<10^{-2}$

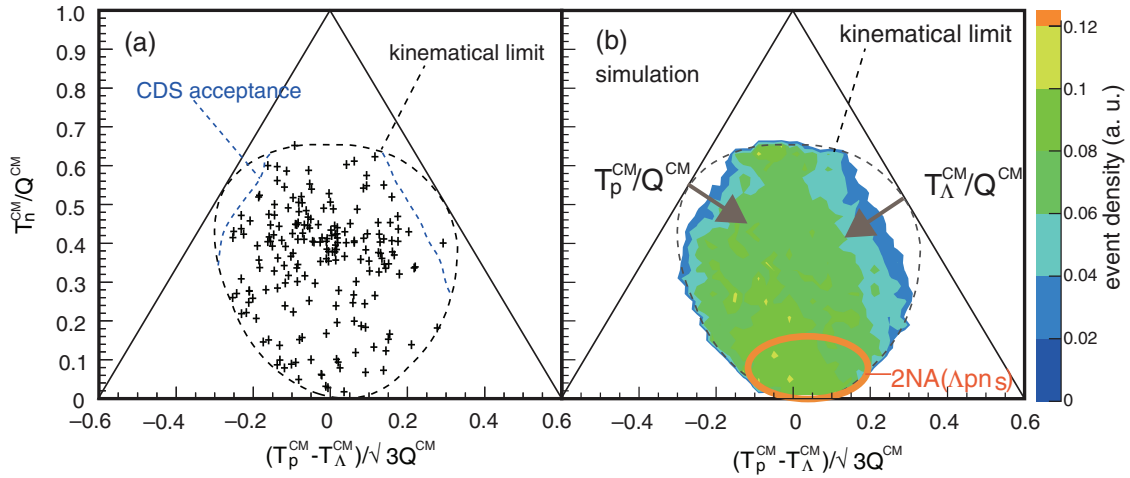


Fig. 5. (a) Dalitz plot of the selected $\Lambda p n$ events in the CM frame. (b) The simulated detector acceptance in a Dalitz plot.

The data events in Fig. 5(a) are widely distributed over the kinematically allowed region, which is consistent with the global fit showing that the major component is coming from the three-nucleon absorption process. In the Dalitz plot, an event concentration was observed at a normalized neutron kinetic energy $T_n^{CM}/Q^{CM} \approx 0.4$, which indicates that the Λp invariant mass will have a structure corresponding to that energy.

The distribution of the Λp invariant mass and the calculated neutron emission angle are shown in Fig. 6(b) and (c), and a scatter plot of the two is given in Fig. 6(a). As indicated in the Dalitz plot, an unexpected peak structure is seen at $M_{\Lambda p} \sim M(K^- + p + p)$ in the Λp invariant mass spectrum, where the global fit is insensitive.

As shown in Fig. 6(c), the neutrons in this peak structure are clearly concentrated at the forward region, where the momentum transfer to the Λp system is minimum, while the wide distribution can be explained by the multi-nucleon absorption processes given by the global fit in the n -window. The values obtained for χ^2/DOF in Fig. 6(b) and (c) are 135/43 and 138/76, respectively. This χ^2 test result indicates that the observed peak structure could not be explained by multi-nucleon absorption processes.

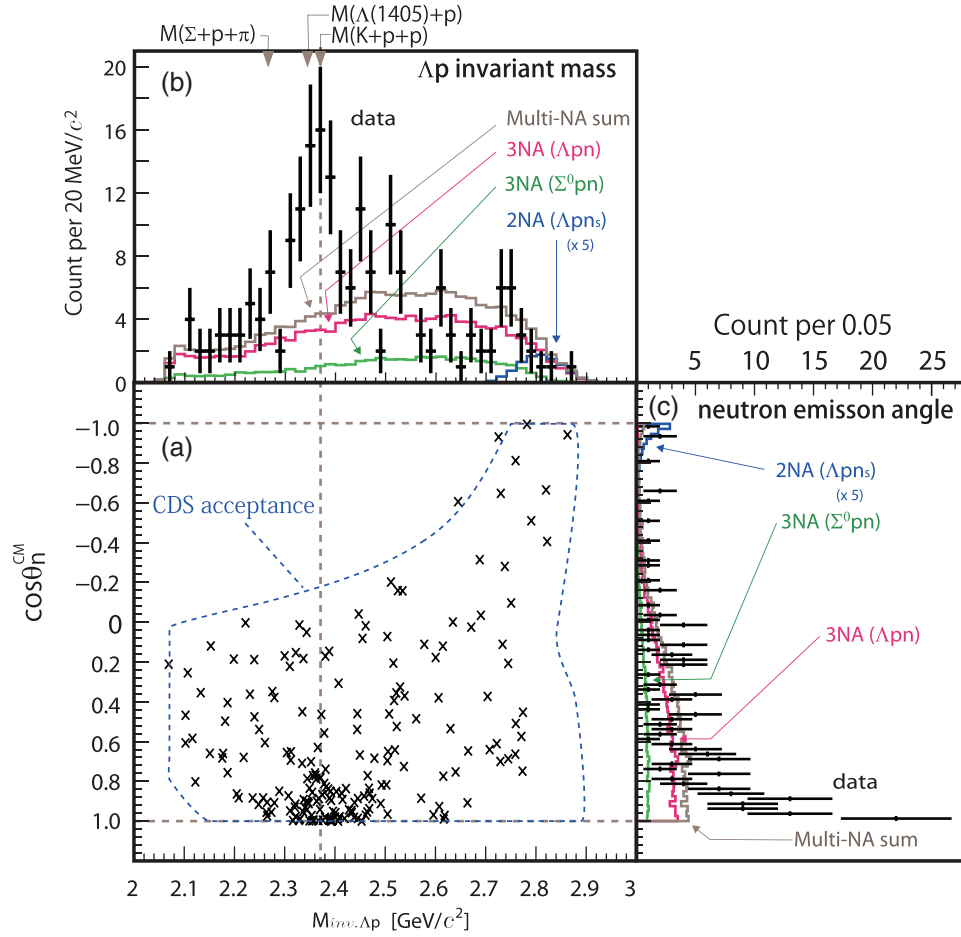


Fig. 6. (a) Two-dimensional distribution of the Λp invariant mass and the emission angle of the missing neutron. The blue dashed lines show a crude CDS acceptance boundary for the Λpn events. (b) Λp invariant mass with simulated spectra obtained by the global fit in the neutron window. (c) Angular distribution of the missing neutron, kinematically reconstructed, as a function of $\cos\theta_n^{\text{CM}}$. The histograms show the contributions of the three remaining channels in the n -window, whose yields are given by the global fit.

5. Discussion To explain the excess, the existence of a simple finite-size S-wave single-pole structure over the three-body phase space which decays to Λp in the final state is assumed. Thus, the formation cross-section (σ_X) of the pole can be written as:

$$\frac{d^2\sigma_X}{dM_{\text{inv},\Lambda p}dq_{\Lambda p}} \propto \rho_3(\Lambda pn) \times \frac{(\Gamma_X/2)^2}{(M_{\text{inv},\Lambda p} - M_X)^2 + (\Gamma_X/2)^2} \times \left| \exp\left(-q_{\Lambda p}^2/2Q_X^2\right) \right|^2, \quad (4)$$

where $M_{\text{inv},\Lambda p}$ is the invariant mass of Λp , $q_{\Lambda p}$ is the momentum transfer of the reaction ($q_{\Lambda p} = |\mathbf{p}_\Lambda + \mathbf{p}_p|$), M_X is the energy, Γ_X is the decay width, and Q_X is the form factor parameter of the pole. The first term in the formula is the three-body Lorentz-invariant phase space of Λpn , the second the Breit–Wigner formula, and the third the square of the form factor, which can also be interpreted as the sticking probability of a plane wave having $q_{\Lambda p}$ to a harmonic oscillator having finite size $\approx \hbar/Q_X$.

We generated events according to Eq. 4 in the simulation. A χ^2 comparative test was made between the experimental data and the simulated pole together with the multi-nucleon absorption processes. Thus, we fitted the spectra keeping the yield of 3NA(Λpn) as a free parameter, because

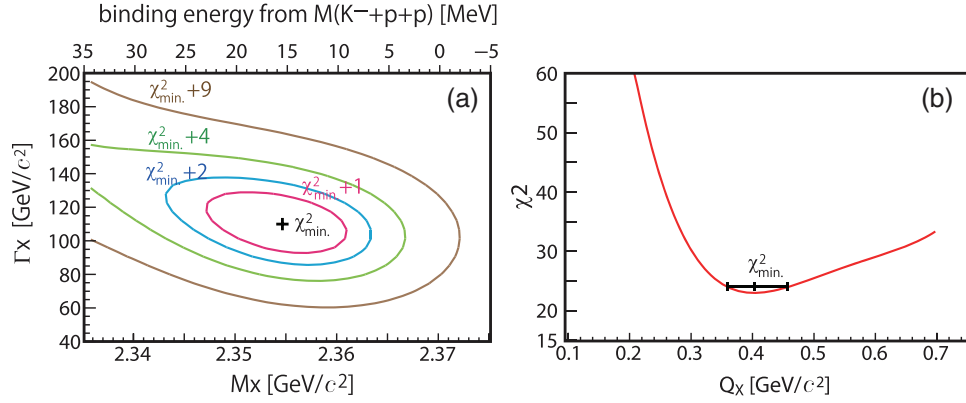


Fig. 7. (a) Two-dimensional χ^2 map of mass and width of pole structure. (b) χ^2 distribution of Q_X .

we introduced a new pole, which decays to the Λp final state. The other two yields for $3\text{NA}(\Sigma^0 pn)$ and $2\text{NA}(\Lambda pn_s)$ are fixed as they are given by the global fit. We first assumed $Q_X = \infty$ and made a two-dimensional χ^2 map on the M_X and Γ_X plane to define the minimum χ^2 for the invariant mass spectra. Then, Q_X is chosen to have minimum χ^2 for the momentum transfer distribution, at a given M_X and Γ_X . This process was iterated until the parameters converged. The two-dimensional χ^2 map as a function of M_X and Γ_X is shown in Fig. 7(a). Figure 7(b) plots χ^2 as a function of momentum transfer. The minimum χ^2 point is at $M_X = 2355^{+6}_{-8}$ (stat.) ± 12 (syst.) MeV/ c^2 , $\Gamma_X = 110^{+19}_{-17}$ (stat.) ± 27 (syst.) MeV/ c^2 , and $Q_X = 400^{+60}_{-40}$ (stat.) MeV/ c . The statistical error is defined as the $\chi^2_{\text{min.}} + 1$ contour. The systematic uncertainties are evaluated by considering the magnetic field strength in the CDS, the likelihood threshold to select the Λp pair, and binning of the invariant mass spectra. The fit results are shown in Fig. 8. Because we simply assumed that the forward neutron emission is due to the form factor as it is given in Eq. 4, we replotted Fig. 6 in terms of $q_{\Lambda p}$ instead of $\cos \theta_n^{\text{CM}}$. The values obtained for χ^2/DOF in Fig. 8(b) and (c) are 68/45 and 23/27, respectively.

To obtain the cross section, one needs to know the detailed event distribution and the acceptance. If we assume that all the angular distributions can be given by the fit results, the acceptance correction can be applied under this assumption. The cross sections can be evaluated as pole: $7 \pm 1 \mu\text{b}$, $3\text{NA}(\Lambda pn)$: $17 \pm 2 \mu\text{b}$, and $2\text{NA}(\Lambda pn_s)$: $0.8^{+2.7}_{-0.8} \mu\text{b}$. If we rely on the global fit and event distribution outside of the neutron window, then we can also determine the $3\text{NA}(\Sigma^0 pn)$ cross section to be $28 \pm 5 \mu\text{b}$ in total. The errors are the quadratic sum of the statistical and the systematic ones, where the systematic uncertainty mainly arises from the target- ^3He and beam-kaon yields.

There could be many interpretations for the pole found in the Λp invariant mass distribution in the $^3\text{He}(K^-, \Lambda p)n_{\text{miss.}}$ reaction channel at $p_K = 1 \text{ GeV}/c$, even if the simplest present assumption of Eq. 4 is valid. A naïve interpretation of the pole would be a $\bar{K} NN$ bound state, since the pole position is located below the $M(K^- + p + p)$ threshold of $2370 \text{ MeV}/c^2$. It could also be a shallow bound or unbound resonance of the $Y^* p$ system, because the $Y^* p$ threshold is located at $2343 \text{ MeV}/c^2$ (assuming $Y^* = \Lambda(1405)$ at $1405 \text{ MeV}/c^2$). However, there is no clear discrimination between the two interpretations given above, if we assume $\Lambda(1405)$ to be a $K^- p$ bound state or penta-quark-like structure.

The Q_X of $\sim 400 \text{ MeV}/c$ is rather large compared to the Fermi motion in light nuclei of about $100 \text{ MeV}/c$, and it implies a short interaction range of about 0.5 fm . Even if we take into account

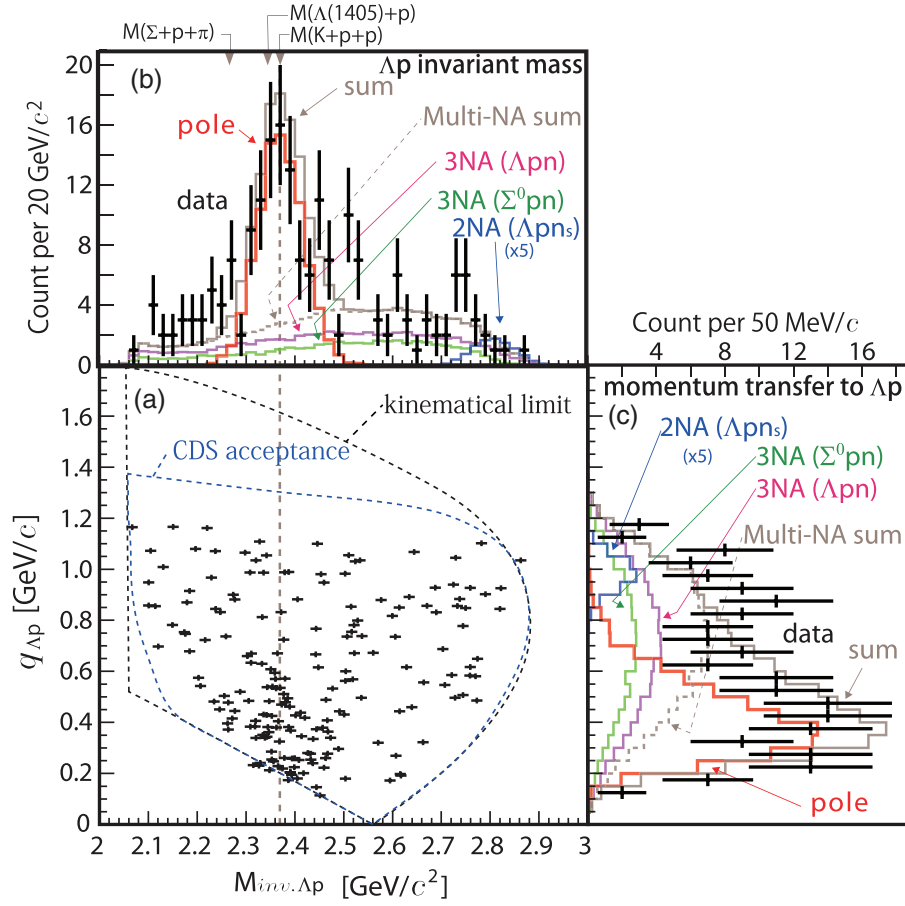


Fig. 8. (a) Two-dimensional distribution of the Λp invariant mass and the momentum transfer to Λp . (b) The fit result of the Λp invariant mass spectrum. (c) Distribution of the momentum transfer to Λp and the fit results of the simulation.

the core motion in the reaction, the effect is as small as about 20%.¹ If we assume a P-wave pole resonance instead of an S-wave, then the reaction with smaller Q_X , namely a longer interaction range, is preferred.

It should be noted that the pole position is close to the two threshold energies, $M(K^- + p + p)$ and $M(Y^* + p)$, and thus the symmetric Breit–Wigner formula could be too simple. For example, the $\bar{K}NN$ -decay channel opens at the K^-pp threshold, so the spectral function observed in the Λp -decay channel could be suppressed above the corresponding thresholds.

There might be a totally different approach to account for the peak structure. One can expect a peak-like structure by assuming $Y^*N \rightarrow \Lambda N$ conversion, after the quasi-free Y^* production by the two-nucleon reaction, namely the $2NA(Y^*np_s)$ reaction followed by $Y^*p_s \rightarrow \Lambda p$ in our notation ($Y^* = \Lambda(1405)$ or $\Sigma(1385)$). In this case, the peak position naturally depends on the momentum transfer as $M_X(q_{\Lambda p})$ (or $\cos \theta_n^{\text{CM}}$). Although the statistics are limited, there is no clear hint of momentum dependence of the peak structure as shown both in Fig. 6 and Fig. 8.

6. Conclusion The ${}^3\text{He}(K^-, \Lambda p)n_{\text{miss.}}$ reaction has been studied with an incident K^- momentum of 1.0 GeV/c for the first time. We found that $17 \pm 2 \mu\text{b}$ of the Λp events in the neutron window

¹ T. Harada, private communication, 2016.

can be explained by a three-nucleon absorption channel, $3\text{NA}(\Lambda pn)$ uniformly spread over the kinematic limit, proportional to the phase space, and without any spectator nucleon. The existence of this channel is already quite interesting, because it seems to follow simple phase space, so that the reaction is “point-like” as in Eq. 1. In contrast, absorption by two nucleons is rather weak compared to three nucleons. The ratio of $2\text{NA}(\Lambda pn_s)/3\text{NA}(\Lambda pn)$ is found to be 0.05 or less.

Apart from the widely distributed 3NA process, a peak structure is observed in the Λp invariant mass around the $\bar{K} NN$ threshold in the low momentum-transfer region. The spectral shape has been examined assuming a single-pole existence whose final state is Λp . Fit results show that the pole has $M_X = 2355_{-8}^{+6}$ (stat.) ± 12 (syst.) MeV/ c^2 , $\Gamma_X = 110_{-17}^{+19}$ (stat.) ± 27 (syst.) MeV/ c^2 , respectively, with a Gaussian form factor parameter $Q_X = 400_{-40}^{+60}$ MeV/ c , having a cross section of about 7 μb . The form factor parameter $Q_X \sim 400$ MeV/ c implies that the range of interaction is about 0.5 fm.

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References

- [1] A. D. Martin, Nucl. Phys. B **179**, 33 (1981).
- [2] M. Bazzi et al., Phys. Lett. B **704**, 113 (2011).
- [3] Y. Akaishi and T. Yamazaki, Phys. Rev. C **65**, 044005 (2002).
- [4] T. Yamazaki and Y. Akaishi, Phys. Lett. B **535**, 70 (2002).
- [5] Y. Ikeda and T. Sato, Phys. Rev. C **79**, 035201 (2009).
- [6] N. V. Shevchenko, A. Gal, and J. Mareš, Phys. Rev. C **76**, 044004 (2007).
- [7] S. Wycech and A. M. Green, Phys. Rev. C **79**, 014001 (2009).
- [8] A. Doté, T. Hyodo, and W. Weise, Phys. Rev. C **79**, 014003 (2009).
- [9] Y. Ikeda and T. Sato, Phys. Rev. C **79**, 035201 (2009).
- [10] S. Maeda, Y. Akaishi, and T. Yamazaki, Proc. Jpn. Acad., Ser. B **89**, 418 (2013).
- [11] J. Revai and N. V. Shevchenko, Phys. Rev. C **90**, 034004 (2014).
- [12] Y. Ikeda, H. Kamano, and T. Sato, Prog. Theor. Phys. **124**, 533 (2010).
- [13] N. Barnea, A. Gal, and E. Z. Liverts, Phys. Lett. B **712**, 132 (2012).
- [14] M. Agnello et al., Phys. Rev. Lett. **94**, 212303 (2005).
- [15] T. Yamazaki et al., Phys. Rev. Lett. **104**, 132502 (2010).
- [16] Y. Ichikawa et al., Prog. Theor. Exp. Phys. **2015**, 021D01 (2015).
- [17] A. O. Tokiyasu et al., Phys. Lett. B **728**, 616 (2014).
- [18] L. Fabbietti et al., Nucl. Phys. A **914**, 60 (2013).
- [19] O. Vázquez Doce et al., [arXiv:1511.04496 [nucl-ex]] [Search INSPIRE].
- [20] A. Filippi and S. Piano, Hyperfine Interact. **233**, 151 (2015).
- [21] T. Hashimoto et al. [E15 Collaboration], Prog. Theor. Exp. Phys. **2015**, 061D01 (2015).
- [22] K. Agari et al., Prog. Theor. Exp. Phys. **2012**, 02B011 (2012).
- [23] M. Iio, S. Ishimoto, M. Sato, S. Enomoto, T. Hashimoto, S. Suzuki, M. Iwasaki, and R. S. Hayano, Nucl. Instrum. Meth. A **687**, 1 (2012).
- [24] S. Agostinelli et al. [GEANT4 Collaboration], Nucl. Instrum. Meth. A **506**, 250 (2003).
- [25] E. Jans et al., Phys. Rev. Lett. **49**, 974 (1982).