

Euler's theorem on the motion of a rigid body

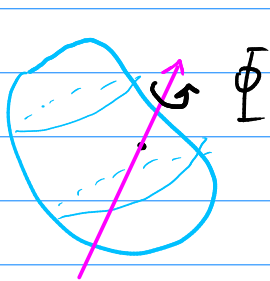
"The general displacement of a rigid body with one point fixed is a rotation about some axis"

GOAL: find an axis at θ and ϕ such that a rotation by ψ will render the actual rotation of the rigid body

We will take the origin of the "body set" of axes (that attached to the body) coinciding with the fixed point - no translation, only rotation

Corollary, Chasles' Theorem:

"The most general displacement of a rigid body is a translation plus a rotation."



$$A = \begin{pmatrix} \cos \phi & \sin \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

ψ

$$A = \begin{bmatrix} \cos \psi \cos \phi - \cos \theta \sin \phi \sin \psi & \cos \psi \sin \phi + \cos \theta \cos \phi \sin \psi & \sin \psi \sin \theta \\ -\sin \psi \cos \phi - \cos \theta \sin \phi \cos \psi & -\sin \psi \sin \phi + \cos \theta \cos \phi \cos \psi & \cos \psi \sin \theta \\ \sin \theta \sin \phi & -\sin \theta \cos \phi & \cos \theta \end{bmatrix}$$

Trace of a transformation matrix is invariant under a similarity transformation.

$$\text{Tr } A' = 1 + 2 \cos \phi = \text{Tr } A$$

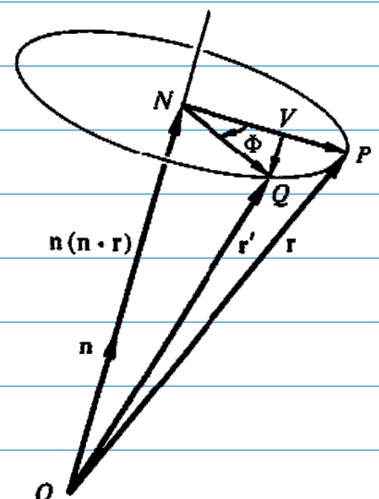
Finite Rotations

GOAL: Find a vectorial transformation that takes r into an r' that is equivalent to a rotation around an axis.

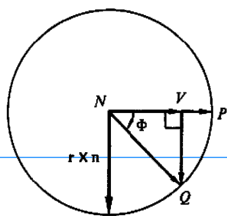
$$\vec{r}' = \vec{ON} + \vec{NV} + \vec{VQ}$$

\vec{ON} : \hat{n} unit vector in the direction of ON

$$|\vec{ON}| = \hat{n} \cdot \vec{r} \Rightarrow \vec{ON} = \hat{n}(\hat{n} \cdot \vec{r})$$



NV:

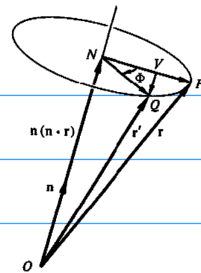


$$|\bar{N}Q| \cos \phi = |\bar{N}V|$$

$$|\bar{N}Q| = |\bar{N}P|$$

$$|\bar{N}V| = |\bar{N}P| \cos \phi$$

$$\Rightarrow \bar{N}V = [\bar{r} - \hat{n}(\hat{n} \cdot \bar{r})] \cos \phi$$



VQ: $|\bar{V}Q| = |\bar{N}Q| \sin \phi = |\bar{r} \times \hat{n}| \sin \phi$

$$\bar{V}Q = (\bar{r} \times \hat{n}) \sin \phi$$

$$\bar{r}' = \hat{n}(\hat{n} \cdot \bar{r}) + [\bar{r} - \hat{n}(\hat{n} \cdot \bar{r})] \cos \phi + (\bar{r} \times \hat{n}) \sin \phi$$

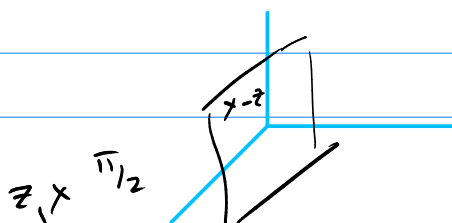
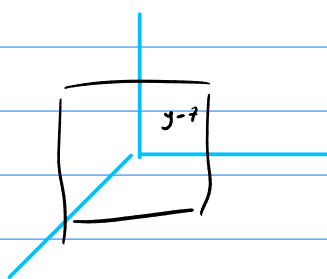
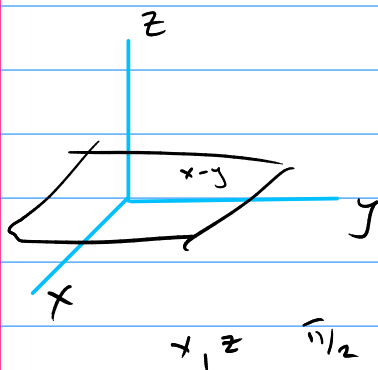
$$\bar{r}' = \bar{r} \cos \phi + \hat{n}(\hat{n} \cdot \bar{r})(1 - \cos \phi) + (\bar{r} \times \hat{n}) \sin \phi$$

This is the "rotation formula" - it is valid for any angle, thus it is a finite-rotation formula (clockwise).

Trace

$$1 + 2 \cos \phi = \text{Tr } A$$

$$\cos \frac{\phi}{2} = \cos \frac{1}{2}(\varphi + \psi) \cos \frac{\theta}{2}$$



$$AB \neq BA$$

Infinitesimal rotations - or in search of the rotation vector through the fixed point

$$x'_i = x_i + \epsilon_{ij} x_j = (\delta_{ij} + \epsilon_{ij}) x_j$$

$$\Rightarrow \boxed{x' = (1 + \epsilon) x}$$

$$1 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

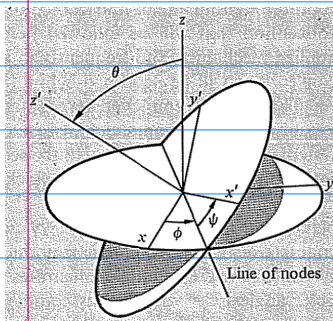
$$(1 + \epsilon_1)(1 + \epsilon_2) = 1 + \epsilon_1 + \epsilon_2 + \epsilon_1 \epsilon_2$$

$$(1 + \epsilon_2)(1 + \epsilon_1) = 1 + \epsilon_2 + \epsilon_1 + \epsilon_2 \epsilon_1$$

so the order of rotations is not important - they commute

Example, use Euler transformation matrix

$$A = \begin{pmatrix} 1 & (d\phi + d\psi) & 0 \\ -(d\phi + d\psi) & 1 & d\theta \\ 0 & -d\theta & 1 \end{pmatrix}$$



$$d\bar{\Omega} = \hat{i} d\theta + \hat{k} (d\phi + d\psi)$$

Inverse operation

$$B = 1 + \epsilon \quad B^{-1} = 1 - \epsilon$$

$$BB^{-1} = 1 - \epsilon^2 = 1$$

Also, because of orthogonality

$$B^{-1} = B^T$$

$$(1 - \epsilon) = 1 + \epsilon^T \Rightarrow -\epsilon = \epsilon^T$$

antisymmetric with zero diagonal elements

$$\begin{pmatrix} a & b & c \\ d & e & f \\ g & h & i \end{pmatrix}$$

$$a = e = i = 0$$

$$b = -d$$

$$c = -g$$

$$h = -f$$

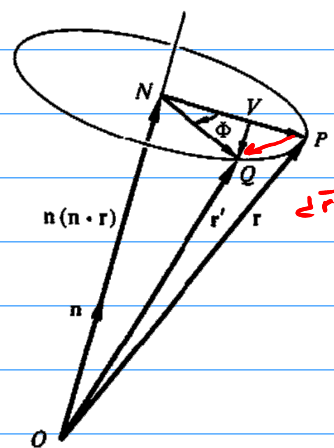
only three independent parameters

$$\epsilon = \begin{pmatrix} 0 & d\Omega_3 & -d\Omega_2 \\ -d\Omega_3 & 0 & d\Omega_1 \\ d\Omega_2 & -d\Omega_1 & 0 \end{pmatrix}$$

$$\mathbf{r}' = (1 + \epsilon)\mathbf{r} \Rightarrow \mathbf{r}' - \mathbf{r} = \epsilon\mathbf{r} \equiv d\mathbf{r}$$

$$\epsilon\mathbf{r} = \begin{pmatrix} 0 & d\Omega_3 & -d\Omega_2 \\ \dots\dots\dots \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix}$$

$$= \begin{pmatrix} x_2 d\Omega_3 - x_3 d\Omega_2 \\ -x_1 d\Omega_3 + x_3 d\Omega_1 \\ x_1 d\Omega_2 - x_2 d\Omega_1 \end{pmatrix} = \begin{pmatrix} dx_1 \\ dx_2 \\ dx_3 \end{pmatrix}$$



$$dx_1 = x_2 d\Omega_3 - x_3 d\Omega_2$$

$$dx_2 = x_3 d\Omega_1 - x_1 d\Omega_3$$

$$dx_3 = x_1 d\Omega_2 - x_2 d\Omega_1$$

elements of a cross product

$$d\mathbf{r} = \mathbf{r} \times d\mathbf{\Omega}$$

$$\mathbf{r}' = \mathbf{r} \cos\phi + \hat{n}(\hat{n} \cdot \mathbf{r})(1 - \cos\phi) + (\mathbf{r} \times \hat{n}) \sin\phi$$

$\phi \rightarrow d\phi$

$$\mathbf{r}' = \mathbf{r} + (\mathbf{r} \times \hat{n}) d\phi \Rightarrow d\mathbf{r} = (\mathbf{r} \times \hat{n}) d\phi$$

$$d\bar{\varphi} = \hat{n} d\phi$$

from now on, we want to treat transformations of vectors in the counter-clockwise sense

$$\bar{r}' = r \cos \phi + \hat{n} (\hat{n} \cdot \bar{r}) (1 - \cos \phi) + (\hat{n} \times \bar{r}) \sin \phi$$

$$d\bar{r} = d\bar{\omega} \times \bar{r}$$

$$(dG)_s = (dG)_{\text{body}} + d\bar{\omega} \times G$$

$$(d\bar{r})_s = (d\bar{r})_r + d\bar{\omega} \times \bar{r}$$

$$\left(\frac{d}{dt}\right)_s = \left(\frac{d}{dt}\right)_r + \bar{\omega} \times$$