

# Initial Boundary Value Problems

$$\frac{\partial u}{\partial t} = f\left(u, \frac{\partial u}{\partial x}, \frac{\partial^2 u}{\partial x^2}, \dots\right)$$

Evolution  
equation

$$u_t = f(u, u_x, u_{xx}, \dots)$$

$$u = u(x, t) \quad x \in [0, 1] \quad t \in [0, T]$$

$$u(x, 0) = \eta(x)$$

$$u(0, t) = \alpha(t)$$

$$u(1, t) = \beta(t)$$

## Diffusion equation

$$u_t = k u_{xx} + f(x)$$



$$u_t = u_{xx}$$

$$u(0, t) = 0$$

$$u(1, t) = 0$$

Exact solution:

$$u(x,t) = \sum_{p=0}^{\infty} \hat{u}_p(t) \sin(p\pi x)$$

$$u(x,0) = \sum_{p=0}^{\infty} \hat{u}_p(0) \sin(p\pi x) = \eta(x)$$

Substitution into the PDE gives:

$$\sum_p \hat{u}'_p(t) \sin(p\pi x) = \sum_p \hat{u}_p(t) \cdot (-p^2 \pi^2) \sin(p\pi x)$$

$$\hat{u}'_p(t) = -p^2 \pi^2 \hat{u}_p(t)$$

$$\hat{u}_p(t) = e^{-p^2 \pi^2 t} \hat{u}_p(0) \Rightarrow \sum_{p=0}^{\infty} \hat{u}_p(0) e^{-p^2 \pi^2 t} \sin(p\pi x)$$

# Discretization (method of lines)

First discretize in space:

$$x_j = jh \quad j=0, 1, \dots, m+1$$

$$h = \frac{1}{m+1}$$

$$U_j(t) \approx u(x_j, t)$$

$$u_{xx}(x_j, t) = \frac{U_{j+1}(t) - 2U_j(t) + U_{j-1}(t)}{h^2}$$

$$U_j'(t) = \frac{U_{j+1}(t) - 2U_j(t) + U_{j-1}(t)}{h^2}$$

Now discretize in time using RK, LM, etc.

$$\begin{bmatrix} U_1 \\ \vdots \\ U_m \end{bmatrix}' = \frac{1}{h^2} \begin{bmatrix} -2 & 1 & & \\ & \ddots & \ddots & \\ & & 1 & -2 \end{bmatrix} \begin{bmatrix} U_1 \\ \vdots \\ U_m \end{bmatrix}$$

$A$

$$U'(t) = AU$$

$$U(t) = e^{tA} U(0).$$

$$A = R \Lambda R^{-1} \text{ where } \lambda_p = \frac{2}{h^2} (\cos(p\pi h) - 1)$$

$$r_{jp} = \sin(p\pi jh)$$

$$U'(t) = R \Lambda R^{-1} U(t)$$

$$R^{-1} U'(t) = \Lambda R^{-1} U(t)$$

$$\hat{U}'(t) = \Lambda \hat{U}(t)$$

$$\hat{U}'_p(t) = \lambda_p \hat{U}_p(t)$$

$$\hat{U}_p(t) = e^{t\lambda_p} \hat{U}_p(0)$$

analogous to

$$\hat{U}_p(t) = e^{-p^2 \pi^2 t} \hat{U}_p(0)$$

The factor  $-p^2 \pi^2$  has been replaced by

$$\lambda_p = \frac{2}{h^2} (\cos(p\pi h) - 1)$$

$$= \frac{2}{h^2} \left( -\frac{p^2 \pi^2 h^2}{2} + \mathcal{O}(h^4) \right)$$

$$= -p^2 \pi^2 + \mathcal{O}(h^2)$$

(for small  $ph$ )

So the numerical "modes" decay at almost the same rate as in the exact solution, for small  $ph$ .



# Stability

Use Euler's method in time:

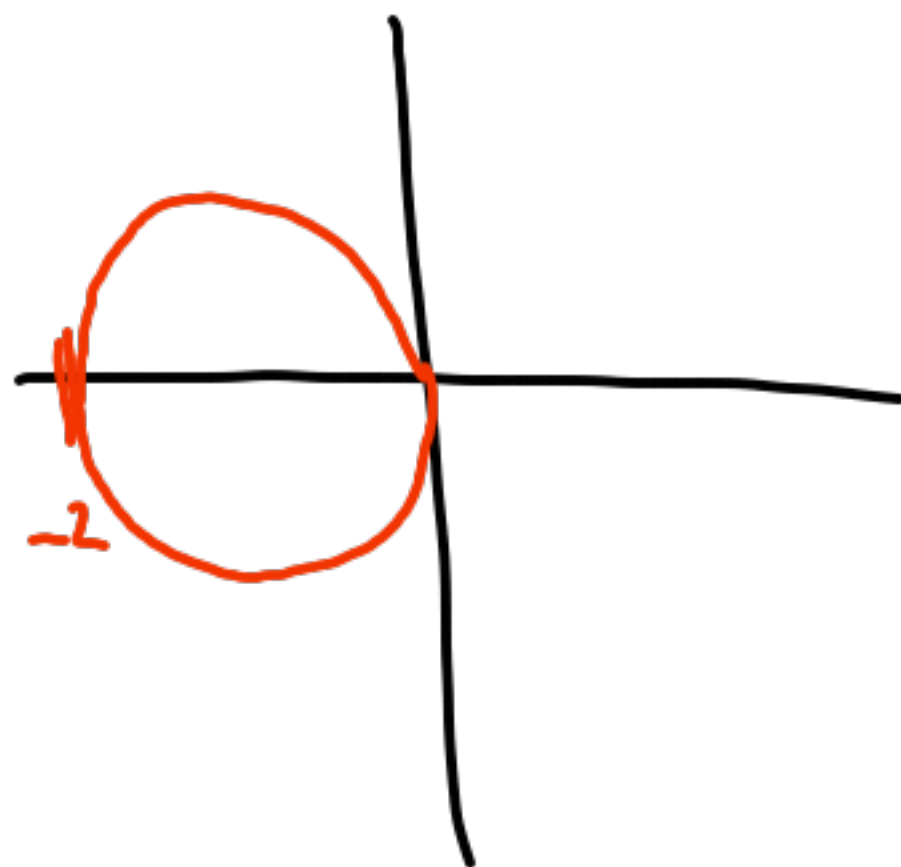
$$U^{n+1} = U^n + \Delta t A U^n$$

For absolute stability,  
We need:

$$-2 \leq K \lambda_p \leq 0$$

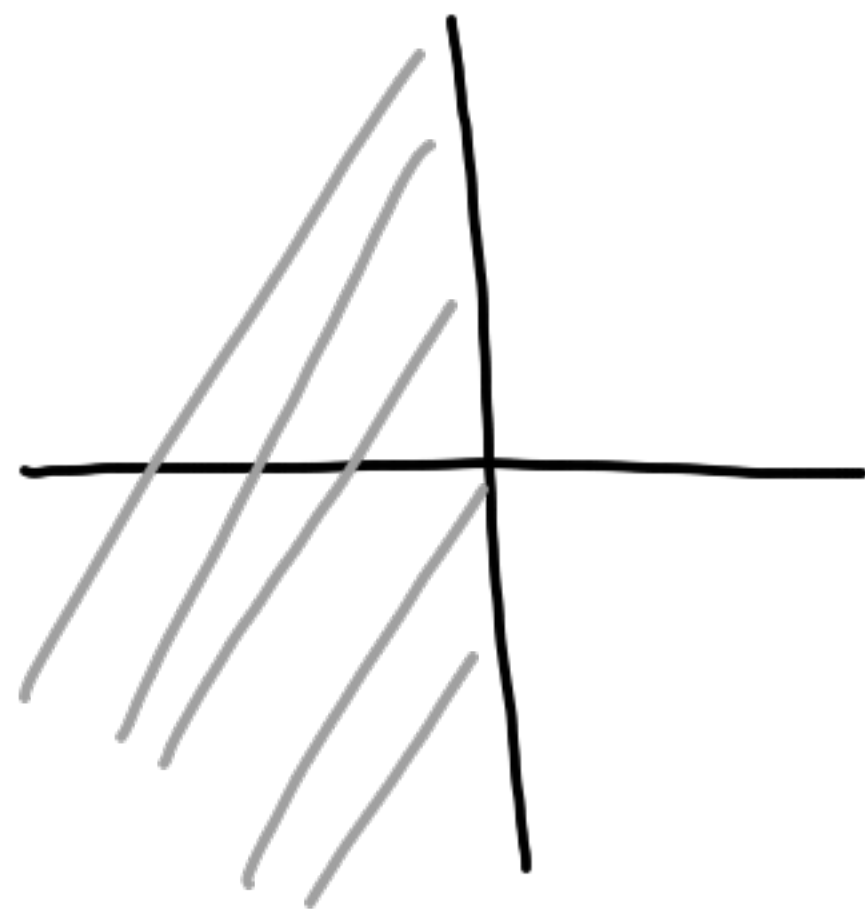
The largest magnitude eigenvalue is  $\lambda_p \approx -\frac{4}{h^2}$

$$-2 \leq -\frac{4K}{h^2} \Rightarrow K \leq \frac{h^2}{2}$$



The original IBVP is infinitely stiff since  $\lim_{p \rightarrow \infty} p^2 \tau^2 = \infty$ , but our discretized problem has finite stiffness.

We should use an A-stable or A( $\alpha$ )-stable method, so we have absolute stability with any step size.



# Diagonally Implicit RK (DIRK)

$$Y_i = U^n + k \sum_{j=1}^i a_{ij} f(Y_j)$$

$$U^{n+1} = U^n + k \sum_{j=1}^r b_j f(Y_j)$$

c	A
	b <sup>r</sup>

A is lower-triangular

More efficient than fully implicit RK methods since we can solve each stage sequentially.

For example:

0	0	0	0
1/2	1/4	1/4	0
1	1/3	1/3	1/3
	1/3	1/3	1/3

A-stable

L-stable

2nd-order accurate

TR-BDF2

Fully-discrete Scheme (Euler):

$$U_j^{n+1} = U_j^n + \frac{K}{h^2} (U_{j+1}^n - 2U_j^n + U_{j-1}^n)$$

To get LTE:

$$U(x_j, t_{n+1}) = U(x_j, t_n) + \frac{K}{h^2} (U(x_{j+1}, t_n) - 2U(x_j, t_n) + U(x_{j-1}, t_n)) + \tau_{j+1}^n$$

After using Taylor series we get

$$\tau_j^n = \frac{K}{2} u_{tt} - \frac{h^2}{12} u_{xxx} + O(K^2) + O(h^4)$$

1st order in time

2nd order in space