So far we've studied: \
-An elliptic PDE:

J2 W=f

- A parabolic PDE: $U_t = \nabla^2 U$

Today: a Hyperbolic PDE

Hyperbolic PDEs model waves:

- Water waves

- Sound (pressure waves)

-EM mones

- Fluid dynamics

Consider flow of a fluid in a channel:

U(x,t): Concentration of Some quantity X, X2 (per unit length)

Total amount in [x,,x2]:

 $\int_{x_1}^{x_2} u(x_3t) dx$

This should only change due to Flow through the endpoints:

 $\frac{d}{dt} \int_{x_1}^{x_2} u dx = f(u(x_1,t)) - f(u(x_2,t))$

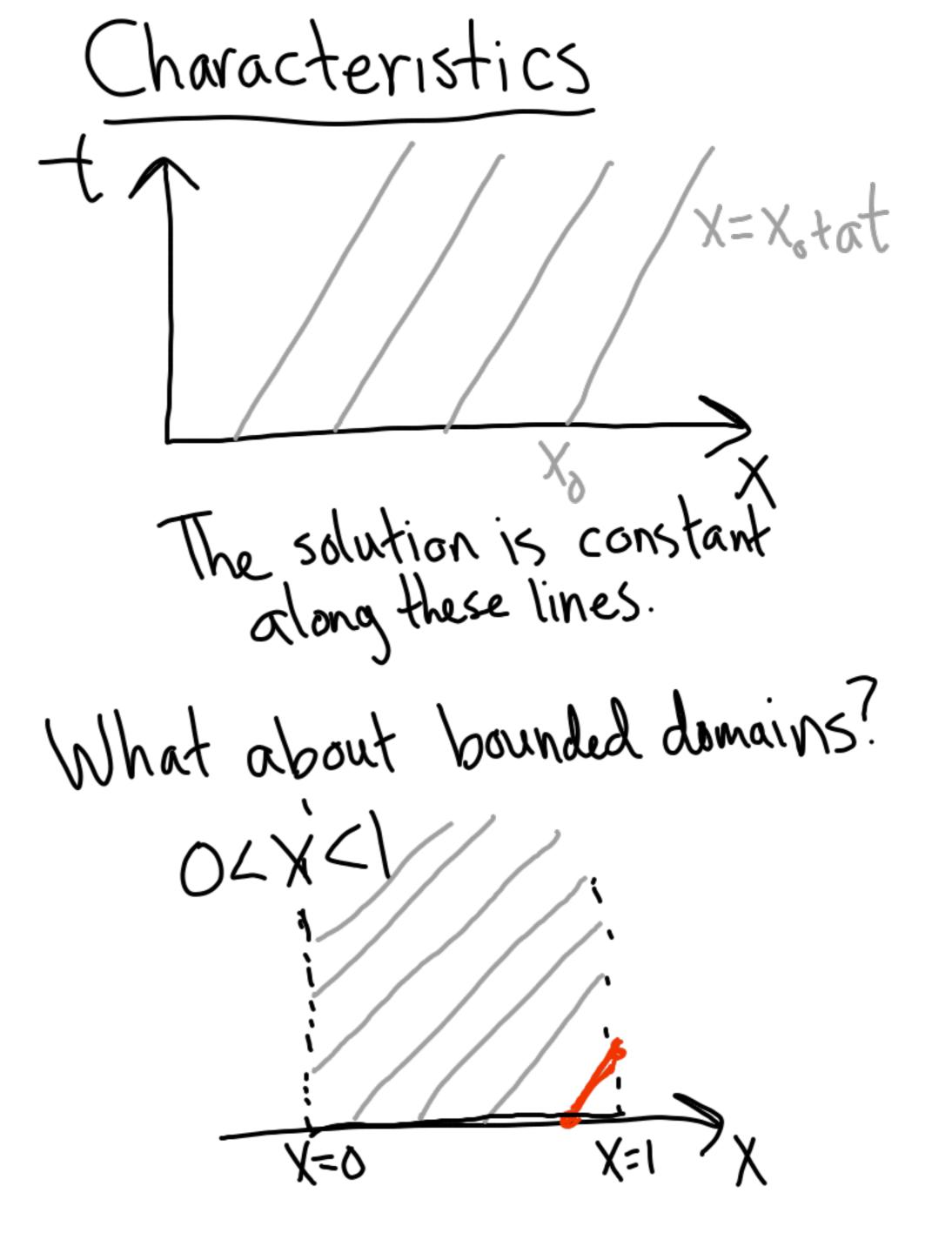
Here f(u) is the flux (rate of flow) If u is smooth enough:

 $\int_{x_{i}}^{x_{i}} \frac{\partial}{\partial t} u(x_{i}t) dx = -\int_{x_{i}}^{x_{i}} \frac{\partial}{\partial x} f(u) dx$

 $\int_{X_1}^{X_2} \left(\frac{\partial U}{\partial t} + \frac{\partial f}{\partial x} \right) dx = 0 \quad \text{law}$

This implies: $U_t + f(u)_x = 0$ (pointwise)

Let's just take a Constant Velocity: f(u) = au(x,t) Ut + aux=0 Advection $U(x)t=0)=\eta(x)$ Solution: $u(x,t) = \eta(x-at)$ Theck it: -an' tan'=0 (For the Cauchy Problem)



We must specify the Solution along each characteristic. Thus we need a boundary Condition u(x=0,t)=x(t) if a>0 $U(x=yt)=B(t) if \alpha < 0.$

Discretization

Centered in space, Euler in time:

 $\int_{j}^{n+1} = \int_{j}^{n} - ka \frac{\int_{j+1}^{n} - \int_{j-1}^{n}}{2h}$

Is this stable?

von Neumann analysis

 $y = a^n e^{ijhg}$

ontheins = oneins-kagneins (ens-eins)

$$g = 1 - \frac{ka}{2h} \left(e^{ih\xi} - e^{-ik\xi} \right)$$

$$2isin(k\xi)$$

$$2isin(k\xi)$$

$$2isin(k\xi)$$

$$Q = 1 - \frac{ka}{h} isin(k\xi)$$

$$Q = 1 - \frac{ka}{h} isin(k\xi)$$
We want $|g| \le 1$ for all ξ .
$$|g|^2 = 1 + \frac{ka}{h} sin(k\xi)^2$$
 Unstable

Method of lines analysis

If we discretize only in space:

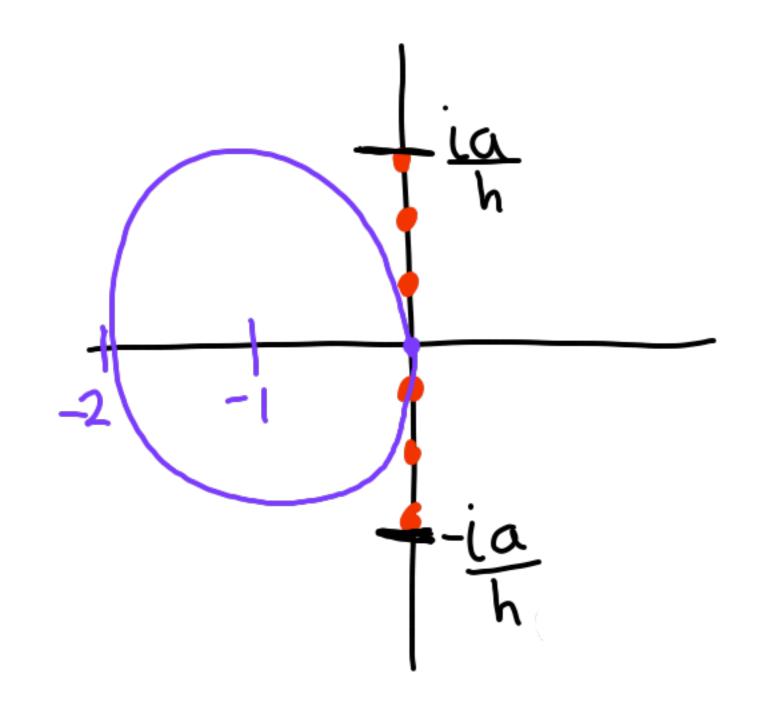
$$U_{j}'(t) = -\frac{a}{2h}(U_{j+1}(t) - U_{j-1}(t))$$

$$U'(t) = -\frac{a}{2h}(U_{j+1}(t) - U_{j-1}(t))$$

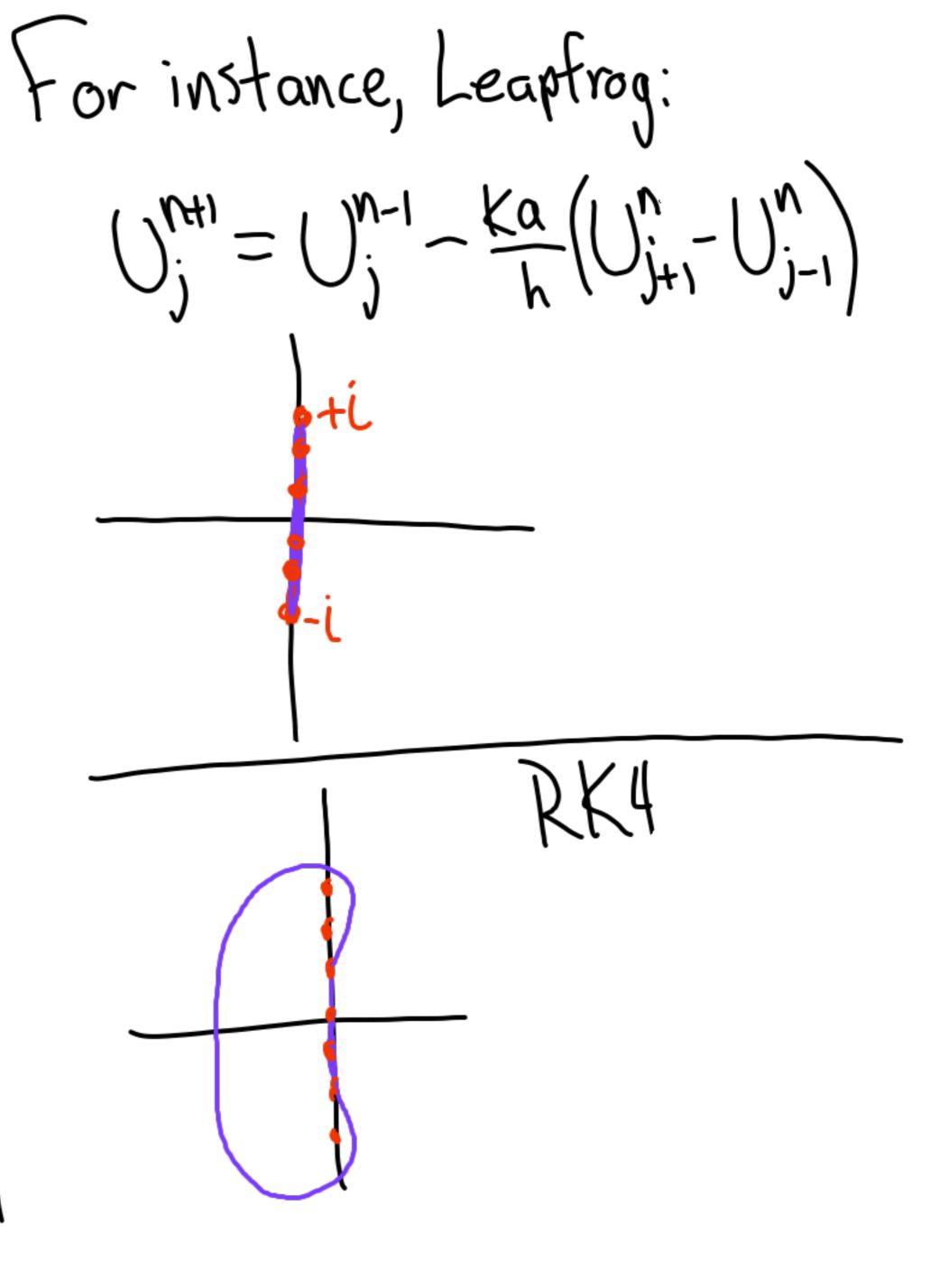
$$U'(t) = -\frac{a}{2h}(U_{j+1}(t) - U_{j-1}(t))$$

Do the eigenvalues lie inside the absolute stability region?

5 Kew- symmetric > imag. eigenvalues/



We should use a method that includes part of the imaginary axis!



$$U_{j}^{n+1} = \frac{U_{j+1}^{n} + U_{j-1}^{n}}{2} - \frac{K\alpha}{2h} \left(U_{j+1}^{n} - U_{j-1}^{n} \right)$$

Observe:
$$U_{j+1}^{n}+U_{j-1}^{n}=U_{j}^{n}+\frac{1}{2}(U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n})$$

So the LF method can be written:

So the LF method can be written:

$$U_{j}^{n+1}-U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n}}{h^{2}}$$
This books like an approximation of $U_{j}^{n}+U_{j}^{n}-U_{j}^{n}+U_{j}^{n}-1$.

$$U_{j}^{n+1}-U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n}}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n}}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n}}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j-1}^{n}}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j}^{n}-1}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j+1}^{n}-1}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j+1}^{n}-1}{h^{2}}$$

$$U_{j}^{n}+U_{j}^{n}=-\frac{a}{2h}(U_{j+1}^{n}-U_{j}^{n}-1)+\frac{h^{2}}{2k}\frac{U_{j+1}^{n}-2U_{j}^{n}+U_{j+1}^{n}-1}{h^{2}}$$

Advection-diffusion