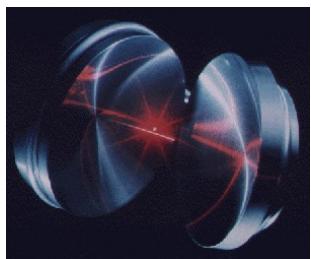
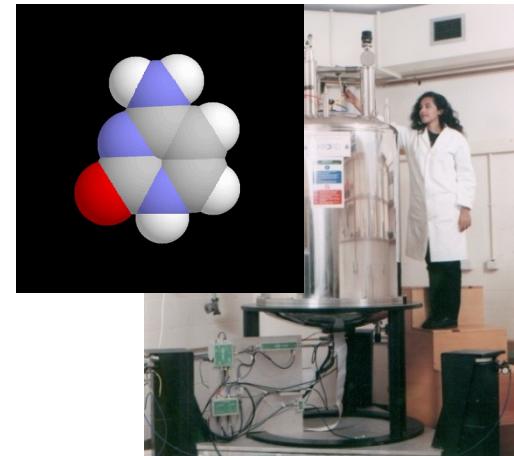
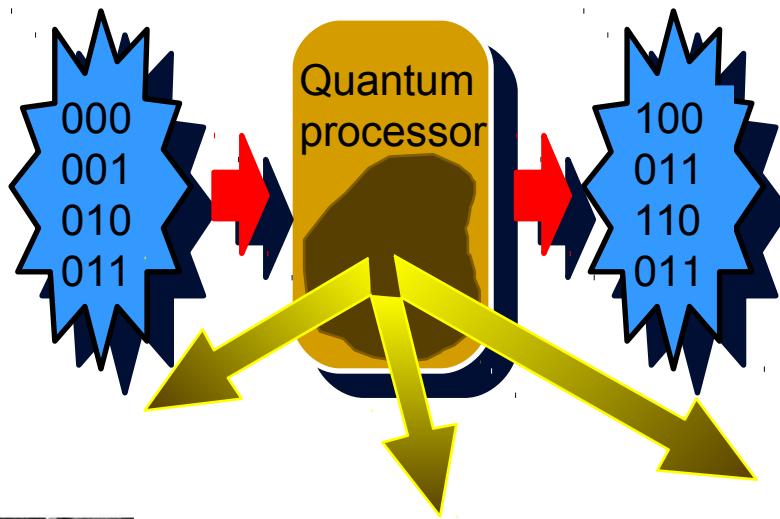


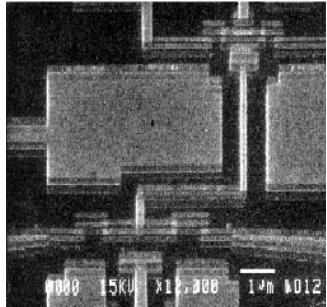
Which technology ?



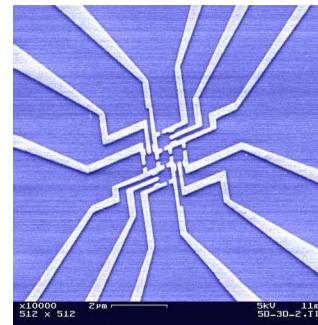
Cavity QED



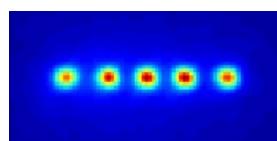
NMR



Superconducting qubits

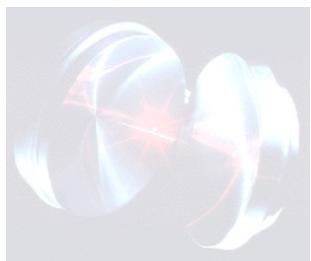


Quantum dots

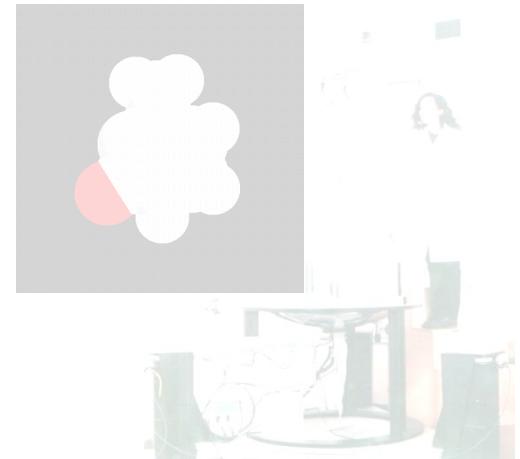
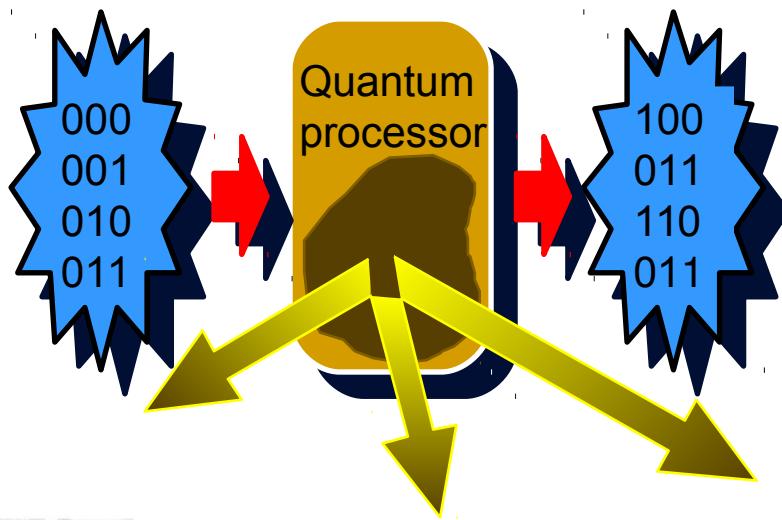


Trapped atoms/ions

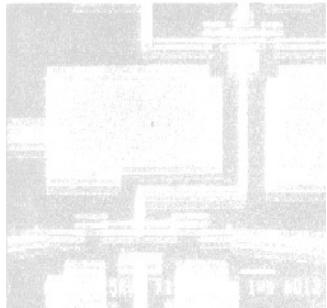
Which technology ?



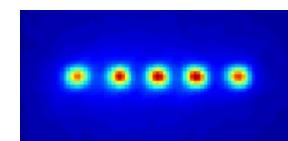
Cavity QED



NMR



Superconducting qubits



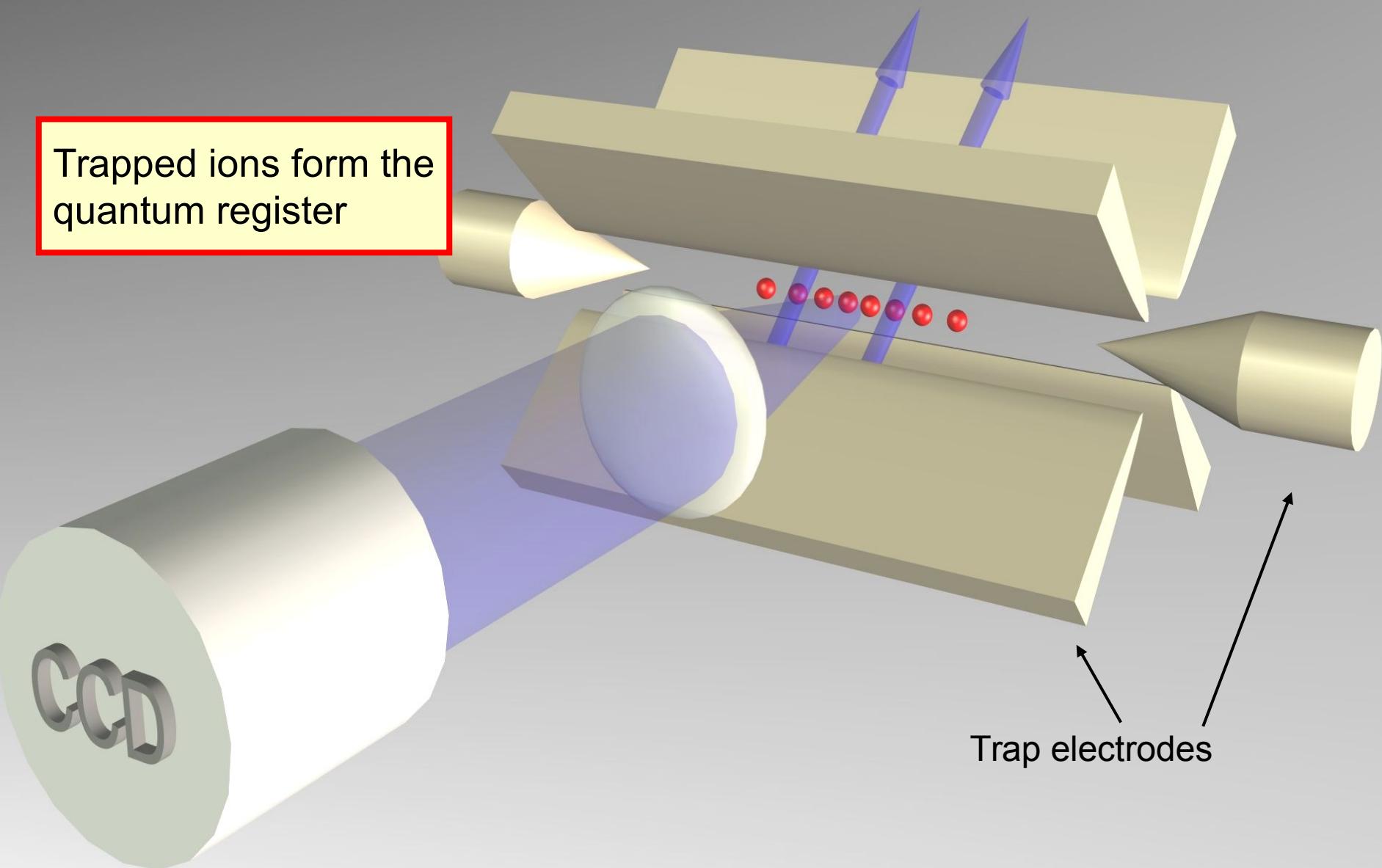
Trapped atoms/ions

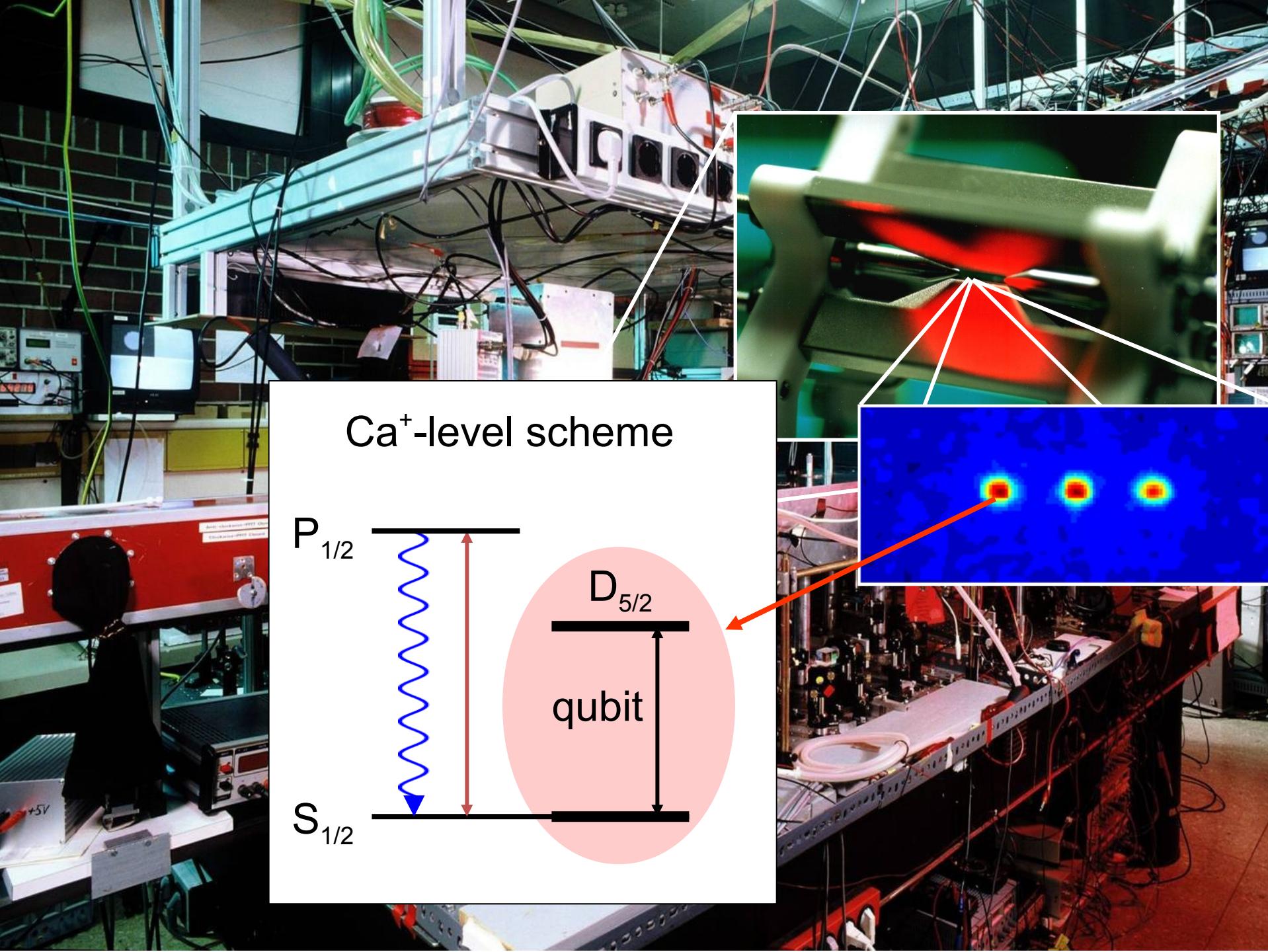


Quantum dots

Ion trap quantum computing

Trapped ions form the quantum register

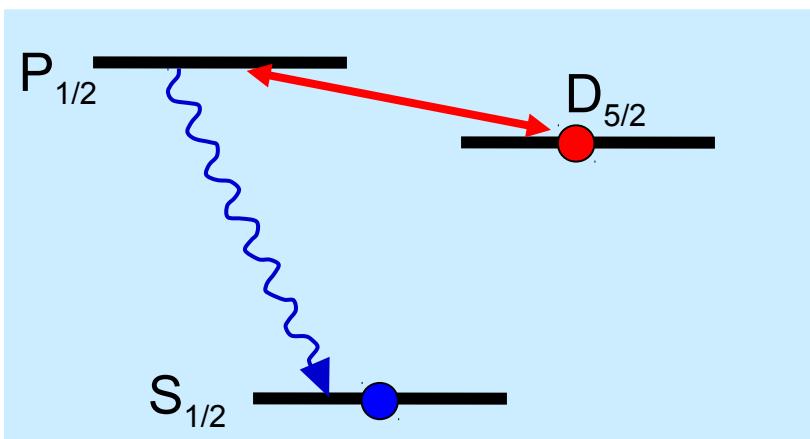




Di Vincenzo criteria

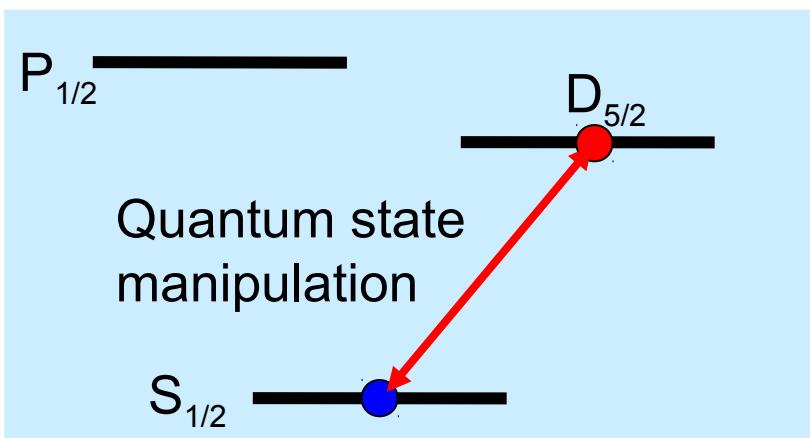
- I. Scalable physical system, well characterized qubits
- II. Ability to initialize the state of the qubits
- III. Long relevant coherence times, much longer than gate operation time
- IV. “Universal” set of quantum gates
- V. Qubit-specific measurement capability

Experimental procedure



1. Initialization in a pure quantum state

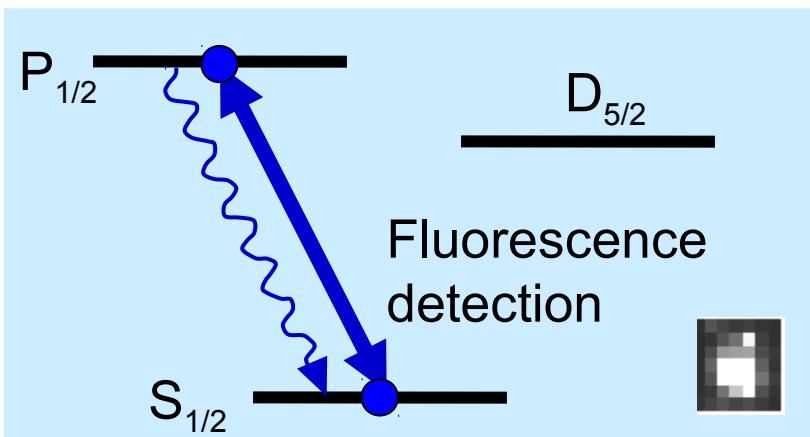
Experimental procedure



1. Initialization in a pure quantum state

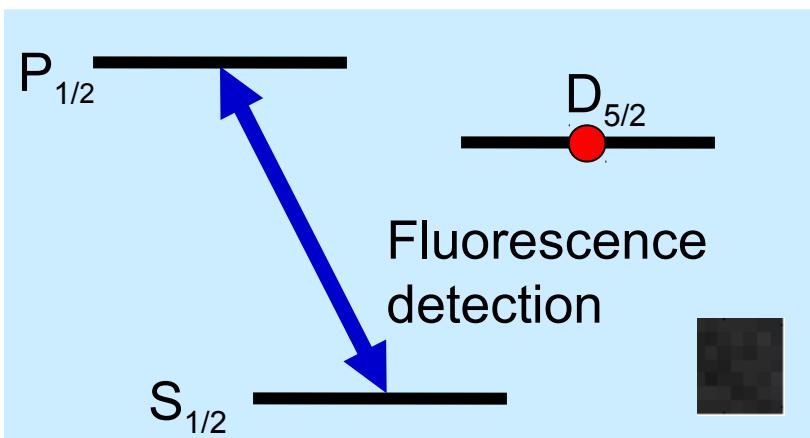
2. Quantum state manipulation on
 $S_{1/2} - D_{5/2}$ transition

Experimental procedure



1. Initialization in a pure quantum state
2. Quantum state manipulation on $S_{1/2} - D_{5/2}$ transition
3. Quantum state measurement by fluorescence detection

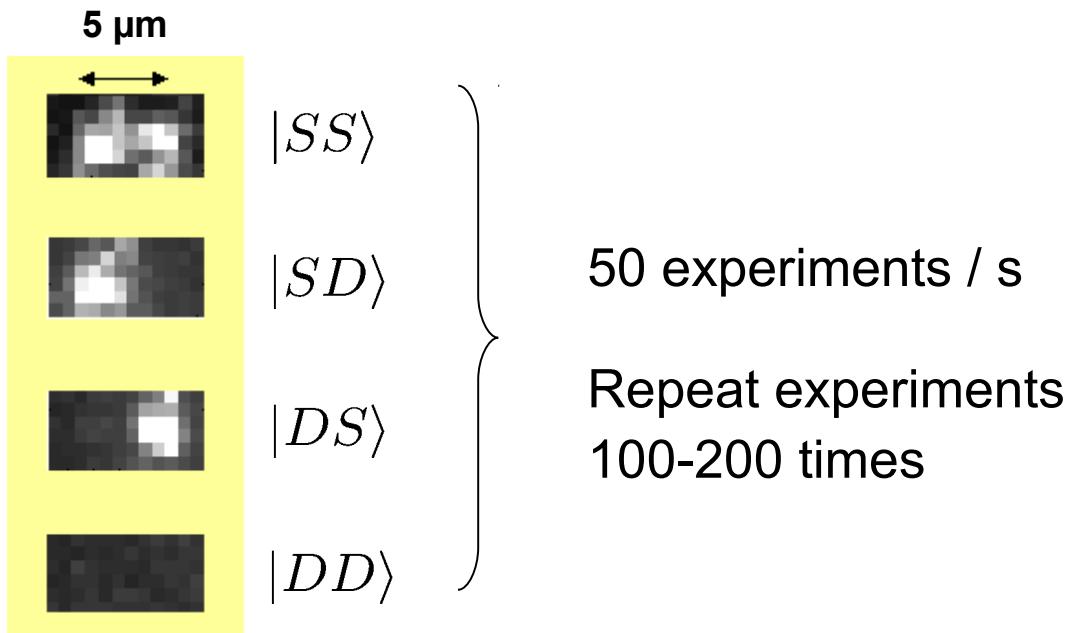
Experimental procedure



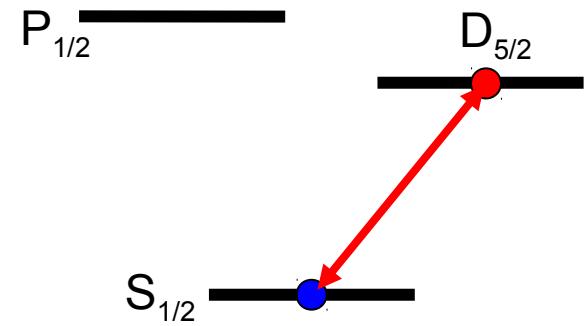
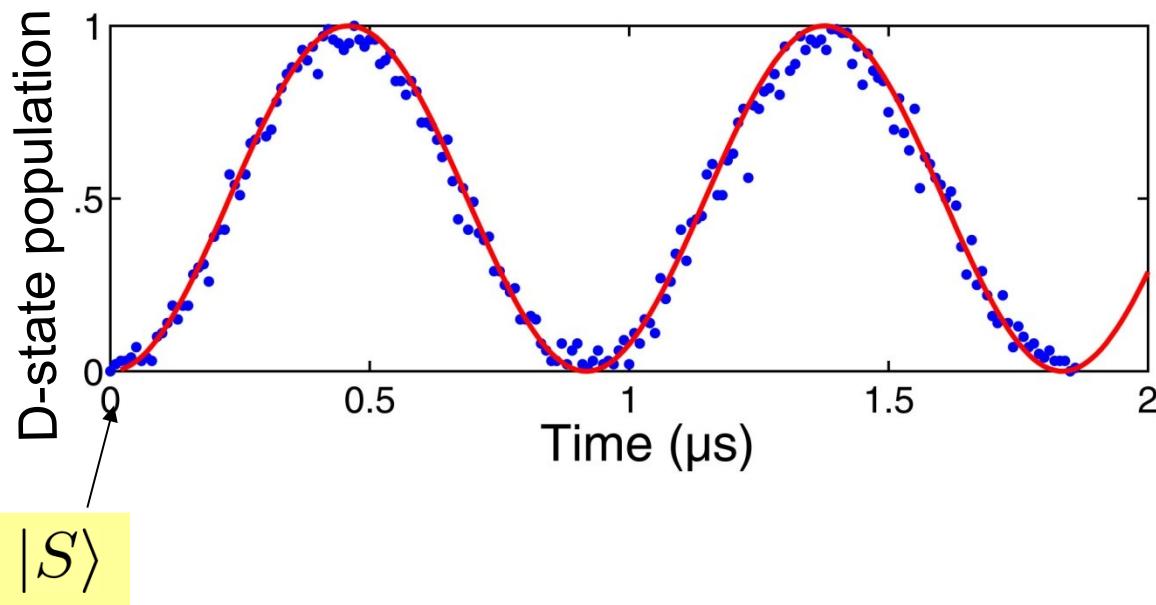
1. Initialization in a pure quantum state
2. Quantum state manipulation on $S_{1/2} - D_{5/2}$ transition
3. Quantum state measurement by fluorescence detection

Two ions:

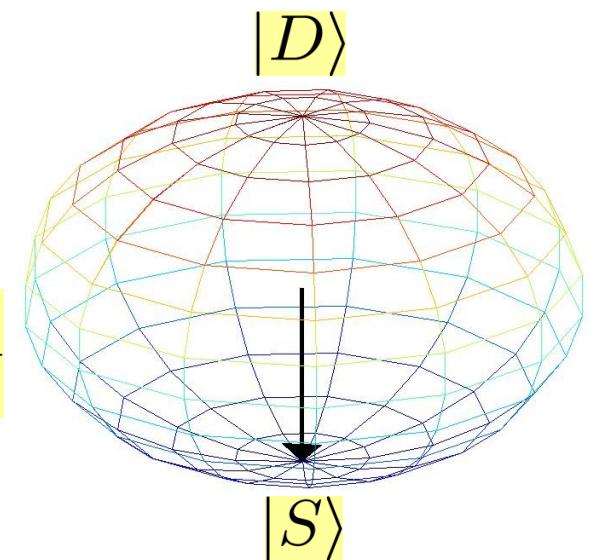
Spatially resolved
detection with
CCD camera



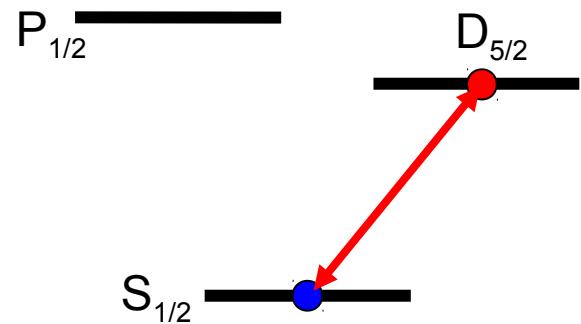
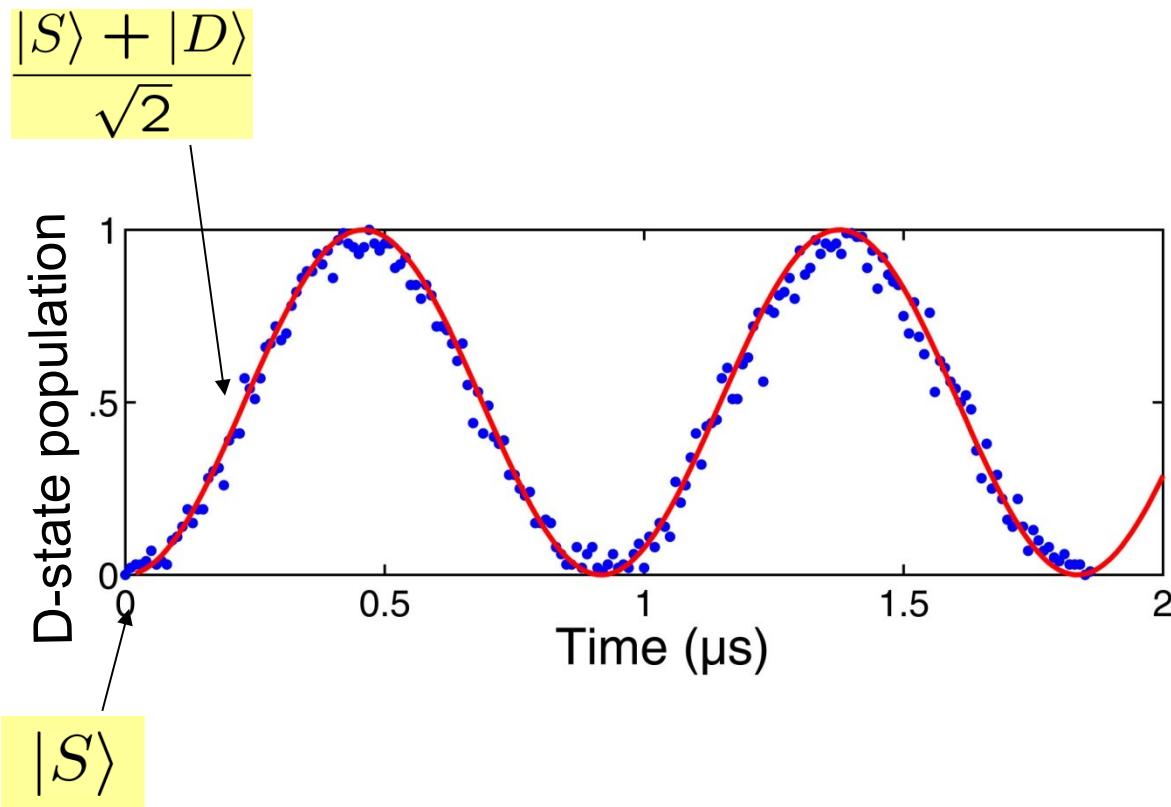
Rabi oscillations



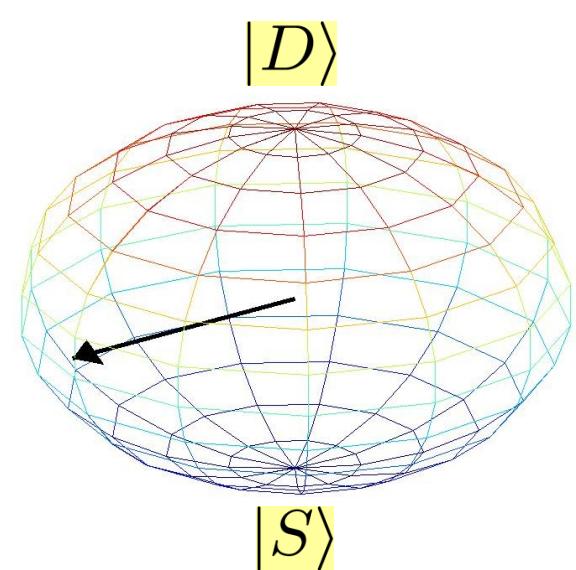
$$\frac{|S\rangle + |D\rangle}{\sqrt{2}}$$



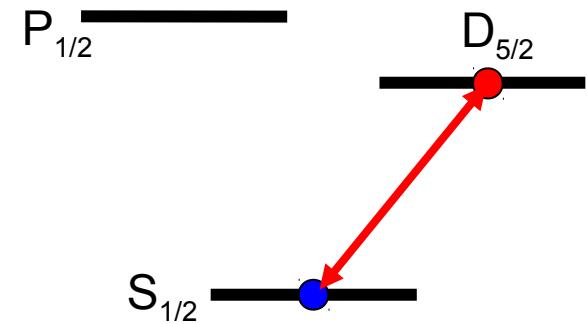
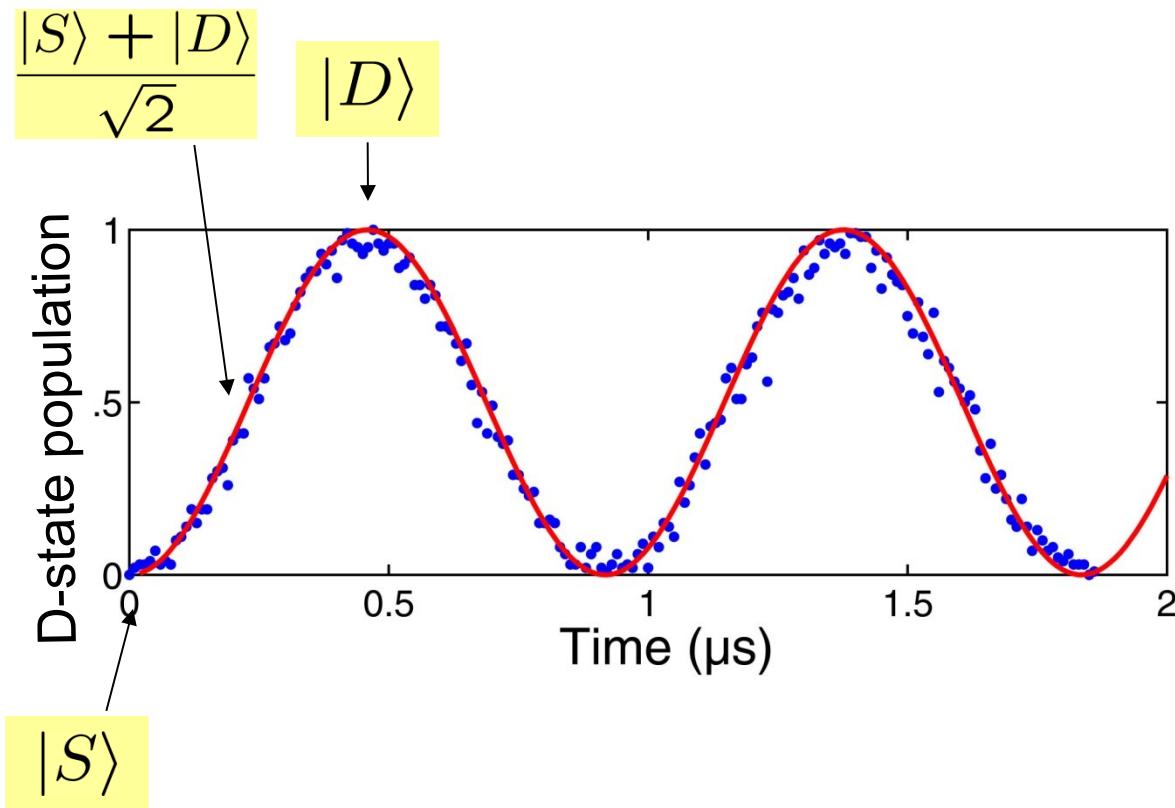
Rabi oscillations



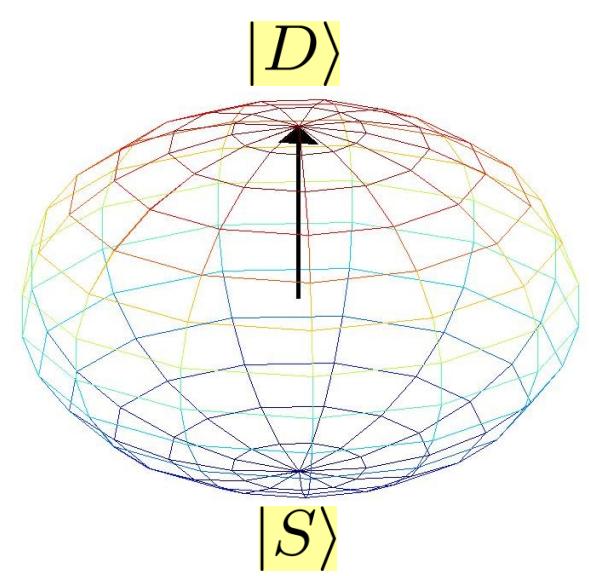
$$\frac{|S\rangle + |D\rangle}{\sqrt{2}}$$



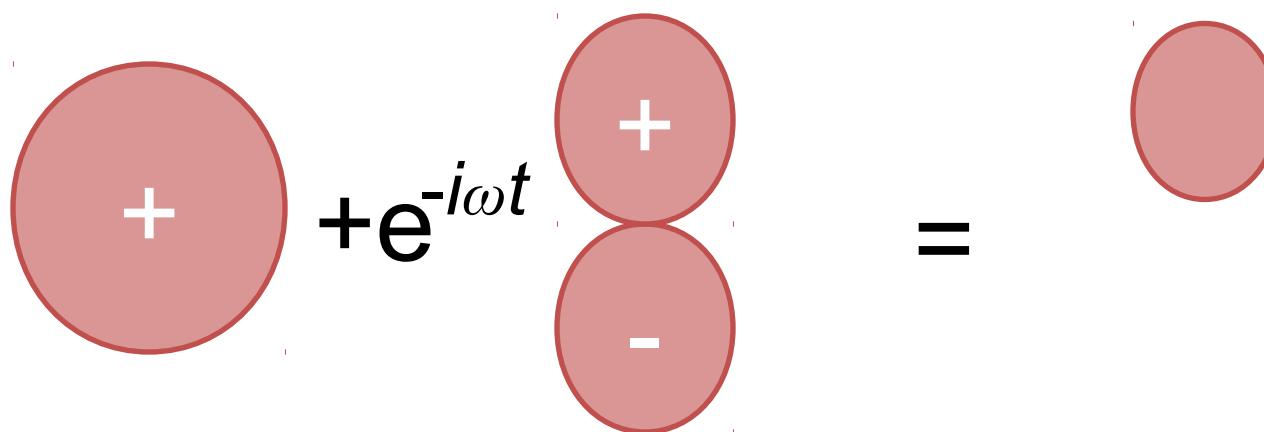
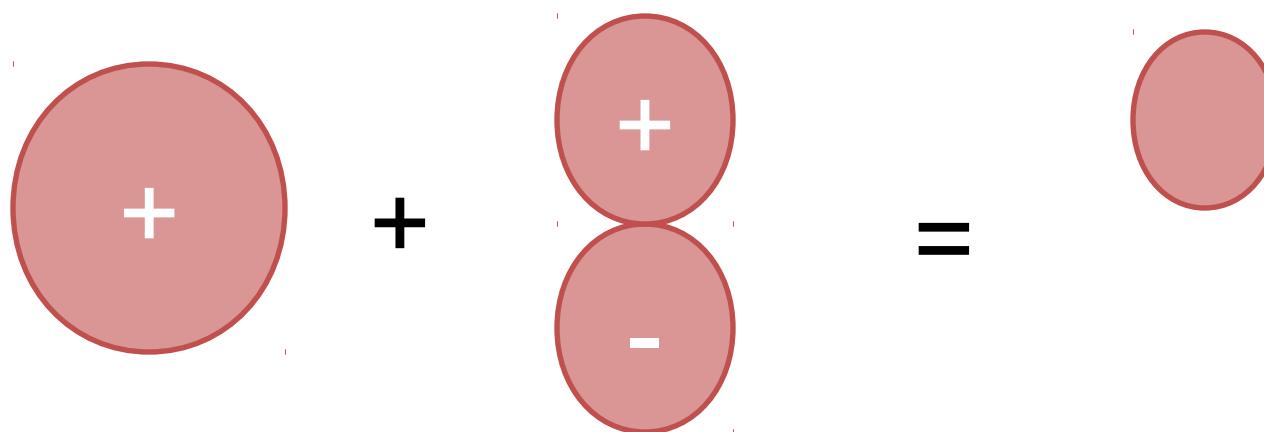
Rabi oscillations



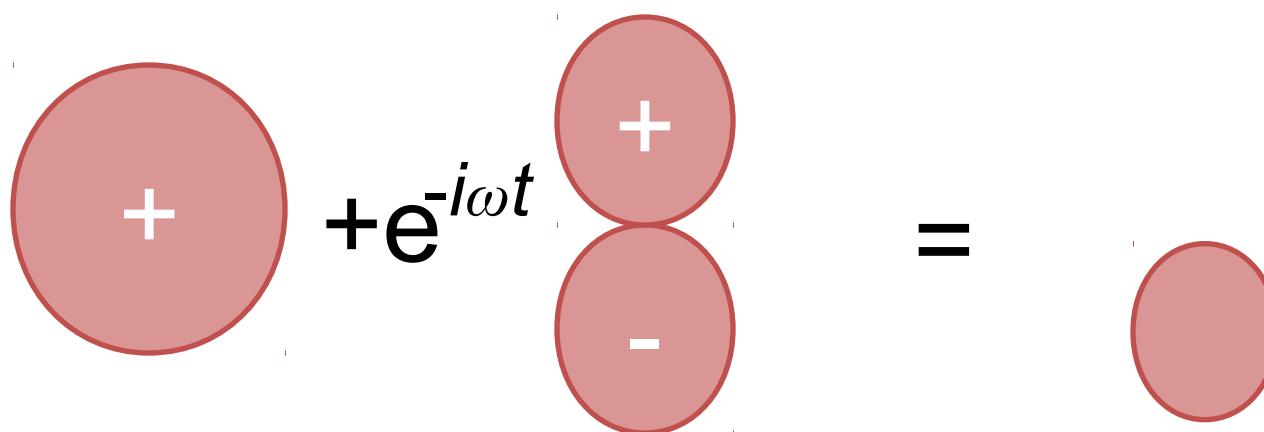
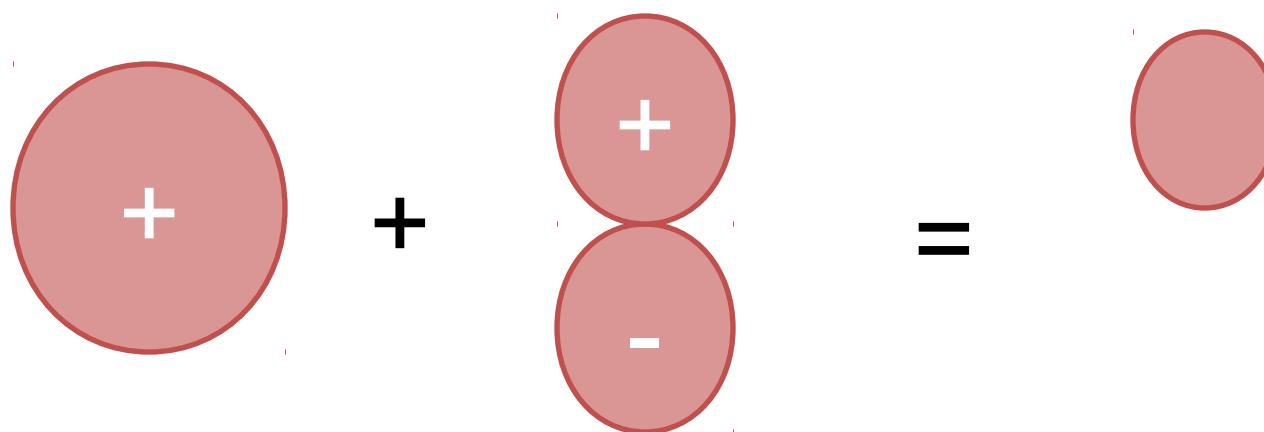
$$\frac{|S\rangle + |D\rangle}{\sqrt{2}}$$



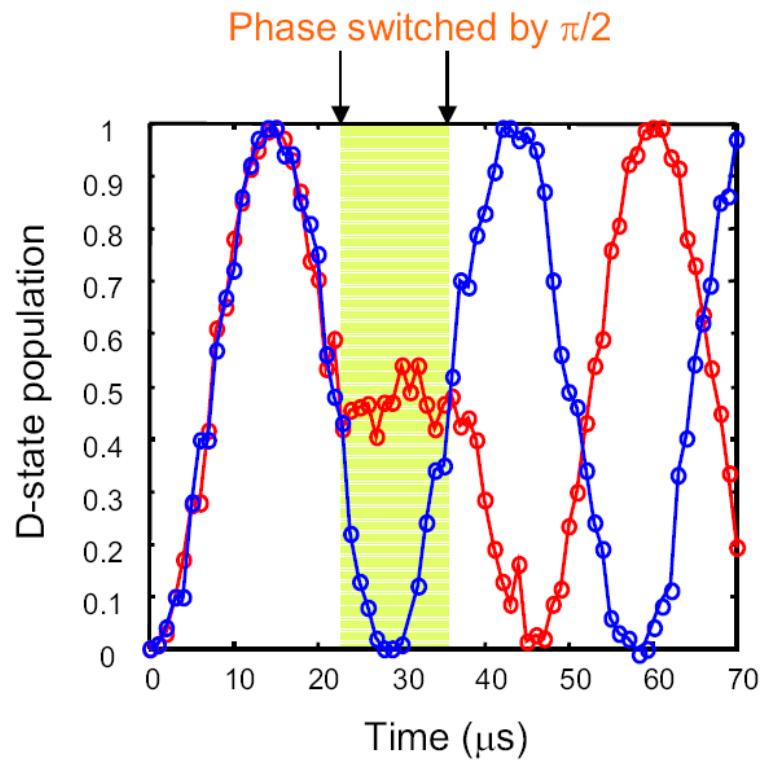
The phase ...



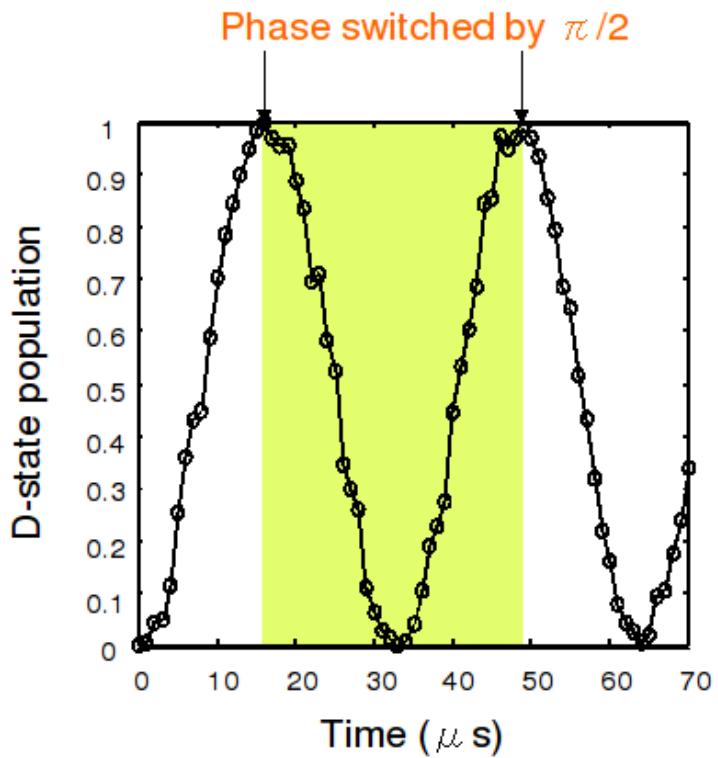
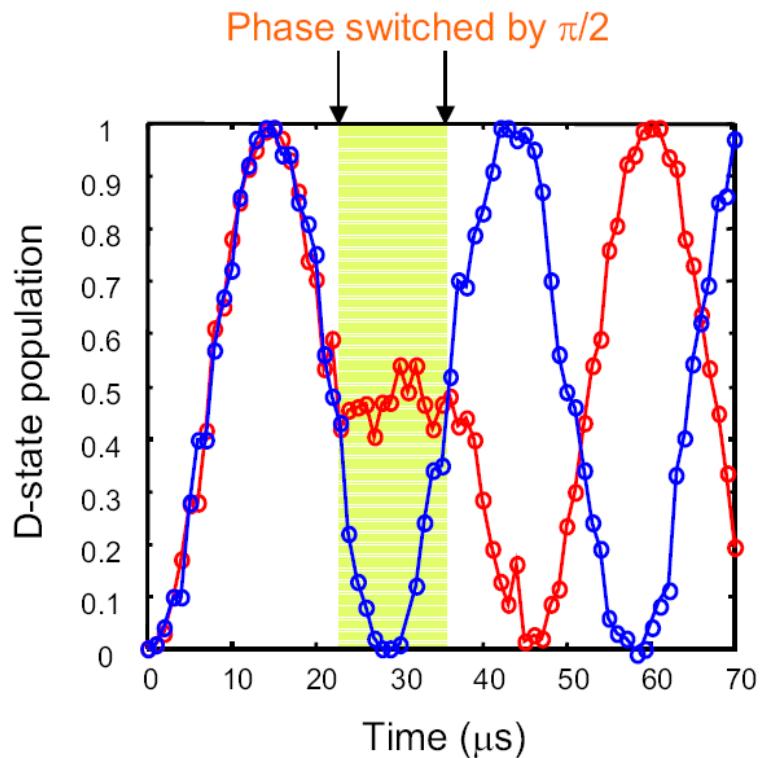
The phase ...



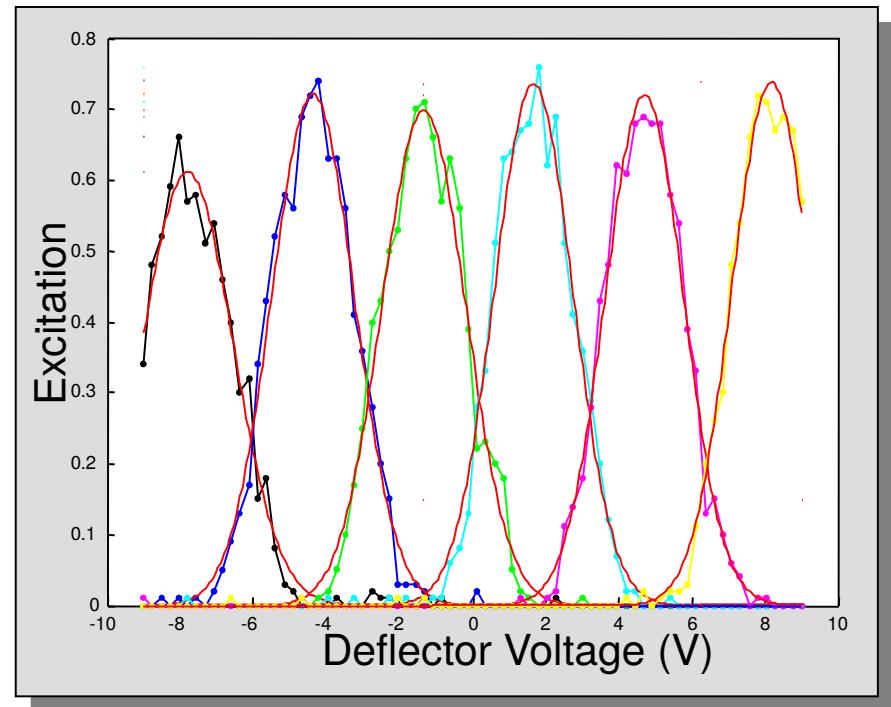
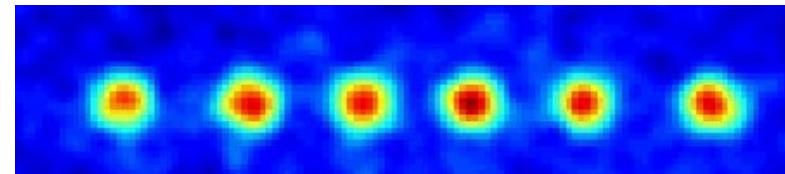
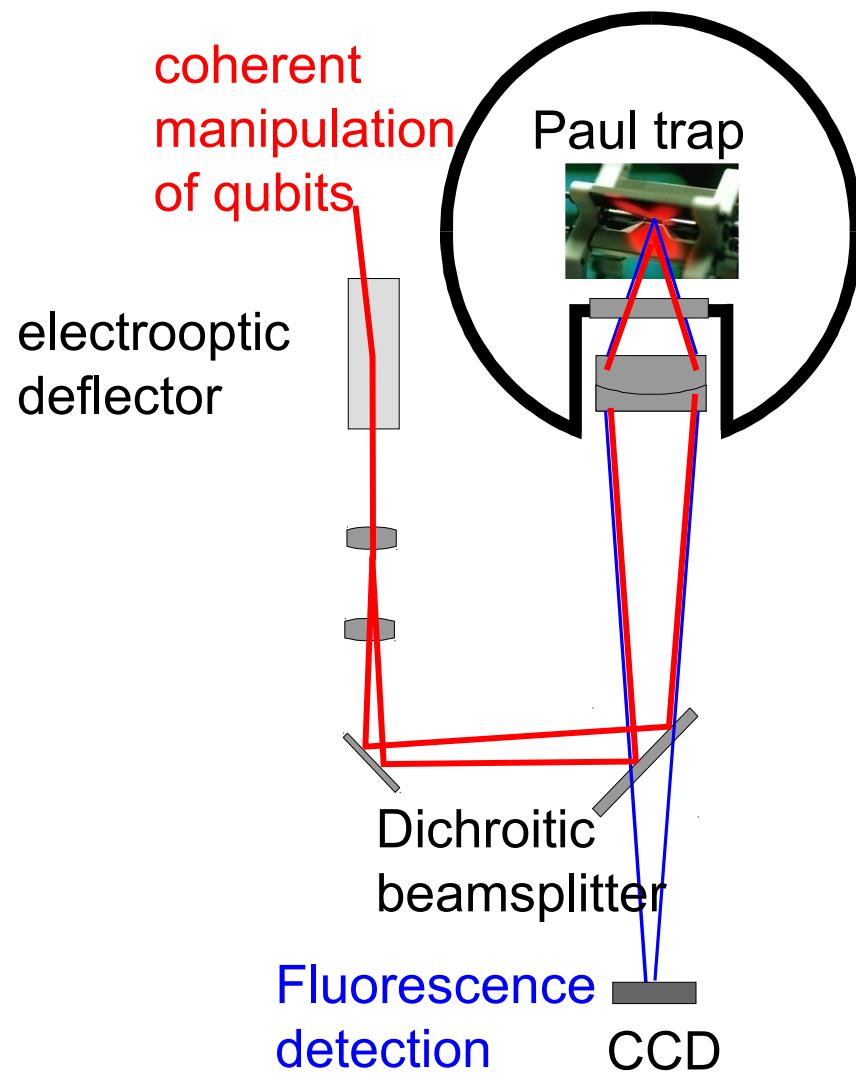
The phase ...



The phase ...



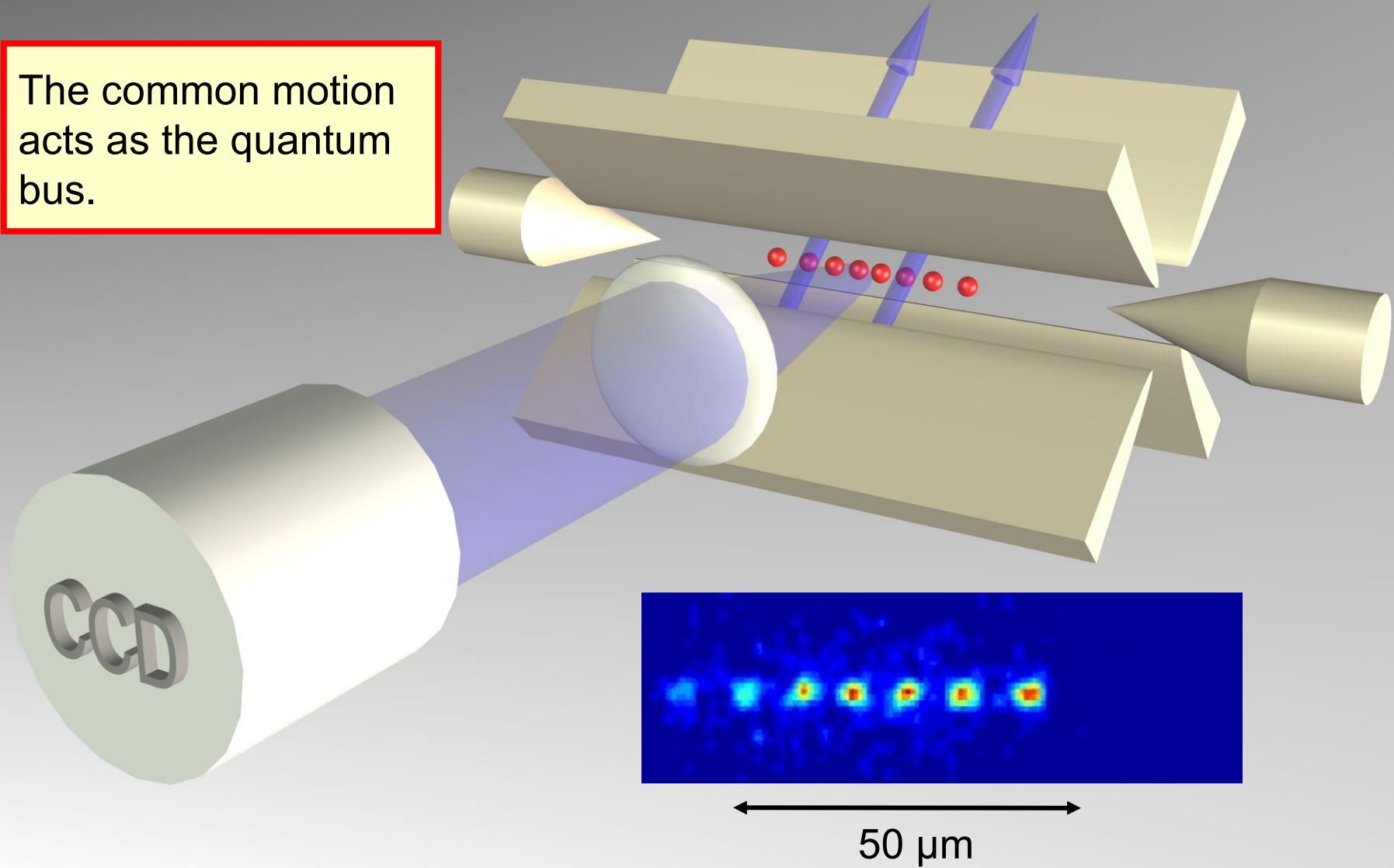
Addressing single qubits



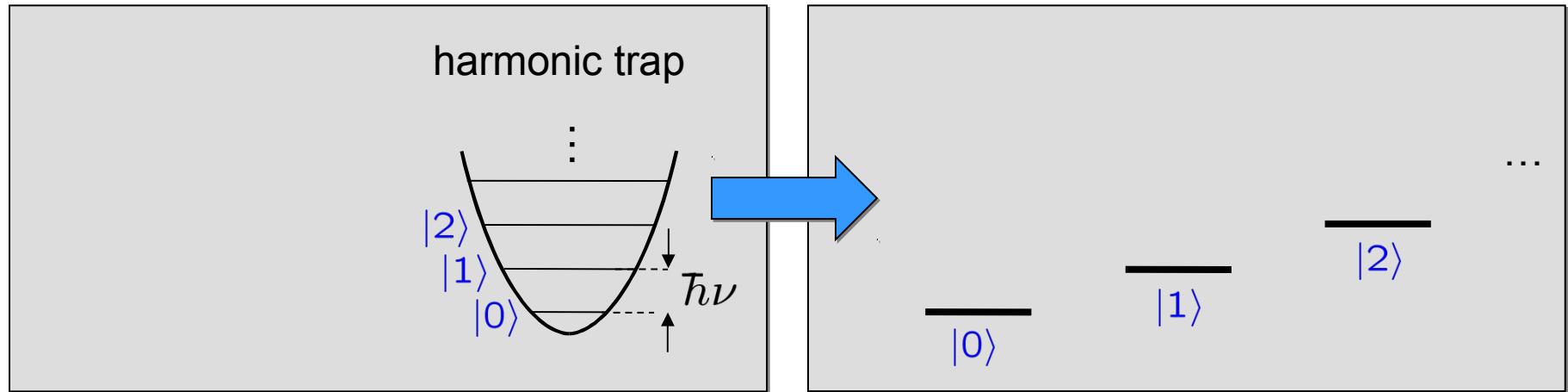
- inter ion distance: $\sim 4 \mu\text{m}$
- addressing waist: $\sim 2 \mu\text{m}$
- < 0.1% intensity on neighbouring ions

Having the qubits interact

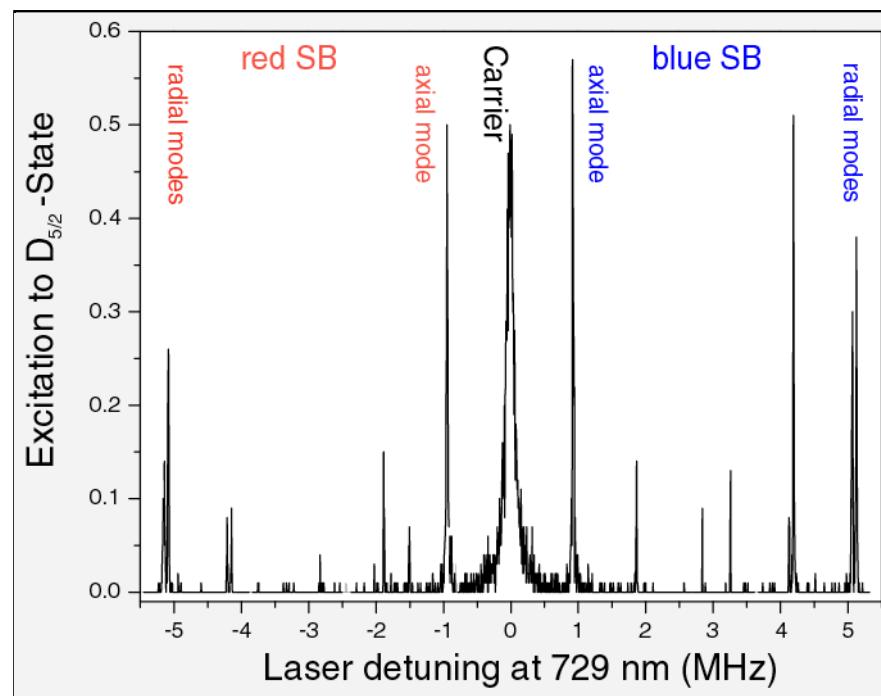
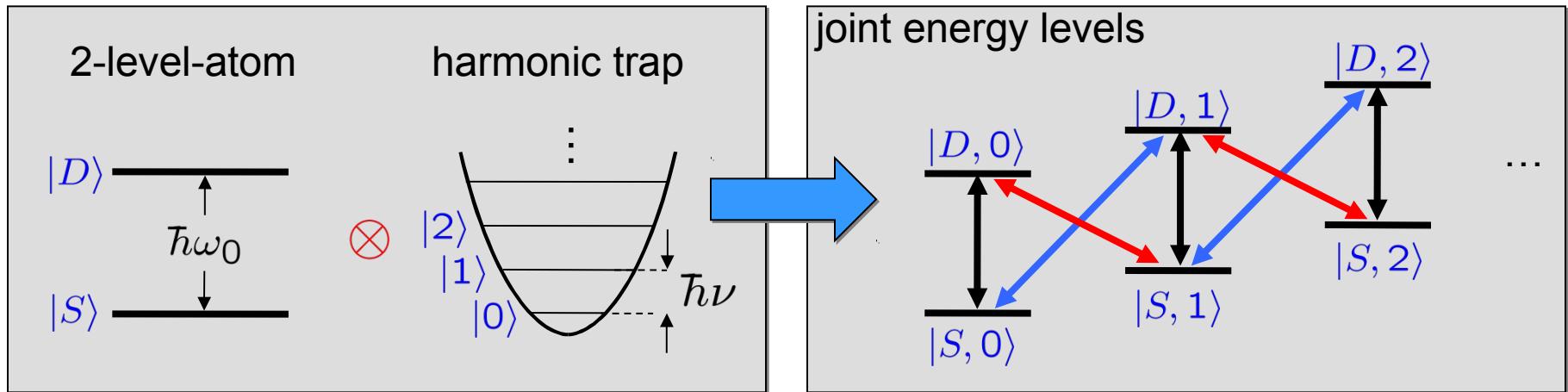
The common motion acts as the quantum bus.



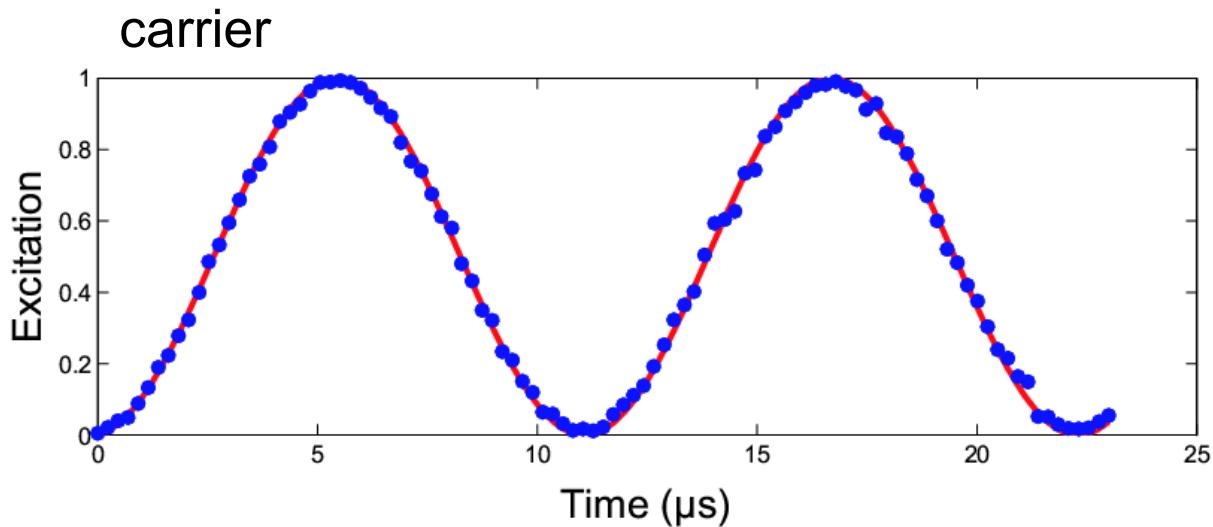
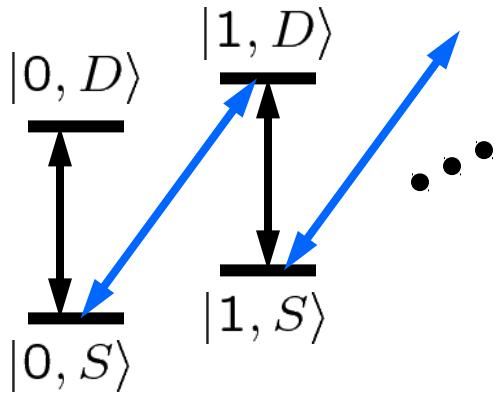
Ion motion



Ion motion

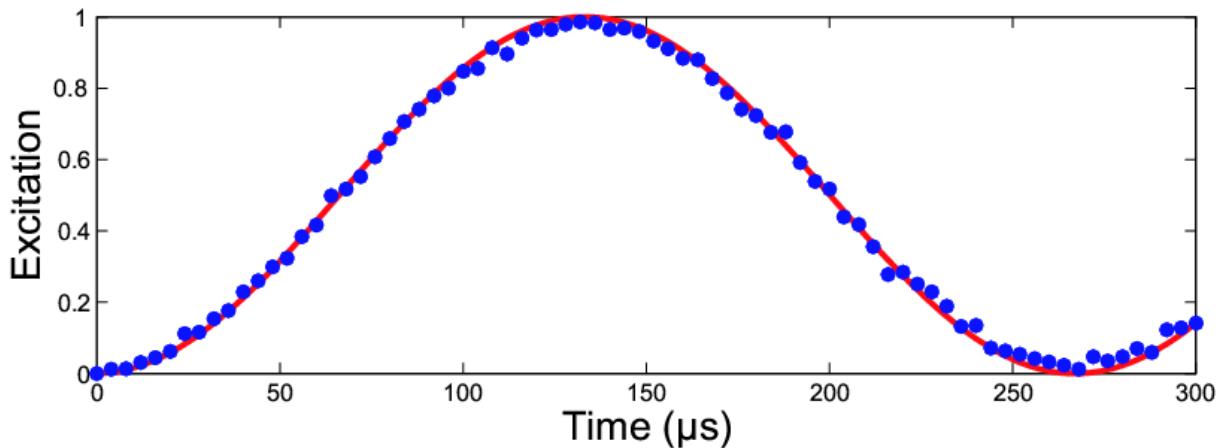


Ion motion



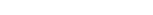
carrier and sideband
Rabi oscillations
with Rabi frequencies

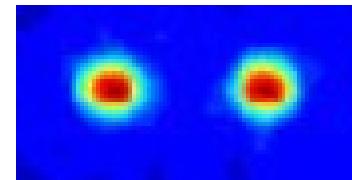
$\Omega, \eta\Omega$



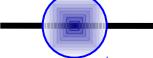
$\eta = kx_0$ Lamb-Dicke parameter

Generation of Bell states

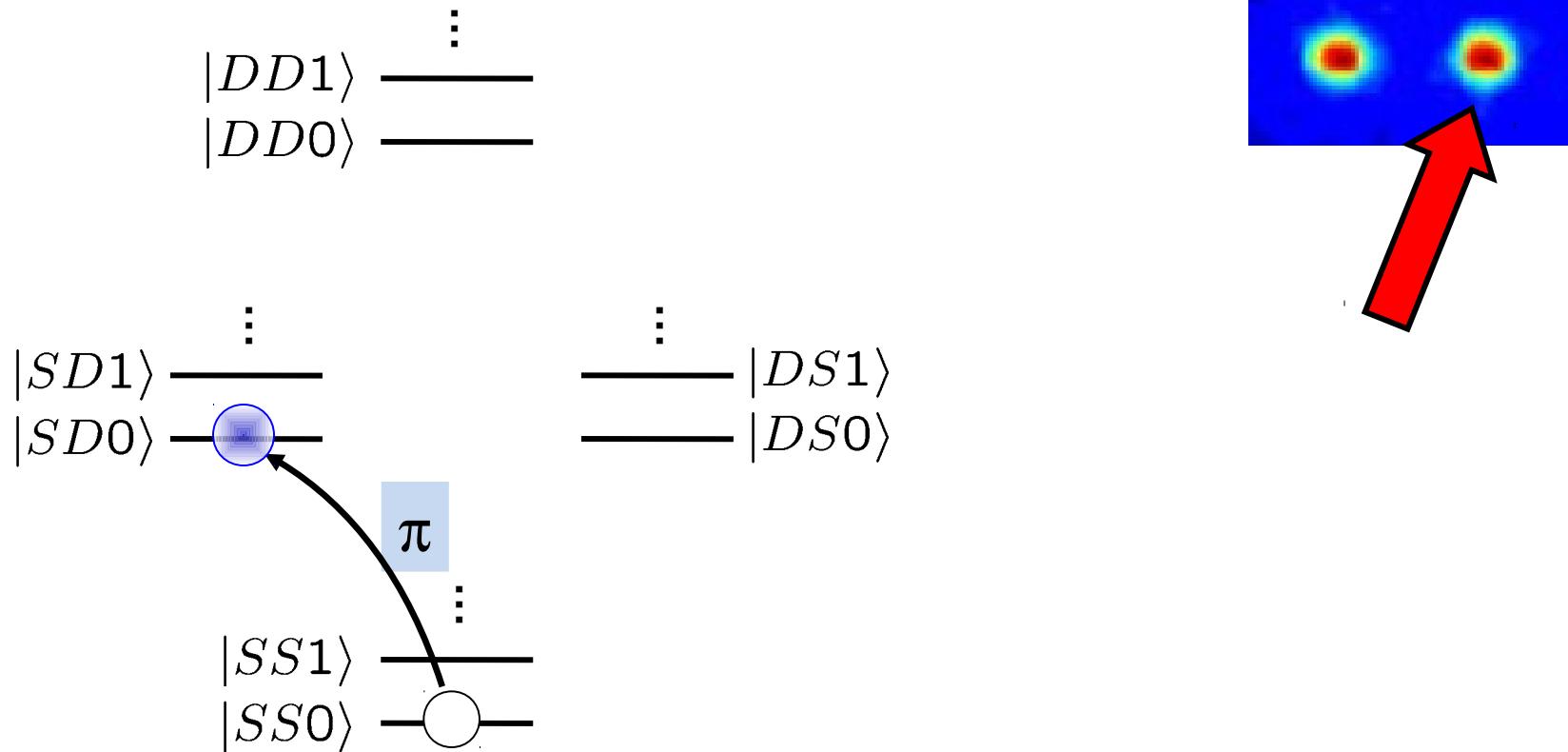
$|DD1\rangle$ 
 $|DD0\rangle$ 



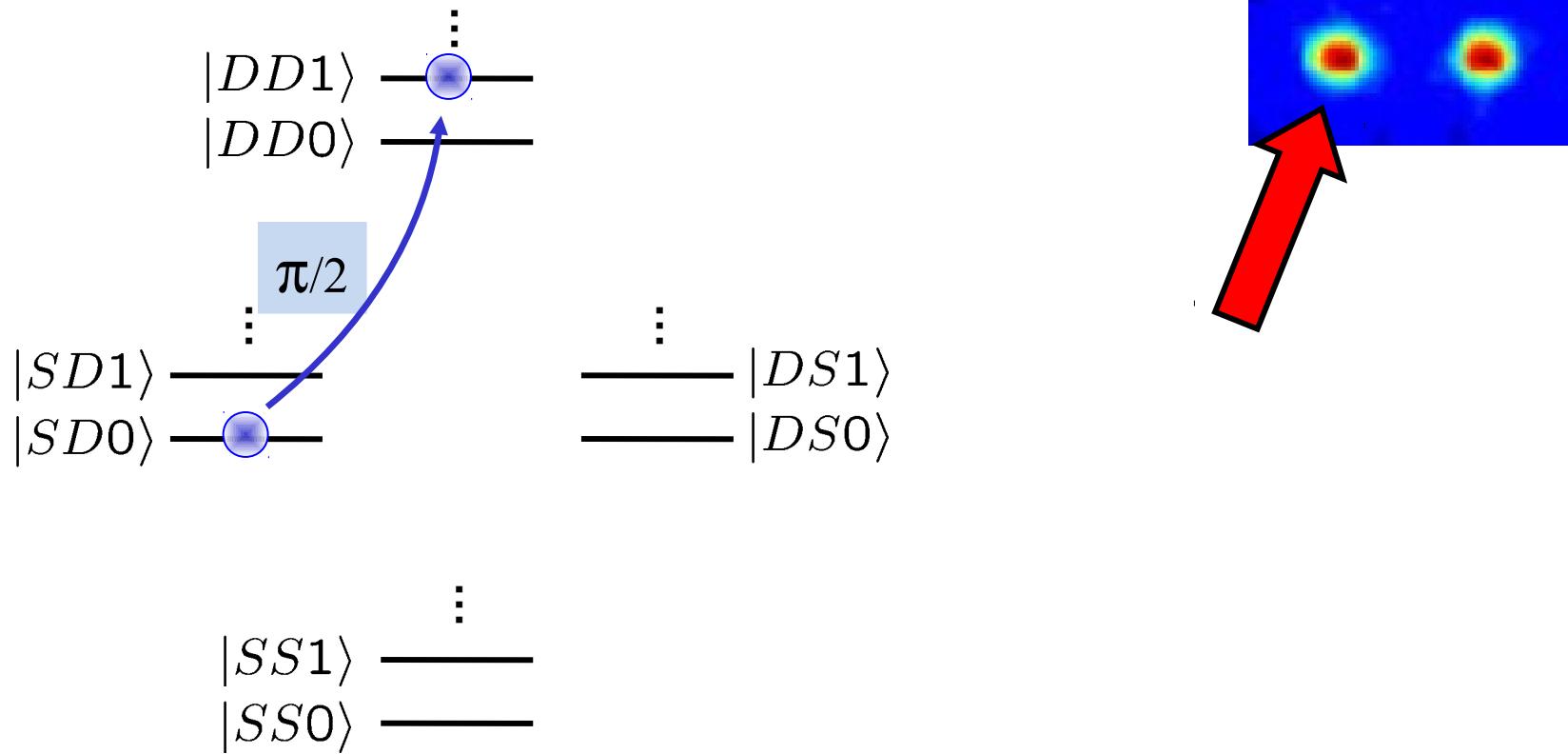
$|SD1\rangle$ 
 $|SD0\rangle$   $|DS1\rangle$
  $|DS0\rangle$

$|SS1\rangle$ 
 $|SS0\rangle$ 

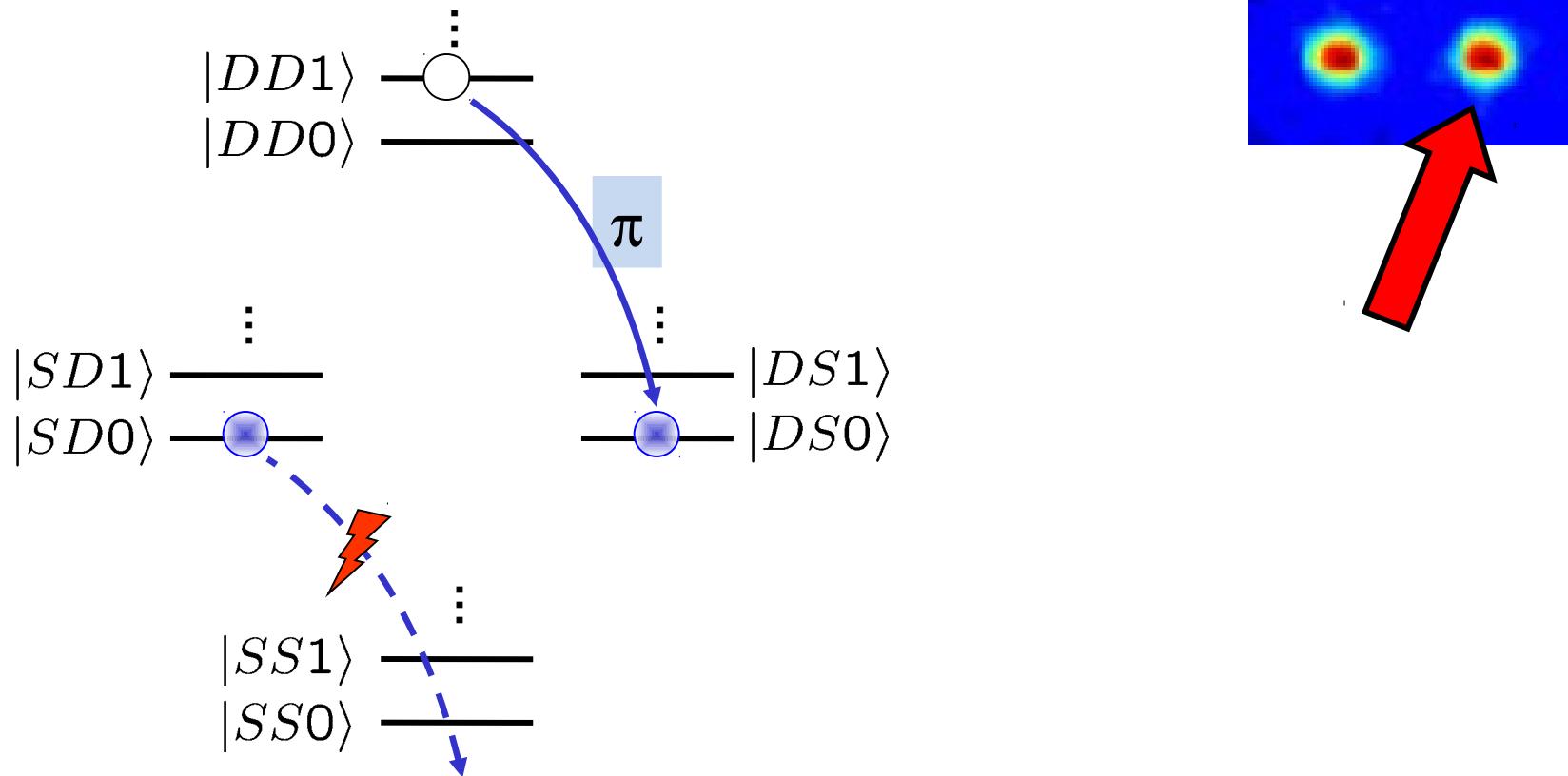
Generation of Bell states



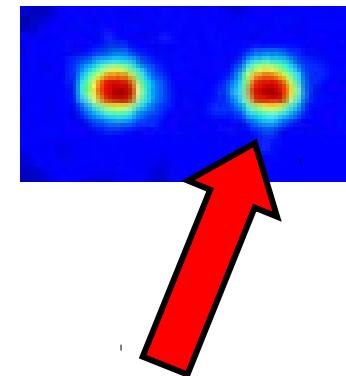
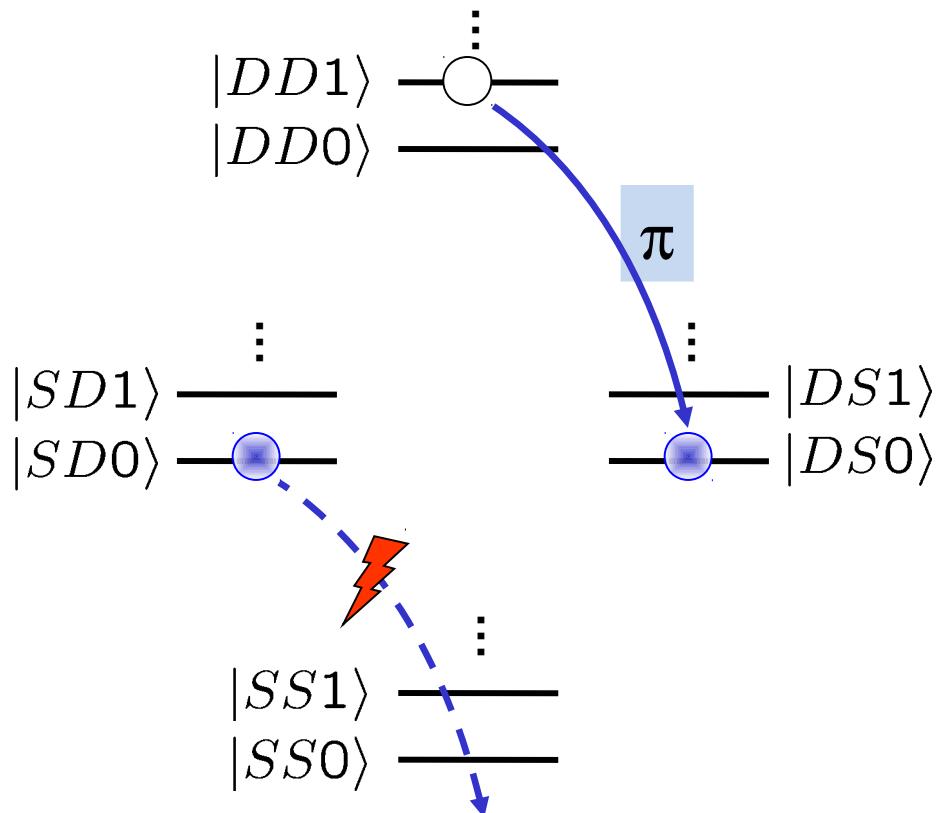
Generation of Bell states



Generation of Bell states



Generation of Bell states



Bell states with atoms

- ${}^9\text{Be}^+$: NIST (fidelity: 97 %)
- ${}^{40}\text{Ca}^+$: Oxford (99.6 %)
- ${}^{111}\text{Cd}^+$: Ann Arbor (79%)
- ${}^{171}\text{Yb}$: Maryland (96%)
- ${}^{25}\text{Mg}^+$: Munich (97%)
- ${}^{40}\text{Ca}^+$: Innsbruck (99.3%)

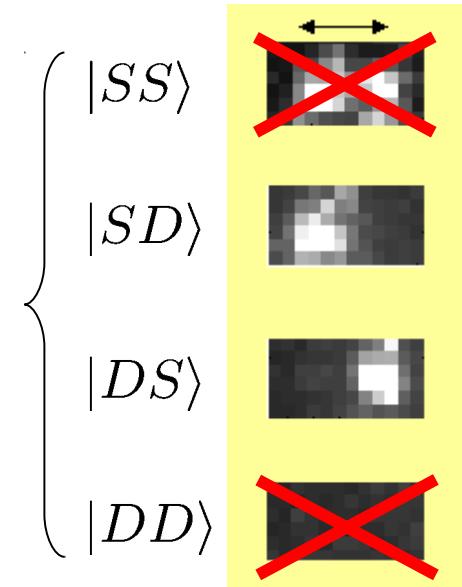
Analysis of Bell states

$$|SD\rangle + |DS\rangle$$

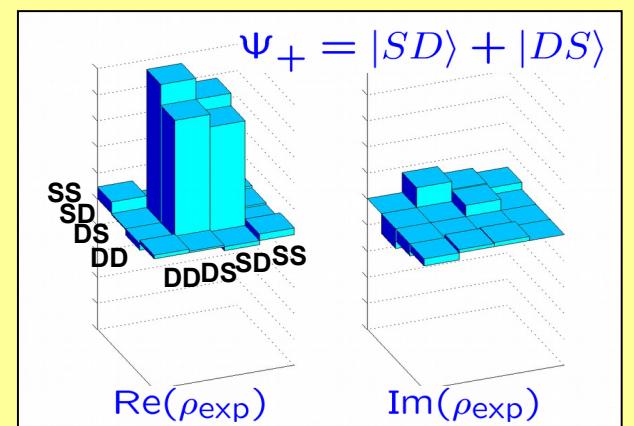
Fluorescence
detection with
CCD camera:

Coherent superposition or incoherent mixture ?

What is the relative phase of the superposition ?



→ Measurement of the density matrix:

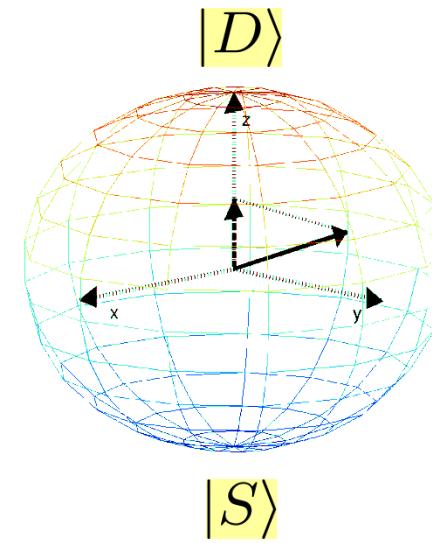


Measuring a density matrix

A measurement yields the z-component of the Bloch vector

=> Diagonal of the density matrix

$$\rho = \begin{pmatrix} P_S & C - iD \\ C + iD & P_D \end{pmatrix}$$

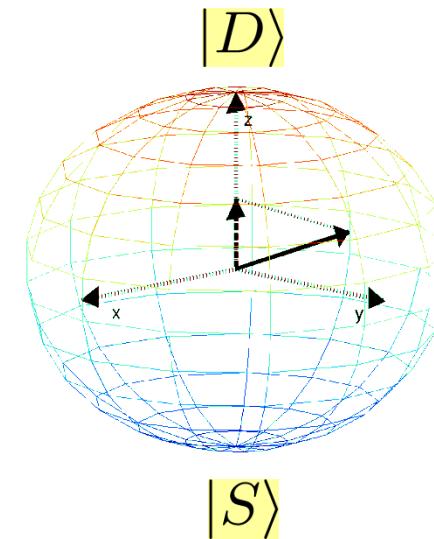


Measuring a density matrix

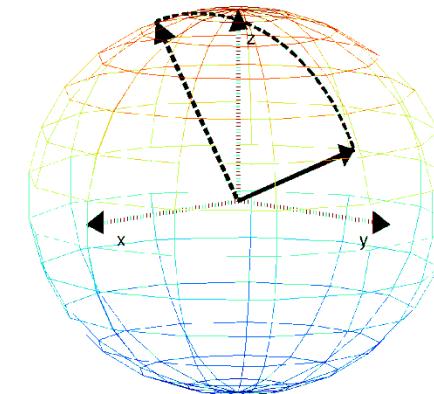
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Rotation around the x- or the y-axis prior to the measurement yields the phase information of the qubit.

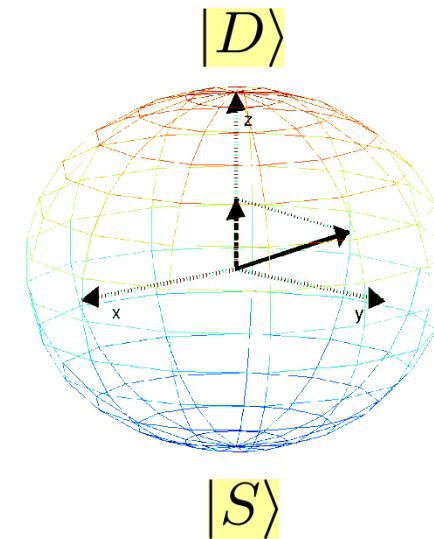


Measuring a density matrix

A measurement yields the z-component of the Bloch vector

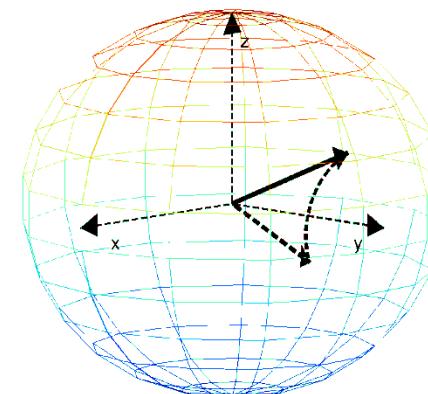
=> Diagonal of the density matrix

$$\rho = \begin{pmatrix} P_S & \mathcal{C} - iD \\ \mathcal{C} + iD & P_D \end{pmatrix}$$

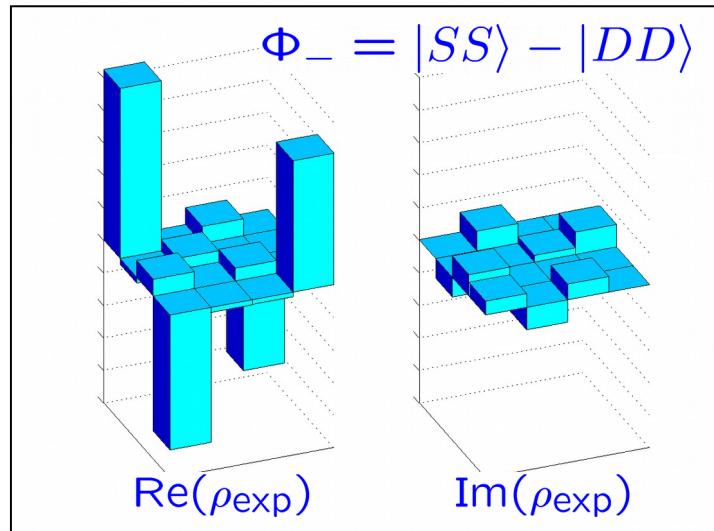
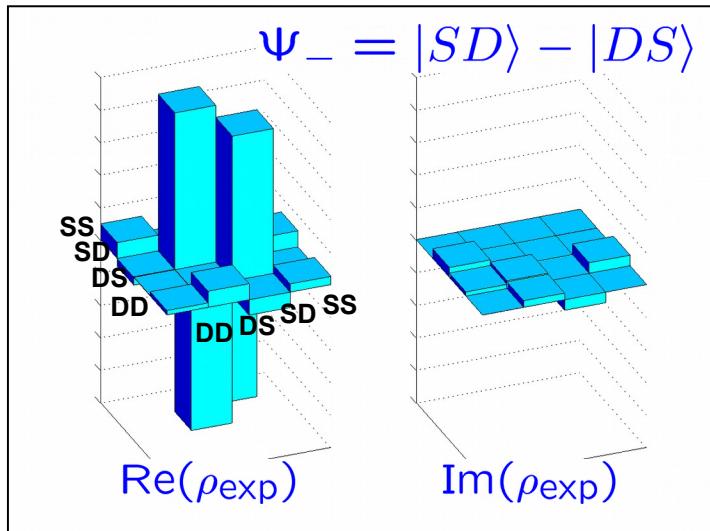
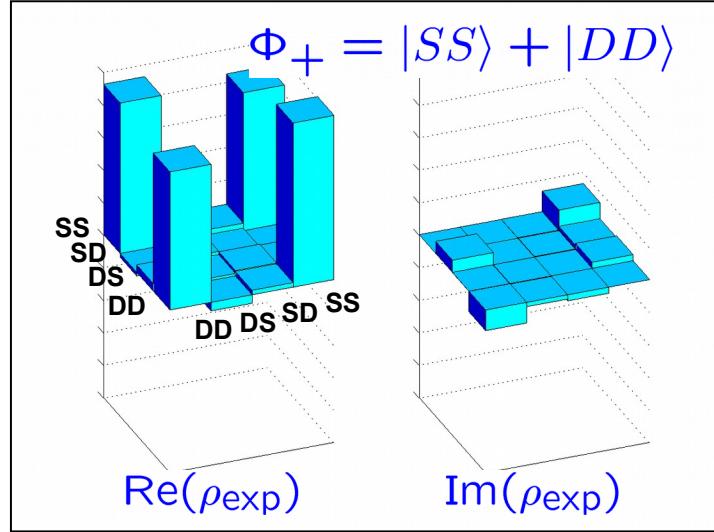
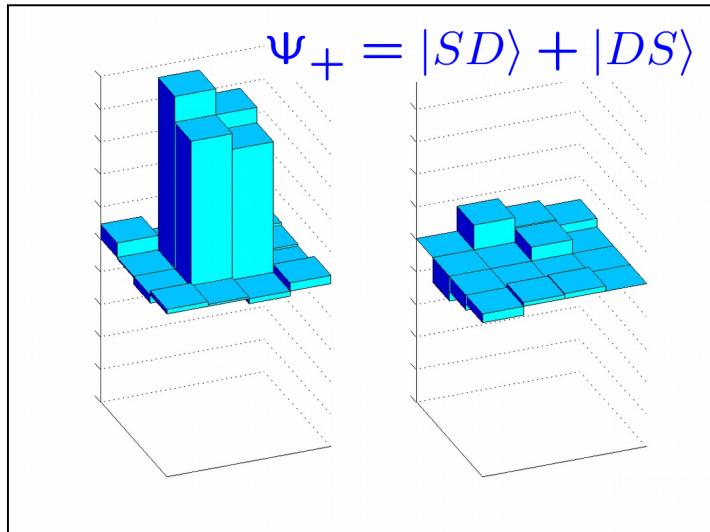


Rotation around the x- or the y-axis prior to the measurement yields the phase information of the qubit.

=> coherences of the density m

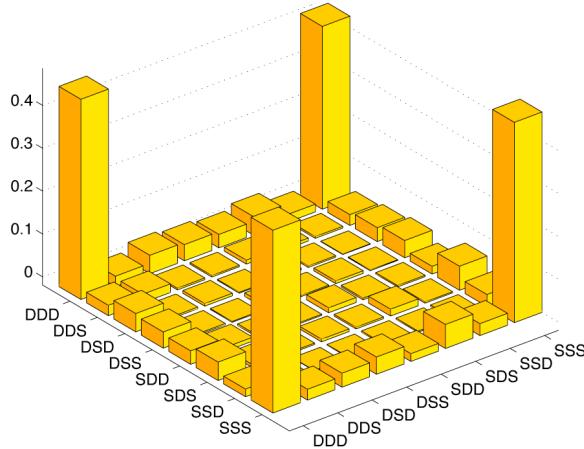


Bell states

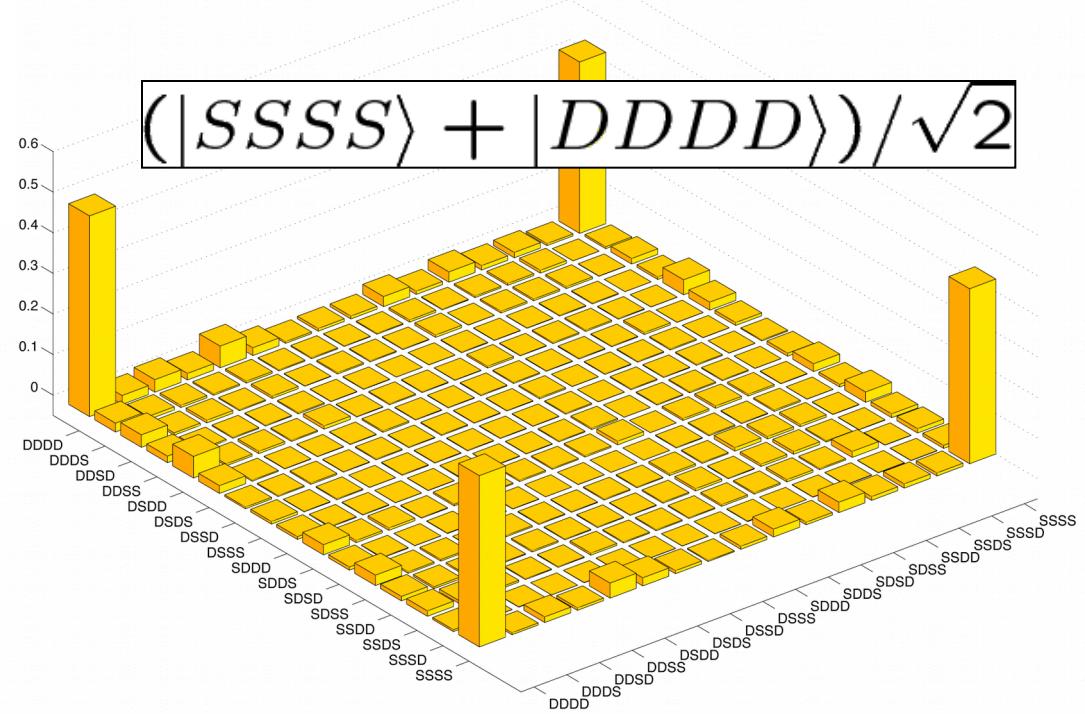


Generalized Bell states

$$(|SSS\rangle + |DDD\rangle)/\sqrt{2}$$

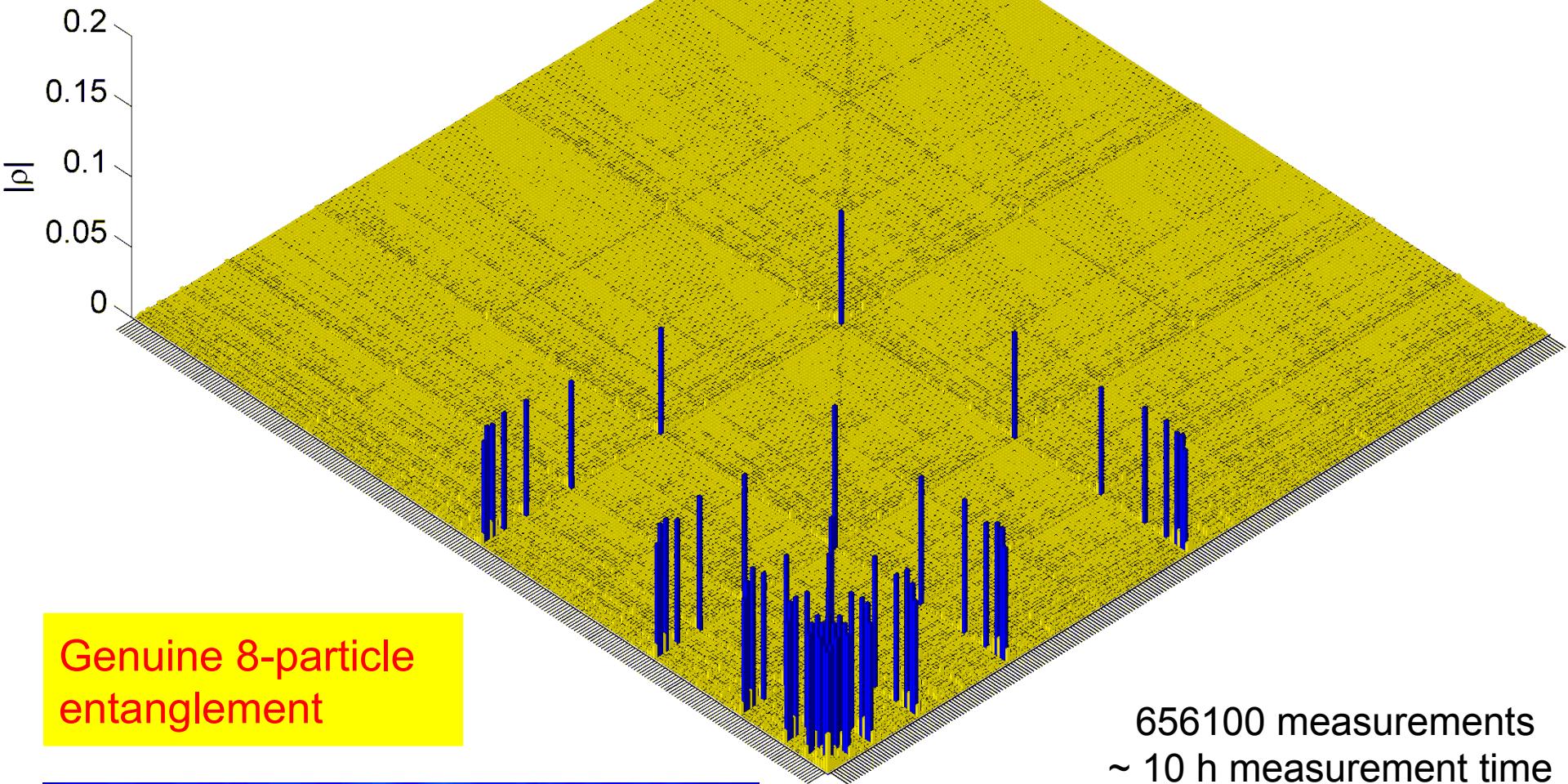


$$(|SSSS\rangle + |DDDD\rangle)/\sqrt{2}$$



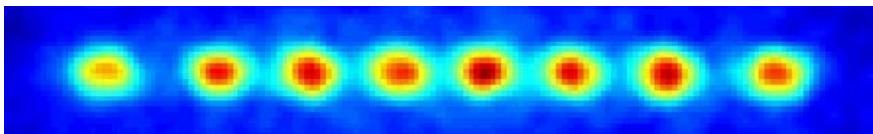
Generalized Bell states

$$\frac{1}{\sqrt{8}}(|DDDDDDDS\rangle + |DDDDDDSD\rangle + \dots + |SDDDDDDD\rangle)$$



Genuine 8-particle
entanglement

656100 measurements
 ~ 10 h measurement time



Häffner et al., Nature 438, 643 (2005)

Universal set of quantum gates ...

Having the qubits interact

VOLUME 74, NUMBER 20

PHYSICAL REVIEW LETTERS

15 MAY 1995

Quantum Computations with Cold Trapped Ions

J. I. Cirac and P. Zoller*

Institut für Theoretische Physik, Universität Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria

(Received 30 November 1994)

A quantum computer can be implemented with cold ions confined in a linear trap and interacting with laser beams. Quantum gates involving any pair, triplet, or subset of ions can be realized by coupling the ions through the collective quantized motion. In this system decoherence is negligible, and the measurement (readout of the quantum register) can be carried out with a high efficiency.

PACS numbers: 89.80.+h, 03.65.Bz, 12.20.Fv, 32.80.Pj

...allows the realization of a
universal quantum computer !

$$|D\rangle|D\rangle \rightarrow |D\rangle|D\rangle$$

$$|D\rangle|S\rangle \rightarrow |D\rangle|S\rangle$$

$$|S\rangle|D\rangle \rightarrow |D\rangle|S\rangle$$

$$|S\rangle|S\rangle \rightarrow |S\rangle|\textcolor{red}{D}\rangle$$

control

target

Having the qubits interact

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$$|S\rangle|S\rangle \rightarrow |S\rangle|\textcolor{red}{D}\rangle$$

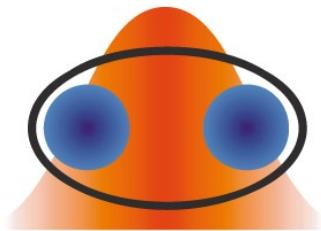
control

target

Most popular gates:

- Cirac-Zoller gate (Schmidt-Kaler et al., Nature **422**, 408 (2003)).
- Geometric phase gate (Leibfried et al., Nature **422**, 412 (2003)).
- Mølmer-Sørensen gate (Sackett et al., Nature **404**, 256 (2000)).

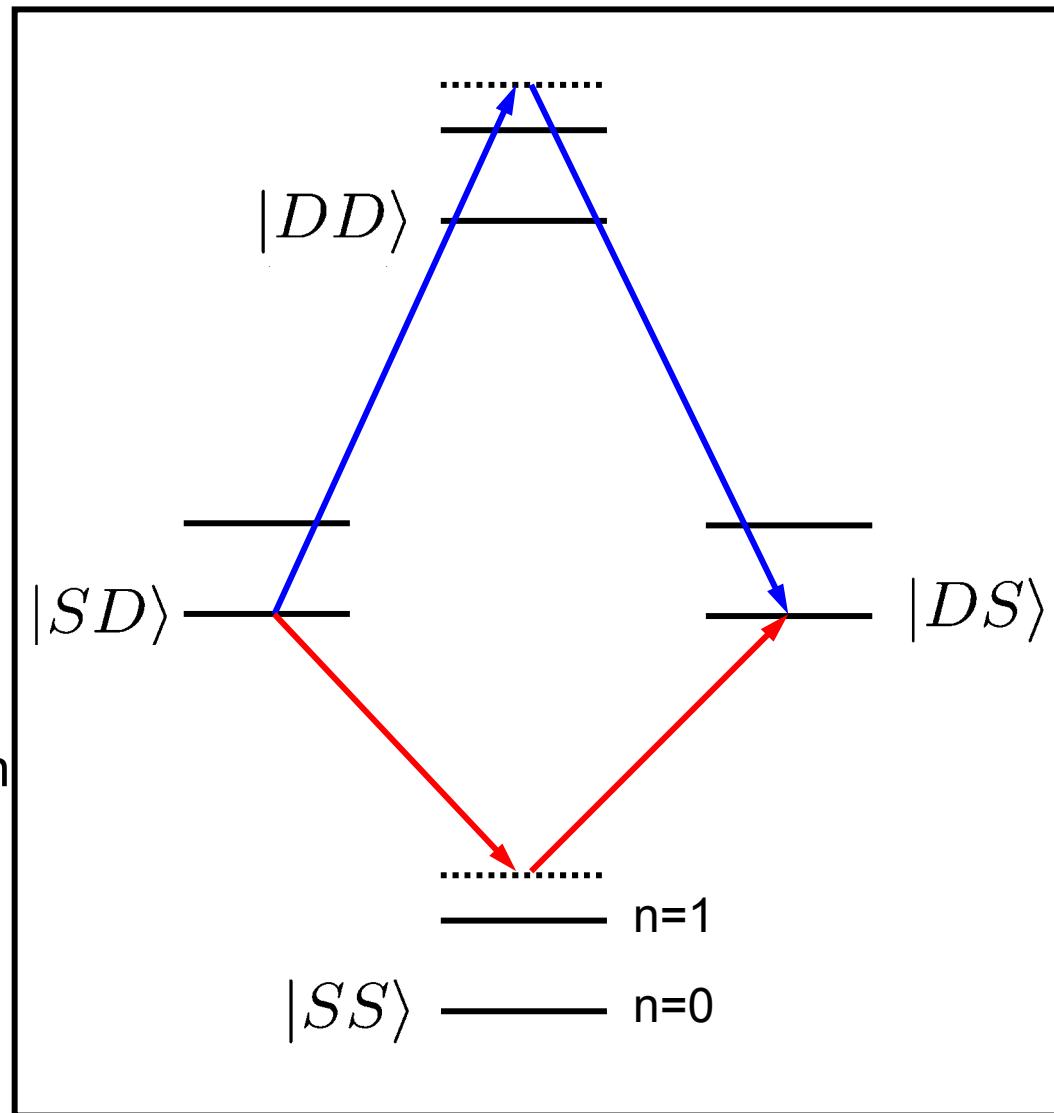
Mølmer-Sørensen gate creates entangled states



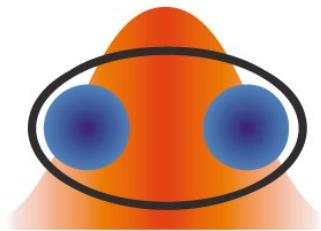
Raman transitions between

$$|SD\rangle \Leftrightarrow |DS\rangle$$

Interaction of two ions via common motion.



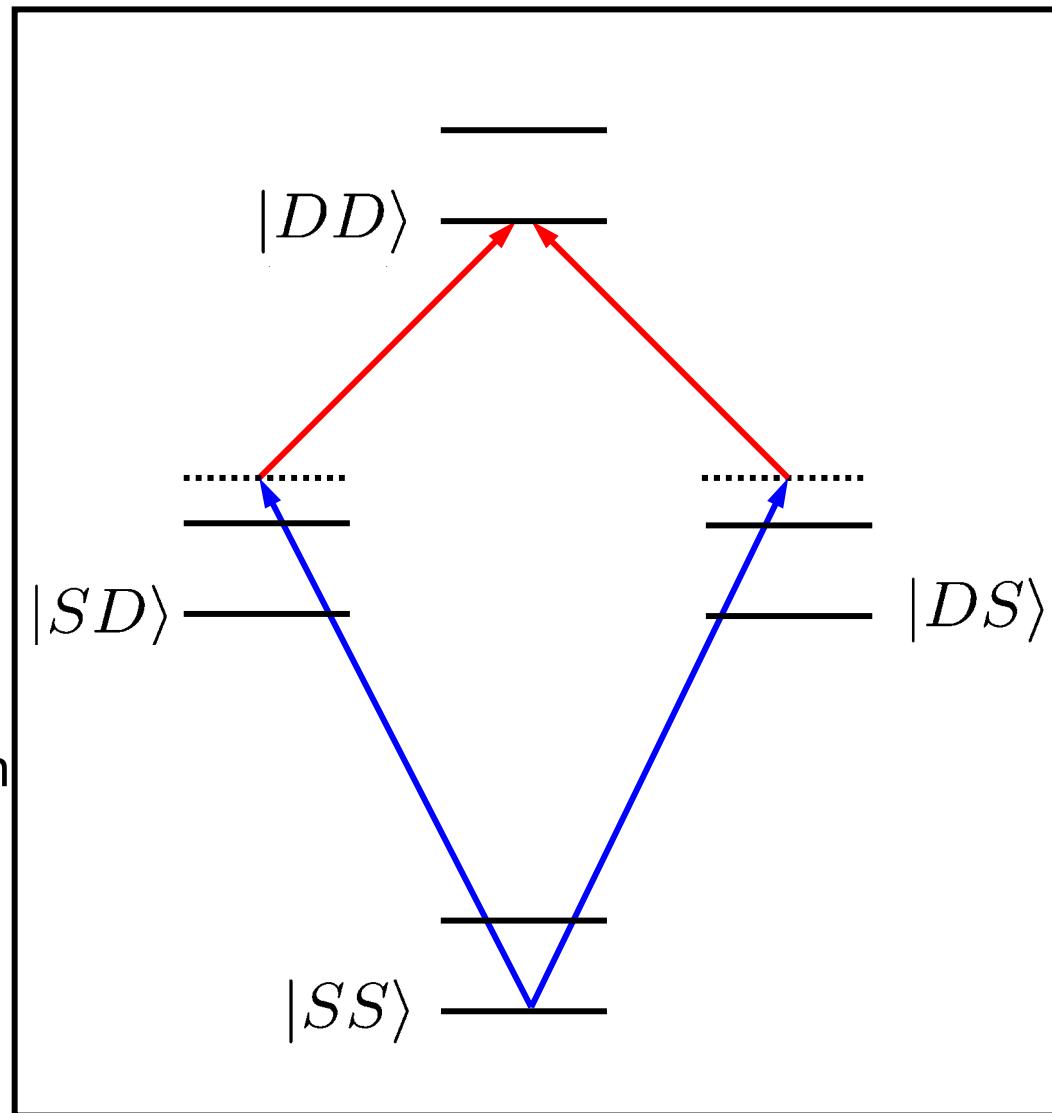
Mølmer-Sørensen gate creates entangled states



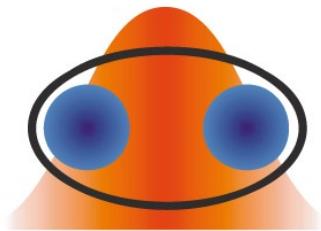
Raman transitions between

$$|SS\rangle \Leftrightarrow |DD\rangle$$

Interaction of two ions via common motion.



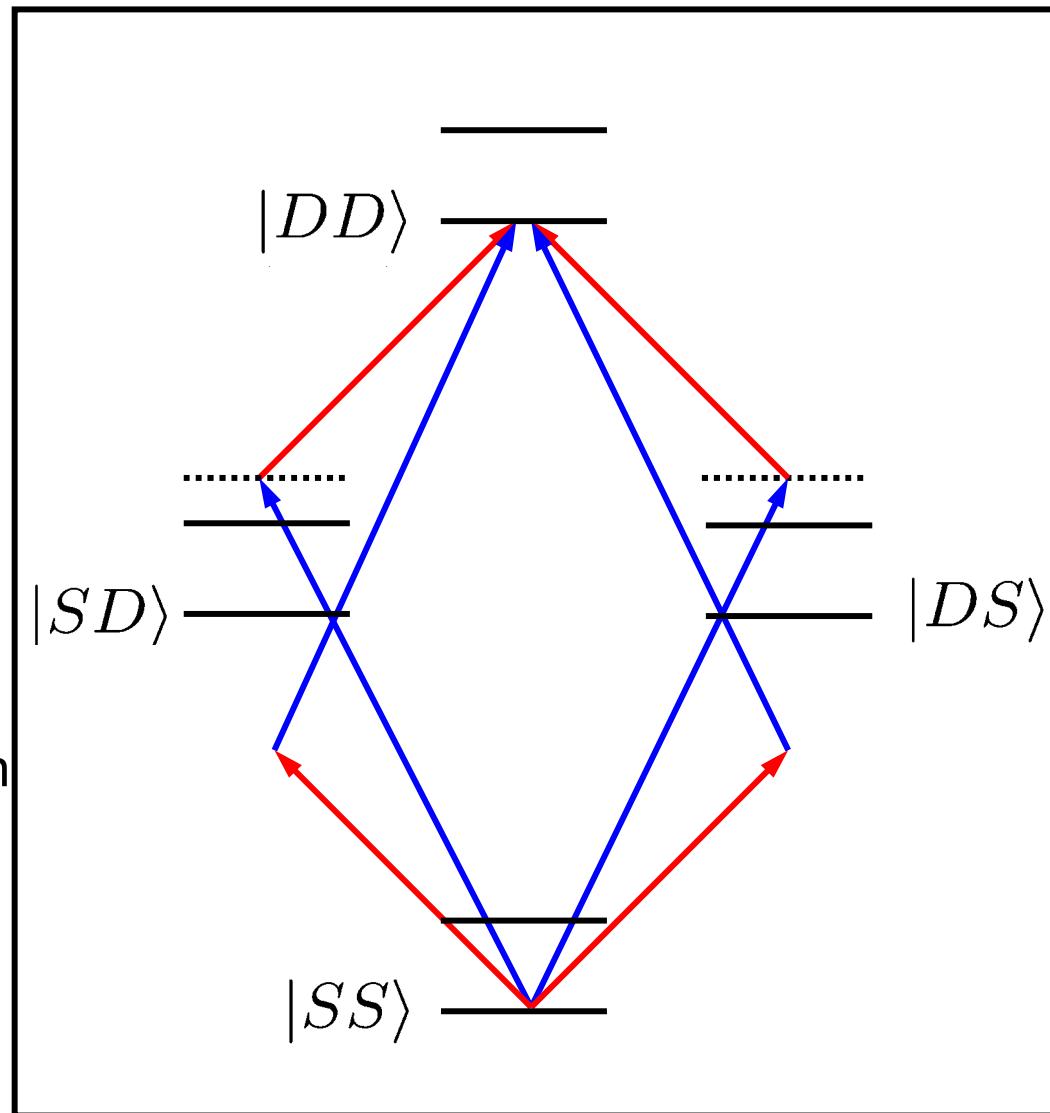
Mølmer-Sørensen gate creates entangled states



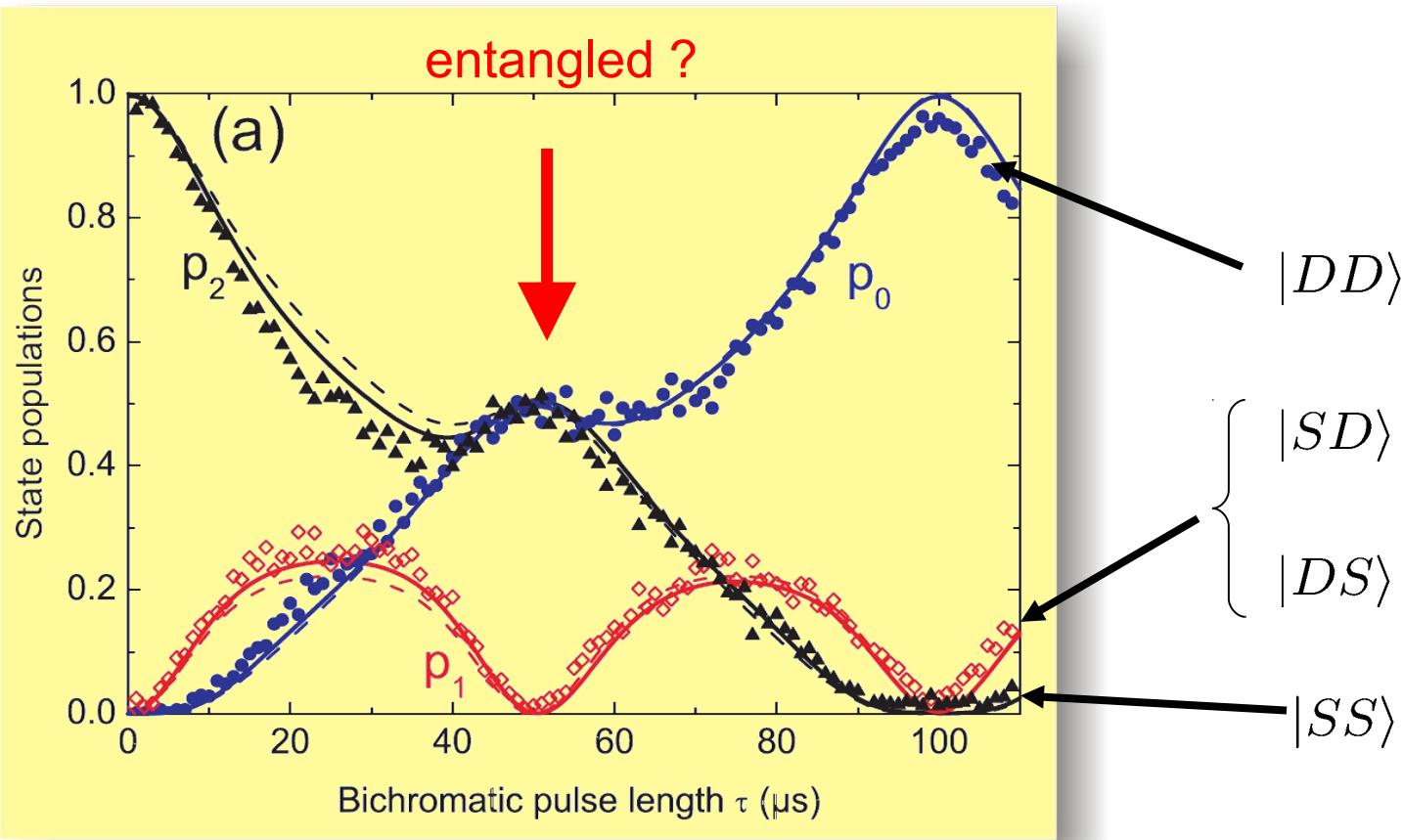
Raman transitions between

$$|SS\rangle \Leftrightarrow |DD\rangle$$

Interaction of two ions via common motion.

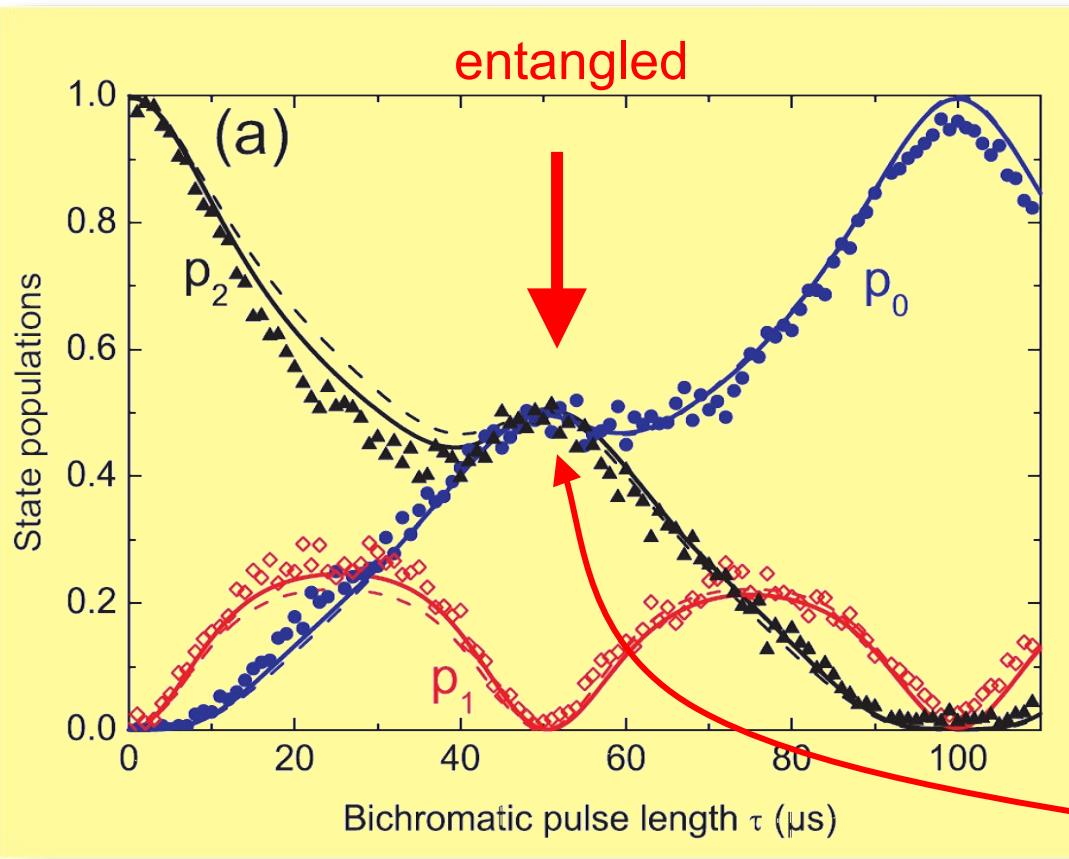


Entangling ions

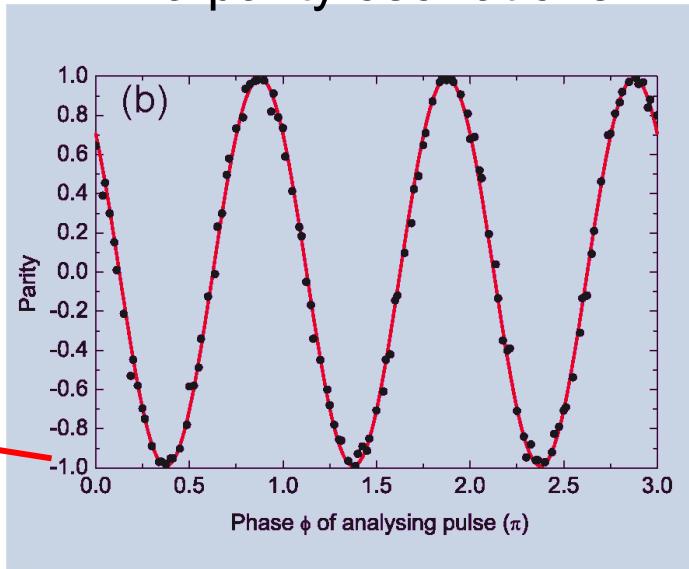


J. Benhelm et al., Nature Physics **4**, 463 (2008)
Theory: C. Roos, NJP **10**, 013002 (2008)

Entangling ions



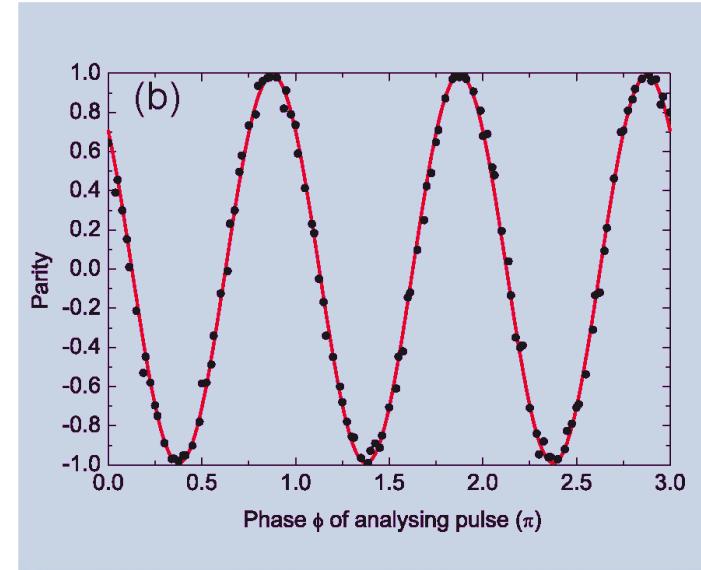
measure entanglement
via parity oscillations



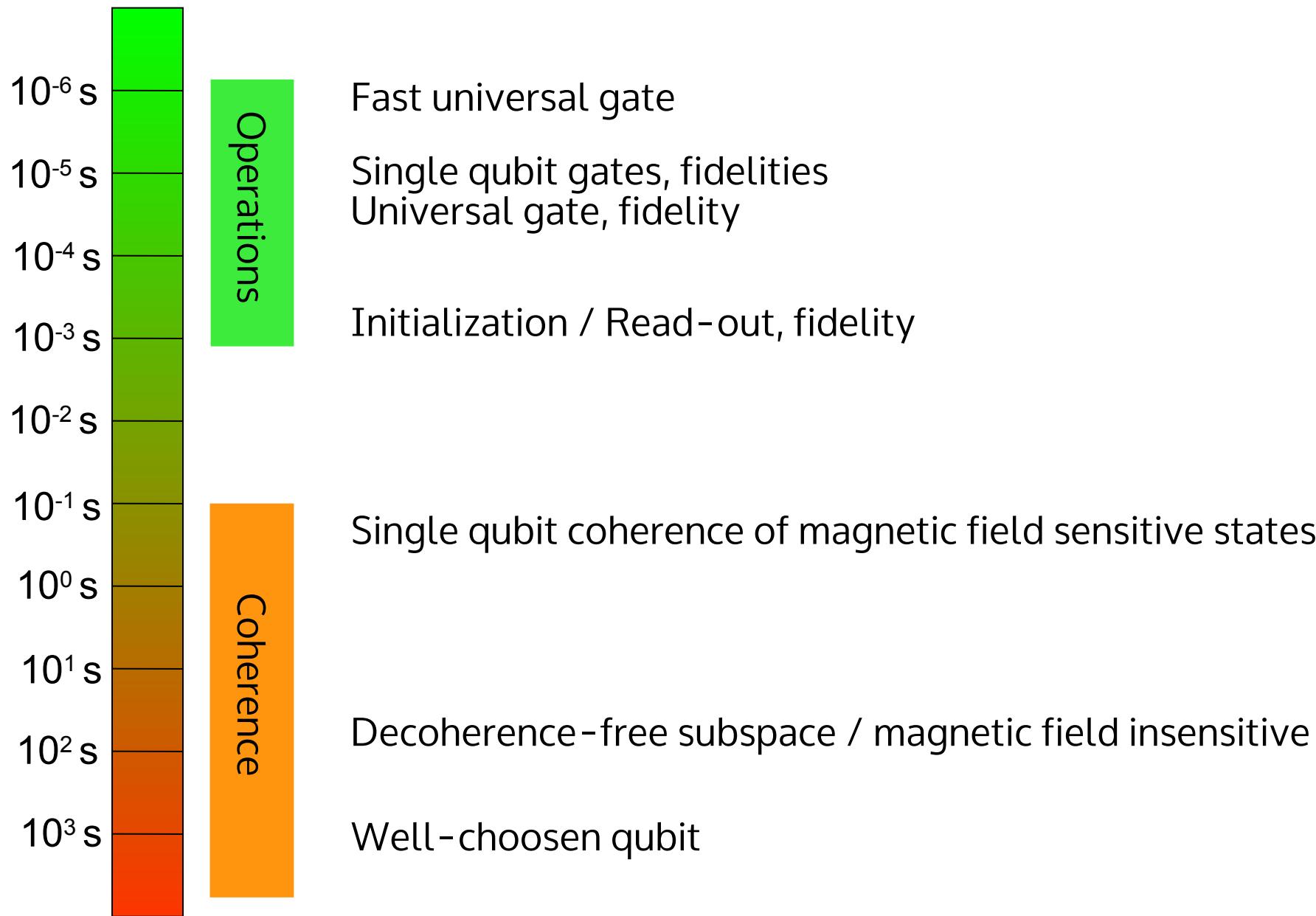
gate duration $51 \mu\text{s}$

average fidelity: 99.3 (2) %

$$\begin{aligned}
|00\rangle + |11\rangle & \xrightarrow{R_2^C(\pi/2, \varphi), R_1^C(\pi/2, \varphi)} \\
& (|0\rangle + ie^{i\varphi}|1\rangle)(|0\rangle + ie^{i\varphi}|1\rangle) + (|1\rangle + ie^{-i\varphi}|0\rangle)(|1\rangle + ie^{-i\varphi}|0\rangle) \\
& = (1 - e^{-2i\varphi})|00\rangle + ie^{i\varphi}(1 + e^{-2i\varphi})|01\rangle \\
& \quad + ie^{i\varphi}(1 + e^{-2i\varphi})|10\rangle + (1 - e^{-2i\varphi})|11\rangle,
\end{aligned}$$



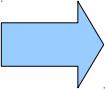
Achieved times scales for ion trap QIP



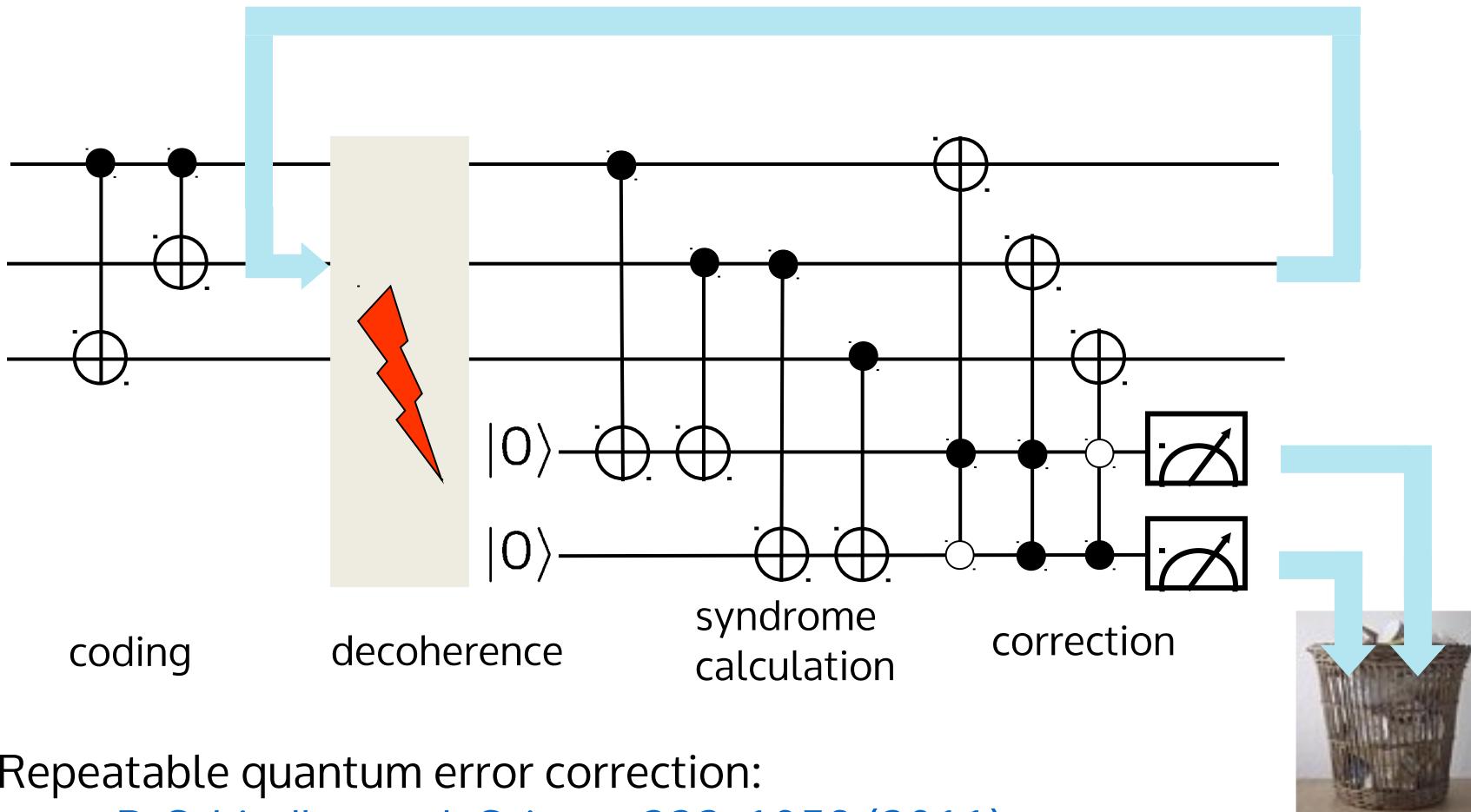
The DiVincenzo criteria for quantum computing

- I. Scalable physical system, well characterized qubits ✓
- II. Ability to initialize the state of the qubits ✓
- III. Long relevant coherence times, much longer than gate operation time ✓
- IV. “Universal” set of quantum gates ✓
- V. Qubit-specific measurement capability ✓

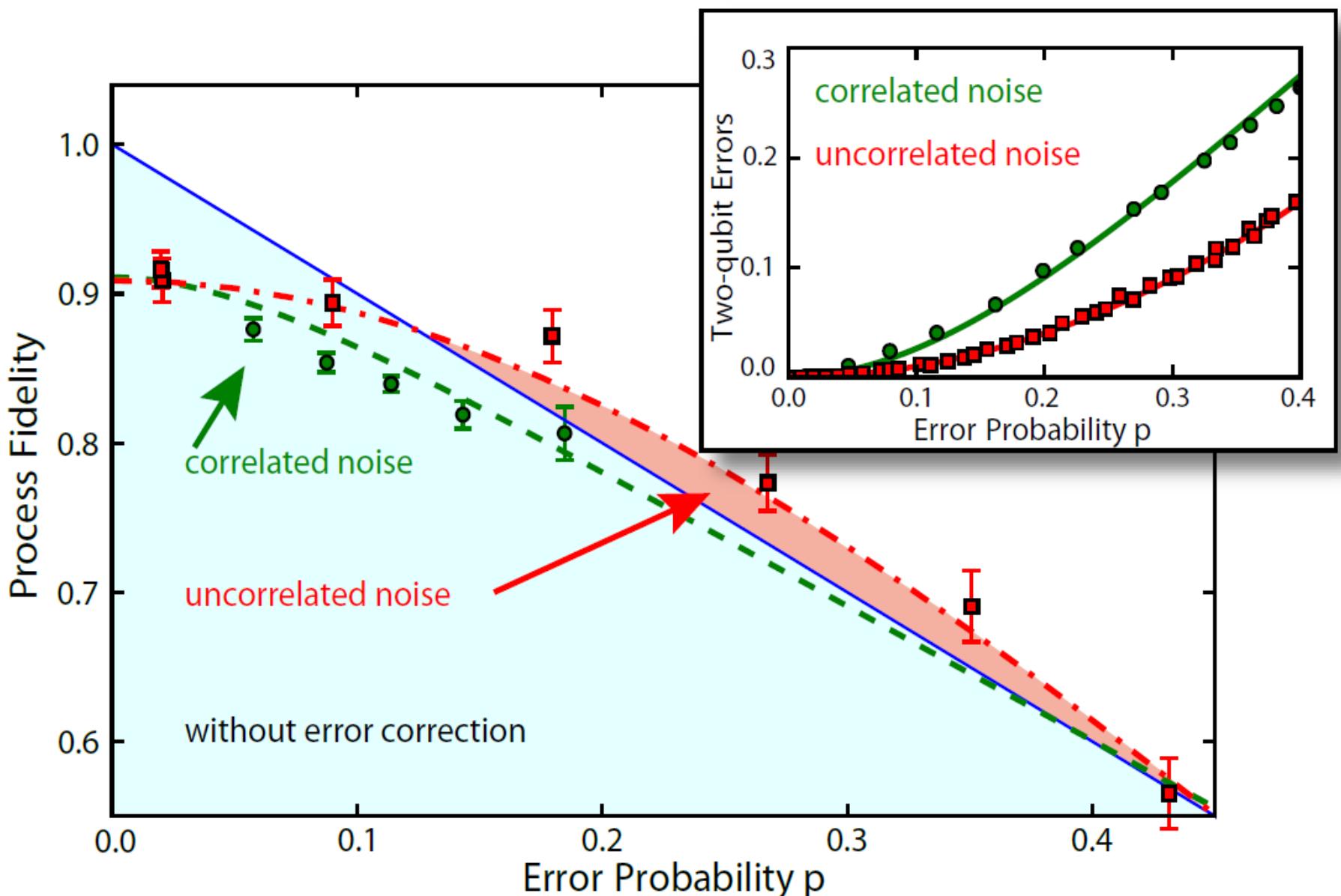
The DiVincenzo criteria for quantum computing

- I. Scalable physical system, well characterized qubits
 - II. Ability to initialize the state of the qubits **with sufficient fidelity**
 - III. Long relevant coherence times, much longer than gate operation time
 - IV. “Universal” set of quantum gates **with sufficient fidelity**
 - V. Qubit-specific measurement capability **with sufficient fidelity**
-  need to beat the fault-tolerant “threshold”

Quantum error correction



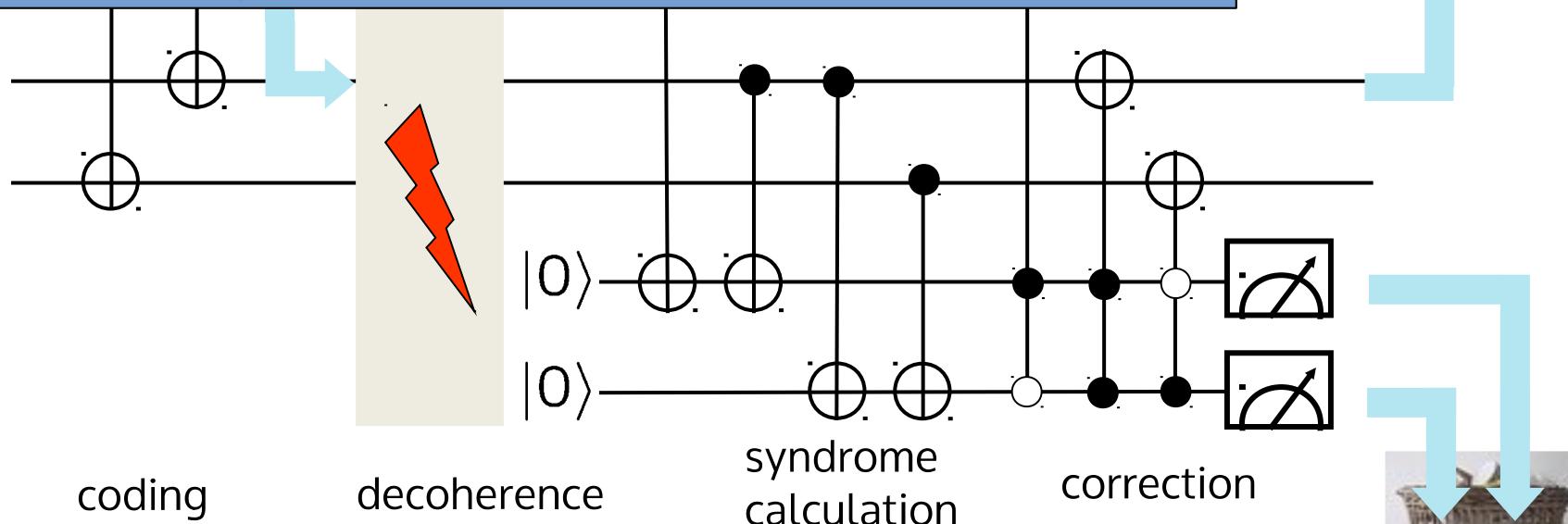
Consequences for QEC



Threshold for quantum computing: 99.99% fidelity / operation

Status of ion-trap QIP:

- Initialization / Read-out: 99.5% – 99.99%
- Single qubit gates: 99% - 99.999%
- Universal gate: 99.9%



Repeatable quantum error correction:

P. Schindler et al, Science 332, 1059 (2011)



The DiVincenzo criteria for quantum computing

- I. Scalable physical system, well characterized qubits
- II. Ability to initialize the state of the qubits with sufficient fidelity
- III. Long relevant coherence times, much longer than gate operation time
- IV. “Universal” set of quantum gates with sufficient fidelity
- V. Qubit-specific measurement capability with sufficient fidelity

The DiVincenzo criteria for quantum computing

- I. Scalable physical system, well characterized qubits
- II. Ability to initialize the state of the qubits with sufficient fidelity
- III. Long relevant coherence times, much longer than gate operation time
- IV. “Universal” set of quantum gates with sufficient fidelity
- V. Qubit-specific measurement capability with sufficient fidelity

Scaling of this approach?

Problems :

- Coupling strength between internal and motional states of a N-ion string decreases as

$$\eta \propto \frac{1}{\sqrt{N}}$$

(momentum transfer from photon to ion string becomes more difficult)

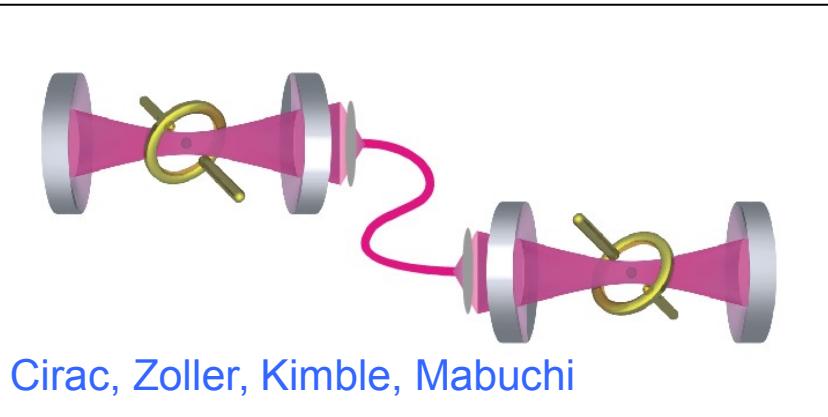
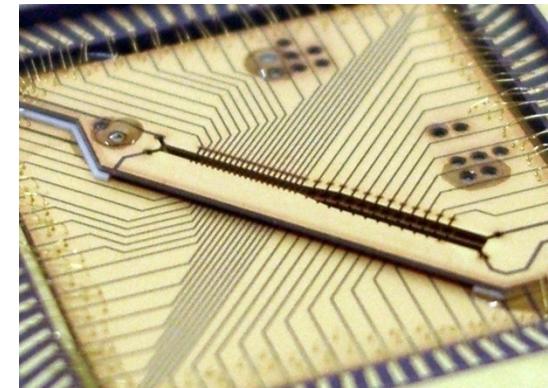
-> Gate operation speed slows down

- More vibrational modes increase risk of spurious excitation of unwanted modes
- Distance between neighbouring ions decreases -> addressing more difficult

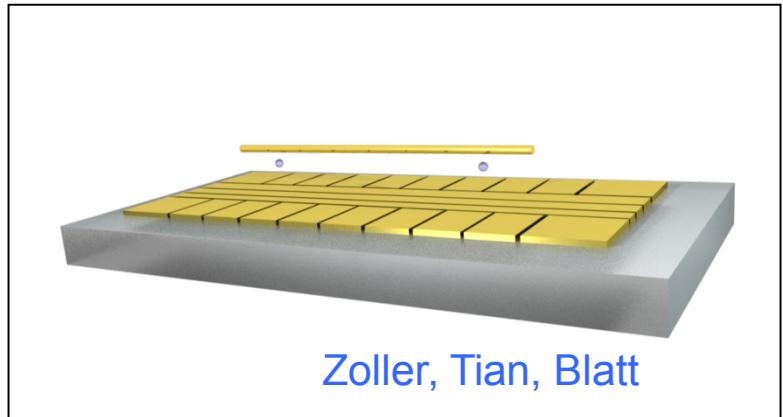
-> Use flexible trap potentials to split long ion string into smaller segments and perform operations on these smaller strings

Scaling of ion trap quantum computers

Kielpinski, Monroe, Wineland



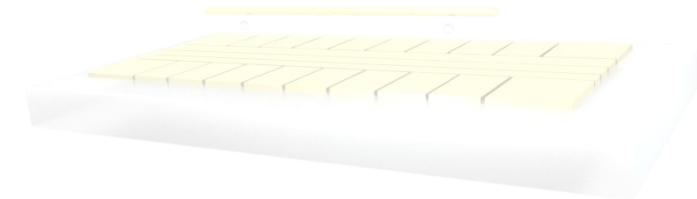
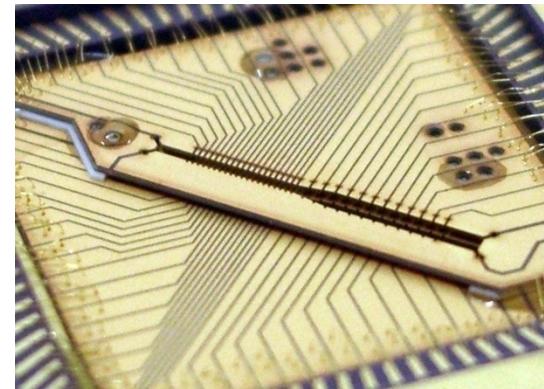
Cirac, Zoller, Kimble, Mabuchi



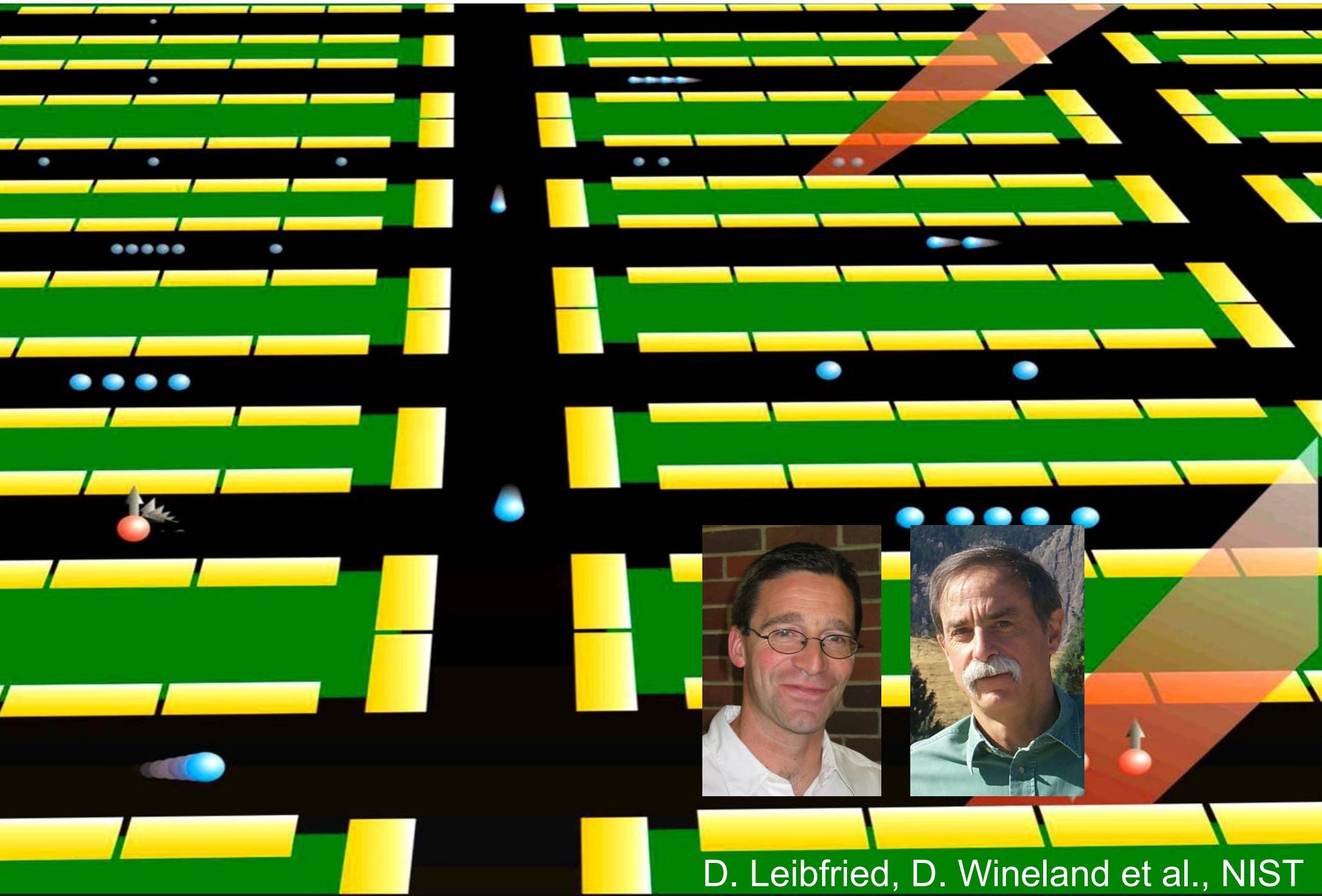
Zoller, Tian, Blatt

Scaling of ion trap quantum computers

Kielpinski, Monroe, Wineland

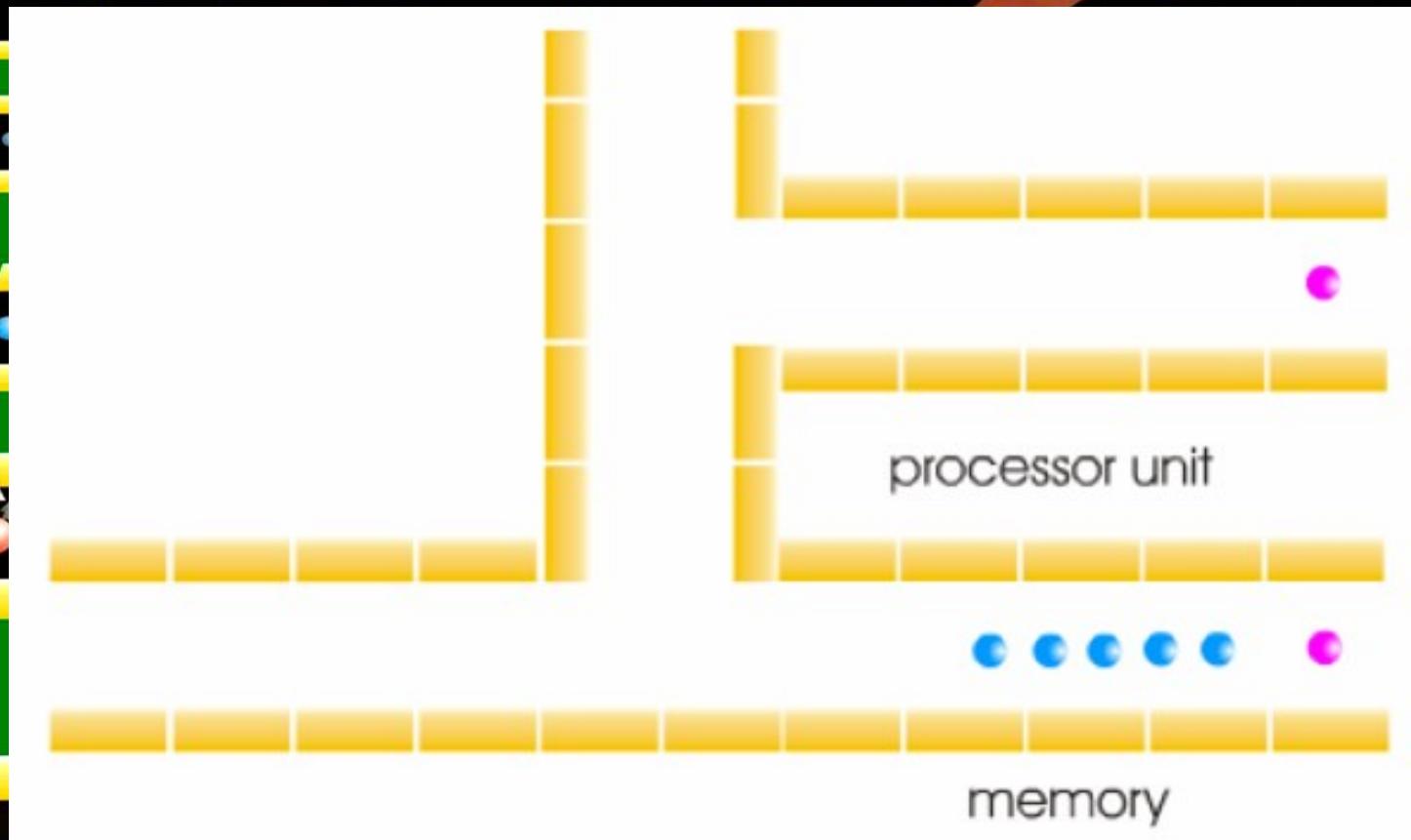


Scaling of ion trap quantum computers



D. Leibfried, D. Wineland et al., NIST

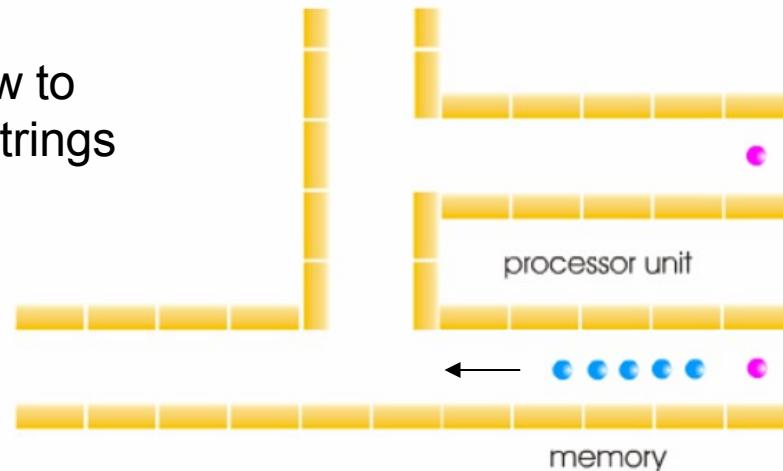
Scaling of ion trap quantum computers



Segmented ion traps as scalable trap architecture

(ideas pioneered by D. Wineland, NIST)

Segmented trap electrode allow to transport ions and to split ion strings



State of the art:

Transport of ions

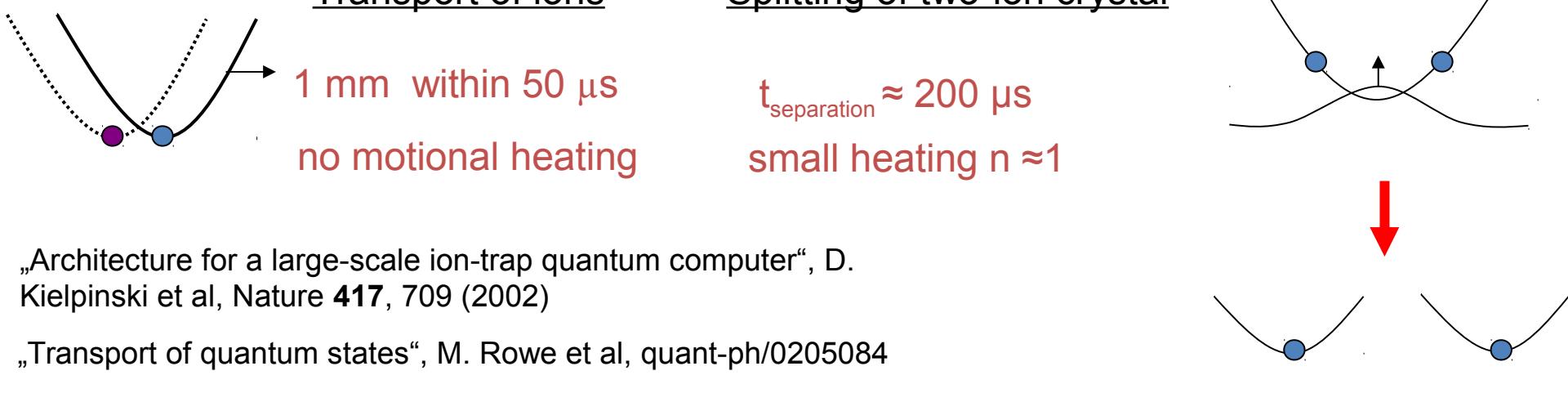
1 mm within 50 μ s

no motional heating

Splitting of two-ion crystal

$t_{\text{separation}} \approx 200 \mu\text{s}$

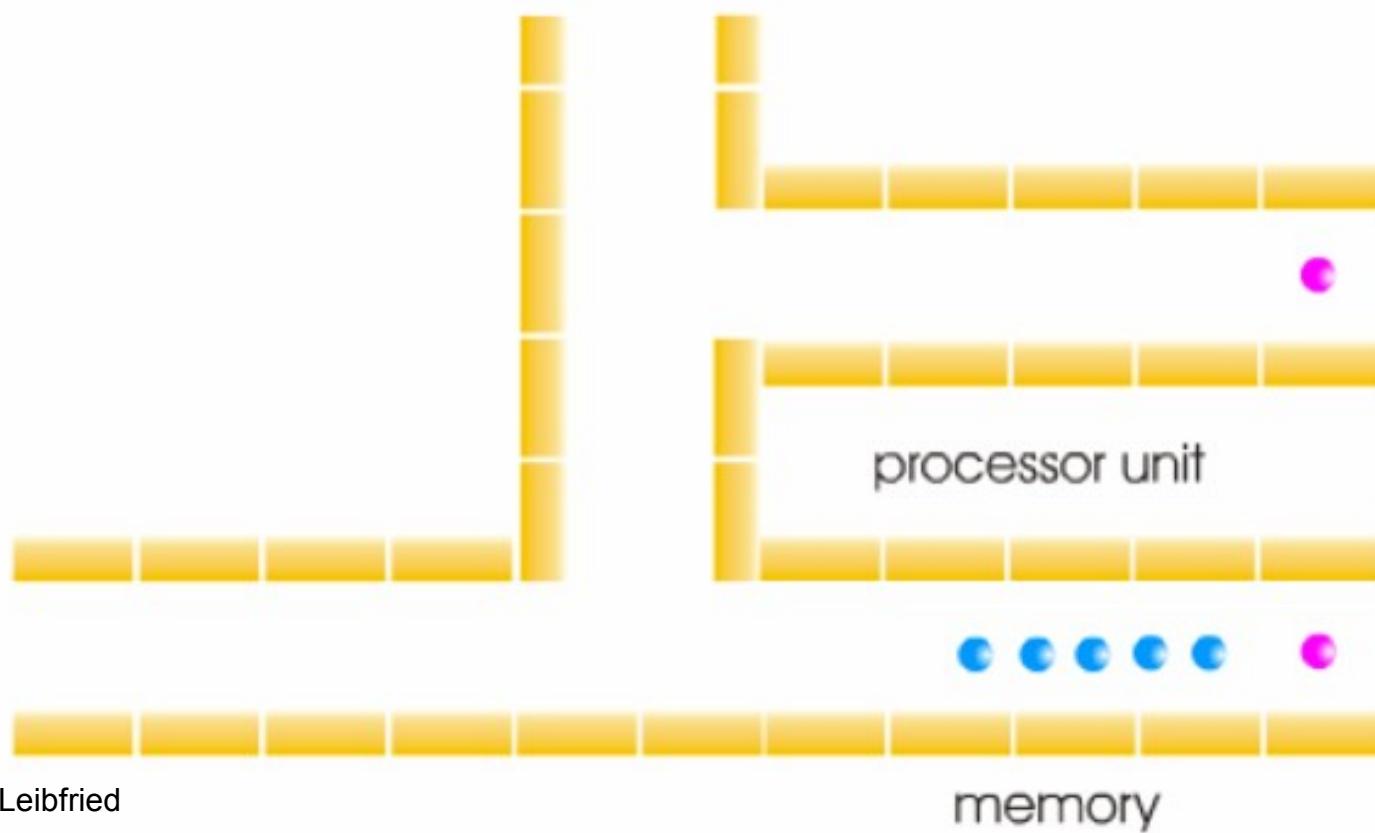
small heating $n \approx 1$



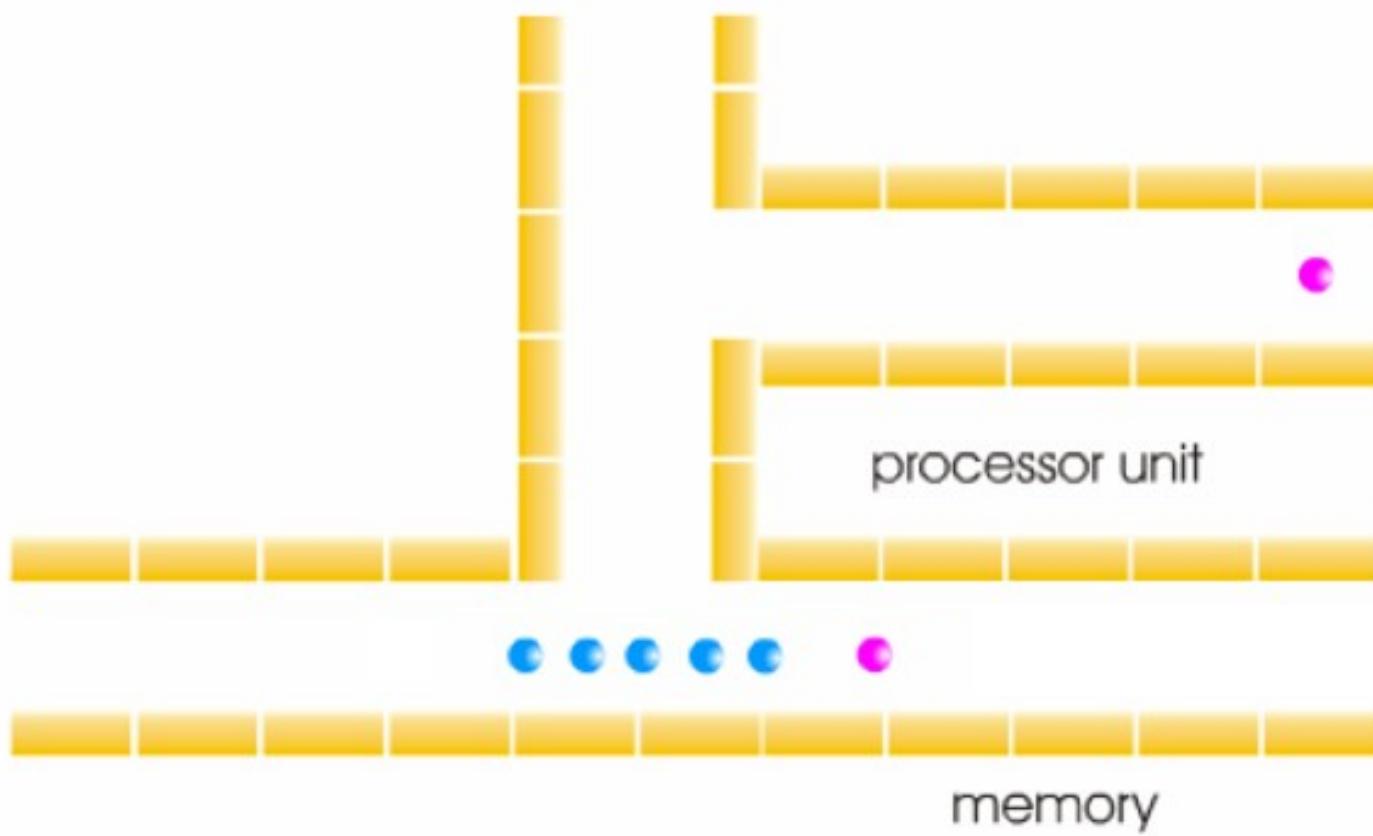
„Architecture for a large-scale ion-trap quantum computer“, D. Kielpinski et al, Nature **417**, 709 (2002)

„Transport of quantum states“, M. Rowe et al, quant-ph/0205084

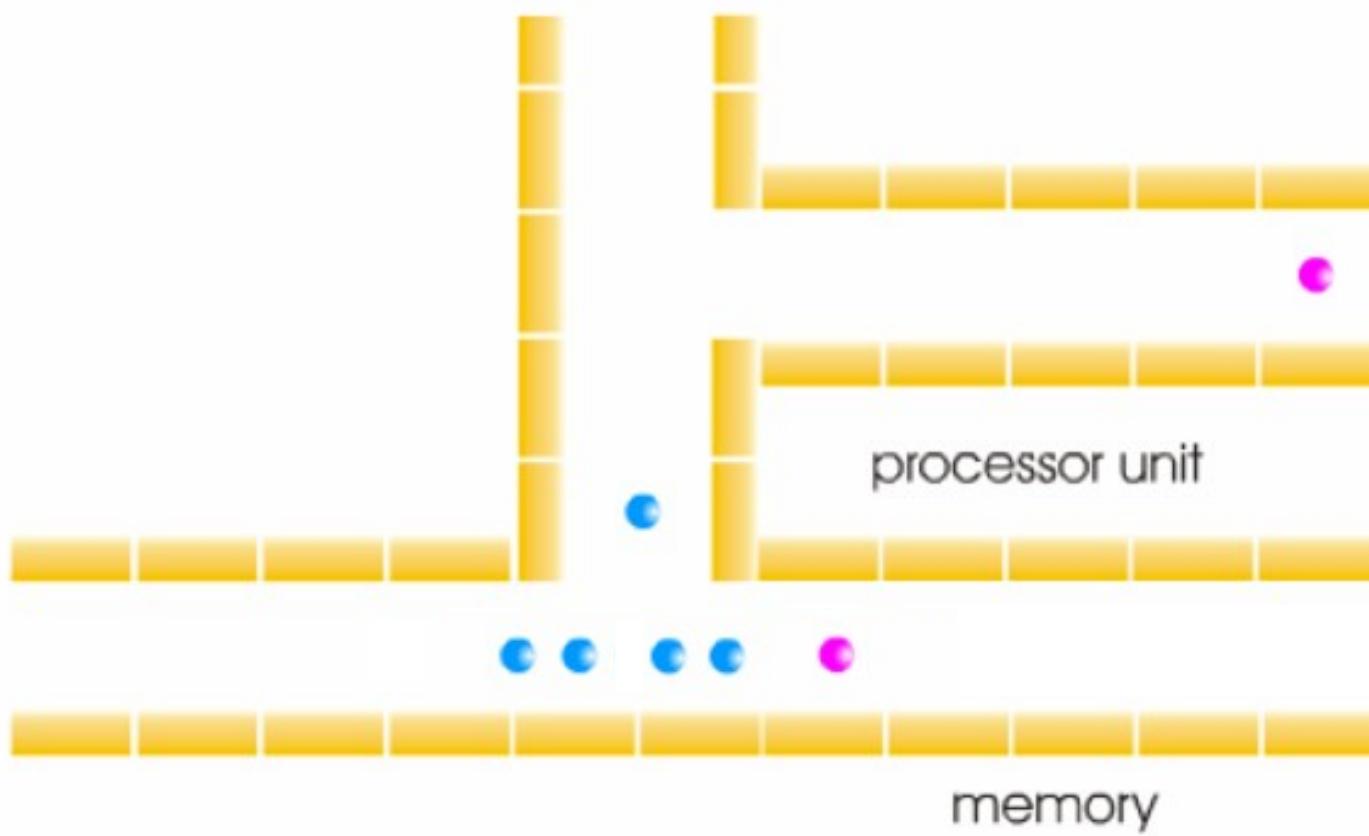
Scaling of ion trap quantum computers



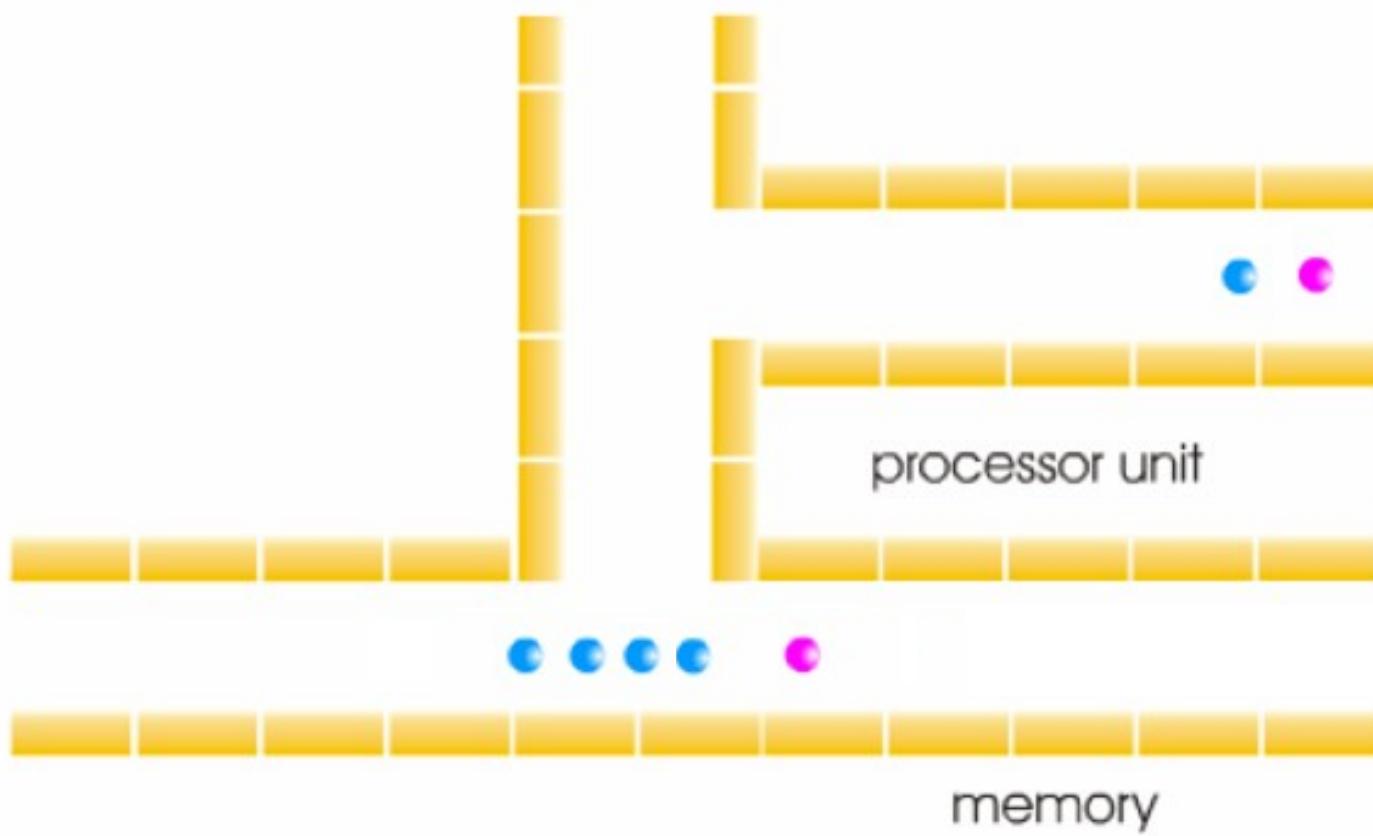
Scaling of ion trap quantum computers



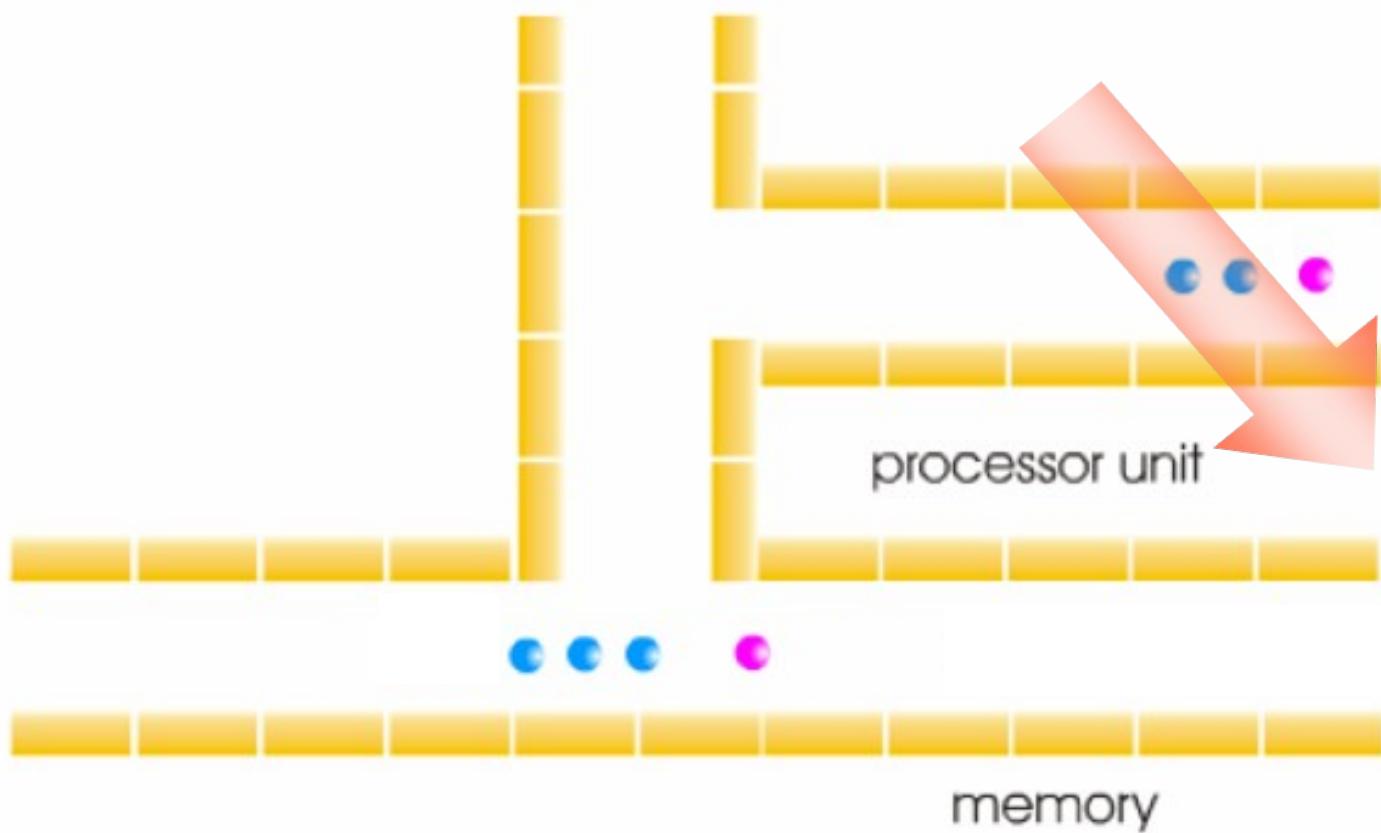
Scaling of ion trap quantum computers



Scaling of ion trap quantum computers

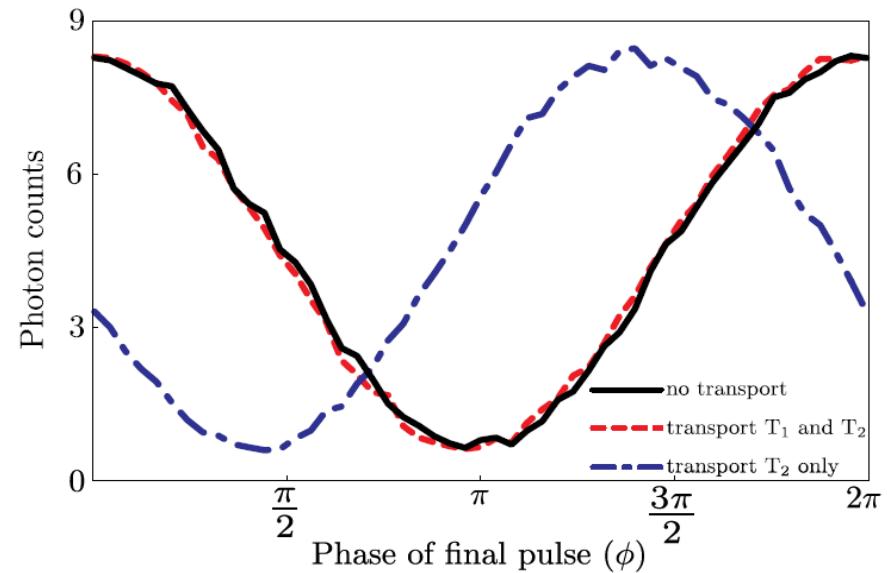
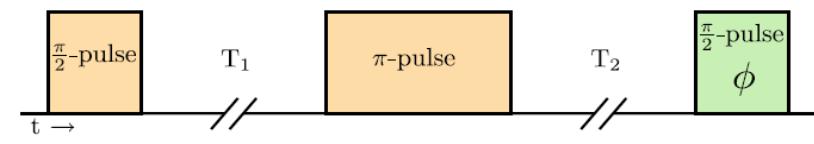
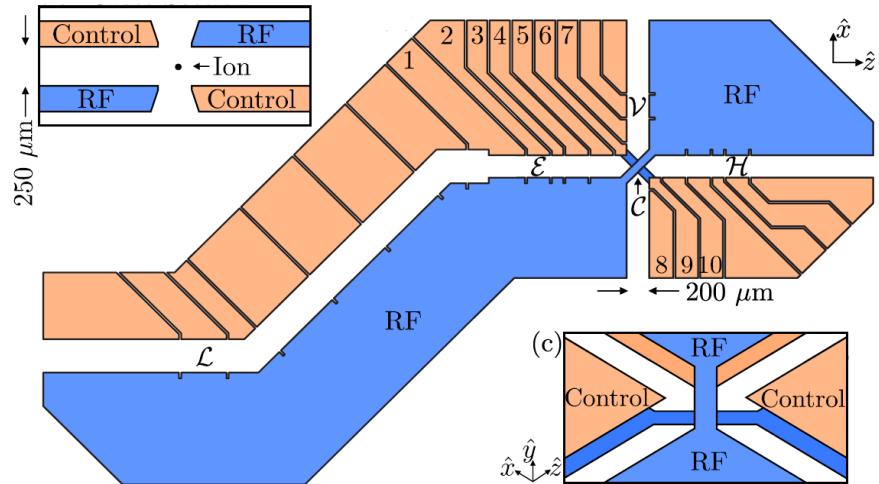


Scaling of ion trap quantum computers



„Architecture for a large-scale ion-trap quantum computer“,
D. Kielpinski et al., Nature **417**, 709 (2002).

Coherent transport through a junction

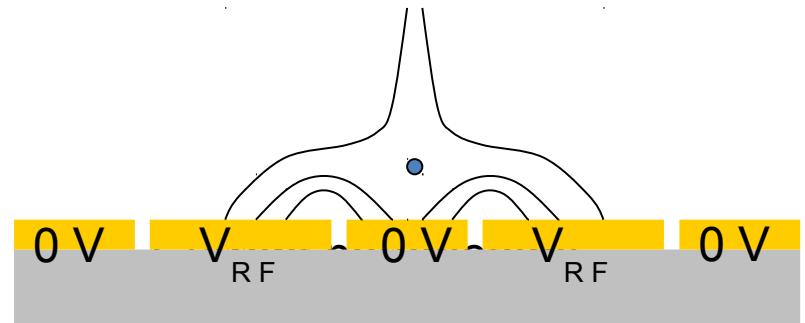
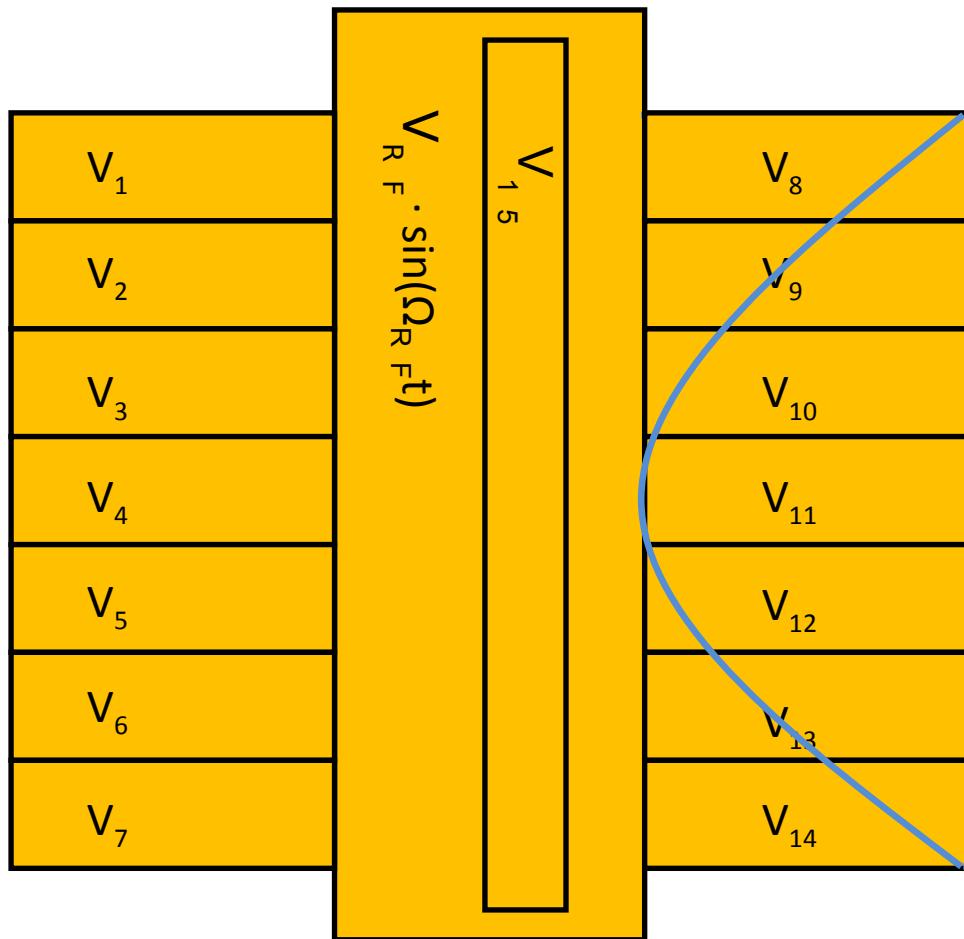


| Transport | Energy Gain (recooling method) | |
|---|-----------------------------------|----------------|
| | quanta/ion | quanta/trip |
| $\mathcal{E}-\mathcal{C}-\mathcal{E}$ | 1 ion | 3.2 ± 1.8 |
| $\mathcal{E}-\mathcal{C}-\mathcal{H}-\mathcal{C}-\mathcal{E}$ | 1 ion | 7.9 ± 1.5 |
| $\mathcal{E}-\mathcal{C}-\mathcal{V}-\mathcal{C}-\mathcal{E}$ | 1 ion | 14.5 ± 2.0 |
| $\mathcal{E}-\mathcal{C}-\mathcal{E}$ | 2 ions | 5.4 ± 1.2 |
| $\mathcal{E}-\mathcal{C}-\mathcal{H}-\mathcal{C}-\mathcal{E}$ | 2 ions | 16.6 ± 1.8 |
| $\mathcal{E}-\mathcal{C}-\mathcal{V}-\mathcal{C}-\mathcal{E}$ | 2 ions | 53.0 ± 1.2 |

NIST:

Blakestad, et al., "High fidelity transport of trapped-ion qubits through an X-junction trap array", arXiv:0901.0533v1

Surface traps



Ion height $\approx 220 \mu\text{m}$

$\Omega_{RF} \approx 2\pi \cdot 15 \text{ MHz}$

$V_{RF} \approx 100 \text{ V}$

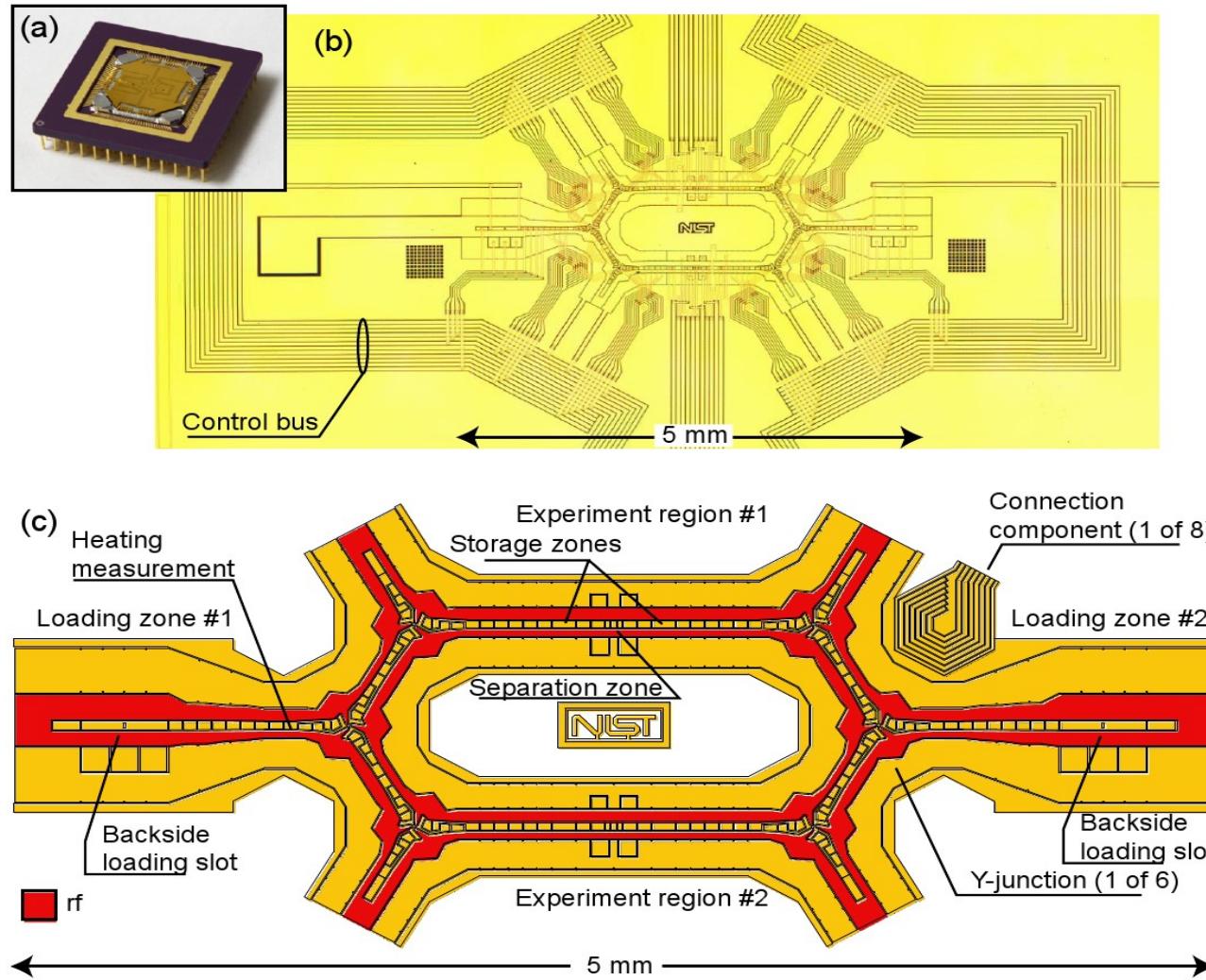
$V_{DC} < 10 \text{ V}$

$\omega_H \approx 2\pi \cdot 1.3 \text{ MHz}$

$\omega_V \approx 2\pi \cdot 1.5 \text{ MHz}$

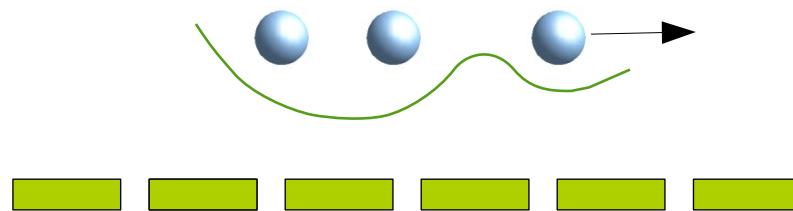
$\omega_A \approx 2\pi \cdot 300 \text{ kHz}$

Surface traps

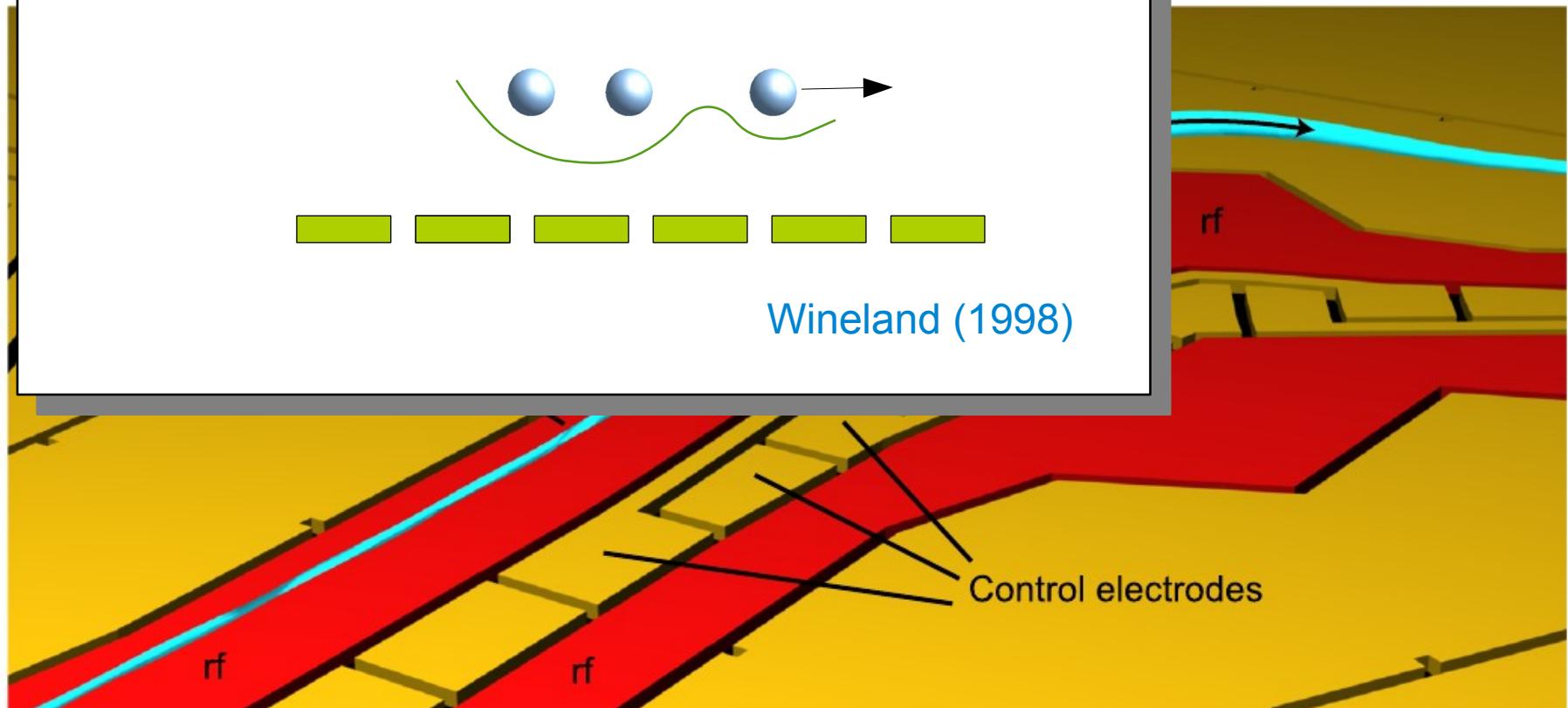


Surface traps

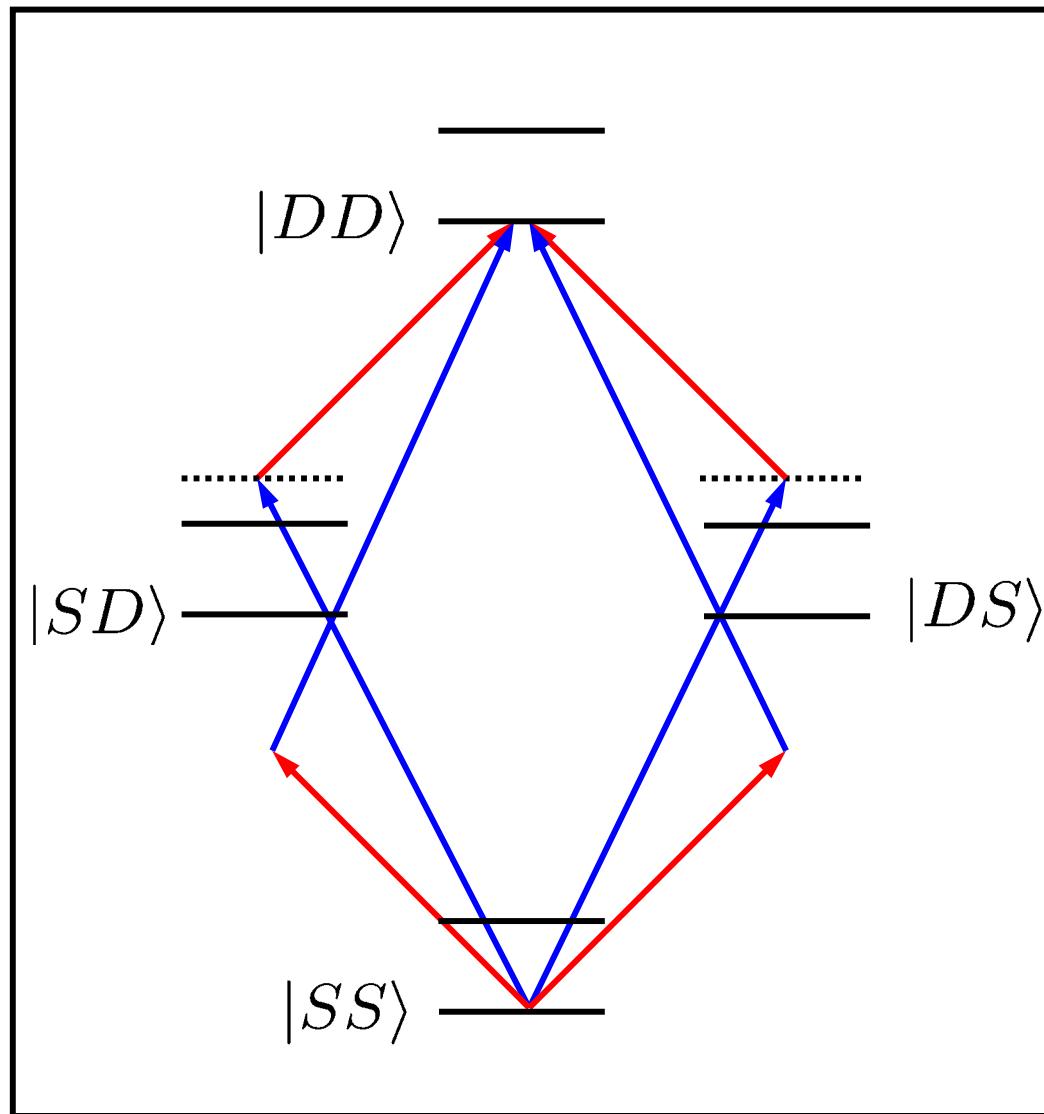
Need to split and merge ion strings fast for multi-qubit gates.



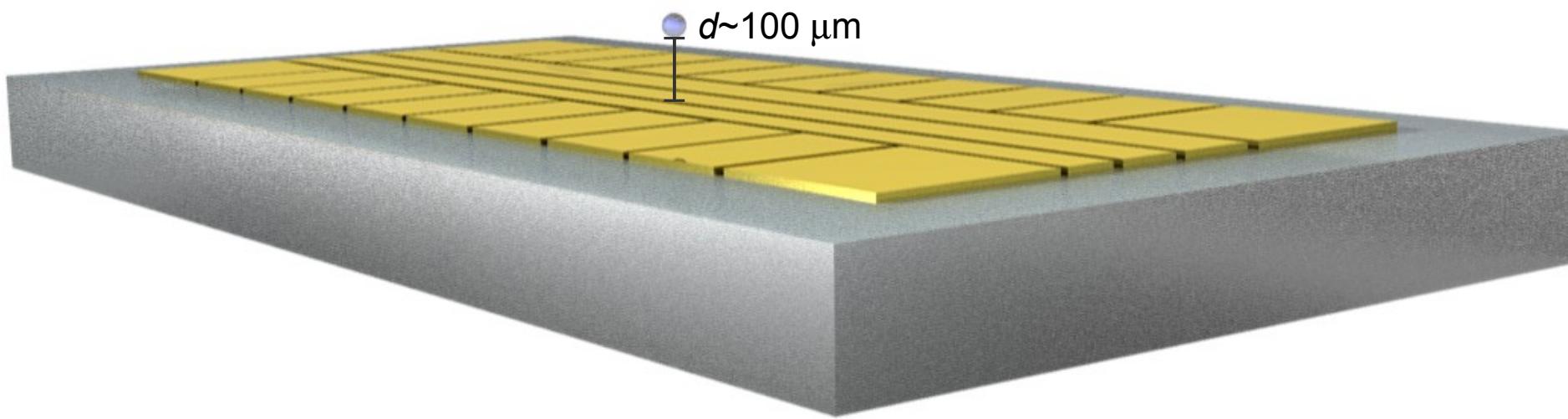
Wineland (1998)



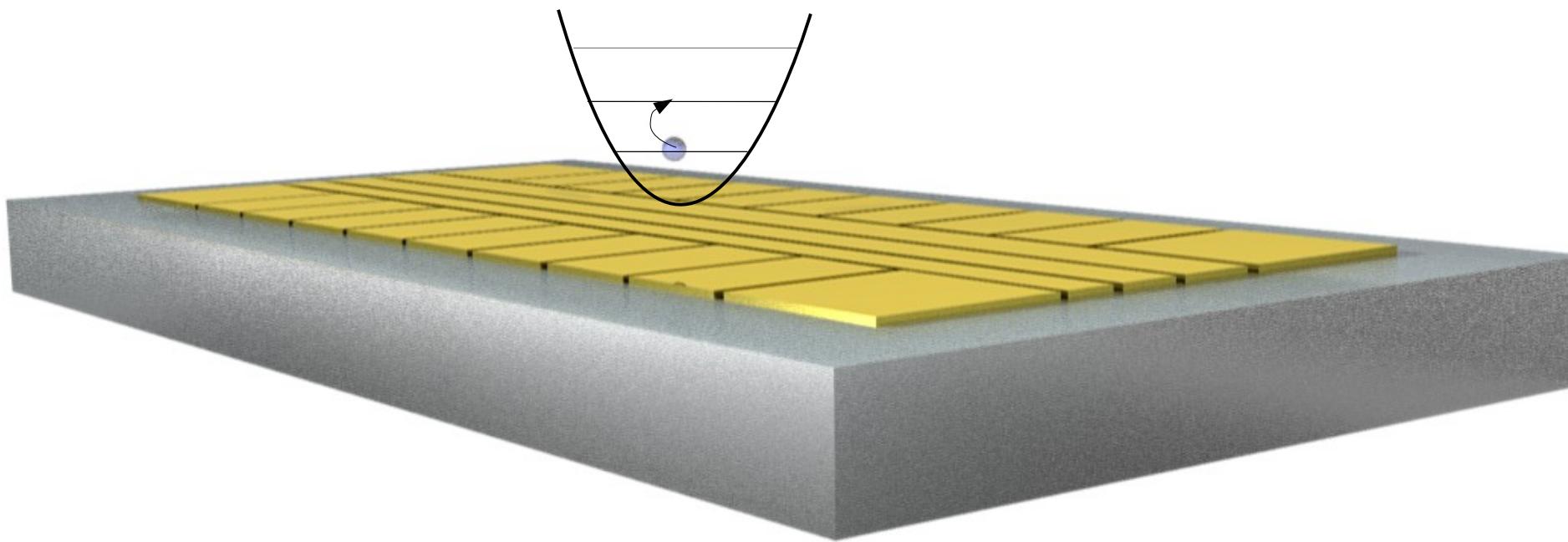
Entangling ions



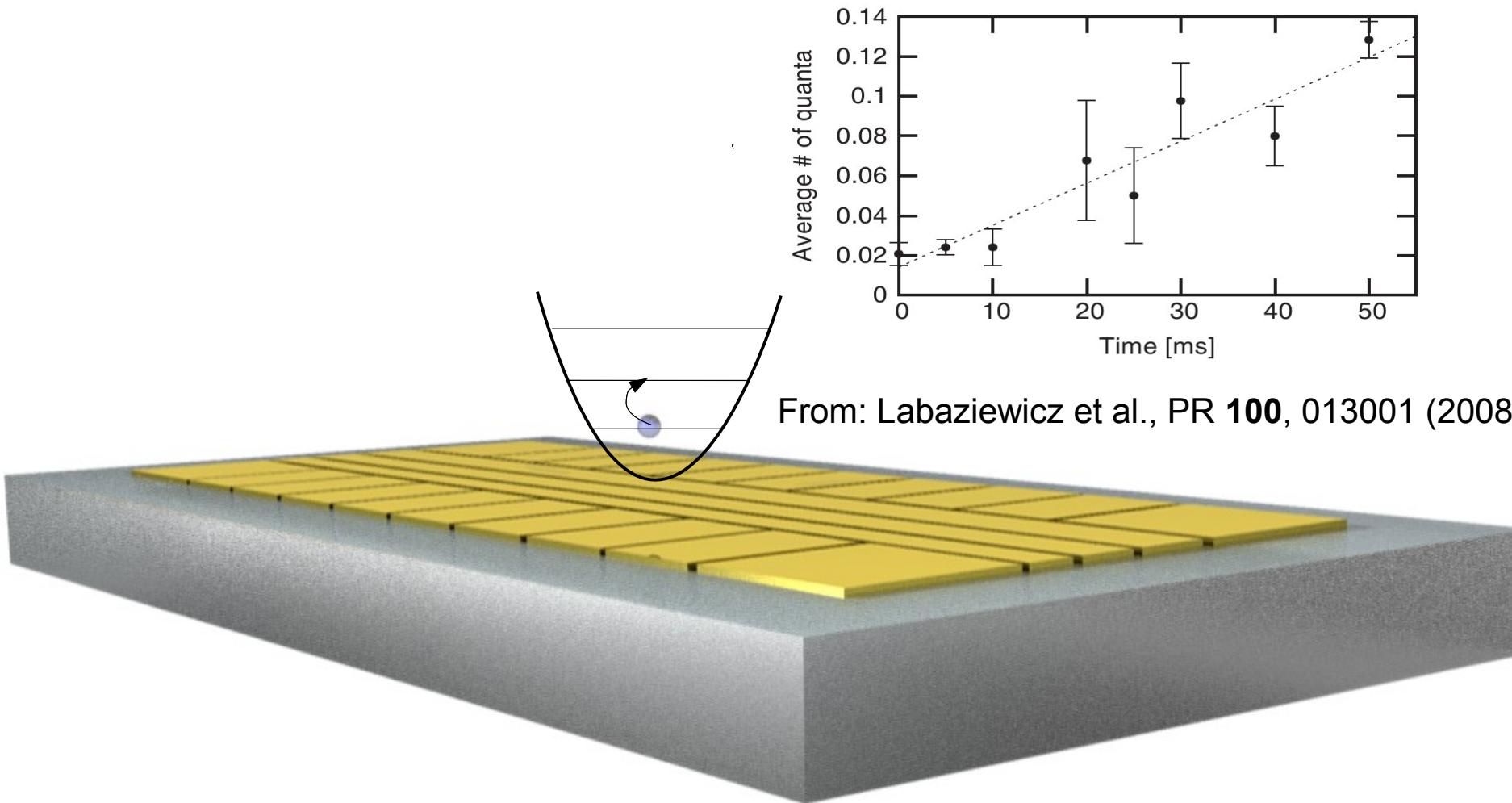
Motional decoherence



Motional decoherence

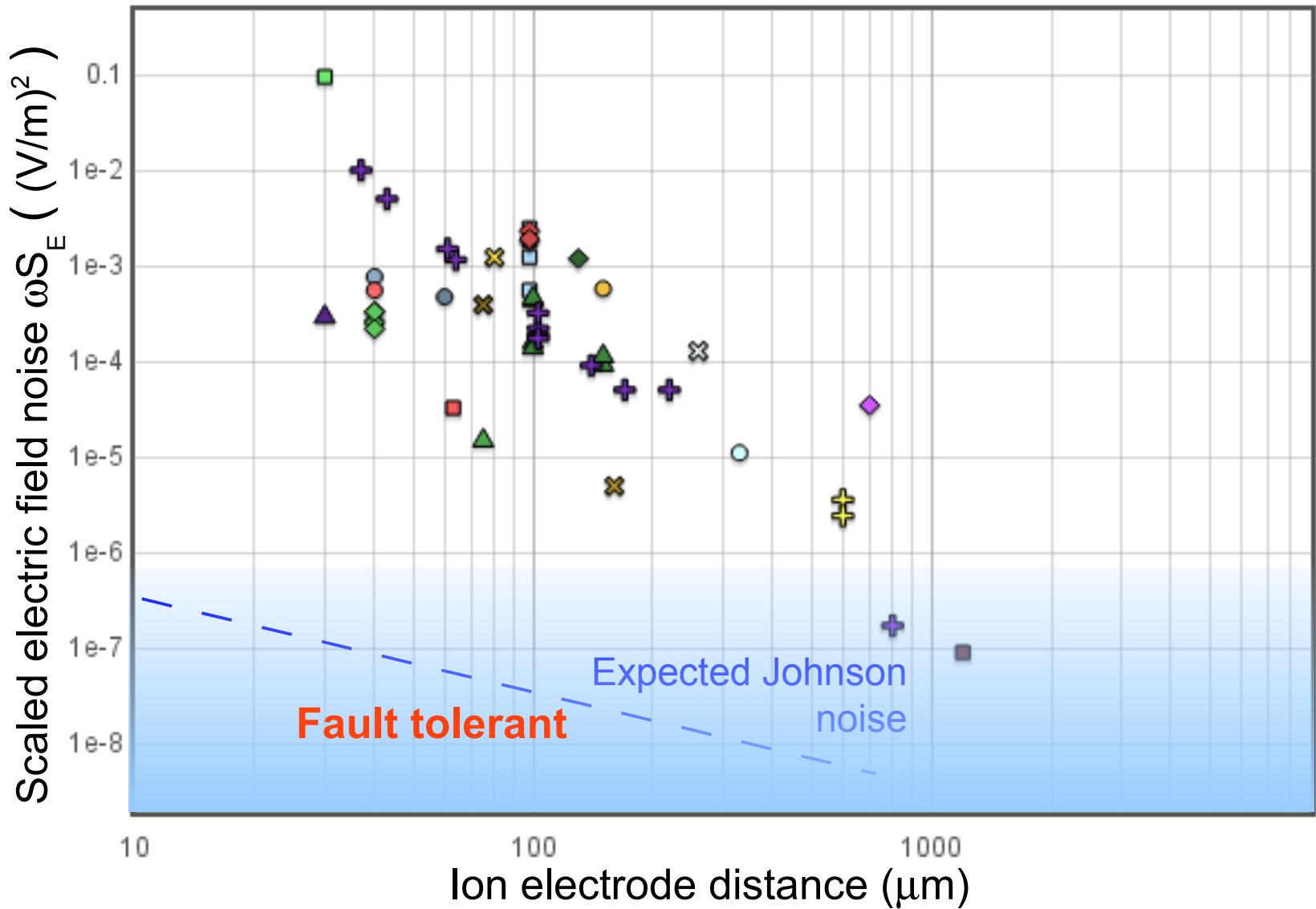


Motional decoherence



Excessive heating in ion traps

From: http://www.quantum.gatech.edu/heating_rate_plot.shtml

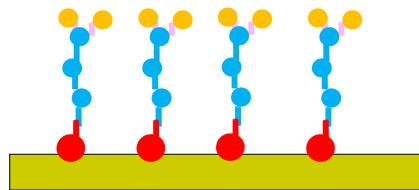


What is causing “the” anomalous heating ?

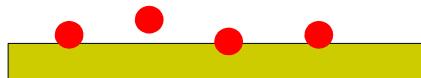
- fluctuating patch potentials, ad-atom diffusion (Wineland 1998)



- independently fluctuating dipoles (Daniilidis 2010)

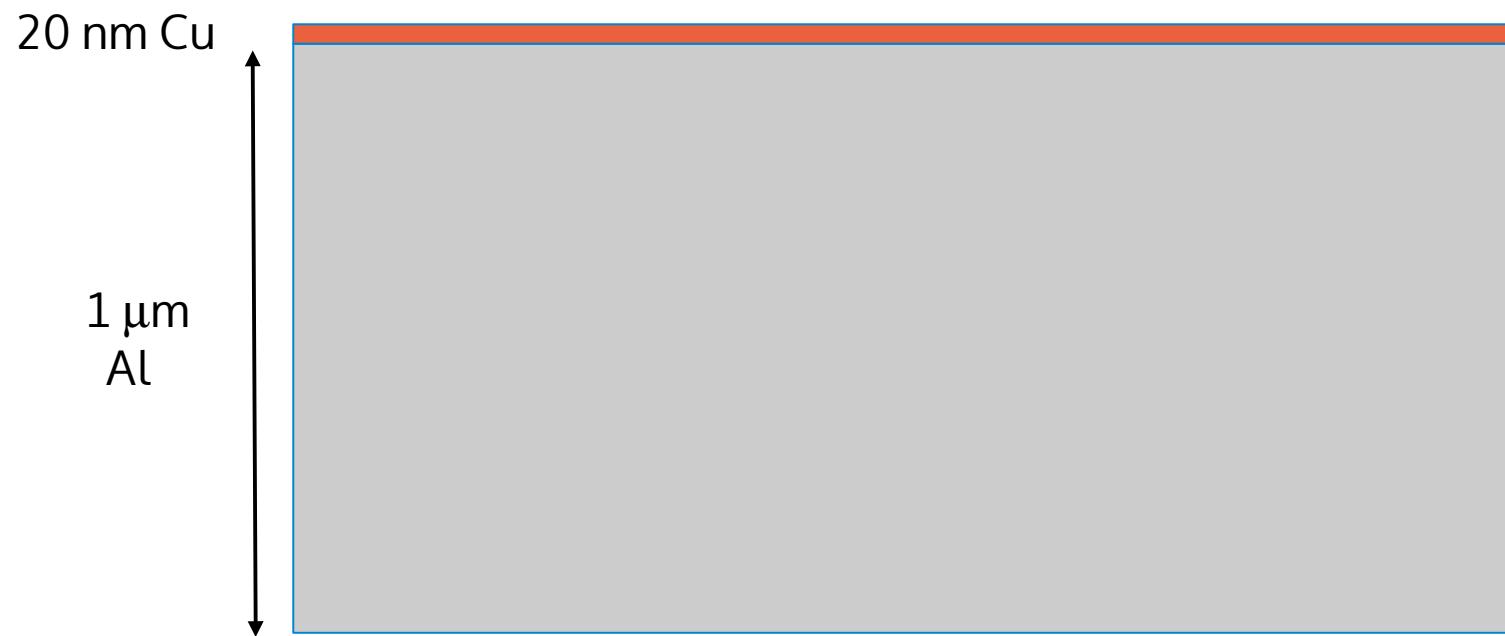


- fluctuating strength of dipoles (Safavi-Naini 2011)

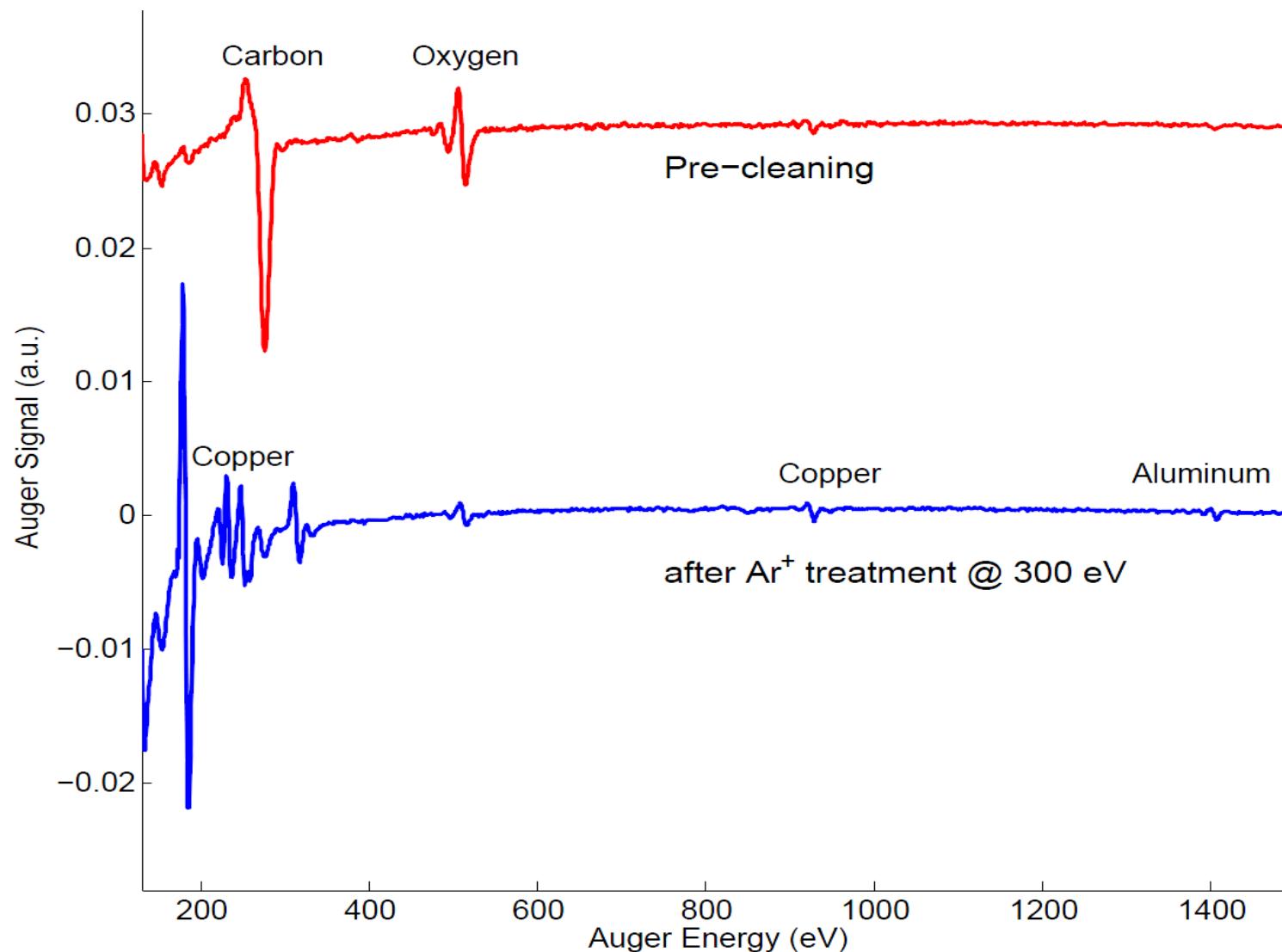


- all of the above and probably something else

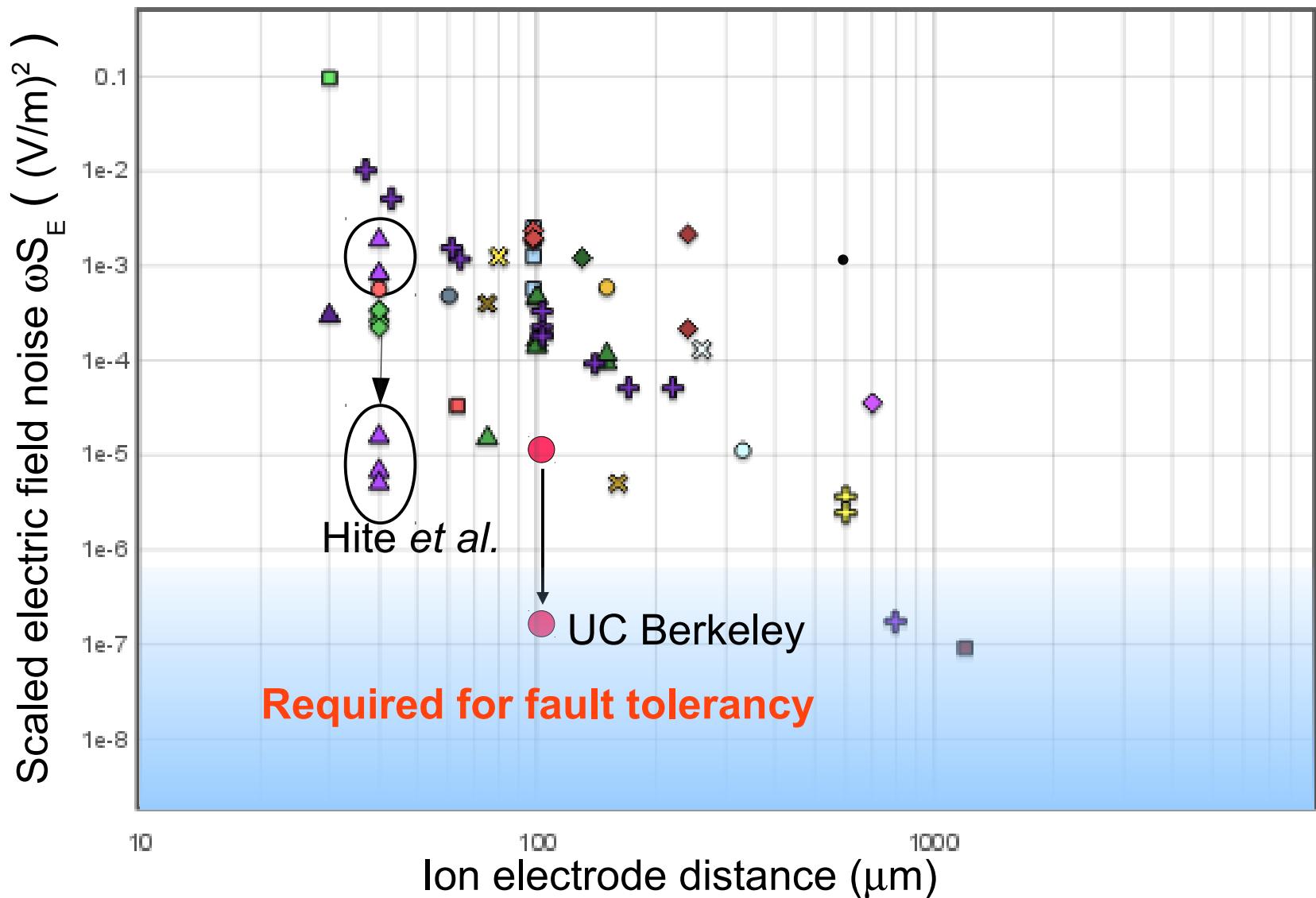
Copper on Aluminum trap



Monitoring the cleaning



Excessive heating in ion traps



More material:

Review on "Quantum computing with trapped ions", H. Häffner, C. F. Roos, R. Blatt, Physics Reports **469**, 155 (2008), <http://xxx.lanl.gov/abs/0809.4368>

Most recent progress:

NIST, Boulder

<http://www.nist.gov/pml/div688/grp10/quantum-logic-and-coherent-control.cfm>

Innsbruck

<http://www.quantumoptics.at/>

University of Maryland:

http://www.iontrap.umd.edu/publications/recent_pubs.html

