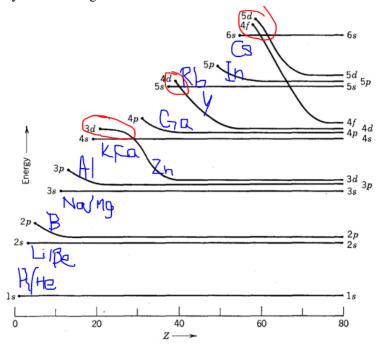
- 1. (a) Presumably because of electron screening. The p and d levels are further than the s, and the electrons in s can screen the electric field from the nucleus. As z increases, the nucleus becomes more positively charged and it reduces the effect of the screening, leading to a plateau.
 - (b) Maybe something like this:



- 2. (a) Because in the experiment, the X-ray emission is coming from the K series lines, which are the same transitions seen by the sodium atom during absorption spectroscopy.
 - (b) Presumably because tungsten is a high Z atom, leading to a larger magnetic moment and a larger energy splitting.
 - (c) Because the energy is proportional to \mathbb{Z}^2 in hydrogen-like states.
 - (d) The K series lines are the lowest ones (\rightarrow 1s). The L series are the next lowest, \rightarrow 2s/2p.
 - (e) No? The upper plot has a higher energy value.
- 3. (a) There are four states in the Slater basis,

$$\begin{split} |\alpha\beta\rangle_1 &= \mathcal{A}\,|\alpha\beta\rangle_a = |\alpha\beta\rangle_a - |\alpha\beta\rangle_{a'} \\ |\alpha\beta\rangle_2 &= \mathcal{A}\,|\alpha\beta\rangle_b = |\alpha\beta\rangle_b - |\alpha\beta\rangle_{b'} \\ |\alpha\beta\rangle_3 &= \mathcal{A}\,|\alpha\beta\rangle_c = |\alpha\beta\rangle_c - |\alpha\beta\rangle_{c'} \\ |\alpha\beta\rangle_4 &= \mathcal{A}\,|\alpha\beta\rangle_d = |\alpha\beta\rangle_d - |\alpha\beta\rangle_{d'} \end{split}$$

(b) The singlet state has an antisymmetric spin function, so the total wavefunction will be

$$(|\alpha\beta\rangle + |\beta\alpha\rangle) (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) = |\alpha\uparrow\rangle |\beta\downarrow\rangle - |\alpha\downarrow\rangle |\beta\uparrow\rangle + |\beta\uparrow\rangle |\alpha\downarrow\rangle - |\beta\downarrow\rangle |\alpha\uparrow\rangle$$
$$= |\alpha\beta\rangle_b - |\alpha\beta\rangle_c + |\alpha\beta\rangle_{c'} - |\alpha\beta\rangle_{b'}$$
$$= |\alpha\beta\rangle_2 - |\alpha\beta\rangle_3$$

For the triplet state, the spinfunction is symmetric, so

$$(|\alpha\beta\rangle - |\beta\alpha\rangle)(|\uparrow\uparrow\rangle) = |\alpha\uparrow\rangle |\beta\uparrow\rangle - |\beta\uparrow\rangle |\alpha\uparrow\rangle$$
$$= |\alpha\beta\rangle_1$$

I'm guessing the other ones will be $|\alpha\beta\rangle_4$. The last one will be

$$(|\alpha\beta\rangle - |\beta\alpha\rangle)(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) = |\alpha\uparrow\rangle |\beta\downarrow\rangle - |\beta\uparrow\rangle |\alpha\downarrow\rangle + |\alpha\downarrow\rangle |\beta\uparrow\rangle - |\beta\downarrow\rangle |\alpha\uparrow\rangle$$
$$= |\alpha\beta\rangle_2 + |\alpha\beta\rangle_3$$

- (c) 4 states, but I don't understand Slater stuff well enough to understand what this question is asking me.
- (d) We need Russell-Saunders coupling because we're dealing with the electron-electron interaction. Treating them separately results in $l = l_1 + l_2 = 2$, $s = s_1 + s_2 = 1$.

For 3 D, l=2 and s=1, $m_l=2,1,0,-1,-2$, $m_s=1,0,-1$, giving 15 unique combinations. For 1 D, $m_s=0$, giving 5 combinations.

For the P states, 9 and 3. For the S states, 3 and 1 combinations.

The sum of all this is 36 states.

(e) This is jj-coupling, right? So this would be needed when there's electron-electron interaction, but for larger atoms with high Z.

Each $j_i=\frac{3}{2},\frac{1}{2}$. This results in a total j=3,2,1,0. For j=3, $m_j=0,\pm 1,\pm 2,\pm 3$ with 7 states. For j=2, it's 5 states. For j=1, it's 3 states. For j=0, there's a single state.

This is a total of 16 states.

4. Omitting the \hbar , then we can factor j_c as

$$J_C = j_c \sqrt{1 + \frac{1}{j_c}}$$

...not actually sure what to do here. Should I be substituting in $j_a + j_b$ for j_c ? Or should I just check the inequality with $j_a(j_a + 1)$ and $j_b(j_b + 1)$?

5. (a) For the p subshell, $\ell = 1$ and the spherical harmonics squared are

$$Y_{10}^2 = \frac{3}{4\pi} \cos^2(\theta)$$

$$Y_{11}^2 = \frac{3}{8\pi} \sin^2(\theta)$$

$$Y_{1(-1)}^2 = \frac{3}{8\pi} \sin^2(\theta)$$

Summing these,

$$\sum_{k} Y_{1k}^{2} = \frac{3}{4\pi} \cos^{2}(\theta) + 2\left(\frac{3}{8\pi}\right) \sin^{2}(\theta)$$
$$= \frac{3}{8\pi}$$

(b) For the d shell,

$$Y_{20}^{2} = \frac{5}{16\pi} (3\cos^{2}(\theta) - 1)^{2}$$

$$Y_{2\pm 1} = \frac{15}{8\pi} (\cos\theta\sin\theta)^{2}$$

$$Y_{2\pm 2} = \frac{15}{32\pi} \sin^{4}(\theta)$$

Summing these in WolframAlpha results in

$$\sum_{k} Y_{2k}^2 = \frac{5}{4\pi}$$

6. (a)
$${}^2S_{3/2}$$
: $s = 1/2$, $l = 0$, $j = 3/2$
 3D_2 : $s = 1$, $l = 2$, $j = 2$
 5P_3 : $s = 2$, $l = 1$, $j = 3$

- (b) $^2S_{3/2}$ is impossible, since the total j exceeds l+s.
- 7. Gonna omit the *j* value to save space on the spectroscopy notation bit,

\overline{L}	S	J	Spect.
5	0	5	1 H
5	1	6, 5, 4	3 H
4	0	4	1 G
4	1	5, 4, 3	3 G
3	0	3	1 F
3	1	4, 3, 2	3 F
2	0	2	1 D
2	1	3, 2, 1	3 D
1	0	1	¹ P
1	1	2, 1, 0	³ P

8. (a) Since the three negative terms ($\epsilon_{321} = \epsilon_{132} = \epsilon_{213}$) squared are just 1. The other terms squared are 1, so summing this results in

$$\epsilon_{ijk}\epsilon_{ijk} = 6$$

- (b) There are only two non-zero components when k = l, resulting in $2\delta_{kl}$.
- 9. Using the determinate form with the Kronecker deltas,

$$\varepsilon_{ijk}\varepsilon_{lmn} = \begin{vmatrix} \delta_{il} & \delta_{im} & \delta_{in} \\ \delta_{jl} & \delta_{jm} & \delta_{jn} \\ \delta_{kl} & \delta_{km} & \delta_{kn} \end{vmatrix}$$

As i = l, then the only non-zero term is

$$\varepsilon_{ijk}\varepsilon_{imn} = \delta_{jm}\delta_{kn} - \delta_{jn}\delta_{km}$$

10. Given the rule provided,

$$\begin{aligned} \left[\mathbf{A} \times \left(\mathbf{B} \times \mathbf{C}\right)\right]_{i} &= \epsilon_{ijk} A_{j} \left(\mathbf{B} \times \mathbf{C}\right)_{k} \\ &= \epsilon_{ijk} A_{j} \epsilon_{klm} B_{l} C_{m} = \epsilon_{kij} \epsilon_{klm} A_{j} B_{l} C_{m} \\ &= \left(\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}\right) A_{j} B_{l} C_{m} \\ &= B_{i} A_{m} B_{m} - C_{i} A_{l} B_{l} \quad \Box \end{aligned}$$