1. (a) The LHS of the equation is the Q of the reaction, it's the available energy that the reaction can yield.

For this reaction, we know the momentum must be conserved, where

$$\mathbf{p}_{\gamma} = \mathbf{p}_{\text{recoil}}.$$

If we assume non-relativistic motion and use nuclear masses,

$$\mathbf{p}_{\gamma} = pc \implies p_{\gamma} = E_{\gamma}/c$$

$$E_D = p^2/2m = \frac{E_{\gamma}^2}{2\mathcal{M}_D c^2}.$$

We can equate the RHS of the equation as

$$E_{\gamma} + E_{\text{recoil}} = E_{\gamma} + \frac{E_{\gamma}^2}{2\mathcal{M}_D c^2}.$$

(b) Atomic masses should be OK to use on the left side for m_p and m_d , as we're dealing with atoms complete with bound electrons.

Using mass values from the internet,

939.565 MeV + 938.272 MeV - 1876 MeV =
$$E_{\gamma}$$
 + $\frac{E_{\gamma}^2}{2 \times 1.875 \text{ GeV}}$
 $E_{\gamma} = 1.84 \text{ MeV}.$

The quadratic term doesn't matter much, as that term is roughly only $900\,\mathrm{eV}$ —far less than the MeV range.

(c) In Williams, $Q_{\alpha} = 7.834 \, \mathrm{MeV}$. We could expect

$$Q_{\alpha} = [213.995186 \,\mathrm{u} - 209.984163 \,\mathrm{u} - 4.001506 \,\mathrm{u}] \,c^2$$

= 8.87 MeV,

which is pretty close to the 7.8 MeV, I guess...

I'm not sure how/what to check if it's consistent with the caption of the cloud chamber figure. We see a long-range (high-energy) decay occurring in the cloud chamber. Since the energy of the decay is higher than the other decay mode, shouldn't it be less common?

- 2. Using Table 5.1 and NIST's Atomic Mass Table, the Q values for each scenario are
 - (a) β^- emission.

$$Q_{\beta^{-}} = (\mathcal{M}(19, 40) - \mathcal{M}(20, 40)) c^{2}$$

= (39.963 998 u - 39.962 590 u) c^{2}
= 1.31 MeV.

(b) β^+ emission.

$$Q_{\beta^{+}} = (\mathcal{M}(19, 40) - \mathcal{M}(18, 40) - 2m_{e}) c^{2}$$

$$= (39.963 998 u - 39.962 383 u - 2 \times 5.485 799 \times 10^{-4} u) c^{2}$$

$$= 482 \text{ keV}.$$

(c) e^- capture.

$$Q_{EC} = (39.963\,998\,\mathrm{u} - 39.962\,383\,\mathrm{u})\,c^2$$

= 1.52 MeV.

3. In this sequence from 238 U to 206 Pb, there are 8 α -decays and 6 β -decays (can we use atomic masses even if there are an uneven number of beta and alpha decays?).

The Q value can be calculated as

$$\begin{split} Q &= (\mathcal{M}(\mathrm{U}, 938) - \mathcal{M}(\mathrm{Pb}, 206) - 8\mathcal{M}(\mathrm{He}, 4)) \, c^2 \\ &= (238.050\, 788\, \mathrm{u} - 205.974\, 46\, \mathrm{u} - 8 \times 4.002\, 603\, \mathrm{u}) \, c^2 \\ &= 51.7\, \mathrm{MeV} = 8.3 \times 10^{-12}\, \mathrm{J}. \end{split}$$

From room temperature, the total change in temperature required is roughly $1135\,\mathrm{K}$. For one gram, this requires an energy of

$$E = 0.12 \,\mathrm{J \cdot g^{-1} \cdot K^{-1}} \times 1135 \,\mathrm{K} = 136.2 \,\mathrm{J}.$$

In terms of decays, this needs 1.644×10^{13} decays, and in a gram of 238 U, there are 2.53×10^{21} atoms. I'm guessing the time required would be solved with

$$1.644 \times 10^{13} = 2.53 \times 10^{21} \exp(-t/4.5 \,\text{Gyr})$$

 $t = 85 \,\text{Gyr}$?

4. If we consider a particle of energy E tunneling through a potential barrier of energy U, its wavefunction is given by

$$\Psi(x) \propto e^{-kx}$$
 where $k = \frac{1}{\hbar} \sqrt{2m(U-E)}$.

Then, we can assume the rate of tunneling across the spatial interval [a, b] is given by its probability, i.e.

tunneling rate
$$\omega \propto |\Psi|^2$$

$$\propto \exp \biggl\{ -\frac{2}{\hbar} \sqrt{2m} \int_a^b \sqrt{U(x)-E} \,\mathrm{d}x \biggr\}.$$

Then, changing the bounds of the integral and taking the log (not sure if I really understand the reasoning of this part...),

$$\ln \omega = -2\sqrt{\frac{2m}{\hbar}} \int_0^{r=c/E} \sqrt{\frac{c}{r} - E} \, \mathrm{d}r$$

Changing variables to x=Er/c,

$$= \dots \frac{c}{\sqrt{E}} \int_0^1 \sqrt{1/x - 1} \, \mathrm{d}x.$$

Using an integration table, we can see that

$$\ln \omega \propto -1/\sqrt{E}$$
,

as depicted in Figure 6.3.