1. If we consider the product of the hermitian operators $\hat{x} \cdot \hat{p}$,

$$\langle \psi | \hat{x} \hat{p} \psi \rangle = \int \psi^* x \frac{\hbar}{i} \frac{\mathrm{d}}{\mathrm{d}x} \psi \, \mathrm{d}x$$
$$= -\int \frac{\hbar}{i} \psi \frac{\mathrm{d}}{\mathrm{d}x} x \psi^*$$
$$= \int (\psi^* \hat{p} \hat{x} \psi \, \mathrm{d}x)^*$$
$$= \langle \psi | \hat{p} \hat{x} \psi \rangle^*$$

This is only hermitian if $\hat{x}\hat{p} = \hat{p}x$.

For the operator $\frac{1}{2}(\hat{x}\hat{p}+\hat{p}\hat{x})$, its expectation is

$$\frac{1}{2} (\langle \psi | \hat{x} \hat{p} \psi \rangle + \langle \psi | \hat{p} \hat{x} \psi \rangle).$$

From the first part of the problem, we can see that this is equivalent to

$$\frac{1}{2} \left(\langle \psi | \hat{p} \hat{x} \psi \rangle^* + \langle \psi | \hat{x} \hat{p} \psi \rangle^* \right) = \frac{1}{2} \left\langle \psi \left| \frac{1}{2} (\hat{x} \hat{p} + \hat{p} \hat{x}) \psi \right\rangle^* \right.$$

Therefore, this is indeed hermitian.

2. (a) Let \hat{A} and \hat{B} be both hermitian operators, then the sum of these are hermitian,

(b) If \hat{Q} is hermitian and $\alpha \in \mathbb{C}$, their product is hermitian when

$$(\alpha \hat{Q})^{\dagger} = \alpha^* \hat{Q}^{\dagger}$$

$$\implies \alpha = \alpha^*$$

$$\implies \alpha \in \mathbb{R}$$

(c) For two hermitian operators \hat{A} and \hat{B} , the expectation of the product is

$$\left\langle f\middle|\hat{A}\hat{B}g\right\rangle = \left\langle \hat{B}\hat{A}f\middle|g\right\rangle$$

For the product to be hermitian, $\hat{A}\hat{B} = \hat{B}\hat{A}$. This was also shown in Problem 1.

(d) The position operator \hat{x} is hermitian as

$$\langle f|\hat{x}g\rangle = \int f^*xg\,\mathrm{d}x$$

As $x \in \mathbb{R}, x^* = x$,

$$\langle f|\hat{x}g\rangle = \int (fx)^* g \, \mathrm{d}x = \langle \hat{x}f|g\rangle$$

Because $\hat{x} = \hat{x}^{\dagger}$, \hat{x} is hermitian.

The Hamiltonian operator is

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(x)$$

Because \hat{p} is hermitian, \hat{p}^2 is also hermitian. From linearity, we can conclude that the Hamiltonian operator is hermitian.

3. (a) The hermitian conjurgate of x is x as it's real.

For i, the hermitian conjurgate is just the complex conjurgate -i.

For $\frac{d}{dx}$, by integration by parts,

$$\left\langle f \middle| \frac{\mathrm{d}}{\mathrm{d}x} g \right\rangle = \int f^* \frac{\mathrm{d}g}{\mathrm{d}x} \, \mathrm{d}x = 0 - \int g \frac{\mathrm{d}}{\mathrm{d}x} f^* \, \mathrm{d}x$$
$$= -\left\langle \frac{\mathrm{d}f^*}{\mathrm{d}x} \middle| g \right\rangle$$

The adjoint is then

$$\frac{\mathrm{d}}{\mathrm{d}x}^{\dagger} = -\frac{\mathrm{d}}{\mathrm{d}x}$$

(b) For the product of two operators $\hat{Q}\hat{R}$, its adjoint is

$$\begin{split} \left(\hat{Q}\hat{R}\right)^{\dagger} &\Longrightarrow \left\langle f\middle|Q^{\dagger}R^{\dagger}g\right\rangle = \left\langle \hat{Q}f\middle|\hat{R}^{\dagger}g\right\rangle \\ &= \left\langle \hat{R}\hat{Q}f\middle|g\right\rangle \end{split}$$

(c) For the raising operator \hat{a}_+ , the adjoint is

$$\hat{a}_{+}^{\dagger} = \frac{1}{\sqrt{2\hbar m\omega}} \left(-i\hat{p} + m\omega x \right) = \hat{a}_{-}$$

4. (a) The expectation of an anti-hermitian operator \hat{Q} is

$$\left\langle f \middle| \hat{Q}g \right\rangle = \left\langle \hat{Q}^{\dagger} f \middle| g \right\rangle = \left\langle -\hat{Q}f \middle| g \right\rangle$$

$$= -\left\langle f \middle| \hat{Q}g \right\rangle^{*}$$

The condition $\left\langle \hat{Q} \right\rangle^* = -\left\langle \hat{Q} \right\rangle$ is only true for imaginary numbers.

- (b) The eigenvalues of an operator are its expectation, so for the eigenvalue $\hat{Q}f = qf$, it must be true that $q^* = -q$, therefore its eigenvalue q is imaginary.
- (c) If we consider two eigenstates of an anti-hermitian operator \hat{Q} ,

$$\hat{Q}\phi_n = a_n\phi_n$$
$$\hat{Q}\phi_m = a_m\phi_m,$$

and consider the inner product

$$\left\langle \phi_n \middle| \hat{Q}\phi_m \right\rangle = \left\langle \phi_n \middle| a_m \phi_m \right\rangle = a_m \left\langle \phi_n \middle| \phi_m \right\rangle$$

and by the anti-hermitian nature of the original inner product,

$$\left\langle \phi_n \middle| \hat{Q}\phi_m \right\rangle = \left\langle -\hat{Q}\phi_n \middle| \phi_m \right\rangle = -a_n \left\langle \phi_n \middle| \phi_m \right\rangle.$$

Equating these expressions together (as they're the same thing),

$$0 = (a_m + a_n) \langle \phi_n | \phi_m \rangle$$

This is only true when ϕ_n and ϕ_m are orthogonal.

(d) For the commutator of two hermitian operators,

$$\begin{split} \left[\hat{A}, \hat{B} \right] &= \left(\hat{A}\hat{B} - \hat{B}\hat{A} \right) \\ \Longrightarrow \left\langle f \middle| \left(\hat{A}\hat{B} - \hat{B}\hat{A} \right) g \right\rangle &= \left\langle f \middle| \hat{A}\hat{B}g \right\rangle - \left\langle f \middle| \hat{B}\hat{A}g \right\rangle \\ &= \left\langle \hat{B}^{\dagger}\hat{A}^{\dagger}f \middle| g \right\rangle - \left\langle \hat{A}^{\dagger}\hat{B}^{\dagger}f \middle| g \right\rangle \\ &= - \left\langle f \middle| [A, B]g \right\rangle \end{split}$$

For anti-hermitian operators, we can follow the same general steps until:

$$\dots = -\left\langle \hat{B}^{\dagger} \hat{A}^{\dagger} f \middle| g \right\rangle + \left\langle \hat{A}^{\dagger} \hat{B}^{\dagger} f \middle| g \right\rangle$$
$$= \left\langle f \middle| [A, B] g \right\rangle$$

The commutator of two anti-hermitian operators is hermitian.

(e) If we let $\hat{Q} = \hat{A} + \hat{B}$, then we can write

$$\hat{A} = \frac{1}{2} \left(\hat{Q} + \hat{Q}^{\dagger} \right)$$

$$\hat{B} = \hat{Q} - \hat{A} \frac{1}{2} \left(\hat{Q} - \hat{Q}^{\dagger} \right)$$