

# Homework 4, High-resolution shock-capturing methods

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## 1 Shallow water with non-horizontal bottom

The shallow water model of HW2 is now extended to a non-horizontal bottom “bathymetry”  $B(x)$ .

$$\begin{pmatrix} h_t \\ hu_t \end{pmatrix} + \begin{pmatrix} h_x u + hu_x \\ hu u_x + gh(h_x + B_x) \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad (1)$$

a) Show that still water ( $u = 0$ ) must have, as it should, a horizontal water level.

If  $u = 0$  for all space and time, then  $u_t$  and  $u_x$  are also zero. Equation (1) becomes

$$\begin{pmatrix} h_t \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ gh(h_x + B_x) \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

This means that  $h_t = 0$  and  $h_x = -B_x$ . We know that  $B(x)_t = 0$  because it isn't a function of time. The total height,  $D(x, t)$ , is  $D(x, t) = h(x, t) + B(x)$ . A horizontal water level means that  $D(x, t) = \text{constant}$  and  $D(x, t)_x = D(x, t)_t = 0$ . Equation (2) shows that derivatives with respect to  $x$  and  $t$  are zero and therefore showing that the total height  $D(x, t)$  is constant.

$$\begin{aligned} D(x, t)_x &= h_x + B_x = h_x + (-h_x) = 0 \\ D(x, t)_t &= h_t + B_t = 0 + 0 = 0 \end{aligned} \quad (2)$$

b) Write the equation in conservation form for  $h$  and  $m = hu$ .

$$\begin{pmatrix} h \\ m \end{pmatrix}_t + \begin{pmatrix} m \\ f_2(h, m) \end{pmatrix}_x = \begin{pmatrix} 0 \\ s(h, m, x) \end{pmatrix} \quad (3)$$

The first equation in Equation (3) is proved using the product rule,  $m_x = h_x u + hu_x$ . The second equation needs to be written out explicitly and solved for the unknowns. Substituting  $m_t = h_t u + hu_t$  into Equation (3)

$$(h_t u + hu_t) + [hu u_x + gh(h_x + B_x)] = h_t u$$

and substituting  $u = m/h$  and  $h_t = -m_x$

$$m_t + [h \left( \frac{m}{h} \right)_x + gh(h_x + B_x)] = (-m_x) \left( \frac{m}{h} \right)$$

and bringing  $x$  and  $t$  derivatives together and putting derivatives other than  $h$  or  $m$  to the source function

$$\begin{aligned} m_t + \underbrace{\left[ m \left( \frac{m}{h} \right)_x + m_x \left( \frac{m}{h} \right) \right]}_{(m^2/h)_x} + \underbrace{ghh_x}_{(gh^2/2)_x} &= -ghB_x \\ m_t + \underbrace{\left( \frac{m^2}{h} + \frac{gh^2}{2} \right)_x}_{f_2(h, m)} &= \underbrace{-ghB_x}_{s(h, m, x)} \end{aligned}$$

Equation (3) written in conservation form becomes,

$$\begin{pmatrix} h \\ m \end{pmatrix}_t + \begin{pmatrix} m \\ \frac{m^2}{h} + \frac{gh^2}{2} \end{pmatrix}_x = \begin{pmatrix} 0 \\ -ghB_x \end{pmatrix}$$

## 2 First order Roe scheme

- (a) Explain how to choose  $m(x,0)$ , to make a single pulse (not two!) based on the linearized problem.

You can make it a single pulse by giving the wave an initial speed proportional to its height in one direction. Since  $m = hv$ , we can set  $m$  to  $hc$ , where the characteristic speed is  $c = \sqrt{gH}$ , and  $H = 0$  since we want a linearized approximation.

- (b) Use this initial data and run the Roe solver with  $n=80, 160, 320$  and the CFL number as large as possible without instability, until the wave has been fully reflected at  $L$ . Plot the wave shapes.

Figure 1 shows the plots using different  $N$  values. The  $N$  value seemed to affect the sharpness of the front edge of the wave. As  $N$  increased, the sharpness increased.

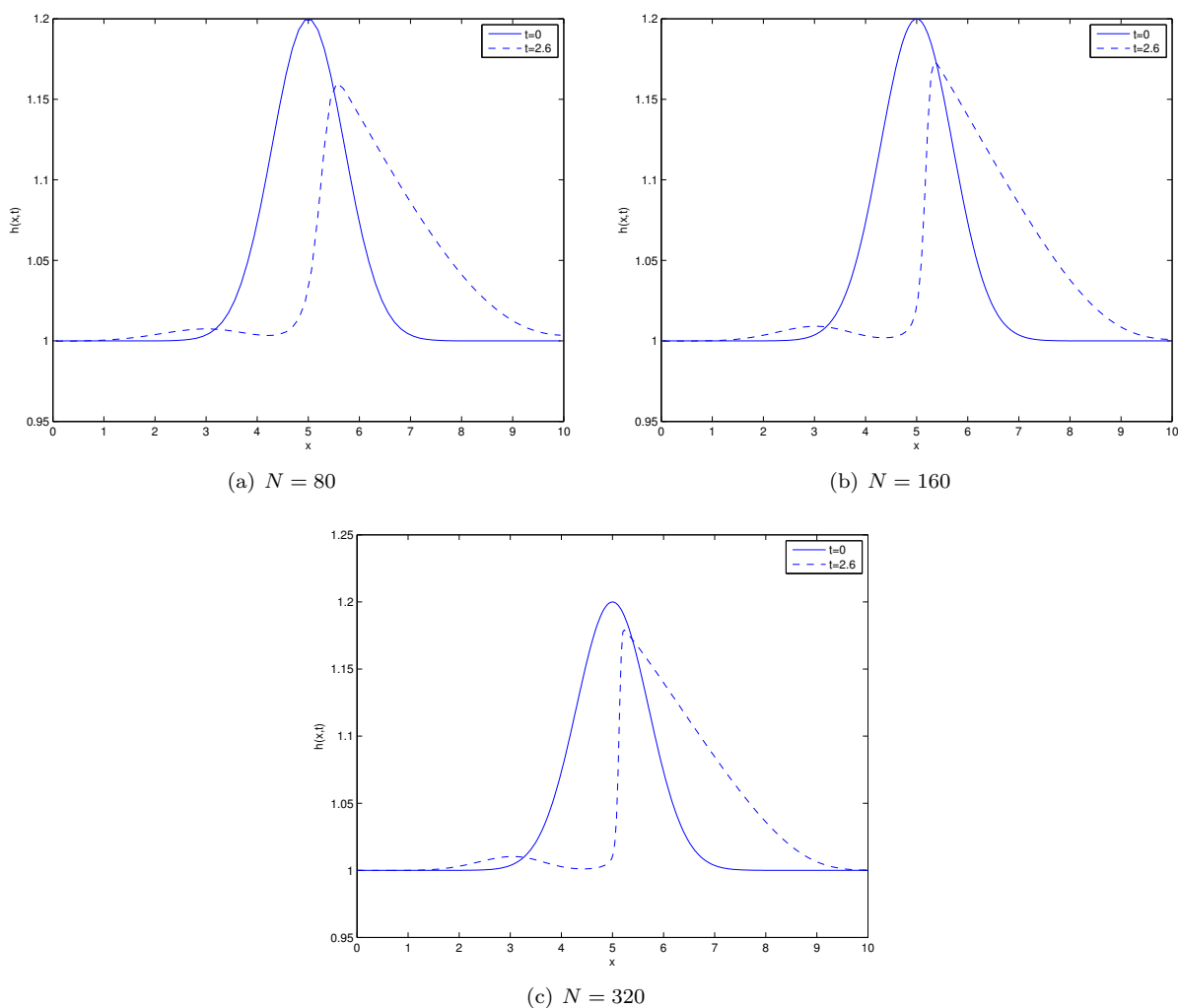


Figure 1: These are the plots with different  $N$  values. The CFL condition was set to 80% of the value that would be ideal for a linear problem. The wave edges get sharper when  $N$  is increased

- (c) Compare with solutions obtained using your Lax-Friedrichs solver from HW2.

The solution obtained here looks very similar to the Lax-Friedrichs solver from HW2 as the shape of the wave forms a sharp edge and looks like a triangle instead of a Gaussian. This solution isn't stable at the CFL condition and we must make the time step a little smaller. The Lax-Friedrichs solution was also not stable at the CFL condition, but I was forced to decrease the time step by a much larger amount in order to see stability for a few seconds.

- (d) Comment about order of accuracy, dissipation and dispersion. The dissipation around the edges decreases as we increase the value of  $N$ . But increasing the value of  $N$  too much will only cause instabilities.
- (e) Then try also a higher pulse, say  $a = 2H$ . Comment on how well the single pulse initial data works.

A pulse with  $a > 1/5H$  gives unstable results using a time step of 80% of the CFL condition, so I had to reduce it to 40% in order to get the results shown in Figure 2. This figure is different from Figure 1 in that the height of the left moving pulse isn't as negligible as it once was. The overlap of the smaller wave is much more noticeable when  $a$  is large compared to  $H$ .

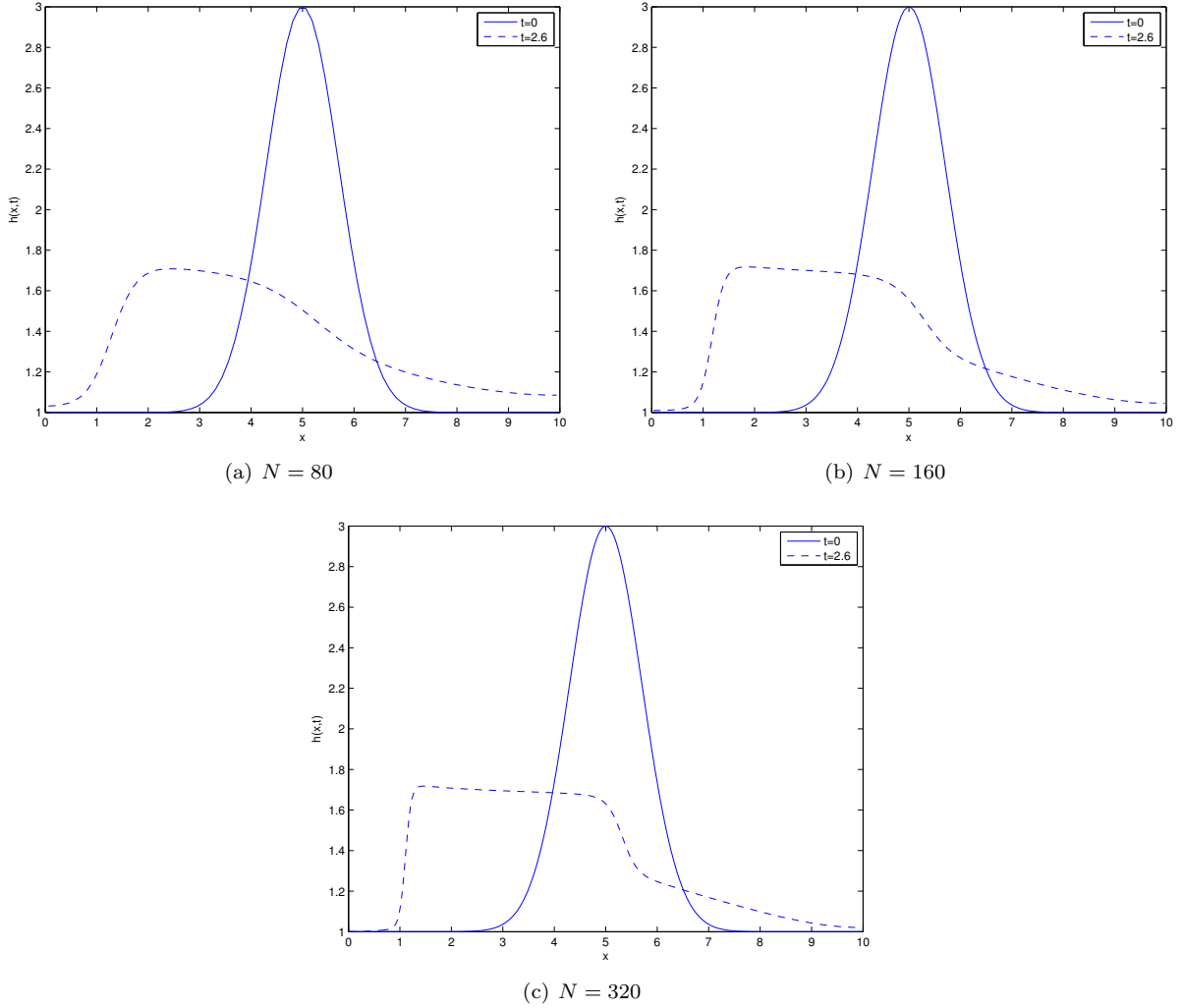


Figure 2: These are the plots with different  $N$  values. The CFL condition was set to 40% of the value that would be ideal for a linear problem. The wave edges get sharper when  $N$  is increased, but changing the value of  $a$  has also changed the characteristic speed and shape of the wave after it recombines with its small wave that moves leftward.

### 3 Boundary conditions

If the flow is sub-critical,  $|u(x)| < \sqrt{gh(x)}$ , then for Dirichlet boundary conditions, we should set  $h(x)$  to be large enough to satisfy the sub-critical condition at the boundaries. The same should be done for super-critical conditions at the boundaries.

## 4 Smooth steady state solutions

Figure 3 shows the plot with

$$B(x) = \begin{cases} B_0 \cos\left(\frac{\pi(x-L/2)}{2r}\right), & \text{if } |x - L/2| < r \\ 0, & \text{if } |x - L/2| > r \end{cases}$$

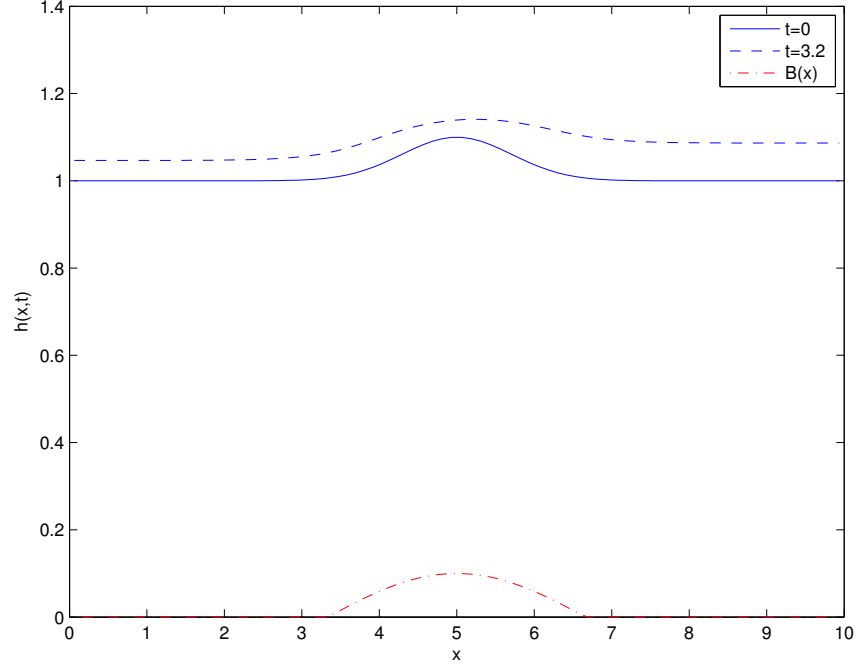


Figure 3: Solution with an uneven bottom that contributes to the wave shape and speed. The depth of the wave is related to the bottom surface as well as its own surface.