Linearly Stabilized Schemes for the Time Integration of Stiff Nonlinear PDEs

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Introduction

- Focus on time stepping for stiff nonlinear PDEs.
 - Stability
 - Accuracy
 - Efficiency
 - Simplicity



Consider the heat equation,

$$u_t = u_{xx}, \quad x \in \Omega, t > 0.$$

Discretize in space:

$$U'=LU, \quad U\in\mathbb{R}^N, t>0.$$

Explicit: $U^{n+1} = G(U^n, U^{n-1}, \dots, LU^n, LU^{n-1}, \dots)$, but $\Delta t \leq Ch^2$. Implicit: $AU^{n+1} = b$; unconditionally stable, but must solve a linear system.

Now compare with

$$u_t = \frac{u_{xx}}{1 + u_x^2} - \frac{1}{u}, \quad x \in \Omega, t > 0.$$

and

$$U' = F(U), \quad U \in \mathbb{R}^N, t > 0.$$

Explicit: $U^{n+1} = G(U^n, U^{n-1}, \dots, F(U^n), F(U^{n-1}), \dots)$, but $\Delta t \leq Ch^2$. Implicit: $AU^{n+1} = b(U^{n+1})$; unconditionally stable, but must solve a nonlinear system because nonlinearity is in the stiff term.

Comparing side-by-side:

$$u_t = u_{xx}, \quad x \in \Omega, t > 0,$$

Explicit: $\Delta t \leq Ch^2$

Implicit: unconditionally stable;

solution to linear system

 $u_t = \frac{u_{xx}}{1 + u_x^2} - \frac{1}{u}, \quad x \in \Omega, t > 0,$

Explicit: $\Delta t \leq Ch^2$

Implicit: unconditionally stable; solution to nonlinear system

Summary: What We Like

Explicit: simple; handles nonlinear terms with no added difficulty.

Implicit: large time steps

Modify the equation,

$$u_{t} = \frac{u_{xx}}{1 + u_{x}^{2}} - \frac{1}{u} - u_{xx} + u_{xx}, \quad x \in \Omega, t > 0,$$

$$\downarrow$$

$$U' = F(U) - LU + LU, \quad U \in \mathbb{R}^{N}, t > 0,$$

and use implicit-explicit time stepping, e.g.

$$\frac{U^{n+1}-U^n}{\Delta t}=F(U^n)-LU^n+LU^{n+1}.$$



Linear Stability

More generally, from U' = F(U), we can modify as

$$U' = \underbrace{F(U) - pLU}_{(\star)} + pLU, \quad p > 0,$$

and apply a time stepping scheme that treats (\star) explicitly. Key question: Is this unconditionally stable?



Scalar test equation

Standard case:

$$U' = F(U)$$

 $\mbox{Linearize} \rightarrow \mbox{Diagonalize} \rightarrow \mbox{Test} \\ \mbox{equation:} \\$

$$u' = \lambda u$$

Apply time stepping method:

With linear modification:

$$U' = F(U) - pLU + pLU$$

Linearize \rightarrow Diagonalize \rightarrow Test equation:

$$u' = \lambda u - p\lambda u + p\lambda u$$
$$= (1 - p)\lambda u + p\lambda u$$

$$u^{n+1} = \xi(\lambda \Delta t) u^n$$

Stability constraint:

$$|\xi(\lambda \Delta t)| \leq 1$$



Implicit-explicit Euler

Applied to the test equation $u' = (1 - p)\lambda u + p\lambda u$, yields

$$\frac{u^{n+1}-u^n}{\Delta t}=(1-p)\lambda u^n+p\lambda u^{n+1}.$$

The amplification factor is

$$\xi_1(\lambda \Delta t) = \frac{1 + (1 - p)\lambda \Delta t}{1 - p\lambda \Delta t}.$$

Impose unconditional stability:

$$|\xi_1(\lambda \Delta t)| \le 1$$
 for all $\lambda \Delta t < 0 \implies p \ge 1/2$.



Explicit-implicit-null (EIN)

Use Richardson extrapolation to get second order. The amplification factor is

$$\xi_{EIN}(\lambda \Delta t) = 2\xi_1^2(\lambda \Delta t/2) - \xi_1(\lambda \Delta t).$$

and

$$|\xi_{EIN}(\lambda \Delta t)| \le 1$$
 for all $\lambda \Delta t < 0 \implies p \ge 2/3$.



Implicit-explicit multistep methods

An alternative for second and higher order methods: IMEX multistep methods.

Order	Method	$p \in$
2	SBDF2 CNAB mCNAB CNLF	$[3/4,\infty)$ $[1,\infty)$ $[8/9,\infty)$ $[1/2,\infty)$
3	SBDF3	[7/8, 2]
4	SBDF4	[11/12, 5/4]

Comparing the methods

Do the methods work as advertised? Examine this with two test problems,

$$u_t = \frac{u_{xx}}{1 + u_x^2} - \frac{1}{u},$$

and

$$u_t = \Delta(u^5).$$

Test Problem 1

First test problem:

$$u_t = \frac{u_{xx}}{1 + u_x^2} - \frac{1}{u}, \quad 0 < x < 10, \quad t > 0,$$

with initial condition

$$u(x,0) = 1 + 0.10\sin\left(\frac{\pi}{5}x\right),$$

and boundary conditions u(0, t) = 1 = u(10, t).



Numerical convergence test

Figure: Numerical convergence of linearly stabilized schemes.

Spatial discretization: Uniform grid, centred differences, N=2048.

Final time: T = 0.35.

Reference solution: Explicit third order Runge-Kutta, $\Delta t = 1.46 \times 10^{-5}$.

Stabilized by adding and subtracting pu_{xx} .

Failure of SBDF3 and SBDF4

How did we choose p? Consider

$$u' = \lambda u - p\lambda u + p\lambda u$$

and

$$U' = F(U) - pLU + pLU.$$

With the test equation, we derived a restriction on p. More generally, the restriction applies to $p\lambda_L/\lambda_F$. For test problem 1 with centred differences, we find

$$\frac{\rho\lambda_L}{\lambda_F}pprox
ho(1+(D_1ar{u}_j^n)^2),$$



Failure of SBDF3 and SBDF4

The selection of p for SBDF3 is dictated by

$$\max_{1 \leq j \leq N} \frac{7}{8} \frac{1}{1 + (D_1 \bar{u}_j^n)^2} \leq p \leq \min_{1 \leq j \leq N} \frac{2}{1 + (D_1 \bar{u}_j^n)^2}.$$

Figure: Development of instabilities using SBDF3, p = 1.625.



Test Problem 2

Second test problem:

$$u_t = \Delta(u^5), \quad (x, y) \in [0, 1]^2, \quad t > 0,$$

with initial and boundary conditions set such that the exact solution is

$$u(x, y, t) = \left(\frac{4}{5}(2t + x + y)\right)^{1/4}.$$

Numerical convergence test

Figure: Numerical convergence of linearly stabilized schemes.

Spatial discretization: second order centred differences; N=2048.

Final time: T = 0.40.

Reference solution: explicit third order Runge-Kutta, $\Delta t = 6.25 \times 10^{-6}$.

Stabilized using $p\Delta u$, and

$$\frac{p\lambda_L}{\lambda_F} \approx \frac{p}{8(1+t)}.$$



Error constant

indicates how the method behaves.

Discretizing $u'=(1-p)\lambda u+p\lambda u$, we observe that the discretization error grows with p.

We examine this behaviour by applying the time stepping method and finding the series expansion of the amplification factor at $\Delta t = 0$. For second order methods, leading order error term is Δt^3 . The coefficient

Table: Coefficient of leading order error term as applied to the test equation.

Method	Coefficient
EIN SBDF2 CNAB mCNAB CNLF	$\frac{\frac{1}{2}(p-p^2)}{p}$

Amplification factor as $\Delta t \to \infty$

Methods performing well at large time step-sizes have small amplication factor as $\Delta t \to \infty$.

Figure: Amplification factor as $\Delta t \to \infty$.



Image Inpainting

In Schönlieb and Bertozzi (2011), the authors proposed the fourth order inpainting model

$$u_t = -\Delta \nabla \cdot \left(\frac{\nabla u}{\sqrt{|\nabla u|^2 + \epsilon^2}} \right) + \lambda (u_0 - u),$$

and numerical solution by the first order accurate method

$$\frac{u^{n+1} - u^n}{\Delta t} = -\Delta \nabla \cdot \left(\frac{\nabla u^n}{\sqrt{|\nabla u^n|^2 + \epsilon^2}}\right) + \lambda (u_0 - u^n) + p_0 \lambda u^n + p_1 \Delta^2 u^n$$
$$- p_0 \lambda u^{n+1} - p_1 \Delta^2 u^{n+1}.$$

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Image inpainting

Figure: A vandalized photograph.

Table: Iteration counts for $\mathsf{TV}\text{-}\mathsf{H}^{-1}$ image restoration.

	Δt	Iterations
SBDF1	0.30	1002
SBDF2	0.54	401
CNAB	0.64	347

Motion by mean curvature

We evolve the level set equation for motion by mean curvature,

$$u_t = \kappa |\nabla u| = |\nabla u| \nabla \cdot \left(\frac{\nabla u}{|\nabla u|}\right)$$

on an initial dumbbell-shaped curve in 3D. We solve to time T=0.75 on a $256\times128\times128$ periodic grid.

On a machine with an Intel[®] CoreTM i5-4570 CPU@3.20GHz running MATLAB 2014b:

- Forward Euler: 3000 time steps \rightarrow over 28 minutes.
- SBDF2: 75 time steps \rightarrow under 100 seconds.

Motion by mean curvature

Figure: Dumbbell-shaped curve under mean curvature flow.

Conclusion

Contributions of this thesis mainly fall into two categories:

- Further developed linearly stabilized schemes.
 - Set out a framework for developing new methods of this type.
 - Identified properties necessary for effective schemes.
- 2 Proposed new methods that outperform existing ones.
 - IMEX multistep methods and exponential time differencing methods.

Future work:

- Oevelopment higher order methods without the deficiencies exhibited by ETD methods.
- Oetailed comparison with popular algorithms for nonlinear stiff PDEs.

