

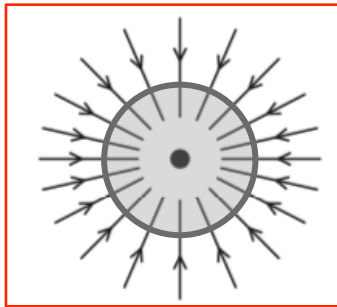
## Chiral anomaly: Second viewpoint (Text 5.4.3)

$$\dot{\mathbf{r}} = \mathbf{v}_g - \dot{\mathbf{k}} \times \boldsymbol{\Omega}, \quad (5.11a)$$

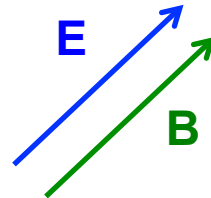
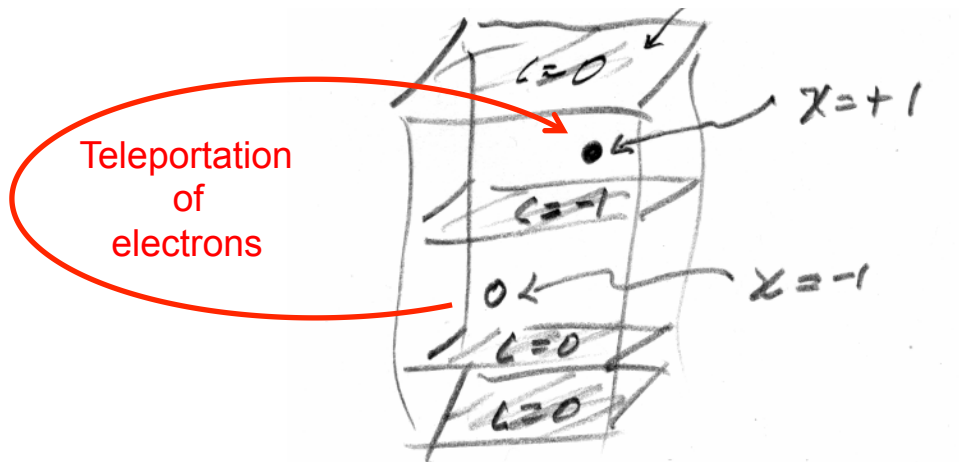
$$\dot{\mathbf{k}} = -\frac{e}{\hbar} \boldsymbol{\mathcal{E}} - \frac{e}{\hbar c} \dot{\mathbf{r}} \times \mathbf{B}, \quad (5.11b)$$

$$\left(1 + \frac{e}{\hbar c} \mathbf{B} \cdot \boldsymbol{\Omega}\right) \dot{\mathbf{r}} = \mathbf{v}_g + \frac{e}{\hbar} \boldsymbol{\mathcal{E}} \times \boldsymbol{\Omega} + \frac{e}{\hbar c} (\mathbf{v}_g \cdot \boldsymbol{\Omega}) \mathbf{B}, \quad (5.39a)$$

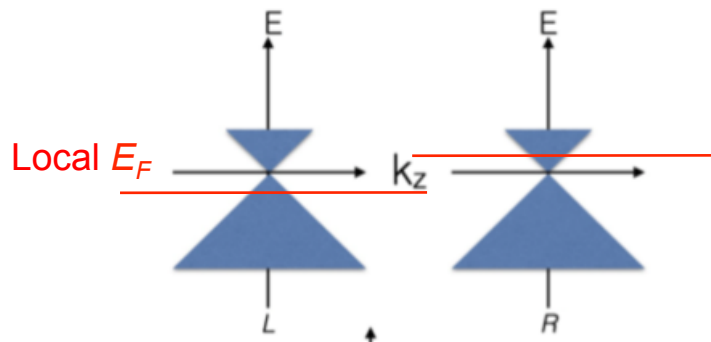
$$\left(1 + \frac{e}{\hbar c} \mathbf{B} \cdot \boldsymbol{\Omega}\right) \dot{\mathbf{k}} = -\frac{e}{\hbar} \boldsymbol{\mathcal{E}} - \frac{e}{\hbar c} \mathbf{v}_g \times \mathbf{B} - \frac{e^2}{\hbar^2 c} (\boldsymbol{\mathcal{E}} \cdot \mathbf{B}) \boldsymbol{\Omega}. \quad (5.39b)$$



$$\begin{aligned} \frac{dn}{dt} &= \frac{-1}{(2\pi)^3} \frac{e^2}{\hbar^2 c} (\boldsymbol{\mathcal{E}} \cdot \mathbf{B}) \int_{S_F} \boldsymbol{\Omega} \cdot \hat{\mathbf{v}}_F d^2k \\ &= \frac{e^2}{h^2 c} (\boldsymbol{\mathcal{E}} \cdot \mathbf{B}) \chi_i. \end{aligned}$$



One consequence:



## Chiral magnetic effect

$$\mathbf{J} = \left( \frac{e^2}{h^2 c} \sum_i \mu_i \chi_i \right) \mathbf{B}$$

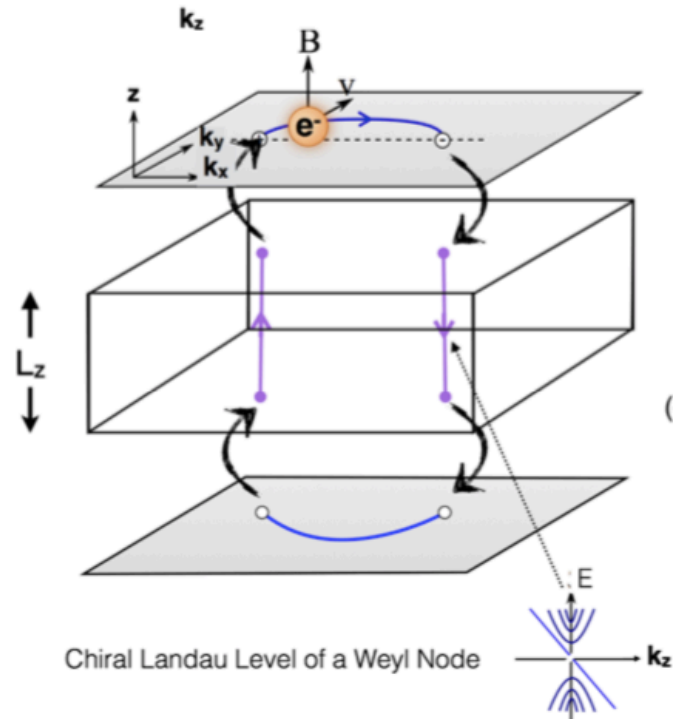
Local  $E_F$   
of  $i^{th}$  Weyl node

### More on Weyl semimetals

See reviews by

- Hasan et al. (2017)
- Yan and Felser (2017)
- Armitage et al. (2018)

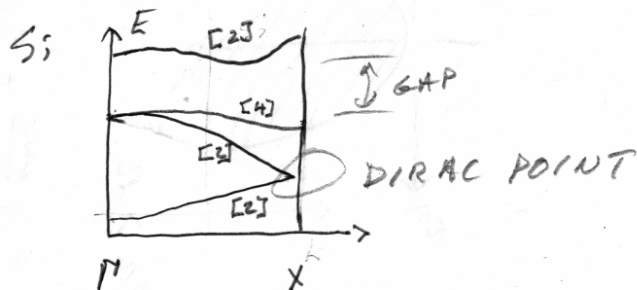
## Landau orbits



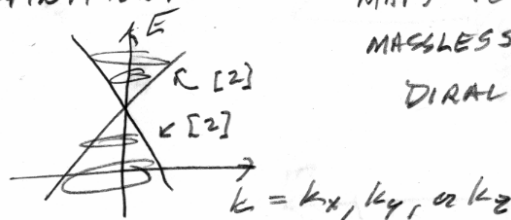
# Dirac semimetals

IF  $I$  AND  $TR$  ARE BOTH GOOD SYMMETRIES,

$$\underline{k} \xrightarrow{I} -\underline{k} \xrightarrow{TR} \underline{k} \quad \text{SO } H_{\underline{k}} = \text{HAM OF TR-INVARIANT SYSTEM FOR ANY } \underline{k}$$



DEFINITION:



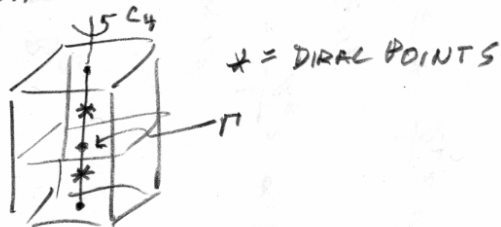
MAPS TO  
MASSLESS  
DIRAC THEORY

CAN IT OCCUR AT EF, AS TOUCHING OF VB & CB?

YES, E.G.,  $Na_3Bi$ ,  $Cd_3As_2$ , ETC.

3 fold rotational axis for  $Na_3Bi$   
4 fold rotational axis for  $Cd_3As_2$

OCCURS ON ROTATIONAL SYMMETRY AXIS:



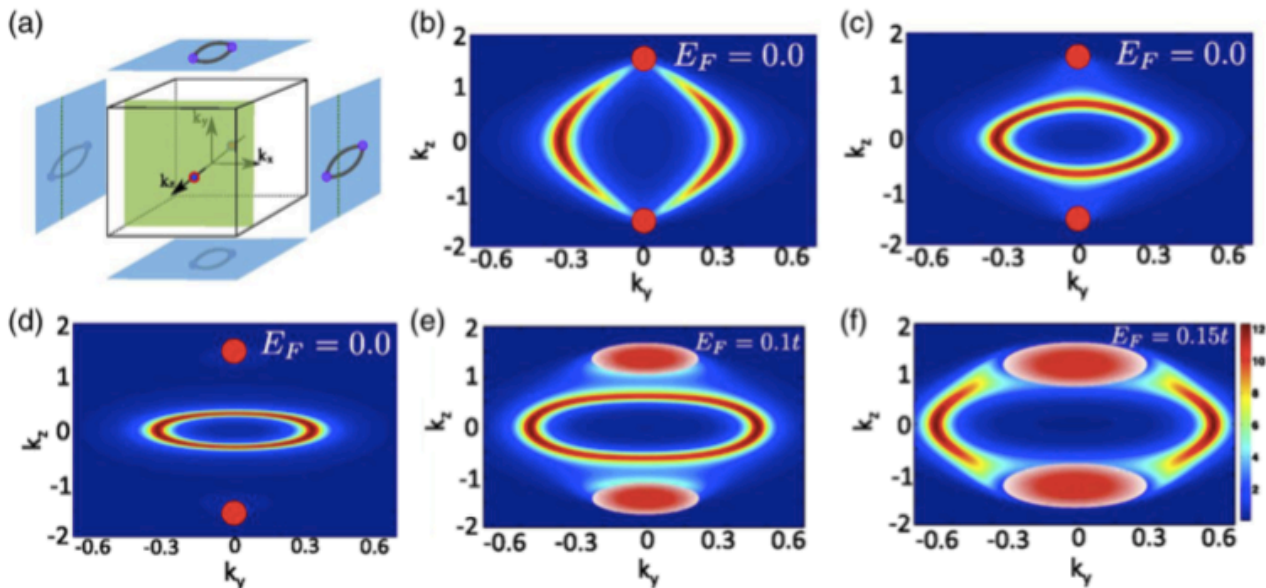


FIG. 13. Fermi arcs on the surface of DSMs. (a) A schematic of a DSM showing Dirac nodes along the  $k_z$  axis in the bulk BZ and double Fermi arcs on the surface BZs. Note that surfaces perpendicular to the  $z$  axis have no arcs. A 2D slice of the bulk BZ perpendicular to the  $k_z$  axis is shown as a shaded (green) plane, which projects to the dashed (green) line on the surface BZ. (b)–(d) A symmetry-allowed mass term at the surface admits backscattering between these branches at the contact point which dissociates the surface band from the projected Dirac point. These surface branches can be deformed but not removed from the time reversal symmetric plane at  $k_z = 0$ . If the chemical potential is not aligned with the bulk Dirac points the surface Fermi arcs disappear by merging with the bulk continuum (e), (f). Adapted from [Kargarian, Randeria, and Lu, 2016](#).

# Topological nodal line semimetals\*

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## 2.2. Nodal lines protected by inversion, time-reversal, and SU(2) spin-rotation symmetries

Here we first assume that all the three symmetries are present in our system. Since SU(2) is a symmetry, we can redefine time-reversal operator, combining it with a  $\pi$  spin rotation about the y-axis,

$$T \rightarrow T e^{i\sigma_y \pi}, \quad (7)$$

after which we have  $T^2 = +1$  instead of  $-1$  for fermions. Since both inversion,  $P$  and  $T$ , reverse the momentum  $\mathbf{k} \rightarrow -\mathbf{k}$ ,  $P*T$  is an anti-unitary symmetry that preserves the momentum. Since  $[P, T] = 0$ , we have

$$(P*T)^2 = P^2 T^2 = 1. \quad (8)$$

Equation (8) dictates that it can be represented as

$$P*T = K, \quad (9)$$

where  $K$  is the complex conjugation, in a proper orbital basis. In this basis,  $P*T$ -symmetry ensures that

$$H(\mathbf{k}) = H^*(\mathbf{k}), \quad (10)$$

or that  $H(\mathbf{k})$  is real at each  $\mathbf{k}$ .

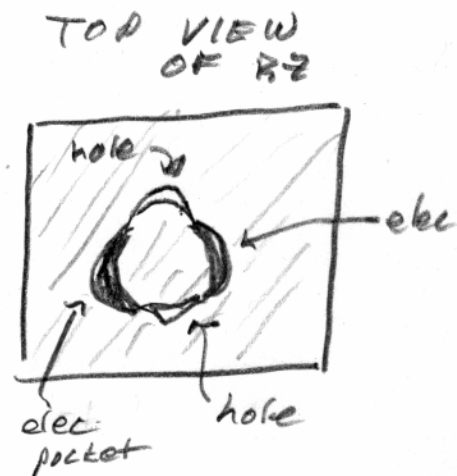
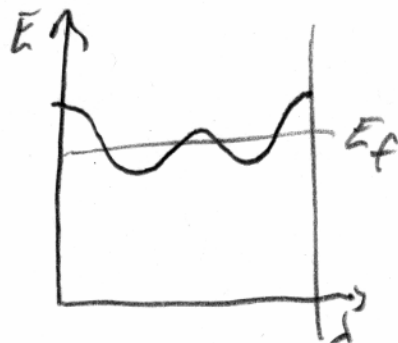
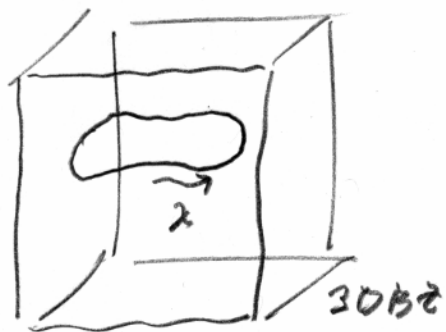
$$H_{2 \times 2} = \begin{pmatrix} a & c \\ c & b \end{pmatrix}$$

$$= \frac{a+b}{2} \mathbb{I} + \frac{a-b}{2} \sigma_3 + c \sigma_1$$

codimension of 2 is enough!

Nodal lines can occur in 3D BZ.

(also can occur due to symmetries)



# Symmorphic Intersecting Nodal Rings in Semiconducting Layers

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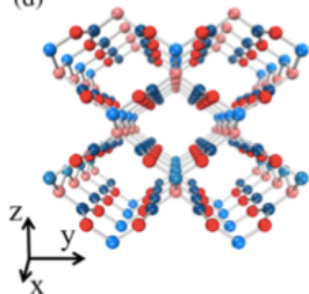
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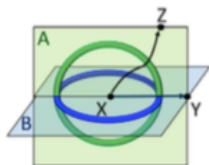
(Received 21 September 2017; revised manuscript received 23 January 2018; published 9 March 2018)

(d)



Hypothetical  
structures of BN,  
AIP, GaP, SiC, BP,  
etc.

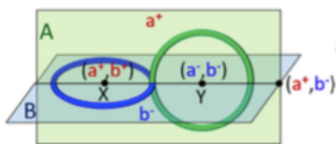
(a)



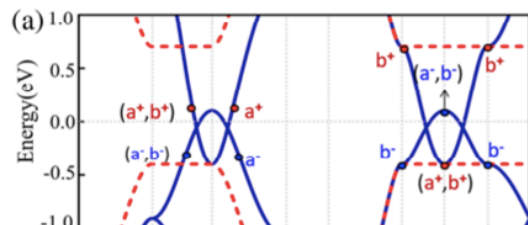
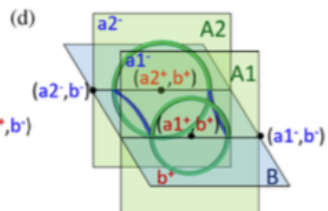
(b)



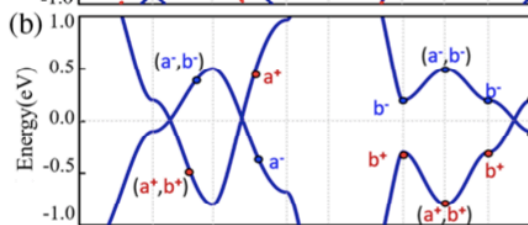
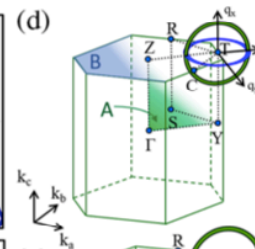
(c)



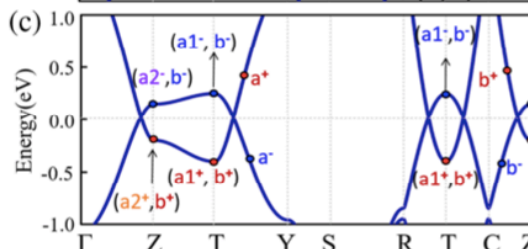
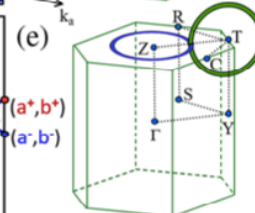
(d)



(d)



(e)



(f)

