1 The Model

A Simple Markov Model Of A Chain with A Single Loop

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In this document we summarize a simple Markov chain describing a chain having N monomers (beads) and one occasion of a loop. Only the first monomer can create a loop with all other j = 2, 3, ..., N. We consider bead N to be fixed. There are several states for the model. Assuming a loop [1, j], indicating bead 1 is connected to bead j, we can form the loops [1, j + 1], [1, j - 1], or detach, indicated by [1, 0].

Using this terminology, we write

$$\begin{split} [1,j] &\to [1,j+1] = \lambda \\ [1,j] &\to [1,j-1] = \mu \\ [1,0] &\to [1,j] = \alpha \\ [1,j] &\to [1,0] = \beta \\ [1,0] &\to [1,0] = 1 - \alpha \end{split}$$

where j = 2, 3, ..., N The α value represent the probability of attaching bead 1 to bead j given that there is no loop at time t. At first approximation we vary α according to the probability that the beads are closer than some distance, given by the function:

$$\alpha_{nm} = \Phi(|R_n - R_m| < \epsilon, n - m) = \left[\frac{3}{2\pi b^2 |n - m|}\right]^{3/2} \int_0^{\epsilon} \exp\left[-\frac{3k^2}{2|n - m|b^2}\right] dk$$

where b>0 is the connection length, R_n is the position of bead n, and $\epsilon>0$ is the maximal distance in which two beads are considered to have met. We indicate the probability to create a loop [1,j] by p_j . Thus, the

master equation is given by:

$$\begin{bmatrix} \dot{p}_0 \\ \dot{p}_2 \\ \vdots \\ \vdots \\ \dot{p}_N \end{bmatrix} = \begin{bmatrix} -N\alpha & \beta & \beta & \dots & \dots & \beta & \beta \\ \alpha & -(\beta+\lambda) & \mu & \dots & \dots & \dots & 0 \\ \alpha & \lambda & -(\beta+\lambda+\mu) & \mu & \dots & \dots & 0 \\ \alpha & 0 & \lambda & -(\beta+\lambda+\mu) & \mu & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ \alpha & 0 & 0 & \dots & \lambda & -(\beta+\mu+\lambda) & \mu \\ \alpha & 0 & 0 & \dots & 0 & \lambda & -(\beta+\mu+\lambda) \end{bmatrix} \begin{bmatrix} p_0 \\ p_2 \\ \vdots \\ p_N \end{bmatrix}$$

in short notation we write

$$[P(t)] = [M] [P(t)].$$

the solution to the equation above is given by

$$[P(t)] = [P(0)] \exp [Mt] = [U] \exp [[D]t) [U^{-1}] [P_0]$$

where [P(0)] is the initial distribution of states, U is a matrix of the eigenvectors in its columns, and D is a diagonal matrix with the corresponding eigenvalues in the diagonal entries.

2 Simulations

Setting $\beta = \mu = \lambda = x$ we get

$$[M] = \begin{bmatrix} -N\alpha & x & x & \dots & \dots & x & x \\ \alpha & -2x & x & \dots & \dots & \dots & 0 \\ \alpha & x & -3x & x & \dots & \dots & 0 \\ \alpha & 0 & x & -3x & x & \dots & 0 \\ \vdots & & & & & & & \\ \vdots & & & & & & \\ \alpha & 0 & 0 & \dots & x & -3x & x \\ \alpha & 0 & 0 & \dots & 0 & x & -2x \end{bmatrix}$$

choosing x = 1/3, d = b/2, and simulating the system until steady-state (note that α remains as defined above), we find, for a chain of N = 64 beads that the probability to find a loop [1, j] at steady-state is proportional to |j - 1| see Figure 2 and Figure 2.

The steady state probability of loops [1, j] is shown in Figure 2

2.1 Introducing 'hot spots' in the chain

We next introduce some beads for which bead 1 have more affinity to, in the sense that once a loop [1,j] is formed,with j being the affine bead to 1, the probability of either detachment (β) , decreasing (μ) or increasing (λ) the loop size are decreased in relation to other beads.

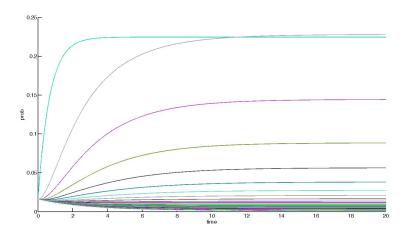


Figure 1: The time evolution of the Markov chain, with x = 1/3. The values of α remained as defined above. The thick cyan curve represents the probability of no loop. The curves are decreasing (at steady states) with bead number

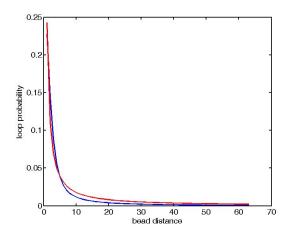


Figure 2: The steady state values of the loop probability (blue curve), with x=1/3, and α as defined above. To test the theoretical results, that the looping probability is proportional to the bead distance, we fit a model of the kind $y=ax^{-b}$ (red curve). The fitted values are a=0.2429, b=1.137. The R-square value of the fit is 0.9687