## 1. The problem setup

Let M be a real n-dimensional manifold and  $x_0 \in M$ . We consider a formal complex-valued function

$$\varphi = \nu^{-1}\varphi_{-1} + \varphi_0 + \nu\varphi_1 + \dots$$

and a formal complex-valued density

$$\rho = \rho_0 + \nu \rho_1 + \dots$$

on M such that  $x_0$  is a nondegenerate critical point of  $\varphi_{-1}$  with zero critical value,  $\varphi_{-1}(x_0) = 0$ , and  $\rho_0(x_0) \neq 0$ . We want to relate to the formal oscillatory integral

$$f \mapsto \nu^{-\frac{n}{2}} \int_{(x_0)} e^{\varphi} f \rho,$$

where f is an amplitude (say,  $f \in C^{\infty}(M)[[\nu]]$ ), a formal distribution supported at  $x_0$ ,

$$f \mapsto \Lambda(f)$$
,

where  $\Lambda = \Lambda_0 + \nu \Lambda_1 + \dots$  The assignment should be based on a number of formal properties of the formal integral (see details in my most recent preprint).

The answer is as follows. Choose local coordinates  $\{x^i\}$  around  $x_0$  such that  $x^i(x_0) = 0$  for all i. The Hessian matrix of  $\varphi_{-1}$  at  $x_0$  is denoted by  $h_{ij}$ ,

$$h_{ij} := \frac{\partial \varphi_{-1}}{\partial x^i \partial x^j}(x_0).$$

It is a complex symmetric nondegenerate matrix with constant coefficients. We denote by  $h^{ij}$  its inverse matrix and introduce the Laplace operator

$$\Delta := -\frac{1}{2}h^{ij}\frac{\partial^2}{\partial x^i \partial x^j}.$$

We use the following model formal Gaussian integral,

$$\nu^{-\frac{n}{2}} \int_{\mathbb{R}^n} e^{\frac{1}{2\nu}h_{ij}x^ix^j} dx^1 \dots dx^n = \pm \sqrt{\frac{(-2\pi)^n}{\det(h_{ij})}} e^{\nu\Delta} f \bigg|_{x=0}.$$

We do not specify the sign on the right-hand side.

Assume that locally

$$\rho = e^u \, dx^1 \dots dx^n,$$

where  $u = u_0 + \nu u_1 + \dots$  Set

$$\chi(x) := \varphi(x) - \frac{1}{2\nu} h_{ij} x^i x^j - \varphi_0(0) + u(x) - u_0(0).$$

Then

$$\nu^{-\frac{n}{2}} \int_{(x_0)} e^{\varphi} f \, \rho = \nu^{-\frac{n}{2}} e^{\varphi_0(0) + u_0(0)} \int_{(x_0)} e^{\frac{1}{2\nu} h_{ij} x^i x^j} (e^{\chi} f) \, dx^1 \dots dx^n =$$

$$\pm \sqrt{\frac{(-2\pi)^n}{\det(h_{ij})}} e^{\varphi_0(0) + u_0(0)} e^{\nu \Delta} (e^{\chi} f)|_{x=0}.$$

This expression is coordinate-independent, which follows, say, from my axiomatic description of formal oscillatory integrals. Consider the functional

$$K(f) := e^{\nu \Delta} (e^{\chi} f)|_{x=0}.$$

It is well-defined (see the Appendix to my preprint). Our task is to identify all such functionals and to recover the phase remainder  $\chi$  from K. In my preprint I proved that the full jet of  $\chi$  at  $x_0 = 0$  is determined uniquely by the functional K. Now I will describe all functionals K without referring to  $\chi$  and recover  $\chi$  from K via a constructive procedure.

## 2. A description of the functional K

Let  $N = N_0 + \nu N_1 + \dots$  be a formal differential operator on a neighborhood U. I call it natural if the order of each differential operator  $N_r$  is not greater than r. The natural operators on U form an algebra. The operators  $\nu^{-1}N$ , where N is natural, form a Lie algebra with respect to the commutator.

We consider formal differential operators acting on formal jets at zero. We introduce a descending filtration on the space of formal jets at zero by assigning filtration degree 2 to  $\nu$  and filtration degree r to a jet at zero (which does not depend on  $\nu$ ) that has zero of order r. This filtration induces a filtration on the space of formal differential operators on U. For example, in local coordinates, the degree of the operator

$$\nu x^1 \frac{\partial}{\partial x^2}$$

is 2+1-1=2. Our main object is the Lie algebra  $\mathfrak{g}$  of formal differential operators on the formal neighborhood of zero of the form  $\nu^{-1}N$ , where N is natural and  $\nu^{-1}N$  has a positive filtration degree. This is a pronilpotent Lie algebra. The multiplication operator by the phase remainder  $\chi = \nu^{-1}\chi_{-1} + \chi_0 + \ldots$  lies in  $\mathfrak{g}$ , because the multiplication operator  $\nu\chi$  is natural, the order of zero of  $\chi_{-1}$  is at least 3, and the order of zero of  $\chi_0$  is at least one. The algebra  $\mathfrak{g}$  acts on the formal

distributions supported at zero from the right:

$$\mathfrak{g} \ni A : u \mapsto \langle u | A,$$

so that

$$\langle u|A|f\rangle = \langle u,Af\rangle,$$

where f is a full jet at zero. Here we use the Dirac bra-ket notation. Denote by  $\mathfrak{b}$  the annihilator of  $\delta(x)$  in  $\mathfrak{g}$ . For  $A \in \mathfrak{b}$ ,

$$\langle \delta | A | f \rangle = (Af)(0) = 0.$$

The subalgebra Lie  $\mathfrak{b}$  does not depend on the choice of local coordinates. Now fix local coordinates around zero. Denote by  $\mathfrak{c}$  the subalgebra of  $\mathfrak{g}$  of formal differential operators with constant coefficients. Clearly,

$$\mathfrak{g}=\mathfrak{b}\oplus\mathfrak{c}.$$

Writing  $A \in \mathfrak{g}$  in the local coordinates in the normal form,  $A = A(x, \frac{\partial}{\partial x})$ , we decompose it according to (1) as

$$A = \left( A\left(x, \frac{\partial}{\partial x}\right) - A\left(0, \frac{\partial}{\partial x}\right) \right) + A\left(0, \frac{\partial}{\partial x}\right).$$

The functional  $f \mapsto \langle \delta | A | f \rangle$  depends only on the  $\mathfrak{c}$ -component of A. Moreover, the  $\mathfrak{c}$ -component of A is completely determined by this functional.

An important remark: An operator  $C \in \mathfrak{c}$  has constant coefficients and positive filtration degree. Also,  $\nu C$  is natural,

$$\nu C = N_0 + \nu N_1 + \nu^2 N_2 + \dots,$$

where  $N_r$  is a differential operator with constant coefficients of order not greater than r. The filtration degree of  $N_r$  is at least -r. Now,

$$C = \nu^{-1} N_0 + N_1 + \nu N_2 + \dots$$

The filtration degree of  $\nu^{r-1}N_r$  is at least 2(r-1)-r=r-2. Since the filtration degree of C is positive, we see that  $N_0=0, N_1=0$ , and  $N_2$  is of order not greater than one.

Now consider the pronilpotent Lie group  $\exp \mathfrak{g}$ . We will show that it can be realized as a group of formal differential operators

$$A = \sum_{r = -\infty}^{\infty} \nu^r A_r$$

of filtration degree zero on the space  $\mathcal{F}$  of formal jets

$$f = \sum_{r = -\infty}^{\infty} \nu^r f_r$$

at zero of finite filtration degree, i.e., for every  $f \in \mathcal{F}$  there exists an integer k such that the order of zero of  $f_r$  is at least k-2r for every  $r \in \mathbb{Z}$ .

First we want to show that any natural operator leaves  $\mathcal{F}$  invariant. Let  $N = N_0 + \nu N_1 + \dots$  be a natural operator and  $f \in \mathcal{F}$ . Assume for the sake of simplicity that the filtration degree of f is zero. Then the filtration degree of  $f_r$  is at least -2r. The coefficient at  $\nu^k$  of Nf is

(2) 
$$\sum_{r=0}^{\infty} N_r f_{k-r}$$

Since the order of  $N_r$  is at most r, its filtration degree is at least -r. The filtration degree of  $f_{k-r}$  is at least -2(k-r) = -2k+2r. Therefore, the filtration degree of  $N_r f_{k-r}$  is at least -r - 2k + 2r = r - 2k. It follows that the series (2) is convergent in the topology induced by the filtration and defines a formal jet at zero. The filtration degree of this formal jet is at least zero, because the filtration degree of a natural operator is at least zero.

If N is a natural operator, then  $\nu^{-1}N$  also leaves  $\mathcal{F}$  invariant.

Given  $f \in \mathcal{F}$  and  $A \in \mathfrak{g}$ ,  $e^A f = \sum_{r=0}^{\infty} \frac{1}{r!} A^r f$  is a convergent series in  $\mathcal{F}$ , because A leaves  $\mathcal{F}$  invariant and the filtration degree of A is positive. The group  $\exp \mathfrak{g}$  respects the filtration on  $\mathcal{F}$ , because its elements have zero filtration degree.

Given  $A \in \mathfrak{g}$ , one can write uniquely

$$e^A = e^B e^C,$$

where  $B \in \mathfrak{b}$  and  $C \in \mathfrak{c}$ . This is a constructive procedure. First decompose  $A = A_1 = B_1 + C_1$ . Then iterate

$$e^{A_{r+1}} = e^{-B_r} e^{A_r} e^{-C_r},$$

where  $B_r \in \mathfrak{b}$ ,  $C_r \in \mathfrak{c}$ , and  $A_r = B_r + C_r$ . The filtration degree of  $A_r$  goes to infinity as  $r \to \infty$ , which means that  $e^{A_r} \to 1$ . We get that

$$e^A = e^{B_1}e^{A_2}e^{C_1} = e^{B_1}e^{B_2}e^{A_3}e^{C_2}e^{C_1} = \dots$$

Therefore,

$$e^A = e^B e^C,$$

where

$$e^{B} = e^{B_{1}}e^{B_{2}}e^{B_{3}}\dots \in \exp \mathfrak{b} \text{ and } e^{C} = \dots e^{C_{3}}e^{C_{2}}e^{C_{1}} \in \mathfrak{c}$$

are convergent products in the topology induced by the filtration.

The formal differential operator  $\exp(\nu\Delta)$  has filtration degree zero and does not lie in  $\exp \mathfrak{g}$ , but it acts upon it by conjugations. Since

 $\chi \in \mathfrak{g}$ , we have

$$A := e^{\nu \Delta} \chi e^{-\nu \Delta} \in \mathfrak{g}.$$

Now, decompose

(3) 
$$e^{\nu\Delta}e^{\chi}e^{-\nu\Delta} = e^A = e^Be^C,$$

where  $B \in \mathfrak{b}$  and  $C \in \mathfrak{c}$ . Denote by  $\mathcal{F}_0$  the subspace of  $\mathcal{F}$  of elements of filtration degree at least zero.

Example:

$$\nu^{-1}x^1x^2 \in \mathcal{F}_0$$
.

Observe that if  $f \in \mathcal{F}_0$ , then  $f(0) \in \mathbb{C}[[\nu]]$ , i.e.,  $\delta : \mathcal{F}_0 \to \mathbb{C}[[\nu]]$ . Denote by  $\mathcal{F}_+$  the space of formal jets at zero of the form

$$f = \sum_{r=0}^{\infty} \nu^r f_r.$$

Then  $\mathcal{F}_+ \subset \mathcal{F}_0$  (a proper subspace). The functional K can be written as

$$K(f) = \langle \delta | (e^{\nu \Delta} e^{\chi} e^{-\nu \Delta}) | e^{\nu \Delta} f \rangle.$$

We want to show that K is defined on  $\mathcal{F}_+$ . The operator  $e^{\nu\Delta}$  leaves  $\mathcal{F}_+$  invariant. We have shown that  $\mathcal{F}_+ \subset \mathcal{F}_0$  and that the element  $e^{\nu\Delta}e^{\chi}e^{-\nu\Delta} \in \exp \mathfrak{g}$  leaves  $\mathcal{F}_0$  invariant. Also,  $\delta: \mathcal{F}_0 \to \mathbb{C}[[\nu]]$ . It follows that  $K: \mathcal{F}_+ \to \mathbb{C}[[\nu]]$ , which means that

$$K: C^{\infty}(M)[[\nu]] \to \mathbb{C}[[\nu]].$$

Using (3), we can write

$$K(f) = \langle \delta | e^{\nu \Delta} e^{\chi} | f \rangle = \langle \delta | e^B e^C e^{\nu \Delta} | f \rangle = \langle \delta | e^{\nu \Delta + C} | f \rangle.$$

Thus, the functional K is given by the formal differential operator with constant coefficients  $\exp(\nu\Delta + C)$ .

According to the remark above, the operator  $\nu\Delta + C$  is of the form

$$\nu A_2 + \nu^2 A_3 + \dots,$$

where the order of  $A_r$  is not greater than r and the principal symbol of  $A_2$  is  $-\frac{1}{2}h^{ij}\xi_i\xi_j$ , where  $h^{ij}$  is nondegenerate. We claim that **any** operator of this format is equal to the operator K for an appropriate phase remainder  $\chi$ .

## 3. The rest of the proof

Now we need to define the action of differential operators on functions **from the right**. This is achieved by passing to the transpose,

$$(x^i)^t = x^i \text{ and } \left(\frac{\partial}{\partial x^i}\right)^t = -\frac{\partial}{\partial x^i},$$

so that

$$\langle f|A:=A^t|f\rangle.$$

For example,

$$\left\langle 1 \middle| x^1 \frac{\partial}{\partial x^1} = \left\langle x^1 \middle| \frac{\partial}{\partial x^1} = -1. \right.$$

This action depends on the choice of coordinates.

Denote by  $\mathfrak{r}$  the subalgebra of  $\mathfrak{g}$  of the operators which annihilate constants from the right and by  $\mathfrak{f}$  the subalgebra of  $\mathfrak{g}$  of the multiplication operators. Then

$$\mathfrak{g}=\mathfrak{r}\oplus\mathfrak{f}.$$

Given  $A \in \mathfrak{g}$ , we expand it according to (4) as follows,

$$A = (A - \langle 1|A) + \langle 1|A,$$

where we interpret the function  $\langle 1|A$  as a multiplication operator.

We will need the following calculation:

$$e^{-\nu\Delta}x^k e^{\nu\Delta} = x^k + \nu h^{kl} \frac{\partial}{\partial x^l}.$$

Given a nondegenerate symmetric complex matrix  $h^{ij}$  with constant coefficients, introduce the function

$$\psi = \frac{1}{2}h_{ij}x^ix^j.$$

Now,

$$e^{\nu^{-1}\psi}e^{-\nu\Delta}x^ke^{\nu\Delta}e^{-\nu^{-1}\psi} = \nu h^{kl}\frac{\partial}{\partial x^l}.$$

Therefore, if  $B \in \mathfrak{b}$ , i.e.,  $B = x^i A_i$ , for some  $A_i \in \mathfrak{g}$ , then

(5) 
$$R := e^{\nu^{-1}\psi} e^{-\nu\Delta} B e^{\nu\Delta} e^{-\nu^{-1}\psi} \in \mathfrak{r}.$$

Assume that, as in (3),

$$e^{\nu\Delta}e^{\chi}e^{-\nu\Delta} = e^Be^C.$$

for some  $B \in \mathfrak{b}$  and  $C \in \mathfrak{c}$ . Then

$$e^{\chi} = e^{-\nu\Delta} e^B e^{\nu\Delta} e^C,$$

from whence we get that

$$e^{\nu^{-1}\psi}e^{\chi}e^{-\nu^{-1}\psi} = \left(e^{\nu^{-1}\psi}e^{-\nu\Delta}e^Be^{\nu\Delta}e^{-\nu^{-1}\psi}\right)e^{\nu^{-1}\psi}e^Ce^{-\nu^{-1}\psi}.$$

Therefore,

(6) 
$$e^{\chi} = e^R e^{\nu^{-1} \psi} e^C e^{-\nu^{-1} \psi},$$

where R is as in (5). Applying (6) to 1 from the right, we get

$$e^{\chi} = \langle 1|e^{\nu^{-1}\psi}e^{C}e^{-\nu^{-1}\psi}.$$

Thus, we have recovered  $\chi$  from an arbitrary  $C \in \mathfrak{c}$ .