

# Duistermat-Heckman localisation formula and locus of vector fields.

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§0

About Two years ago (summer 2012) Sasha Belavin explained how to calculate an integral

$$Z(t) = \int e^{t d_K \omega} \quad (0.1)$$

( $\omega$ -1 form,  $d_K = d + L_K$ ). He ~~exp~~ showed first that

this integral does not depend on  $t$ , then showed that it is localised at zeros of vector field  $K$ :

$$Z(t) \sim \frac{1}{\sqrt{\det \frac{\partial K}{\partial x}}} \Big|_{K=0}. \quad (0.2)$$

It is typical localisation formula.

I tried to revive these calculations. ~~Here~~ On one hand they are leading to Duistermat-Heckman formula in more less general case.

On the other hand ~~we may discuss~~ it is interesting to analyze geometrical meaning of answer.

§1.

localisation

Two words about Duistermat-Heckman formula (DHL)-formula.

Let  $M$  be compact manifold

$(M^{2n}, \Omega)$  be compact symplectic manifold

Let  $H$  be an Hamiltonian, such that the vector field

$$K = D_H: \Omega \rightarrow D_H = -dH$$

is a compact vector field  
i.e. it generates compact subgroup  $E^{1k}$   
in the group of diffeomorphisms

Then

$$\int \Omega^n e^{iH}$$

is localized at zero locus  
of vector field  $K$   
of  $e^{iH(x_i)}$

$$\sum_{x_i: K(x_i)=0} \int \Omega^n e^{iH} = \int \Omega^n e^{iH}$$

$$\text{Vol}(Hess H(x_i))$$

(we suppose that  $K(x_i)$  are not-degenerate)

This is famous Duistermaat-Hackman formula.

We will consider here a special but very illuminating case of this formula.

[See in more detail the next file.]

We consider now the following set up:

Let  $\omega$  be 1-form on  $M$  ( $\dim M = 2n$ )

such that  $\Omega = d\omega$  defines symplectic structure.  
(of course condition  $\Omega = d\omega$  is in contradiction with compactness of  $M$ :  $\int \Omega^n \neq 0$ , but we ignore now this.  
E.g. we suppose that  $M$  is not compact)

Let  $K$  be a vector field such that

$$\mathcal{L}_K \omega = d\omega \lrcorner K + d(\omega \lrcorner K) = 0$$

Then it is evident that  $K$  is

Hamiltonian vector field of  $H = \omega \lrcorner K$

$$\Omega \lrcorner K = d\omega \lrcorner K = -d(\omega \lrcorner K) = -dH.$$

$$\begin{array}{ccc} & \omega & \\ \swarrow \Omega = d\omega & & \searrow H = \omega \lrcorner K \\ \Omega & & H_1 = \Omega + H \end{array} \quad \begin{array}{l} d_K \omega = \\ = (d + \mathcal{L}_K)\omega = \end{array}$$

We see that

$$\begin{aligned} \int \Omega^n e^{iH} &= \int e^{iH + i\Omega} = \\ &= \int e^{i d_K \omega} \end{aligned}$$

We come to integral (0)

Calculation of

$$\int e^{ikx} dx$$

Consider

$$Z(t) = \int_{-M}^M e^{ikx} dx$$

Show that  $Z(t)$  does not depend on  $t$ .

$$\frac{dZ(t)}{dt} = i \int_{-M}^M dx \, x \, e^{ikx} = 0$$

$$\left( \int_{-M}^M dx \, x \, e^{ikx} \right) = 0 \quad (2.1)$$

(under some technical conditions)

$\int_{-M}^M dx \, x \, e^{ikx} = 0$  since form  $kx$  has rank  $\leq 2n-1$

We see that  $Z(t)$  does not depend on  $t$ .

Hence we can calculate  $Z(t)$  at  $t \rightarrow \infty$ .

$$\int e^{ikx} dx = \int e^{ik(x+H)} dx$$

$$dx = d, \quad x+H = H, \quad (dx=0)$$

$$= \int_{-M}^M \frac{1}{\sqrt{2\pi}} e^{ikH} dx = \frac{1}{\sqrt{2\pi}} \int_{-M}^M e^{ikH} dx \quad (2.2)$$

$$(dm = 2m)$$

Calculate using stationary phase method:

$$dH = d(u \cdot k) = -du \cdot k \quad (2.3)$$

locus of  $dH = \text{locus of } k$



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We see that at stationary point  $dH=0$   
Hessian is:

$$\begin{aligned} \frac{\partial^2}{\partial x^i \partial x^r} H \Big|_{K=0} &= \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^r} (W_r K^r) \Big|_{K=0} = \frac{\partial}{\partial x^i} (\Omega_{rp} K^p) = \\ &= \Omega_{ip} \frac{\partial K^p}{\partial x^r} \quad (d\omega(K) = -d\omega(K)) \\ &\quad (H(x_i) = 0 \text{ for } K(x_d) = 0) \end{aligned}$$

Hence .

$$\begin{aligned} \int \Omega^n e^{i t H} &\sim \frac{\frac{1}{\det \Omega} e^{i t H(x_0)}}{\sqrt{\det(\Omega \cdot \frac{\partial K}{\partial x})}} \sim \\ &\sim \int_{x_0} \frac{e^{i t H(x_0)}}{\sqrt{\det \frac{\partial K}{\partial x}}} \end{aligned}$$

Note:  $\frac{\partial K}{\partial x}$  is linear operator at points where  $K(x) \neq 0$

$$L_K = \frac{\partial K}{\partial x}: \quad L_K u = -[K, u].$$

We see that answer does not depend on  
choise of W.

$$\int \sim \frac{1}{\sqrt{\det \frac{\partial K}{\partial x}}}$$

Our formula is a special case of DTL formula. (In particular  $H(x)=0$ ). On the other hand this formula emphasises the role of vector field  $K$ : It states that

$$\int_C \langle \dot{\gamma}, L(\gamma) \rangle = \frac{\int_C \sqrt{\det g_K}}{C}$$

depends only on  $K$  at  $\gamma$  in the case if  $W$  is an "arbitrary"  $K$ -invariant 1-form (of course  $dx$  is not - dependent).

It is useful to study DTL formula

in the supermetric manifestation. I.A. Neressian. "Anchors and localisation of paths, integrals."

(see I.A. Schwarz, O. Zakharenko "Supersymmetry and localization"

JETP Lett. 58:1 (1993) — CMP (1995 or 1996)

(see for details next slide)

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