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1 Surfaces in \mathbf{E}^3

Now equipped by the knowledge of Riemannian geometry we consider surfaces in \mathbf{E}^3 . We reconsider again conceptions of Shape (Weingarten) opera-

tor, Gaussian and mean curvatures, focusing attention on the the fact what properties are internal and what properties are external.

1.1 Induced metric on surfaces.

Recall here again induced metric (see for detail subsection 1.4)

If surface $M: \mathbf{r} = \mathbf{r}(u, v)$ is embedded in \mathbf{E}^3 then induced Riemannian metric G_M is defined by the formulae

$$\langle \mathbf{X}, \mathbf{Y} \rangle = G_M(\mathbf{X}, \mathbf{Y}) = G(\mathbf{X}, \mathbf{Y}), \quad (1.1)$$

where G is Euclidean metric in \mathbf{E}^3 :

$$\begin{aligned} G_M = dx^2 + dy^2 + dz^2|_{\mathbf{r}=\mathbf{r}(u,v)} &= \sum_{i=1}^3 (dx^i)^2|_{\mathbf{r}=\mathbf{r}(u,v)} = \sum_{i=1}^3 \left(\frac{\partial x^i}{\partial u^\alpha} du^\alpha \right)^2 \\ &= \frac{\partial x^i}{\partial u^\alpha} \frac{\partial x^i}{\partial u^\beta} du^\alpha du^\beta, \end{aligned}$$

i.e.

$$G_M = g_{\alpha\beta} du^\alpha, \text{ where } g_{\alpha\beta} = \frac{\partial x^i}{\partial u^\alpha} \frac{\partial x^i}{\partial u^\beta} du^\alpha du^\beta.$$

We use notations x, y, z or x^i ($i = 1, 2, 3$) for Cartesian coordinates in \mathbf{E}^3 , u, v or u^α ($\alpha = 1, 2$) for coordinates on the surface. We usually omit summation symbol over dummy indices. For coordinate tangent vectors

$$\underbrace{\frac{\partial}{\partial u_\alpha}}_{\text{Internal observer}} = \underbrace{\mathbf{r}_\alpha = \frac{\partial x^i}{\partial u^\alpha} \frac{\partial}{\partial x^i}}_{\text{External observer}}$$

We have already plenty examples in the subsection 1.4. In particular for scalar product

$$g_{\alpha\beta} = \left\langle \frac{\partial}{\partial u_\alpha}, \frac{\partial}{\partial u_\beta} \right\rangle = x^i \alpha x^i \beta \cdot \langle \mathbf{r}_\alpha, \mathbf{r}_\beta \rangle. \quad (1.2)$$

1.2 Induced connection

We already studied before induced connection on manifolds embedded in \mathbf{E}^n (see the subsection 2.2).

We do it again in a more detail for surfaces in \mathbf{E}^3 .

Let $\nabla^{\text{can.flat}}$ be a canonical flat connection in Euclidean space \mathbf{E}^3 , i.e. a connection such that its Christoffel symbols vanish in cartesian coordinates, i.e.

$$\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y} = \partial_{\mathbf{X}} \mathbf{Y} = \mathbf{X}^m(x) \frac{\partial}{\partial x^m} Y^i(x) \frac{\partial}{\partial x^i} \quad (1.3)$$

in Cartesian coordinates. (Of course in arbitrary curvilinear coordinates Christoffel symbols of canonical flat connection may have non-zero components, but in Cartesian coordinates Christoffel symbols vanish. It is why we sometimes use notation $\partial_{\mathbf{X}}$ instead $\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y}$.)

Induced connection ∇ on the surface can be defined by the relation

$$\nabla_{\mathbf{X}} \mathbf{Y} = (\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y})_{\text{tangent}} \quad (1.4)$$

(See also the subsection (2.2)) One can show that this formula indeed defines the connection on M .

Calculate Christoffel symbols of this connection using the formulae above and the expressions for basic vectors $\partial_{\alpha} = \frac{\partial}{\partial u^{\alpha}}$, $\partial_{\beta} = \frac{\partial}{\partial u^{\beta}}$ for external observer.

We have

$$\nabla_{\alpha} \partial_{\beta} = \Gamma_{\alpha\beta}^{\gamma} \partial_{\gamma} = (\nabla_{\partial_{\alpha}}^{\text{can.flat}} \partial_{\beta})_{\text{tangent}} = (\nabla_{\partial_{\alpha}}^{\text{can.flat}} \mathbf{r}_{\beta})_{\text{tangent}} \quad (1.5)$$

i.e.

$$\Gamma_{\alpha\beta}^{\gamma} \mathbf{r}_{\gamma} = \left(\frac{\partial^2 \mathbf{r}}{\partial u^{\alpha} \partial u^{\beta}} \right)_{\text{tangent}}, \quad (1.6)$$

Before continuing calculations note that we see already that the induced connection is *symmetric* one.

Now continue to calculate. Any vector attached at the surface can be decomposed as a sum of tangent vector and the vector orthogonal to the surface:

$$\frac{\partial^2 \mathbf{r}}{\partial u^{\alpha} \partial u^{\beta}} = \left(\frac{\partial^2 \mathbf{r}}{\partial u^{\alpha} \partial u^{\beta}} \right)_{\text{tangent}} + \mathbf{N}$$

Hence

$$\Gamma_{\alpha\beta}^{\gamma} \mathbf{r}_{\gamma} = \left(\frac{\partial^2 \mathbf{r}}{\partial u^{\alpha} \partial u^{\beta}} \right)_{\text{tangent}} = \frac{\partial^2 \mathbf{r}}{\partial u^{\alpha} \partial u^{\beta}} - \mathbf{N} = \mathbf{r}_{\alpha\beta} - \mathbf{N} \quad (1.7)$$

Take scalar product of both parts of this equation on the tangent vector vector \mathbf{r}_π , we come to:

$$\Gamma_{\alpha\beta}^\gamma \langle \mathbf{r}_\gamma, \mathbf{r}_\pi \rangle = \langle \mathbf{r}_{\alpha\beta}, \mathbf{r}_\pi \rangle + \langle \mathbf{N}, \mathbf{r}_\pi \rangle.$$

We have that $\mathbf{N} \perp \mathbf{r}_\pi$, hence $\langle \mathbf{N}, \mathbf{r}_\pi \rangle = 0$. Using the fact that $\langle \mathbf{r}_\gamma, \mathbf{r}_\pi \rangle = g_{\gamma\pi}$ are entries of the matrix $||G_M||$ of the induced Riemannian metric (see the equation (1.2)) we come to

$$\Gamma_{\alpha\beta}^\gamma g_{\gamma\pi} = \langle \mathbf{r}_{\alpha\beta}, \mathbf{r}_\pi \rangle = x_{\alpha\beta}^i x_\pi^i.$$

Multiplying both sides of this equation on inverse matrix we come to :

$$\Gamma_{\alpha\beta}^\gamma g_{\gamma\pi} g^{\pi\rho} = \Gamma_{\alpha\beta}^\gamma \delta_\gamma^\rho = \Gamma_{\alpha\beta}^\rho = x_{\alpha\beta}^i x_\pi^i g^{\pi\rho}.$$

We come to the explicit formula for the Christoffel symbols of the induced connection in terms of parametric equations $x^i(u^\alpha)$ and induced metric $g_{\alpha\beta}$:

$$\Gamma_{\alpha\beta}^\gamma = x_{\alpha\beta}^i x_\pi^i g^{\pi\gamma}. \quad (1.8)$$

1.2.1 Induced connection—the Levi Civita connection of induced metric.

Now you may ask a question: is there another way to obtain induced connection (may be different one) on the manifold $x^i = x^i(u^\alpha)$. Recall that one can consider Levi-Civita connection of Riemannian metric $G_M = g_{\alpha\beta} du^\alpha du^\beta$ with Christoffel symbols:

$$\Gamma_{\alpha\beta}^\gamma = \frac{1}{2} g^{\gamma\pi} \left(\frac{\partial g_{\pi\alpha}}{\partial u^\beta} + \frac{\partial g_{\pi\beta}}{\partial u^\alpha} - \frac{\partial g_{\alpha\beta}}{\partial u^\pi} \right) \quad (1.9)$$

(see the subsection about Levi-Civita connection)

Natural question arises: may be these connections are the same?—Yes, it is

Proposition *The connection ∇ induced through the flat canonical connection on the surface M coincides with Levi-Civita connection. In other words Christoffel symbols (1.2) coincide with Christoffel symbols (1.9).. The proof of this Proposition follows from the following claim*

Claim *The induced connection ∇ is symmetric connection which is compatible with metric.*

Indeed we proved that Levi-Civita connection (1.9) is a unique symmetric connection which is compatible with metric. Hence every connection which is compatible with metric equals to Levi-Civita connection (1.9).

It remains to prove the Claim. When calculating Christoffel symbols of induced connection we already noted that its Christoffel symbols are symmetric. It remains to prove that induced connection is compatible with induced metric, i.e. covariant derivative preserves scalar product of tangent vectors. To prove this note that covariant derivative with respect to flat canonical metric preserves metric in Euclidean space:

$$\partial_{\mathbf{X}}\langle \mathbf{Y}, \mathbf{Z} \rangle = \langle \nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y}, \mathbf{Z} \rangle + \langle \mathbf{Y}, \nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Z} \rangle$$

On the other hand for an arbitrary tangent vector \mathbf{R}

$$\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{R} = (\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{R})_{\text{tangent}} + \text{vector orthogonal to the surface}$$

and $(\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{R})_{\text{tangent}} = \nabla_{\mathbf{X}} \mathbf{R}$, where ∇ is the induced connection. Hence

$$\partial_{\mathbf{X}}\langle \mathbf{Y}, \mathbf{Z} \rangle = \langle \nabla_{\mathbf{X}} \mathbf{Y} + \dots, \mathbf{Z} \rangle + \langle \mathbf{Y}, \nabla_{\mathbf{X}} \mathbf{Z} + \dots \rangle = \langle \nabla_{\mathbf{X}} \mathbf{Y}, \mathbf{Z} \rangle + \langle \mathbf{Y}, \nabla_{\mathbf{X}} \mathbf{Z} \rangle,$$

where we denote by \dots vectors orthogonal to the surface. Thus ∇ preserves induced metric. The claim is proved.

Note that we come to two different definitions for induced connections and prove that they are the same. The second connection is totally described in terms of metric on the surface, i.e. in terms of *Internal Observer*. The first connection is explicitly dependent on the embedding functions. It is described in terms of *External Observer*.

We continue to consider different geometrical structures from the point of view of External and Internal Observers.

We will try to prove the remarkable theorem Gauß's *Theorema Egregium* that the Gaussian curvature of the surface (which you already learned in the course of geometry) in fact depends only on induced metric...

1.3 Weingarten (shape) operator and second quadratic form

1.3.1 Recalling Weingarten operator

Continue to play with formulae ¹.

¹In some sense differential geometry it is when we write down the formulae expressing the geometrical facts, differentiate these formulae then reveal the geometrical meaning of

Recall the Weingarten (shape) operator which acts on tangent vectors:

$$S\mathbf{X} = -\partial_{\mathbf{X}}\mathbf{n}, \quad (1.10)$$

where we denote by \mathbf{n} -unit normal vector field at the points of the surface M : $\langle \mathbf{n}, \mathbf{r}_\alpha \rangle = 0$, $\langle \mathbf{n}, \mathbf{n} \rangle = 1$.

Now we realise that the derivative $\partial_{\mathbf{X}}\mathbf{R}$ of vector field with respect to another vector field is not a well-defined object: we need a connection. The formula $\partial_{\mathbf{X}}\mathbf{R}$ in Cartesian coordinates, is nothing but the derivative with respect to flat canonical connection: If we work only in Cartesian coordinates we do not need to distinguish between $\partial_{\mathbf{X}}\mathbf{R}$ and $\nabla_{\mathbf{X}}^{\text{can.flat}}\mathbf{R}$. Sometimes with some abuse of notations we will use $\partial_{\mathbf{X}}\mathbf{R}$ instead $\nabla_{\mathbf{X}}^{\text{can.flat}}\mathbf{R}$, but never forget: this can be done only in Cartesian coordinates where Christoffel symbols of flat canonical connection vanish:

$$\partial_{\mathbf{X}}\mathbf{R} = \nabla_{\mathbf{X}}^{\text{can.flat}}\mathbf{R} \quad \text{in Cartesian coordinates.}$$

So the rigorous definition of Weingarten operator is

$$S\mathbf{X} = -\nabla_{\mathbf{X}}^{\text{can.flat}}\mathbf{n}, \quad (1.11)$$

but we often use the former one (equation (1.10)) just remembering that this can be done only in Cartesian coordinates.

Recall that the fact that Weingarten operator S maps tangent vectors to tangent vectors follows from the property: $\langle \mathbf{n}, \mathbf{X} \rangle = 0 \Rightarrow \mathbf{X}$ is tangent to the surface.

Indeed:

$$0 = \partial_{\mathbf{X}}\langle \mathbf{n}, \mathbf{n} \rangle = 2\langle \partial_{\mathbf{X}}\mathbf{n}, \mathbf{n} \rangle = -2\langle S\mathbf{X}, \mathbf{n} \rangle = 0 \Rightarrow S\mathbf{X} \text{ is tangent to the surface}$$

Recall also that normal unit vector is defined up to a sign, $\mathbf{n} \rightarrow -\mathbf{n}$. On the other hand if \mathbf{n} is chosen then \mathbf{S} is defined uniquely.

1.3.2 Second quadratic form

We define now the new object: *second quadratic form*

Definition. For two tangent vectors \mathbf{X}, \mathbf{Y} $A(\mathbf{X}, \mathbf{Y})$ is defined such that

$$(\nabla_{\mathbf{X}}^{\text{can.flat}}\mathbf{Y})_{\perp} = A(\mathbf{X}, \mathbf{Y})\mathbf{n} \quad (1.12)$$

the new obtained formulae e.t.c.

i.e. we take orthogonal component of the derivative of \mathbf{Y} with respect to \mathbf{X} .

This definition seems to be very vague: to evaluate covariant derivative we have to consider not a vector \mathbf{Y} at a given point but the vector field. In fact one can see that $A(\mathbf{X}, \mathbf{Y})$ does depend only on the value of \mathbf{Y} at the given point.

Indeed it follows from the definition of second quadratic form and from the properties of Weingarten operator that

$$\begin{aligned} A(\mathbf{X}, \mathbf{Y}) &= \langle (\nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y})_{\perp}, \mathbf{n} \rangle = \langle \nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{Y}, \mathbf{n} \rangle = \\ &= \partial_{\mathbf{X}} \langle \mathbf{Y}, \mathbf{n} \rangle - \langle \mathbf{Y}, \nabla_{\mathbf{X}}^{\text{can.flat}} \mathbf{n} \rangle = \langle S(\mathbf{X}), \mathbf{Y} \rangle \end{aligned} \quad (1.13)$$

We proved that second quadratic form depends only on value of vector field \mathbf{Y} at the given point and we established the relation between second quadratic form and Weingarten operator.

Proposition *The second quadratic form $A(\mathbf{X}, \mathbf{Y})$ is symmetric bilinear form on tangent vectors \mathbf{X}, \mathbf{Y} in a given point.*

$$A: T_{\mathbf{p}}M \otimes T_{\mathbf{p}}M \rightarrow \mathbf{R}, \quad A(\mathbf{X}, \mathbf{Y}) = A(\mathbf{Y}, \mathbf{X}) = \langle S\mathbf{X}, \mathbf{Y} \rangle. \quad (1.14)$$

In components

$$A = A_{\alpha\beta} du^{\alpha} du^{\beta} = \langle \mathbf{r}_{\alpha\beta}, \mathbf{n} \rangle = \frac{\partial^2 x^i}{\partial u^{\alpha} \partial u^{\beta}} n^i. \quad (1.15)$$

and

$$S_{\beta}^{\alpha} = g^{\alpha\pi} A_{\pi\beta} = g^{\alpha\pi} x_{\pi\beta}^i n^i, \quad (1.16)$$

i.e.

$$A = GS, S = G^{-1}A.$$

Remark The normal unit vector field is defined up to a sign.

1.3.3 Gaussian and mean curvatures

Recall that Gaussian curvature

$$K = \det S$$

and mean curvature

$$H = \text{Tr } S$$

It is easy to see that for Gaussian curvature

$$K = \det S = \det(G^{-1}A) = \frac{\det A}{\det G}.$$

We know already the geometrical meaning of Gaussian and mean curvatures from the point of view of the External Observer:

Gaussian curvature K equals to the product of principal curvatures, and mean curvatures equals to the sum of principal curvatures.

Now we ask a principal question: what about internal observer, "aunt" living on the surface?

We will show that Gaussian curvature can be expressed in terms of induced Riemannian metric, i.e. it is an internal characteristic of the surface, invariant of isometries.

It is not the case with mean curvature: cylinder is isometric to the plane but it has non-zero mean curvature.

1.3.4 Examples of calculation of Weingarten operator, Second quadratic forms, curvatures for cylinder, cone and sphere.

Cylinder

We already calculated induced Riemannian metric on the cylinder (see (??)).

Cylinder is given by the equation $x^2 + y^2 = R^2$. One can consider the following parameterisation of this surface:

$$\mathbf{r}(h, \varphi): \quad \begin{cases} x = R \cos \varphi \\ y = R \sin \varphi \\ z = h \end{cases}, \quad \mathbf{r}_h = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \mathbf{r}_\varphi = \begin{pmatrix} -R \sin \varphi \\ R \cos \varphi \\ 0 \end{pmatrix}, \quad (1.17)$$

$$G_{cylinder} = (dx^2 + dy^2 + dz^2) \big|_{x=R \cos \varphi, y=R \sin \varphi, z=h} =$$

$$= (-R \sin \varphi d\varphi)^2 + (R \cos \varphi d\varphi)^2 + dh^2 = R^2 d\varphi^2 + dh^2, \quad ||g_{\alpha\beta}|| = \begin{pmatrix} 1 & 0 \\ 0 & R^2 \end{pmatrix}.$$

Normal unit vector $\mathbf{n} = \pm \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix}$. Choose $\mathbf{n} = \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix}$. Weingarten

operator

$$\begin{aligned}
S\partial_h &= -\nabla_{\mathbf{r}_h}^{\text{can.flat}} \mathbf{n} = -\partial_{\mathbf{r}_h} \mathbf{n} = -\partial_h \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix} = 0, \\
S\partial_\varphi &= -\nabla_{\mathbf{r}_\varphi}^{\text{can.flat}} \mathbf{n} = -\partial_{\mathbf{r}_\varphi} \mathbf{n} = -\partial_\varphi \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} \sin \varphi \\ -\cos \varphi \\ 0 \end{pmatrix} = -\frac{\partial_\varphi}{R}. \\
S \begin{pmatrix} \partial_h \\ \partial_\varphi \end{pmatrix} &= \begin{pmatrix} 0 \\ -\frac{\partial_\varphi}{R} \end{pmatrix}, \quad S = \begin{pmatrix} 0 & 0 \\ 0 & -\frac{1}{R} \end{pmatrix}. \tag{1.18}
\end{aligned}$$

Calculate second quadratic form: $\mathbf{r}_{hh} = \partial_h \mathbf{r}_h = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$, $\mathbf{r}_{h\varphi} = \mathbf{r}_{\varphi h} =$

$$\partial_h \begin{pmatrix} -R \sin \varphi \\ R \cos \varphi \\ 0 \end{pmatrix} = 0, \quad \mathbf{r}_{\varphi\varphi} = \partial_\varphi \begin{pmatrix} -R \sin \varphi \\ R \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -R \cos \varphi \\ -R \sin \varphi \\ 0 \end{pmatrix} = -R\mathbf{n}.$$

We have

$$A_{\alpha\beta} = \langle \mathbf{r}_{\alpha\beta}, \mathbf{n} \rangle, \quad A = \begin{pmatrix} \langle \mathbf{r}_{hh}, \mathbf{n} \rangle & \langle \mathbf{r}_{h\varphi}, \mathbf{n} \rangle \\ \langle \mathbf{r}_{\varphi h}, \mathbf{n} \rangle & \langle \mathbf{r}_{\varphi\varphi}, \mathbf{n} \rangle \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & -R \end{pmatrix}, \tag{1.19}$$

$$A = GS = \begin{pmatrix} 1 & 0 \\ 0 & R^2 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & -R \end{pmatrix},$$

For Gaussian and mean curvatures we have

$$K = \det S = \frac{\det A}{\det G} = \det \begin{pmatrix} 0 & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = 0, \tag{1.20}$$

and mean curvature

$$H = \text{Tr } S = \text{Tr} \begin{pmatrix} 0 & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = -\frac{1}{R}. \tag{1.21}$$

Mean curvature is define up to a sign. If we change $\mathbf{n} \rightarrow -\mathbf{n}$ mean curvature $H \rightarrow \frac{1}{R}$ and Gaussian curvature will not change.

Cone

We already calculated induced Riemannian metric on the cone (see (??)).

Cone is given by the equation $x^2 + y^2 - k^2 z^2 = 0$. One can consider the following parameterisation of this surface:

$$\mathbf{r}(h, \varphi): \begin{cases} x = kh \cos \varphi \\ y = kh \sin \varphi \\ z = h \end{cases}, \quad \mathbf{r}_h = \begin{pmatrix} k \cos \varphi \\ k \sin \varphi \\ 1 \end{pmatrix}, \quad \mathbf{r}_\varphi = \begin{pmatrix} -kh \sin \varphi \\ kh \cos \varphi \\ 0 \end{pmatrix}, \quad (1.22)$$

$$\begin{aligned} G_{\text{cone}} &= (dx^2 + dy^2 + dz^2) \Big|_{x=kh \cos \varphi, y=kh \sin \varphi, z=h} = \\ &= (-kh \sin \varphi d\varphi + k \cos \varphi dh)^2 + (kh \cos \varphi d\varphi + k \sin \varphi dh)^2 + dh^2 = \\ &= k^2 h^2 d\varphi^2 + (k^2 + 1) dh^2, \quad ||g_{\alpha\beta}|| = \begin{pmatrix} k^2 + 1 & 0 \\ 0 & k^2 h^2 \end{pmatrix}. \end{aligned}$$

One can see that $\mathbf{N} = \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix}$ is orthogonal to the surface: $\mathbf{N} \perp \mathbf{r}_h, \mathbf{r}_\varphi$.

Hence normal unit vector $\mathbf{n} = \pm \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix}$. Choose $\mathbf{n} = \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix}$.

Weingarten operator

$$\begin{aligned} S\partial_h &= -\nabla_{\mathbf{r}_h}^{\text{can.flat}} \mathbf{n} = -\partial_{\mathbf{r}_h} \mathbf{n} = -\partial_h \left(\frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix} \right) = 0, \\ S\partial_\varphi &= -\nabla_{\mathbf{r}_\varphi}^{\text{can.flat}} \mathbf{n} = -\partial_{\mathbf{r}_\varphi} \mathbf{n} = -\partial_\varphi \left(\frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix} \right) = \\ &= \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = -\frac{\partial_\varphi}{kh\sqrt{k^2+1}}. \\ S \begin{pmatrix} \partial_h \\ \partial_\varphi \end{pmatrix} &= \begin{pmatrix} 0 \\ -\frac{\partial_\varphi}{kh\sqrt{k^2+1}} \end{pmatrix}, \quad S = \begin{pmatrix} 0 & 0 \\ 0 & \frac{-1}{kh\sqrt{k^2+1}} \end{pmatrix}. \quad (1.23) \end{aligned}$$

Calculate second quadratic form: $\mathbf{r}_{hh} = \partial_h \mathbf{r}_h = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$, $\mathbf{r}_{h\varphi} = \mathbf{r}_{\varphi h} =$

$$\partial_h \begin{pmatrix} -kh \sin \varphi \\ kh \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -k \sin \varphi \\ k \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{r}_{\varphi\varphi} = \partial_\varphi \begin{pmatrix} -kh \sin \varphi \\ kh \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -kh \cos \varphi \\ -kh \sin \varphi \\ 0 \end{pmatrix}.$$

We have

$$A_{\alpha\beta} = \langle \mathbf{r}_{\alpha\beta}, \mathbf{n} \rangle, \quad A = \begin{pmatrix} \langle \mathbf{r}_{hh}, \mathbf{n} \rangle & \langle \mathbf{r}_{h\varphi}, \mathbf{n} \rangle \\ \langle \mathbf{r}_{\varphi h}, \mathbf{n} \rangle & \langle \mathbf{r}_{\varphi\varphi}, \mathbf{n} \rangle \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & -\frac{kh}{\sqrt{1+k^2}} \end{pmatrix}, \quad (1.24)$$

$$A = GS = \begin{pmatrix} k^2 + 1 & 0 \\ 0 & k^2 h^2 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & \frac{-1}{kh\sqrt{k^2+1}} \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & \frac{-kh}{\sqrt{k^2+1}} \end{pmatrix},$$

For Gaussian and mean curvatures we have

$$K = \det S = \frac{\det A}{\det G} = \det \begin{pmatrix} 0 & 0 \\ 0 & \frac{-1}{kh\sqrt{k^2+1}} \end{pmatrix} = 0, \quad (1.25)$$

and mean curvature

$$H = \text{Tr } S = \text{Tr} \begin{pmatrix} 0 & 0 \\ 0 & \frac{-1}{kh\sqrt{k^2+1}} \end{pmatrix} = \frac{-1}{kh\sqrt{k^2+1}}. \quad (1.26)$$

Mean curvature is define up to a sign. If we change $\mathbf{n} \rightarrow -\mathbf{n}$ mean curvature $H \rightarrow \frac{1}{R}$ and Gaussian curvature will not change.

Sphere

Sphere is given by the equation $x^2 + y^2 + z^2 = a^2$. Consider the parametrisation of sphere in spherical coordinates

$$\mathbf{r}(\theta, \varphi): \quad \begin{cases} x = R \sin \theta \cos \varphi \\ y = R \sin \theta \sin \varphi \\ z = R \cos \theta \end{cases} \quad (1.27)$$

We already calculated induced Riemannian metric on the sphere (see (1.3.4)). Recall that

$$\mathbf{r}_\theta = \begin{pmatrix} R \cos \theta \cos \varphi \\ R \cos \theta \sin \varphi \\ -R \sin \theta \end{pmatrix}, \quad \mathbf{r}_\varphi = \begin{pmatrix} -R \sin \theta \sin \varphi \\ R \sin \theta \cos \varphi \\ 0 \end{pmatrix}$$

and

$$\begin{aligned}
G_{S^2} &= (dx^2 + dy^2 + dz^2) \Big|_{x=R \sin \theta \cos \varphi, y=R \sin \theta \sin \varphi, z=R \cos \theta} = \\
&= (R \cos \theta \cos \varphi d\theta - R \sin \theta \sin \varphi d\varphi)^2 + (R \cos \theta \sin \varphi d\theta + R \sin \theta \cos \varphi d\varphi)^2 + \\
&= (-R \sin \theta d\theta)^2 = R^2 \cos^2 \theta d\theta^2 + R^2 \sin^2 \theta d\varphi^2 + R^2 \sin^2 \theta d\theta^2 = \\
&= R^2 d\theta^2 + R^2 \sin^2 \theta d\varphi^2, \quad ||g_{\alpha\beta}|| = \begin{pmatrix} R^2 & 0 \\ 0 & R^2 \sin^2 \theta \end{pmatrix}.
\end{aligned}$$

For the sphere $\mathbf{r}(\theta, \varphi)$ is orthogonal to the surface. Hence normal unit vector $\mathbf{n}(\theta, \varphi) = \pm \frac{\mathbf{r}(\theta, \varphi)}{R} = \pm \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix}$. Choose $\mathbf{n} = \frac{\mathbf{r}}{R} = \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix}$.

Weingarten operator

$$\begin{aligned}
S\partial_\theta &= -\nabla_{\mathbf{r}_\theta}^{\text{can.flat}} \mathbf{n} = -\partial_\theta \mathbf{n} = -\partial_\theta \left(\frac{\mathbf{r}}{R} \right) = -\frac{\mathbf{r}_\theta}{R}, \\
S\partial_\varphi &= -\nabla_{\mathbf{r}_\varphi}^{\text{can.flat}} \mathbf{n} = -\partial_\varphi \mathbf{n} = -\partial_\varphi \left(\frac{\mathbf{r}}{R} \right) = -\frac{\mathbf{r}_\varphi}{R}. \\
S \begin{pmatrix} \partial_\theta \\ \partial_\varphi \end{pmatrix} &= \begin{pmatrix} -\frac{\partial_\theta}{R} \\ -\frac{\partial_\varphi}{R} \end{pmatrix}, \quad S = - \begin{pmatrix} \frac{1}{R} & 0 \\ 0 & \frac{1}{R} \end{pmatrix}. \tag{1.28}
\end{aligned}$$

For second quadratic form: $\mathbf{r}_{\theta\theta} = \partial_\theta \mathbf{r}_\theta = \begin{pmatrix} -R \sin \theta \cos \varphi \\ -R \sin \theta \sin \varphi \\ -R \cos \theta \end{pmatrix}$, $\mathbf{r}_{\theta\varphi} = \mathbf{r}_{\varphi\theta} =$

$$\partial_\theta \begin{pmatrix} -R \sin \theta \sin \varphi \\ R \sin \theta \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -R \cos \theta \sin \varphi \\ R \cos \theta \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{r}_{\varphi\varphi} = \partial_\varphi \begin{pmatrix} -R \sin \theta \sin \varphi \\ R \sin \theta \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -R \sin \theta \cos \varphi \\ -R \sin \theta \sin \varphi \\ 0 \end{pmatrix}.$$

We have

$$A_{\alpha\beta} = \langle \mathbf{r}_{\alpha\beta}, \mathbf{n} \rangle, \quad A = \begin{pmatrix} \langle \mathbf{r}_{\theta\theta}, \mathbf{n} \rangle & \langle \mathbf{r}_{\theta\varphi}, \mathbf{n} \rangle \\ \langle \mathbf{r}_{\varphi\theta}, \mathbf{n} \rangle & \langle \mathbf{r}_{\varphi\varphi}, \mathbf{n} \rangle \end{pmatrix} = \begin{pmatrix} -R & 0 \\ 0 & -R \sin^2 \theta \end{pmatrix}, \tag{1.29}$$

$$A = GS = \begin{pmatrix} R^2 & 0 \\ 0 & R^2 \sin^2 \theta \end{pmatrix} \begin{pmatrix} -\frac{1}{R} & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = -R \begin{pmatrix} 1 & 0 \\ 0 & \sin^2 \theta \end{pmatrix},$$

For Gaussian and mean curvatures we have

$$K = \det S = \frac{\det A}{\det G} = \det \begin{pmatrix} -\frac{1}{R} & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = \frac{1}{R^2}, \tag{1.30}$$

and mean curvature

$$H = \text{Tr } S = \text{Tr} \begin{pmatrix} -\frac{1}{R} & 0 \\ 0 & -\frac{1}{R} \end{pmatrix} = -\frac{2}{R}, \quad (1.31)$$

Mean curvature is define up to a sign. If we change $\mathbf{n} \rightarrow -\mathbf{n}$ mean curvature $H \rightarrow \frac{1}{R}$ and Gaussian curvature will not change.

We see that for the sphere Gaussian curvature is not equal to zero, whilst for cylinder and cone Gaussian curvature equals to zero.

1.4 Parallel transport along closed curves on surfaces. Derivation formulae and *Theorema Egregium*

1.4.1 Formulation of the main result

Let M be a surface in Euclidean space \mathbf{E}^3 . Consider a closed curve C on M , $M: \mathbf{r} = \mathbf{r}(u, v)$, $C: \mathbf{r} = \mathbf{r}(u(t), v(t))$, $0 \leq t \leq t_1$, $\mathbf{x}(0) = \mathbf{x}(t_1)$. ($u(t), v(t)$ are internal coordinates of the curve C .)

Consider the parallel transport of an arbitrary tangent \mathbf{X} vector along the closed curve C :

$$\mathbf{X}(t) = \underbrace{X^\alpha(t) \frac{\partial}{\partial u_\alpha} \Big|_{u^\alpha(t)}}_{\text{Internal observer}} = \underbrace{X^\alpha(t) \mathbf{r}_\alpha \Big|_{\mathbf{r}(u(t), v(t))}}_{\text{External observer}}, \quad \left(\mathbf{r}_\alpha = \frac{\partial x^i}{\partial u^\alpha} \frac{\partial}{\partial x^i} \right).$$

$$\mathbf{X}(t): \quad \frac{\nabla \mathbf{X}(t)}{dt} = 0, \quad 0 \leq t \leq t_1,$$

i.e.

$$\frac{dX^\alpha(t)}{dt} + X^\beta(t) \Gamma_{\beta\gamma}^\alpha(u(t)) \frac{du^\gamma(t)}{dt} = 0, \quad 0 \leq t \leq t_1, \quad (1.32)$$

where ∇ is the connection induced on the surface M by canonical flat connection (see (1.5)), i.e. the Levi-Civita connection of the induced Riemannian metric on the surface M and $\Gamma_{\beta\gamma}^\alpha$ its Christoffel symbols:

$$\Gamma_{\alpha\beta}^\gamma = \frac{1}{2} g^{\gamma\pi} \left(\frac{\partial g_{\pi\alpha}}{\partial u^\beta} + \frac{\partial g_{\pi\beta}}{\partial u^\alpha} - \frac{\partial g_{\alpha\beta}}{\partial u^\pi} \right), \quad \text{where } g_{\alpha\beta} = \langle \mathbf{r}_\alpha, \mathbf{r}_\beta \rangle = \frac{\partial x^i}{\partial u^\alpha} \frac{\partial x^i}{\partial u^\beta} \quad (1.33)$$

are components of induced Riemannian metric $G_M = g_{\alpha\beta} du^\alpha$ (see the sub-section 4.2.1)

Let $\mathbf{r}(0) = \mathbf{p}$ be a starting (and ending) point of the curve C : $\mathbf{r}(0) = \mathbf{r}(t_1) = \mathbf{p}$. The differential equation (1.32) defines the linear operator

$$R_C: T_{\mathbf{p}}M \longrightarrow T_{\mathbf{p}}M \quad (1.34)$$

For any vector $\mathbf{X} \in T_{\mathbf{p}}M$, its image the vector $R_C\mathbf{X}$ as the solution of the differential equation (1.32) with initial condition $\mathbf{X}(t)|_{t=0} = \mathbf{X}$.

On the other hand we know that parallel transport of the vector does not change its length (see (??) in the subsection 3.2.1):

$$\langle \mathbf{X}, \mathbf{X} \rangle = \langle R_C\mathbf{X}, R_C\mathbf{X} \rangle \quad (1.35)$$

We see that \mathbf{R}_C is an orthogonal operator in the 2-dimensional vector space $T_{\mathbf{p}}M$. We know that orthogonal operator preserving orientation is the operator of rotation on some angle Φ .

We see that if \mathbf{R}_C preserves orientation² then the action of operator \mathbf{R}_C on vectors is rotation on the angle, i.e. the result of parallel transport along closed curve is rotation on the angle

$$\Delta\Phi = \Delta\Phi(C) \quad (1.36)$$

which depends on the curve.

The very beautiful question arises: How to calculate this angle $\Delta\Phi(C)$

Theorem *Let M be a surface in Euclidean space \mathbf{E}^3 . Let C be a closed curve C on M such that C is a boundary of a compact oriented domain $D \subset M$. Consider the parallel transport of an arbitrary tangent vector along the closed curve C . As a result of parallel transport along this closed curve any tangent vector rotates through the angle*

$$\Delta\Phi = \angle(\mathbf{X}, \mathbf{R}_C\mathbf{X}) = \int_D K d\sigma, \quad (1.37)$$

where K is the Gaussian curvature and $d\sigma = \sqrt{\det g} du dv$ is the area element induced by the Riemannian metric on the surface M , i.e. $d\sigma = \sqrt{\det g} du dv$.

Remark One can show that the angle of rotation does not depend on initial point of the curve.

²We consider the case if the operator R_C preserves orientation. In our considerations we consider only the case if the closed curve C is a boundary of a compact oriented domain $D \subset M$. In this case one can see that operator R_C preserves an orientation.

Example Consider the closed curve, "latitude" C_{θ_0} on the sphere of the radius R . Calculations show that

$$\Delta\Phi(C_{\theta_0}) = 2\pi(1 - \cos \theta_0) \quad (1.38)$$

(see the Homework 6). On the other hand the latitude C_{θ_0} is the boundary of the segment D with area $2\pi RH$ where $H = R(1 - \cos \theta_0)$. Hence

$$\angle(\mathbf{X}, \mathbf{R}_C \mathbf{X}) = \frac{2\pi RH}{R^2} = \frac{1}{R^2} \cdot \text{area of the segment} = \int_D K d\sigma$$

since Gaussian curvature is equal to $\frac{1}{R^2}$

The proof of this Theorem is the content of the subsections above.

Before show one of the remarkable corollaries of this Theorem.

Let D be a small domain around a given point \mathbf{p} , let C its boundary and $\Delta\Phi(D)$ the angle of rotation. Denote by $S(D)$ an area of this domain. Applying the Theorem for the case when area of the domain D tends to zero we come to the statement that

$$\text{if } S(D) \rightarrow 0 \text{ then } \Delta\Phi(D) = \int_D K d\sigma \rightarrow K(\mathbf{p})S(D), \text{ i.e.}$$

$$K(\mathbf{p}) = \lim_{S(D) \rightarrow 0} \frac{\Delta\Phi(D)}{S(D)}, \text{ i.e.} \quad (1.39)$$

Now notice that left hand side of this equation defining Gaussian curvature $K(\mathbf{p})$ depends only on Riemannian metric on the surface C . Indeed numerator of LHS is defined by the solution of differential equation (1.32) which depends on Levi-Civita connection depending on the induced Riemannian metric, and denominator is an area depending on Riemannian metric too.

Thus we come to very important

Corollary *Gauß Egregium Theorema*

Gaussian curvature of the surface is invariant of isometries.

In next subsections we develop the technique which itself is very interesting. One of the applications of this technique is the proof of the Theorem (1.37). This we will prove Theorema Egregium too.

Later in the fifth section we will give another proof of the Theorema Egregium.

1.4.2 Derivation formulae

Let M be a surface embedded in Euclidean space \mathbf{E}^3 , $M: \mathbf{r} = \mathbf{r}(u, v)$.

Let $\mathbf{e}, \mathbf{f}, \mathbf{n}$ be three vector fields defined on the points of this surface such that they form an orthonormal basis at any point, so that the vectors \mathbf{e}, \mathbf{f} are tangent to the surface and the vector \mathbf{n} is orthogonal to the surface³. Vector fields $\mathbf{e}, \mathbf{f}, \mathbf{n}$ are functions on the surface M :

$$\mathbf{e} = \mathbf{e}(u, v), \mathbf{f} = \mathbf{f}(u, v), \mathbf{n} = \mathbf{n}(u, v).$$

Consider 1-forms $d\mathbf{e}, d\mathbf{f}, d\mathbf{n}$:

$$d\mathbf{e} = \frac{\partial \mathbf{e}}{\partial u} du + \frac{\partial \mathbf{e}}{\partial v} dv, d\mathbf{f} = \frac{\partial \mathbf{f}}{\partial u} du + \frac{\partial \mathbf{f}}{\partial v} dv, d\mathbf{n} = \frac{\partial \mathbf{n}}{\partial u} du + \frac{\partial \mathbf{n}}{\partial v} dv$$

These 1-forms take values in the vectors in \mathbf{E}^3 , i.e. they are *vector valued* 1-forms. Any vector in \mathbf{E}^3 attached at an arbitrary point of the surface can be expanded over the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$. Thus vector valued 1-forms $d\mathbf{e}, d\mathbf{f}, d\mathbf{n}$ can be expanded in a sum of 1-forms with values in basic vectors $\mathbf{e}, \mathbf{f}, \mathbf{n}$. E.g. for $d\mathbf{e} = \frac{\partial \mathbf{e}}{\partial u} du + \frac{\partial \mathbf{e}}{\partial v} dv$ expanding vectors $\frac{\partial \mathbf{e}}{\partial u}$ and $\frac{\partial \mathbf{e}}{\partial v}$ over basis vectors we come to

$$\frac{\partial \mathbf{e}}{\partial u} = A_1(u, v)\mathbf{e} + B_1(u, v)\mathbf{f} + C_1(u, v)\mathbf{n}, \quad \frac{\partial \mathbf{e}}{\partial v} = A_2(u, v)\mathbf{e} + B_2(u, v)\mathbf{f} + C_2(u, v)\mathbf{n}$$

thus

$$\begin{aligned} d\mathbf{e} &= \frac{\partial \mathbf{e}}{\partial u} du + \frac{\partial \mathbf{e}}{\partial v} dv = (A_1\mathbf{e} + B_1\mathbf{f} + C_1\mathbf{n}) du + (A_2\mathbf{e} + B_2\mathbf{f} + C_2\mathbf{n}) dv = \\ &= \underbrace{(A_1 du + A_2 dv)}_{M_{11}} \mathbf{e} + \underbrace{(B_1 du + B_2 dv)}_{M_{12}} \mathbf{f} + \underbrace{(C_1 du + C_2 dv)}_{M_{13}} \mathbf{n}, \end{aligned} \quad (1.40)$$

i.e.

$$d\mathbf{e} = M_{11}\mathbf{e} + M_{12}\mathbf{f} + M_{13}\mathbf{n},$$

where M_{11}, M_{12} and M_{13} are 1-forms on the surface M defined by the relation (1.40).

In the same way we do the expansions of vector-valued 1-forms $d\mathbf{f}$ and $d\mathbf{n}$ we come to

$$\begin{aligned} d\mathbf{e} &= M_{11}\mathbf{e} + M_{12}\mathbf{f} + M_{13}\mathbf{n} \\ d\mathbf{f} &= M_{21}\mathbf{e} + M_{22}\mathbf{f} + M_{23}\mathbf{n} \\ d\mathbf{n} &= M_{31}\mathbf{e} + M_{32}\mathbf{f} + M_{33}\mathbf{n} \end{aligned}$$

³One can say that $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ is an orthonormal basis in $T_{\mathbf{p}}\mathbf{E}^3$ at every point of surface $\mathbf{p} \in M$ such that $\{\mathbf{e}, \mathbf{f}\}$ is an orthonormal basis in $T_{\mathbf{p}}\mathbf{E}^3$ at every point of surface $\mathbf{p} \in M$.

This equation can be rewritten in the following way:

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} \quad (1.41)$$

Proposition The matrix M in the equation (1.41) is antisymmetrical matrix, i.e.

$$\begin{aligned} M_{11} &= M_{22} = M_{33} = 0 \\ M_{12} &= -M_{21} = a \\ M_{13} &= -M_{31} = b \\ M_{23} &= -M_{32} = -b \end{aligned} \quad (1.42)$$

i.e.

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix}, \quad (1.43)$$

where a, b, c are 1-forms on the surface M .

Formulae (2.5) are called *derivation formulae*.

Prove this Proposition. Recall that $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ is orthonormal basis, i.e. at every point of the surface

$$\langle \mathbf{e}, \mathbf{e} \rangle = \langle \mathbf{f}, \mathbf{f} \rangle = \langle \mathbf{n}, \mathbf{n} \rangle = 1, \text{ and } \langle \mathbf{e}, \mathbf{f} \rangle = \langle \mathbf{e}, \mathbf{n} \rangle = \langle \mathbf{f}, \mathbf{n} \rangle = 0$$

Now using (1.41) we have

$$\langle \mathbf{e}, \mathbf{e} \rangle = 1 \Rightarrow d\langle \mathbf{e}, \mathbf{e} \rangle = 0 = 2\langle \mathbf{e}, d\mathbf{e} \rangle = \langle \mathbf{e}, M_{11}\mathbf{e} + M_{12}\mathbf{f} + M_{13}\mathbf{n} \rangle =$$

$$M_{11}\langle \mathbf{e}, \mathbf{e} \rangle + M_{12}\langle \mathbf{e}, \mathbf{f} \rangle + M_{13}\langle \mathbf{e}, \mathbf{n} \rangle = M_{11} \Rightarrow M_{11} = 0.$$

Analogously

$$\langle \mathbf{f}, \mathbf{f} \rangle = 1 \Rightarrow d\langle \mathbf{f}, \mathbf{f} \rangle = 0 = 2\langle \mathbf{f}, d\mathbf{f} \rangle = \langle \mathbf{f}, M_{21}\mathbf{e} + M_{22}\mathbf{f} + M_{23}\mathbf{n} \rangle = M_{22} \Rightarrow M_{22} = 0,$$

$$\langle \mathbf{n}, \mathbf{n} \rangle = 1 \Rightarrow d\langle \mathbf{n}, \mathbf{n} \rangle = 0 = 2\langle \mathbf{n}, d\mathbf{n} \rangle = \langle \mathbf{n}, M_{31}\mathbf{e} + M_{32}\mathbf{f} + M_{33}\mathbf{n} \rangle = M_{33} \Rightarrow M_{33} = 0.$$

We proved already that $M_{11} = M_{22} = M_{33} = 0$. Now prove that $M_{12} = -M_{21}$, $M_{13} = -M_{31}$ and $M_{23} = -M_{32}$.

$$\langle \mathbf{e}, \mathbf{f} \rangle = 0 \Rightarrow d\langle \mathbf{e}, \mathbf{f} \rangle = 0 = \langle \mathbf{e}, d\mathbf{f} \rangle + \langle d\mathbf{e}, \mathbf{f} \rangle =$$

$$\langle \mathbf{e}, M_{21}\mathbf{e} + M_{22}\mathbf{f} + M_{23}\mathbf{n} \rangle + \langle M_{11}\mathbf{e} + M_{12}\mathbf{f} + M_{13}\mathbf{n}, \mathbf{f} \rangle = M_{21} + M_{12} = 0.$$

Analogously

$$\langle \mathbf{e}, \mathbf{n} \rangle = 0 \Rightarrow d\langle \mathbf{e}, \mathbf{n} \rangle = 0 = \langle \mathbf{e}, d\mathbf{n} \rangle + \langle d\mathbf{e}, \mathbf{n} \rangle =$$

$$\langle \mathbf{e}, M_{31}\mathbf{e} + M_{32}\mathbf{f} + M_{33}\mathbf{n} \rangle + \langle M_{11}\mathbf{e} + M_{12}\mathbf{f} + M_{13}\mathbf{n}, \mathbf{n} \rangle = M_{31} + M_{13} = 0$$

and

$$\langle \mathbf{f}, \mathbf{n} \rangle = 0 \Rightarrow d\langle \mathbf{f}, \mathbf{n} \rangle = 0 = \langle \mathbf{f}, d\mathbf{n} \rangle + \langle d\mathbf{f}, \mathbf{n} \rangle =$$

$$\langle \mathbf{f}, M_{31}\mathbf{e} + M_{32}\mathbf{f} + M_{33}\mathbf{n} \rangle + \langle M_{21}\mathbf{e} + M_{22}\mathbf{f} + M_{23}\mathbf{n}, \mathbf{n} \rangle = M_{32} + M_{23} = 0.$$

Remark This proof could be performed much more shortly in condensed notations. Derivation formulae (2.5) in condensed notations are

$$d\mathbf{e}_i = M_{ik}\mathbf{e}_k \tag{1.44}$$

Orthonormality condition means that $\langle \mathbf{e}_i, \mathbf{e}_k \rangle = \delta_{ik}$. Hence

$$d\langle \mathbf{e}_i, \mathbf{e}_k \rangle = 0 = \langle d\mathbf{e}_i, \mathbf{e}_k \rangle + \langle \mathbf{e}_i, d\mathbf{e}_k \rangle = \langle M_{im}\mathbf{e}_m, \mathbf{e}_k \rangle + \langle \mathbf{e}_i, M_{kn}\mathbf{e}_n \rangle = M_{ik} + M_{ki} = 0 \quad \blacksquare \tag{1.45}$$

Much shorter, is not it?

1.4.3 Gauss condition (structure equations)

Derive the relations between 1-forms a, b and c in derivation formulae.

Recall that a, b, c are 1-forms, $\mathbf{e}, \mathbf{f}, \mathbf{n}$ are vector valued functions (0-forms) and $d\mathbf{e}, d\mathbf{f}, d\mathbf{n}$ are vector valued 1-forms. (We use the simple identity that $ddf = 0$ and the fact that for 1-form $\omega \wedge \omega = 0$.) We have from derivation formulae (2.5) that

$$d^2\mathbf{e} = 0 = d(a\mathbf{f} + b\mathbf{n}) = da\mathbf{f} - a \wedge d\mathbf{f} + db\mathbf{n} - b \wedge d\mathbf{n} =$$

$$da\mathbf{f} - a \wedge (-a\mathbf{e} + c\mathbf{n}) + db\mathbf{n} - b \wedge (-b\mathbf{e} - c\mathbf{f}) =$$

$$(da + b \wedge c)\mathbf{f} + (a \wedge a + b \wedge b)\mathbf{e} + (db - a \wedge c)\mathbf{n} = (da + b \wedge c)\mathbf{f} + (db + c \wedge a)\mathbf{n} = 0.$$

We see that

$$(da + b \wedge c)\mathbf{f} + (db + c \wedge a)\mathbf{n} = 0 \tag{1.46}$$

Hence components of the left hand side equal to zero:

$$(da + b \wedge c) = 0 \quad (db + c \wedge a) = 0. \tag{1.47}$$

Analogously

$$\begin{aligned}
d^2\mathbf{f} = 0 &= d(-a\mathbf{e} + c\mathbf{n}) = -da\mathbf{e} + a \wedge d\mathbf{e} + dc\mathbf{n} - c \wedge d\mathbf{n} = \\
&= -da\mathbf{e} + a \wedge (a\mathbf{f} + b\mathbf{n}) + dc\mathbf{n} - c \wedge (-b\mathbf{e} - c\mathbf{f}) = \\
&= (-da + c \wedge b)\mathbf{e} + (dc + a \wedge b)\mathbf{n} = 0.
\end{aligned}$$

Hence we come to structure equations:

$$\begin{aligned}
da + b \wedge c &= 0 \\
db + c \wedge a &= 0 \\
dc + a \wedge b &= 0
\end{aligned} \tag{1.48}$$

1.4.4 Geometrical meaning of derivation formulae. Weingarten operator and second quadratic form in terms of derivation formulae.

Let M be a surface in \mathbf{E}^3 .

Let $\mathbf{e}, \mathbf{f}, \mathbf{n}$ be three vector fields defined on the points of this surface such that they form an orthonormal basis at any point, so that the vectors \mathbf{e}, \mathbf{f} are tangent to the surface and the vector \mathbf{n} is orthogonal to the surface. Note that in generally these vectors are not coordinate vectors.

Describe Riemannian geometry on the surface M in terms of this basis and derivation formulae (2.5).

Induced Riemannian metric

If G is the Riemannian metric induced on the surface M then since \mathbf{e}, \mathbf{f} is orthonormal basis at every tangent space $T_{\mathbf{p}}M$ then

$$G(\mathbf{e}, \mathbf{e}) = G(\mathbf{f}, \mathbf{f}) = 1, \quad G(\mathbf{e}, \mathbf{f}) = G(\mathbf{f}, \mathbf{e}) = 0 \tag{1.49}$$

The matrix of the Riemannian metric in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$G = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} b(\mathbf{e}) & b(\mathbf{f}) \\ c(\mathbf{e}) & c(\mathbf{f}) \end{pmatrix} \tag{1.50}$$

Induced connection Let ∇ be the connection induced by the canonical flat connection on the surface M .

Then according equations (1.4) and derivation formulae (2.5) for every tangent vector \mathbf{X}

$$\nabla_{\mathbf{X}}\mathbf{e} = (\partial_{\mathbf{X}}\mathbf{e})_{\text{tangent}} = (d\mathbf{e}(\mathbf{X}))_{\text{tangent}} = (a(\mathbf{X})\mathbf{f} + b(\mathbf{X})\mathbf{n})_{\text{tangent}} = a(\mathbf{X})\mathbf{f}. \quad (1.51)$$

and

$$\nabla_{\mathbf{X}}\mathbf{f} = (\partial_{\mathbf{X}}\mathbf{f})_{\text{tangent}} = (d\mathbf{f}(\mathbf{X}))_{\text{tangent}} = (-a(\mathbf{X})\mathbf{e} + c(\mathbf{X})\mathbf{n})_{\text{tangent}} = -a(\mathbf{X})\mathbf{e}. \quad (1.52)$$

In particular

$$\begin{aligned} \nabla_{\mathbf{e}}\mathbf{e} &= a(\mathbf{e})\mathbf{f} & \nabla_{\mathbf{f}}\mathbf{e} &= a(\mathbf{f})\mathbf{f} \\ \nabla_{\mathbf{e}}\mathbf{f} &= -a(\mathbf{e})\mathbf{e} & \nabla_{\mathbf{f}}\mathbf{f} &= -a(\mathbf{f})\mathbf{e} \end{aligned} \quad (1.53)$$

We know that the connection ∇ is Levi-Civita connection of the induced Riemannian metric (1.51) (see the subsection 4.2.1)⁴.

Second Quadratic form Let $A(\mathbf{X}, \mathbf{Y})$ be second quadratic form. Then according to (1.12) and derivation formulae (2.5) we have

$$\begin{aligned} A(\mathbf{e}, \mathbf{e}) &= \langle \partial_{\mathbf{e}}\mathbf{e}, \mathbf{n} \rangle = \langle d\mathbf{e}(\mathbf{e}), \mathbf{n} \rangle = \langle a(\mathbf{e})\mathbf{f} + b(\mathbf{e})\mathbf{n}, \mathbf{n} \rangle = b(\mathbf{e}), \\ A(\mathbf{f}, \mathbf{e}) &= \langle \partial_{\mathbf{f}}\mathbf{e}, \mathbf{n} \rangle = \langle d\mathbf{e}(\mathbf{f}), \mathbf{n} \rangle = \langle a(\mathbf{f})\mathbf{f} + b(\mathbf{f})\mathbf{n}, \mathbf{n} \rangle = b(\mathbf{f}), \\ A(\mathbf{e}, \mathbf{f}) &= \langle \partial_{\mathbf{e}}\mathbf{f}, \mathbf{n} \rangle = \langle d\mathbf{f}(\mathbf{e}), \mathbf{n} \rangle = \langle -a(\mathbf{e})\mathbf{f} + c(\mathbf{e})\mathbf{n}, \mathbf{n} \rangle = c(\mathbf{e}), \\ A(\mathbf{f}, \mathbf{f}) &= \langle \partial_{\mathbf{f}}\mathbf{f}, \mathbf{n} \rangle = \langle d\mathbf{f}(\mathbf{f}), \mathbf{n} \rangle = \langle -a(\mathbf{f})\mathbf{f} + c(\mathbf{f})\mathbf{n}, \mathbf{n} \rangle = c(\mathbf{f}), \end{aligned}$$

The matrix of the second quadratic form in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$A = \begin{pmatrix} A(\mathbf{e}, \mathbf{e}) & A(\mathbf{f}, \mathbf{e}) \\ A(\mathbf{e}, \mathbf{f}) & A(\mathbf{f}, \mathbf{f}) \end{pmatrix} = \begin{pmatrix} b(\mathbf{e}) & b(\mathbf{f}) \\ c(\mathbf{e}) & c(\mathbf{f}) \end{pmatrix} \quad (1.55)$$

This is symmetrical matrix (see the subsection 4.3.2):

$$A(\mathbf{f}, \mathbf{e}) = b(\mathbf{f}) = A(\mathbf{e}, \mathbf{f}) = c(\mathbf{e}). \quad (1.56)$$

Weingarten operator

⁴In particular this implies that this is symmetric connection, i.e.

$$\nabla_{\mathbf{f}}\mathbf{e} - \nabla_{\mathbf{e}}\mathbf{f} - [\mathbf{f}, \mathbf{e}] = a(\mathbf{f})\mathbf{f} + a(\mathbf{e})\mathbf{e} - [\mathbf{f}, \mathbf{e}] = 0. \quad (1.54)$$

Let S be Weingarten operator: $S\mathbf{X} = -\partial_{\mathbf{X}}\mathbf{n}$ (see the subsection 4.3.1). Then it follows from derivation formulae that

$$S\mathbf{X} = -\partial_{\mathbf{X}}\mathbf{n} = -d\mathbf{n}(\mathbf{X}) = -(-b(\mathbf{X})\mathbf{e} - c(\mathbf{X})\mathbf{f}) = b(\mathbf{X})\mathbf{e} + c(\mathbf{X})\mathbf{f}$$

In particular

$$S(\mathbf{e}) = b(\mathbf{e})\mathbf{e} + c(\mathbf{e})\mathbf{f}, S(\mathbf{f}) = b(\mathbf{f})\mathbf{e} + c(\mathbf{f})\mathbf{f}$$

and the matrix of the Weingarten operator in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$S = \begin{pmatrix} b(\mathbf{e}) & c(\mathbf{e}) \\ b(\mathbf{f}) & c(\mathbf{f}) \end{pmatrix} \quad (1.57)$$

According to the condition (1.56) the matrix S is symmetrical.

The relations $A = GS, S = G^{-1}A$ for Weingarten operator, Riemannian metric and second quadratic form are evidently obeyed for matrices of these operators in the basis \mathbf{e}, \mathbf{f} where $G = 1, A = S$.

1.4.5 Gaussian and mean curvature in terms of derivation formulae

Now we are equipped to express Gaussian and mean curvatures (see the subsection 4.3.3) in terms of derivation formulae. Using (1.57) we have for Gaussian curvature

$$K = \det S = b(\mathbf{e})c(\mathbf{f}) - c(\mathbf{e})b(\mathbf{f}) = (b \wedge c)(\mathbf{e}, \mathbf{f}) \quad (1.58)$$

and for mean curvature

$$H = \text{Tr } S = b(\mathbf{e}) + c(\mathbf{f}) \quad (1.59)$$

What next? We will study in more detail formula (1.58) later.

Now consider some examples of calculation derivation formulae, Weingarten operator, e.t.c. for some examples using derivation formulae.

1.4.6 Examples of calculations of derivation formulae for cylinder, cone and sphere

In the subsection 4.3.4 we calculated Weingarten operator, second quadratic form and curvatures for cylinder, cone and sphere. Now we do the same but in terms of derivation formulae.

Cylinder

We have to define three vector fields $\mathbf{e}, \mathbf{f}, \mathbf{n}$ on the points of the cylinder

$$\mathbf{r}(h, \varphi): \begin{cases} x = R \cos \varphi \\ y = R \sin \varphi \\ z = h \end{cases} \quad \text{such that they form an orthonormal basis at any}$$

point, so that the vectors \mathbf{e}, \mathbf{f} are tangent to the surface and the vector \mathbf{n} is orthogonal to the surface.

We already calculated coordinate vector fields $\mathbf{r}_h, \mathbf{r}_\varphi$ and normal unit vector field \mathbf{n} (see (1.17)):

$$\mathbf{r}_h = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \mathbf{r}_\varphi = \begin{pmatrix} -R \sin \varphi \\ R \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{n} = \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix}. \quad (1.60)$$

Vectors $\mathbf{r}_h, \mathbf{r}_\varphi$ and \mathbf{n} are orthogonal to each other but not all of them have unit length. One can choose

$$\mathbf{e} = \mathbf{r}_h = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \mathbf{f} = \frac{\mathbf{r}_\varphi}{R} = \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{n} = \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix} \quad (1.61)$$

These vectors form an orthonormal basis and \mathbf{e}, \mathbf{f} form an orthonormal basis in tangent space.

Derive for this basis derivation formulae (2.5). For vector fields $\mathbf{e}, \mathbf{f}, \mathbf{n}$ in (1.61) we have

$$d\mathbf{e} = 0, \quad d\mathbf{f} = d \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -\cos \varphi \\ -\sin \varphi \\ 0 \end{pmatrix} d\varphi = -\mathbf{n} d\varphi,$$

$$d\mathbf{n} = d \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} d\varphi = \mathbf{f} d\varphi,$$

i.e.

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -d\varphi \\ 0 & d\varphi & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix}, \quad (1.62)$$

i.e. in derivation formulae $a = b = 0$, $c = -d\varphi$.

The matrix of Weingarten operator in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$S = \begin{pmatrix} b(\mathbf{e}) & c(\mathbf{e}) \\ b(\mathbf{f}) & c(\mathbf{f}) \end{pmatrix} = \begin{pmatrix} 0 & -d\varphi(\mathbf{e}) \\ 0 & -d\varphi(\mathbf{f}) \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & -\frac{1}{R} \end{pmatrix}$$

According to (1.58) and (1.59) Gaussian curvature $K = b(\mathbf{e})c(\mathbf{f}) - b(\mathbf{f})c(\mathbf{e}) = 0$ and mean curvature

$$H = b(\mathbf{e}) + c(\mathbf{f}) = -d\varphi(\mathbf{f}) = -d\varphi\left(\frac{\mathbf{r}_\varphi}{R}\right) = -\frac{1}{R}$$

(Compare with calculations in the subsection 4.3.4)

Cone

For cone:

$$\mathbf{r}(h, \varphi): \begin{cases} x = kh \cos \varphi \\ y = kh \sin \varphi \\ z = h \end{cases},$$

$$\mathbf{r}_h = \begin{pmatrix} k \cos \varphi \\ k \sin \varphi \\ 1 \end{pmatrix}, \quad \mathbf{r}_\varphi = \begin{pmatrix} -kh \sin \varphi \\ kh \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{n} = \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix}$$

Tangent vectors $\mathbf{r}_h, \mathbf{r}_\varphi$ are orthogonal to each other. The length of the vector \mathbf{r}_h equals to $\sqrt{1+k^2}$ and the length of the vector \mathbf{r}_φ equals to kh . Hence we can choose orthonormal basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ such that vectors \mathbf{e}, \mathbf{f} are unit vectors in the directions of the vectors $\mathbf{r}_h, \mathbf{r}_\varphi$:

$$\mathbf{e} = \frac{\mathbf{r}_h}{\sqrt{1+k^2}} = \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} k \cos \varphi \\ k \sin \varphi \\ 1 \end{pmatrix}, \quad \mathbf{f} = \frac{\mathbf{r}_\varphi}{hk} = \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{n} = \frac{1}{\sqrt{1+k^2}} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix}$$

Calculate $d\mathbf{e}, d\mathbf{f}$ and $d\mathbf{n}$:

$$d\mathbf{e} = d \begin{pmatrix} k \cos \varphi \\ k \sin \varphi \\ 1 \end{pmatrix} = \frac{kd\varphi}{\sqrt{1+k^2}} \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = \frac{kd\varphi}{\sqrt{1+k^2}} \mathbf{f},$$

$$d\mathbf{f} = d \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = \begin{pmatrix} -\cos \varphi \\ -\sin \varphi \\ 0 \end{pmatrix} d\varphi =$$

$$\frac{-k}{1+k^2} \begin{pmatrix} k \cos \varphi \\ k \sin \varphi \\ 1 \end{pmatrix} d\varphi - \frac{d\varphi}{1+k^2} \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix} = \frac{-kd\varphi}{\sqrt{1+k^2}} \mathbf{e} - \frac{d\varphi}{\sqrt{1+k^2}} \mathbf{n},$$

and

$$d\mathbf{n} = \frac{1}{\sqrt{1+k^2}} d \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ -k \end{pmatrix} = \frac{d\varphi}{\sqrt{1+k^2}} \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix}.$$

We come to

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & \frac{kd\varphi}{\sqrt{1+k^2}} & 0 \\ -\frac{kd\varphi}{\sqrt{1+k^2}} & 0 & \frac{-d\varphi}{\sqrt{1+k^2}} \\ 0 & \frac{d\varphi}{\sqrt{1+k^2}} & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix}, \quad (1.63)$$

i.e. in derivation formulae for 1-forms $a = \frac{kd\varphi}{\sqrt{1+k^2}}$, $b = 0$ and $c = -\frac{d\varphi}{\sqrt{1+k^2}}$.

The matrix of Weingarten operator in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$S = \begin{pmatrix} b(\mathbf{e}) & c(\mathbf{e}) \\ b(\mathbf{f}) & c(\mathbf{f}) \end{pmatrix} = S = \begin{pmatrix} 0 & \frac{-d\varphi(\mathbf{e})}{\sqrt{1+k^2}} \\ 0 & \frac{-d\varphi(\mathbf{f})}{\sqrt{1+k^2}} \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & \frac{-1}{kh\sqrt{1+k^2}} \end{pmatrix}.$$

since $d\varphi(\mathbf{f}) = d\varphi\left(\frac{\mathbf{r}_\varphi}{kh}\right) = \frac{1}{kh}d\varphi(\partial_\varphi) = \frac{1}{kh}$.

According to (1.58), (1.59) Gaussian curvature

$$K = b(\mathbf{e})c(\mathbf{f}) - b(\mathbf{f})c(\mathbf{e}) = 0$$

and mean curvature

$$H = b(\mathbf{e}) + c(\mathbf{f}) = -d\varphi(\mathbf{f}) = -d\varphi\left(\frac{\mathbf{r}_\varphi}{R}\right) = -\frac{1}{R}$$

(Compare with calculations in the subsection 4.3.4)

Sphere

For sphere

$$\mathbf{r}(\theta, \varphi): \begin{cases} x = R \sin \theta \cos \varphi \\ y = R \sin \theta \sin \varphi \\ z = R \cos \theta \end{cases} \quad (1.64)$$

$$\mathbf{r}_\theta(\theta, \varphi) = \frac{\partial \mathbf{r}}{\partial \theta} = \begin{pmatrix} R \cos \theta \cos \varphi \\ R \cos \theta \sin \varphi \\ -R \sin \theta \end{pmatrix}, \quad \mathbf{r}_\varphi(\theta, \varphi) = \frac{\partial \mathbf{r}}{\partial \varphi} = \begin{pmatrix} -R \sin \theta \sin \varphi \\ R \sin \theta \cos \varphi \\ 0 \end{pmatrix},$$

$$\mathbf{n}(\theta, \varphi) = \frac{\mathbf{r}}{R} = \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix}.$$

Tangent vectors $\mathbf{r}_\theta, \mathbf{r}_\varphi$ are orthogonal to each other. The length of the vector \mathbf{r}_θ equals to R and the length of the vector \mathbf{r}_φ equals to $R \sin \theta$. Hence we can choose orthonormal basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ such that vectors \mathbf{e}, \mathbf{f} are unit vectors in the directions of the vectors $\mathbf{r}_\theta, \mathbf{r}_\varphi$:

$$\mathbf{e}(\theta, \varphi) = \frac{\mathbf{r}_\theta}{R} = \begin{pmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{pmatrix}, \quad \mathbf{f}(\theta, \varphi) = \frac{\mathbf{r}_\varphi}{R \sin \theta} = \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix}, \quad \mathbf{n}(\theta, \varphi) = \frac{\mathbf{r}}{R} = \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix}.$$

Calculate $d\mathbf{e}, d\mathbf{f}$ and $d\mathbf{n}$:

$$\begin{aligned} d\mathbf{e} &= d \begin{pmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{pmatrix} = \\ &= \begin{pmatrix} -\cos \theta \sin \varphi \\ \cos \theta \cos \varphi \\ 0 \end{pmatrix} d\varphi - \begin{pmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{pmatrix} d\theta = \cos \theta d\varphi \mathbf{f} - d\theta \mathbf{n}, \\ d\mathbf{f} &= d \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = - \begin{pmatrix} \cos \varphi \\ \sin \varphi \\ 0 \end{pmatrix} d\varphi = \\ &= -\cos \theta d\varphi \begin{pmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{pmatrix} - \sin \theta d\varphi \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix} = -\cos \theta d\varphi \mathbf{e} - \sin \theta d\varphi \mathbf{n}, \\ d\mathbf{n} &= d \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix} = \begin{pmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{pmatrix} d\theta + \begin{pmatrix} -\sin \theta \sin \varphi \\ \sin \theta \cos \varphi \\ 0 \end{pmatrix} d\varphi \\ &= d\theta \mathbf{e} + \sin \theta d\varphi \mathbf{f}. \end{aligned}$$

i.e.

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & \cos \theta d\varphi & -d\theta \\ -\cos \theta d\varphi & 0 & -\sin \theta d\varphi \\ d\theta & \sin \theta d\varphi & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix}, \quad (1.65)$$

i.e. in derivation formulae $a = \cos \theta d\varphi$, $b = -d\theta$, $c = -\sin \theta d\varphi$.

The matrix of Weingarten operator in the basis $\{\mathbf{e}, \mathbf{f}\}$ is

$$S = \begin{pmatrix} b(\mathbf{e}) & c(\mathbf{e}) \\ b(\mathbf{f}) & c(\mathbf{f}) \end{pmatrix} = \begin{pmatrix} -d\theta(\mathbf{e}) & -\sin \theta d\varphi(\mathbf{e}) \\ -d\theta(\mathbf{f}) & -\sin \theta d\varphi(\mathbf{f}) \end{pmatrix} = \begin{pmatrix} \frac{-1}{R} & 0 \\ 0 & -\frac{1}{R} \end{pmatrix}$$

since $d\theta(\mathbf{e}) = d\theta\left(\frac{\partial_\theta}{R}\right) = \frac{1}{R}d\theta(\partial_\theta) = \frac{1}{R}$, $d\varphi(\mathbf{e}) = d\varphi\left(\frac{\partial_\theta}{R}\right) = \frac{1}{R}d\varphi(\partial_\theta) = 0$.

According to (1.58) and (1.59) Gaussian curvature

$$K = b(\mathbf{e})c(\mathbf{f}) - b(\mathbf{f})c(\mathbf{e}) = \frac{1}{R^2}$$

and mean curvature

$$H = b(\mathbf{e}) + c(\mathbf{f}) = -\frac{2}{R}$$

Notice that for calculation of Weingarten operator and curvatures we used only 1-forms b and c , i.e. the derivation equation for $d\mathbf{n}$, ($d\mathbf{n} = d\theta\mathbf{e} + \sin \theta d\varphi\mathbf{f}$). (Compare with calculations in the subsection 4.3.4)

Mean curvature is define up to a sign. If we change $\mathbf{n} \rightarrow -\mathbf{n}$ mean curvature $H \rightarrow \frac{1}{R}$ and Gaussian curvature will not change.

We see that for the sphere Gaussian curvature is not equal to zero, whilst for cylinder and cone Gaussian curvature equals to zero.

1.4.7 Proof of the Theorem of parallel transport along closed curve.

We are ready now to prove the Theorem. Recall that the Theorem states following:

If C is a closed curve on a surface M such that C is a boundary of a compact oriented domain $D \subset M$, then during the parallel transport of an arbitrary tangent vector along the closed curve C the vector rotates through the angle

$$\Delta\Phi = \angle(\mathbf{X}, \mathbf{R}_C\mathbf{X}) = \int_D K d\sigma, \quad (1.66)$$

where K is the Gaussian curvature and $d\sigma = \sqrt{\det g} du dv$ is the area element induced by the Riemannian metric on the surface M , i.e. $d\sigma = \sqrt{\det g} du dv$.

(see (1.37)).

Recall that for derivation formulae (2.5) we obtained structure equations

$$\begin{aligned} da + b \wedge c &= 0 \\ db + c \wedge a &= 0 \\ dc + a \wedge b &= 0 \end{aligned} \tag{1.67}$$

We need to use only one of these equations, the equation

$$da + b \wedge c = 0. \tag{1.68}$$

This condition sometimes is called *Gauß condition*.

Let as always $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ be an orthonormal basis in $T_{\mathbf{p}}\mathbf{E}^3$ at every point of surface $\mathbf{p} \in M$ such that $\{\mathbf{e}, \mathbf{f}\}$ is an orthonormal basis in $T_{\mathbf{p}}\mathbf{E}^3$ at every point of surface $\mathbf{p} \in M$. Then the Gauß condition (1.68) and equation (1.58) mean that for Gaussian curvature on the surface M can be expressed through the 2-form da and base vectors $\{\mathbf{e}, \mathbf{f}\}$:

$$K = b \wedge c(\mathbf{e}, \mathbf{f}) = -da(\mathbf{e}, \mathbf{f}) \tag{1.69}$$

We use this formula to prove the Theorem.

Now calculate the parallel transport of an arbitrary tangent vector over the closed curve C on the surface M .

Let $\mathbf{r} = \mathbf{r}(u, v) = \mathbf{r}(u^\alpha)$ ($\alpha = 1, 2$, $(u, v) = (u^1, v^1)$) be an equation of the surface M .

Let $u^\alpha = u^\alpha(t)$ ($\alpha = 1, 2$) be the equation of the curve C . Let $\mathbf{X}(t)$ be the parallel transport of vector field along the closed curve C , i.e. $\mathbf{X}(t)$ is tangent to the surface M at the point $u(t)$ of the curve C and vector field $\mathbf{X}(t)$ is covariantly constant along the curve:

$$\frac{\nabla \mathbf{X}(t)}{dt} = 0$$

To write this equation in components we usually expanded the vector field in the coordinate basis $\{\mathbf{r}_u = \partial_u, \mathbf{r}_v = \partial_v\}$ and used Christoffel symbols of the connection $\Gamma_{\beta\gamma}^\alpha: \nabla_\beta \partial_\gamma = \Gamma_{\beta\gamma}^\alpha \partial_\alpha$.

Now we will do it in different way: *instead coordinate basis $\{\mathbf{r}_u = \partial_u, \mathbf{r}_v = \partial_v\}$ we will use the basis $\{\mathbf{e}, \mathbf{f}\}$* . In the subsection 3.4.4 we obtained that the connection ∇ has the following appearance in this basis

$$\nabla_{\mathbf{v}} \mathbf{e} = a(\mathbf{v}) \mathbf{f}, \quad \nabla_{\mathbf{v}} \mathbf{f} = -a(\mathbf{v}) \mathbf{e} \tag{1.70}$$

(see the equations (1.51) and (1.52))

Let

$$\mathbf{X} = \mathbf{X}(u(t)) = X^1(t)\mathbf{e}(u(t)) + X^2(t)\mathbf{f}(u(t))$$

Let be an expansion of tangent vector field $\mathbf{X}(t)$ over basis $\{\mathbf{e}, \mathbf{f}\}$. Let \mathbf{v} be velocity vector of the curve C . Then the equation of parallel transport $\frac{\nabla \mathbf{X}(t)}{dt} = 0$ will have the following appearance:

$$\begin{aligned} \frac{\nabla \mathbf{X}(t)}{dt} = 0 &= \nabla_{\mathbf{v}} (X^1(t)\mathbf{e}(u(t)) + X^2(t)\mathbf{f}(u(t))) = \\ &= \frac{dX^1(t)}{dt}\mathbf{e}(u(t)) + X^1(t)\nabla_{\mathbf{v}}\mathbf{e}(u(t)) + \frac{dX^2(t)}{dt}\mathbf{f}(u(t)) + X^2(t)\nabla_{\mathbf{v}}\mathbf{f}(u(t)) = \\ &= \frac{dX^1(t)}{dt}\mathbf{e}(u(t)) + X^1(t)a(\mathbf{v})\mathbf{f}(u(t)) + \frac{dX^2(t)}{dt}\mathbf{f}(u(t)) - X^2(t)a(\mathbf{v})\mathbf{e}(u(t)) = \\ &= \left(\frac{dX^1(t)}{dt} - X^2(t)a(\mathbf{v}) \right) \mathbf{e}(u(t)) + \left(\frac{dX^2(t)}{dt} + X^1(t)a(\mathbf{v}) \right) \mathbf{f}(u(t)) = 0. \end{aligned}$$

Thus we come to equation:

$$\begin{cases} \dot{X}^1(t) - a(\mathbf{v}(t))X^2 = 0 \\ \dot{X}^2(t) + a(\mathbf{v}(t))X^1 = 0 \end{cases}$$

There are many ways to solve this equation. It is very convenient to consider complex variable

$$Z(t) = X^1(t) + iX^2(t)$$

We see that

$$\dot{Z}(t) = \dot{X}^1(t) + i\dot{X}^2(t) = a(\mathbf{v}(t))X^2 - ia(\mathbf{v}(t))X^1 = -ia(\mathbf{v})Z(t),$$

i.e.

$$\frac{dZ(t)}{dt} = -ia(\mathbf{v}(t))Z(t) \quad (1.71)$$

The solution of this equation is:

$$Z(t) = Z(0)e^{-i \int_0^t a(\mathbf{v}(\tau))d\tau} \quad (1.72)$$

Calculate $\int_0^{t_1} a(\mathbf{v}(\tau))d\tau$ for closed curve $u(0) = u(t_1)$. Due to Stokes Theorem:

$$\int_0^{t_1} a(\mathbf{v}(t))dt = \int_C a = \int_D da$$

Hence using Gauss condition (1.68) we see that

$$\int_0^{t_1} a(\mathbf{v}(t))dt = \int_C a = \int_D da = - \int_D b \wedge c$$

Claim

$$\int_D b \wedge c = - \int_D da = \int K d\sigma. \quad (1.73)$$

Theorem follows from this claim:

$$Z(t_1) = Z(0)e^{-i \int_C a} = Z(0)e^{i \int_D b \wedge c} \quad (1.74)$$

Denote the integral $i \int_D b \wedge c$ by $\Delta\Phi$: $\Delta\Phi = i \int_D b \wedge c$. We have

$$Z(t_1) = X^1(t_1) + iX^2(t_1) = (X^1(0) + iX^2(0)) e^{i\Delta\Phi} = \quad (1.75)$$

It remains to prove the claim. The induced volume form $d\sigma$ is 2-form. Its value on two orthogonal unit vector \mathbf{e}, \mathbf{f} equals to 1:

$$d\sigma(\mathbf{e}, \mathbf{f}) = 1 \quad (1.76)$$

(In coordinates u, v volume form $d\sigma = \sqrt{\det g} du \wedge dv$).

The value of the form $b \wedge c$ on vectors $\{\mathbf{e}, \mathbf{f}\}$ equals to Gaussian curvature according to (1.69). We see that

$$b \wedge c(\mathbf{e}, \mathbf{f}) = -da(\mathbf{e}, \mathbf{f}) = K d\sigma(\mathbf{e}, \mathbf{f})$$

Hence 2-forms $b \wedge c$, $-da$ and volume form $d\sigma$ coincide. Thus we prove (1.73).

2 Curvature tensor

2.1 Definition

Let $\mathbf{X}, \mathbf{Y}, \mathbf{Z}$ be an arbitrary vector fields on the manifold equipped with affine connection ∇ .

Consider the following operation which assigns to the vector fields \mathbf{X}, \mathbf{Y} and \mathbf{Z} the new vector field.

$$\mathcal{R}(\mathbf{X}, \mathbf{Y})\mathbf{Z} = (\nabla_{\mathbf{X}}\nabla_{\mathbf{Y}} - \nabla_{\mathbf{Y}}\nabla_{\mathbf{X}} - \nabla_{[\mathbf{X}, \mathbf{Y}]})\mathbf{Z} \quad (2.1)$$

This operation is obviously linear over the scalar coefficients and it is $C^\infty(M)$ with respect to vector fields $\mathbf{X}, \mathbf{Y}, \mathbf{Z}$, i.e. for an arbitrary functions f, g, h

$$\mathcal{R}(f\mathbf{X}, g\mathbf{Y})(h\mathbf{Z}) = fgh\mathcal{R}(\mathbf{X}, \mathbf{Y})\mathbf{Z}, \quad (2.2)$$

i.e. it defines the tensor field of the type $\begin{pmatrix} 1 \\ 3 \end{pmatrix}$: If $\mathbf{X} = X^i\partial_i, \mathbf{Y} = Y^n\partial_n, \mathbf{Z} = Z^r\partial_r$ then according to (2.2)

$$\mathcal{R}(\mathbf{X}, \mathbf{Y})\mathbf{Z} = \mathcal{R}(X^m\partial_m, Y^n\partial_n)(Z^r\partial_r) = Z^r R_{rmn}^i X^m Y^n \partial_i$$

where we denote by R_{rmn}^i the components of the tensor \mathcal{R} in the coordinate basis ∂_i

$$R_{rmn}^i \partial_i = \mathcal{R}(\partial_m, \partial_n)\partial_r \quad (2.3)$$

Express components of the curvature tensor in terms of Christoffel symbols of the connection. If $\nabla_m \partial_n = \Gamma_{mn}^r \partial_r$ then according to the (2.1) we have:

$$\begin{aligned} R_{rmn}^i \partial_i &= \mathcal{R}(\partial_m, \partial_n)\partial_r = \nabla_{\partial_m} \nabla_{\partial_n} \partial_r - \nabla_{\partial_n} \nabla_{\partial_m} \partial_r = \\ &= \nabla_{\partial_m} (\Gamma_{nr}^p \partial_p) - \nabla_{\partial_n} (\Gamma_{mr}^p \partial_p) = \\ &= \partial_m \Gamma_{nr}^i + \Gamma_{mp}^i \Gamma_{nr}^p - (m \leftrightarrow n) \end{aligned} \quad (2.4)$$

The proof of the property (2.2) can be given just by straightforward calculations: Consider e.g. the case $f = g = 1$, then

$$\begin{aligned} \mathcal{R}(\mathbf{X}, \mathbf{Y})(h\mathbf{Z}) &= \nabla_{\mathbf{X}} \nabla_{\mathbf{Y}}(h\mathbf{Z}) - \nabla_{\mathbf{Y}} \nabla_{\mathbf{X}}(h\mathbf{Z}) - \nabla_{[\mathbf{X}, \mathbf{Y}]}(h\mathbf{Z}) = \\ &= \nabla_{\mathbf{X}} (\partial_{\mathbf{Y}} h\mathbf{Z} + h\nabla_{\mathbf{Y}} \mathbf{Z}) - \nabla_{\mathbf{Y}} (\partial_{\mathbf{X}} h\mathbf{Z} + h\nabla_{\mathbf{X}} \mathbf{Z}) - \partial_{[\mathbf{X}, \mathbf{Y}]} h\mathbf{Z} - h\nabla_{[\mathbf{X}, \mathbf{Y}]} \mathbf{Z} = \\ &= \partial_{\mathbf{X}} \partial_{\mathbf{Y}} h\mathbf{Z} + \partial_{\mathbf{Y}} h\nabla_{\mathbf{X}} \mathbf{Z} + \partial_{\mathbf{X}} h\nabla_{\mathbf{Y}} \mathbf{Z} + h\nabla_{\mathbf{X}} \nabla_{\mathbf{Y}} \mathbf{Z} - \\ &= \partial_{\mathbf{Y}} \partial_{\mathbf{X}} h\mathbf{Z} - \partial_{\mathbf{X}} h\nabla_{\mathbf{Y}} \mathbf{Z} - \partial_{\mathbf{Y}} h\nabla_{\mathbf{X}} \mathbf{Z} + h\nabla_{\mathbf{Y}} \nabla_{\mathbf{X}} \mathbf{Z} - \\ &= \partial_{[\mathbf{X}, \mathbf{Y}]} h\mathbf{Z} - h\nabla_{[\mathbf{X}, \mathbf{Y}]} \mathbf{Z} = \\ &= h [\nabla_{\mathbf{X}} \nabla_{\mathbf{Y}} \mathbf{Z} - \nabla_{\mathbf{Y}} \nabla_{\mathbf{X}} \mathbf{Z}] + [\partial_{\mathbf{X}} \partial_{\mathbf{Y}} h - \partial_{\mathbf{Y}} \partial_{\mathbf{X}} h - \partial_{[\mathbf{X}, \mathbf{Y}]} h] \mathbf{Z} = \\ &= h\nabla_{\mathbf{X}} \nabla_{\mathbf{Y}} \mathbf{Z} - \nabla_{\mathbf{Y}} \nabla_{\mathbf{X}} \mathbf{Z} - \nabla_{[\mathbf{X}, \mathbf{Y}]} \mathbf{Z} = h\mathcal{R}(\mathbf{X}, \mathbf{Y})\mathbf{Z}, \end{aligned}$$

since $\partial_{\mathbf{X}} \partial_{\mathbf{Y}} h - \partial_{\mathbf{Y}} \partial_{\mathbf{X}} h - \partial_{[\mathbf{X}, \mathbf{Y}]} h = 0$.

2.2 Curvature of surfaces in \mathbf{E}^3 . *Theorema Egregium* again

Express Riemannian curvature of surfaces in \mathbf{E}^3 in terms of derivation formulae (2.5).

Consider derivation formulae for the orthonormal basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ adjusted to the surface M :

$$d \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix} = \begin{pmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{pmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \\ \mathbf{n} \end{pmatrix}, \quad (2.5)$$

where as usual $\mathbf{e}, \mathbf{f}, \mathbf{n}$ vector fields of unit length which are orthogonal to each other and \mathbf{n} is orthogonal to the surface M . As we know the induced connection on the surface M is defined by the formulae (1.51) and (1.52):

$$\nabla_{\mathbf{Y}} \mathbf{e} = (d\mathbf{e}(\mathbf{Y}))_{\text{tangent}} = a(\mathbf{X})\mathbf{f}, \quad \nabla_{\mathbf{Y}} \mathbf{f} = (d\mathbf{f}(\mathbf{Y}))_{\text{tangent}} = -a(\mathbf{X})\mathbf{e}, \quad (2.6)$$

According to the definition of curvature calculate

$$R(\mathbf{e}, \mathbf{f})\mathbf{e} = \nabla_{\mathbf{e}} \nabla_{\mathbf{f}} \mathbf{e} - \nabla_{\mathbf{f}} \nabla_{\mathbf{e}} \mathbf{e} - \nabla_{[\mathbf{e}, \mathbf{f}]} \mathbf{e}.$$

Note that since the induced connection is symmetrical connection then:

$$\nabla_{\mathbf{e}} \mathbf{f} - \nabla_{\mathbf{f}} \mathbf{e} - [\mathbf{e}, \mathbf{f}] = 0. \quad (2.7)$$

hence due to (2.6)

$$[\mathbf{e}, \mathbf{f}] = \nabla_{\mathbf{e}} \mathbf{f} - \nabla_{\mathbf{f}} \mathbf{e} = a(\mathbf{e})\mathbf{e} + a(\mathbf{f})\mathbf{f} \quad (2.8)$$

Thus we see that $R(\mathbf{e}, \mathbf{f})\mathbf{e} =$

$$\begin{aligned} \nabla_{\mathbf{e}} \nabla_{\mathbf{f}} \mathbf{e} - \nabla_{\mathbf{f}} \nabla_{\mathbf{e}} \mathbf{e} - \nabla_{[\mathbf{e}, \mathbf{f}]} \mathbf{e} &= \nabla_{\mathbf{e}} (a(\mathbf{f})\mathbf{f}) - \nabla_{\mathbf{f}} (a(\mathbf{e})\mathbf{f}) - \nabla_{a(\mathbf{e})\mathbf{e} + a(\mathbf{f})\mathbf{f}} \mathbf{e} = \\ \partial_{\mathbf{e}} a(\mathbf{f})\mathbf{f} + a(\mathbf{f})\nabla_{\mathbf{e}} \mathbf{f} - \partial_{\mathbf{f}} a(\mathbf{e})\mathbf{f} - a(\mathbf{e})\nabla_{\mathbf{f}} \mathbf{f} - a(\mathbf{e})\nabla_{\mathbf{e}} \mathbf{e} - a(\mathbf{f})\nabla_{\mathbf{f}} \mathbf{e} &= \\ \partial_{\mathbf{e}} a(\mathbf{f})\mathbf{f} - a(\mathbf{f})a(\mathbf{e})\mathbf{e} - \partial_{\mathbf{f}} a(\mathbf{e})\mathbf{f} + a(\mathbf{e})a(\mathbf{f})\mathbf{e} - a(\mathbf{e})a(\mathbf{e})\mathbf{f} - a(\mathbf{f})a(\mathbf{f})\mathbf{f} &= \\ [\partial_{\mathbf{e}} a(\mathbf{f})\mathbf{f} - \partial_{\mathbf{f}} a(\mathbf{e})\mathbf{f} - a((\mathbf{e})\mathbf{e} - a(\mathbf{f})\mathbf{f})]\mathbf{f} &= da(\mathbf{e}, \mathbf{f})\mathbf{f}. \end{aligned}$$

Recall that we established in 1.69 that for Gaussian curvature K

$$K = b \wedge c(\mathbf{e}, \mathbf{f}) = -da(\mathbf{e}, \mathbf{f})$$

Hence we come to the relation:

$$R(\mathbf{e}, \mathbf{f})\mathbf{e} = da(\mathbf{e}, \mathbf{f}) = -K\mathbf{f}. \quad (2.9)$$

This means that

$$R_{112}^2 = -K$$

(in the basis \mathbf{e}, \mathbf{f}), i.e. the scalar curvature

$$R = 2R_{1212} = 2K$$

We come to fundamental relation which claims that the Gaussian curvature (the magnitude defined in terms of External observer) equals to the scalar curvature, the magnitude defined in terms of Internal observer. This gives us the straightforward proof of Theorema Egregium.