Introduction to Geometry

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1 Euclidean space

1.1 Recollection of vector space and Euclidean vector space

We recall here important notions from linear algebra of vector space and Euclidean vector space.

1.1.1 Vector space.

Vector space V on real numbers is a set of vectors with operations " + "—addition of vector and " \cdot "—multiplication of vector Lon real number (sometimes called coefficients, scalars).

1

Remark We denote by 0 real number 0 and *vector* $\mathbf{0}$. Sometimes we have to be careful to distinguish between zero vector $\mathbf{0}$ and number zero.

1.1.2 Basic example of (n-dimensional) vector space— \mathbb{R}^n

A basic example of vector space (over real numbers) is a space of ordered n-tuples of real numbers.

 \mathbf{R}^2 is a space of pairs of real numbers. $\mathbf{R}^2 = \{(x,y), x,y \in \mathbf{R}\}$

- $\forall \mathbf{a}, \mathbf{b} \in V, \mathbf{a} + \mathbf{b} \in V,$
- $\forall \lambda \in \mathbf{R}, \forall \mathbf{a} \in V, \lambda \mathbf{a} \in V$.
- $\forall \mathbf{a}, \mathbf{ba} + \mathbf{b} = \mathbf{b} + \mathbf{a}$ (commutativity)
- $\forall \mathbf{a}, \mathbf{b}, \mathbf{c}, \ \mathbf{a} + (\mathbf{b} + \mathbf{c}) = (\mathbf{a} + \mathbf{b}) + \mathbf{c}$ (associativity)
- \exists **0** such that \forall **a**, **a** + **0** = **a**
- $\forall \mathbf{a}$ there exists a vector $-\mathbf{a}$ such that $\mathbf{a} + (-\mathbf{a}) = 0$.
- $\forall \lambda \in \mathbf{R}, \lambda(\mathbf{a} + \mathbf{b}) = \lambda \mathbf{a} + \lambda \mathbf{b}$
- $\forall \lambda, \mu \in \mathbf{R}(\lambda + \mu)\mathbf{a} = \lambda \mathbf{a} + \mu \mathbf{a}$
- $(\lambda \mu) \mathbf{a} = \lambda(\mu \mathbf{a})$
- $1\mathbf{a} = \mathbf{a}$

¹These operations obey the following axioms

 \mathbf{R}^3 is a space of triples of real numbers. $\mathbf{R}^3 = \{(x,y,z), \ x,y,z \in \mathbf{R}\}$ \mathbf{R}^4 is a space of quadruples of real numbers. $\mathbf{R}^4 = \{(x,y,z,t), \ x,y,z,t,\in \mathbf{R}\}$ and so on...

 \mathbf{R}^n —is a space of *n*-typles of real numbers:

$$\mathbf{R}^{n} = \{ (x^{1}, x^{2}, \dots, x^{n}), \ x^{1}, \dots, x^{n} \in \mathbf{R} \}$$
 (1.1)

If $\mathbf{x}, \mathbf{y} \in \mathbf{R}^n$ are two vectors, $\mathbf{x} = (x^1, \dots, x^n), \mathbf{y} = (y^1, \dots, y^n)$ then

$$\mathbf{x} + \mathbf{y} = (x_1 + y_1, \dots, x_n + y_n).$$

and multiplication on scalars is defined as

$$\lambda \mathbf{x} = \lambda \cdot (x^1, \dots, x^n) = (\lambda x^1, \dots, \lambda x^n), \quad (\lambda \in \mathbf{R}).$$

1.1.3 Linear dependence of vectors

We often consider linear combinations in vector space:

$$\sum_{i} \lambda_{i} \mathbf{x}_{i} = \lambda_{1} \mathbf{x}_{1} + \lambda_{2} \mathbf{x}_{2} + \dots + \lambda_{m} \mathbf{x}_{m}, \qquad (1.2)$$

where $\lambda_1, \lambda_2, \ldots, \lambda_m$ are coefficients (real numbers), $\mathbf{x}_1, \mathbf{x}_2, \ldots, \mathbf{x}_m$ are vectors from vector space V. We say that linear combination (1.2) is *trivial* if all coefficients $\lambda_1, \lambda_2, \ldots, \lambda_m$ are equal to zero.

$$\lambda_1 = \lambda_2 = \cdots = \lambda_m = 0$$
.

We say that linear combination (1.2) is *not trivial* if at least one of coefficients $\lambda_1, \lambda_2, \ldots, \lambda_m$ is not equal to zero:

$$\lambda_1 \neq 0, \text{ or } \lambda_2 \neq 0, \text{ or } \dots \text{ or } \lambda_m \neq 0.$$

Recall definition of linearly dependent and linearly independent vectors:

Definition The vectors $\{\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m\}$ in vector space V are linearly dependent if there exists a non-trivial linear combination of these vectors such that it is equal to zero.

In other words we say that the vectors $\{\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m\}$ in vector space V are *linearly dependent* if there exist coefficients $\mu_1, \mu_2, \dots, \mu_m$ such that at least one of these coefficients is not equal to zero and

$$\mu_1 \mathbf{x}_1 + \mu_2 \mathbf{x}_2 + \dots + \mu_m \mathbf{x}_m = 0. \tag{1.3}$$

Respectively vectors $\{\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m\}$ are linearly independent if they are not linearly dependent. This means that an arbitrary linear combination of these vectors which is equal zero is trivial.

In other words vectors $\{\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_m\}$ are linearly independent if the condition

$$\mu_1 \mathbf{x}_1 + \mu_2 \mathbf{x}_2 + \dots + \mu_m \mathbf{x}_m = 0$$

implies that $\mu_1 = \mu_2 = \cdots = \mu_m = 0$.

Very useful and workable

Proposition Vectors $\{\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m\}$ in vector space V are linearly dependent if and only if at least one of these vectors is expressed via linear combination of other vectors:

$$\mathbf{x}_i = \sum_{j \neq i} \lambda_j \mathbf{x}_j$$
 .

1.1.4 Dimension of vector space. Basis in vector space.

Definition Vector space V has a dimension n if there exist n linearly independent vectors in this vector space, and any n+1 vectors in V are linearly dependent.

In the case if in the vector space V for an arbitrary N there exist N linearly independent vectors then the space V is *infinite-dimensional*. An example of infinite-dimensional vector space is a space V of all polynomials of an arbitrary order. One can see that for an arbitrary N polynomials $\{1, x, x^2, x^3, \ldots, x^N\}$ are linearly idependent. (Try to prove it!). This implies V is infinite-dimensional vector space.

Rasis

Definition Let V be n-dimensional vector space. The ordered set $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ of n linearly independent vectors in V is called a basis of the vector space V.

Remark We say 'a basis', not 'the basis' since there are many bases in the vector space (see also Homeworks 1.2).

Remark Focus your attention: basis is *an ordered* set of vectors, not just a set of vectors².

The following Proposition is very useful:

²See later on orientation of vector spaces, where the ordering of vectors of basis will be highly important.

Proposition Let $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ be an arbitrary basis in n-dimensional vector space V. Then any vector $\mathbf{x} \in V$ can be expressed as a linear combination of vectors $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ in a unique way, i.e. for every vector $\mathbf{x} \in V$ there exists an ordered set of coefficients $\{x^1, \ldots, a^n\}$ such that

$$\mathbf{x} = x^1 \mathbf{e}_1 + \dots + x^n \mathbf{e}_n \tag{1.4}$$

and if

$$\mathbf{x} = a^1 \mathbf{e}_1 + \dots + a^n \mathbf{e}_n = b^1 \mathbf{e}_1 + \dots + b^n \mathbf{e}_n, \tag{1.5}$$

then $a^1 = b^1, a^2 = b^2, \ldots, a^n = b^n$. In other words for any vector $\mathbf{x} \in V$ there exists an ordered n-tuple (x^1, \ldots, x^n) of coefficients such that $\mathbf{x} = \sum_{i=1}^n x^i \mathbf{e}_i$ and this n-tuple is unique.

In other words:

Basis is a set of linearly independent vectors in vector space V which span (generate) vector space V.

Recall that we say that vector space V is spanned by vectors $\{\mathbf{x}_1, \dots, \mathbf{x}_n\}$ (or vectors vectors $\{\mathbf{x}_1, \dots, \mathbf{x}_n\}$ span vector space V) if any vector $\mathbf{a} \in V$ can be expresses as a linear combination of vectors $\{\mathbf{x}_1, \dots, \mathbf{x}_n\}$.

Definition Coefficients $\{a^1, \ldots, a^n\}$ are called *components of the vector* \mathbf{x} in the basis $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ or just shortly *components of the vector* \mathbf{x} .

Example Canonical basis in \mathbb{R}^n

We considered above the basic example of vector space—a space of ordered n-tuples of real numbers: $\mathbf{R}^n = \{(x^1, x^2, \dots, x^n), x^i \in \mathbf{R}\}$ (see (1.1)). One can see that it is n-dimensional vector space. Consider vectors $\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n \in \mathbf{R}^n$:

Then for an arbitrary vector $\mathbf{R}^n \ni \mathbf{a} = (a^1, a^2, a^3, \dots, a^n),$

$$\mathbf{a} = a^1(1,0,0\ldots,0,0) + a^2(0,1,0\ldots,0,0) + a^3(0,0,1,0\ldots,0,0) + \cdots + a^n(0,1,0\ldots,0,1) = a^n(0,0,0,0) + a^n(0,0,0$$

$$\sum_{i=1}^{m} a^{i} \mathbf{e}_{i} = a^{i} \mathbf{e}_{i} \qquad \text{(we will use sometimes condensed notations } \mathbf{x} = x^{i} \mathbf{e}_{i}\text{)}$$

For every vector $\mathbf{a} \in \mathbf{R}^n$ we have unique expansion via the vectors $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ is a basis of \mathbf{R}^n . The basis (1.6) is the distinguished basis. Sometimes it is called *canonical basis in* \mathbf{R}^n . One can find another basis in \mathbf{R}^n -just take an arbitrary ordered set of n linearly independent vectors. (See exercises in Homework 0).

1.1.5 Scalar product. Euclidean space

In vector space one have additional structure: scalar product of vectors.

Definition Scalar product in a vector space V is a function $B(\mathbf{x}, \mathbf{y})$ on a pair of vectors which takes real values and satisfies the the following conditions:

$$B(\mathbf{x}, \mathbf{y}) = B(\mathbf{y}, \mathbf{x})$$
 (symmetricity condition)
 $B(\lambda \mathbf{x} + \mu \mathbf{x}', \mathbf{y}) = \lambda B(\mathbf{x}, \mathbf{y}) + \mu B(\mathbf{x}', \mathbf{y})$ (linearity condition) (1.7)
 $B(\mathbf{x}, \mathbf{x}) \ge 0$, $B(\mathbf{x}, \mathbf{x}) = 0 \Leftrightarrow \mathbf{x} = 0$ (positive-definiteness condition)

Definition Euclidean space is a vector space equipped with a scalar product.

One can easy to see that the function $B(\mathbf{x}, \mathbf{y})$ is bilinear function, i.e. it is linear function with respect to the second argument also. This follows from previous axioms:

$$B(\mathbf{x}, \lambda \mathbf{y} + \mu \mathbf{y}') \underbrace{=}_{\text{symm.}} B(\lambda \mathbf{y} + \mu \mathbf{y}', \mathbf{x}) \underbrace{=}_{\text{linear.}} \lambda B(\mathbf{y}, \mathbf{x}) + \mu B(\mathbf{y}', \mathbf{x}) \underbrace{=}_{\text{symm.}} \lambda B(\mathbf{x}, \mathbf{y}) + \mu B(\mathbf{x}, \mathbf{y}').$$

A bilinear function $B(\mathbf{x}, \mathbf{y})$ on pair of vectors is called sometimes bilinear form on vector space. Bilinear form $B(\mathbf{x}, \mathbf{y})$ which satisfies the symmetricity condition is called symmetric bilinear form. Scalar product is nothing but symmetric bilinear form on vectors which is positive-definite: $B(\mathbf{x}, \mathbf{x}) \geq 0$ and is non-degenerate $((\mathbf{x}, \mathbf{x}) = 0 \Rightarrow \mathbf{x} = 0$.

Example We considered the vector space \mathbb{R}^n , the space of *n*-tuples (see the subsection 1.2). One can consider the vector space \mathbb{R}^n as Euclidean space provided by the scalar product

$$B(\mathbf{x}, \mathbf{y}) = x^1 y^1 + \dots + x^n y^n \tag{1.8}$$

This scalar product sometimes is called *canonical scalar product*.

Exercise Check that it is indeed scalar product.

Example We consider in 2-dimensional vector space V with basis $\{\mathbf{e}_1, \mathbf{e}_2\}$ and $B(\mathbf{X}, \mathbf{Y})$ such that $B(\mathbf{e}_1, \mathbf{e}_1) = 3$, $B(\mathbf{e}_2, \mathbf{e}_2) = 5$ and $B(\mathbf{e}_1, \mathbf{e}_2) = 0$. Then for every two vectors $\mathbf{X} = x^1\mathbf{e}_1 + x^2\mathbf{e}_2$ and $\mathbf{Y} = y^1\mathbf{e}_1 + y^2\mathbf{e}_2$ we have that

$$B(\mathbf{X}, \mathbf{Y}) = (\mathbf{X}, \mathbf{Y}) = (x^1 \mathbf{e}_1 + x^2 \mathbf{e}_2, y^1 \mathbf{e}_1 + y^2 \mathbf{e}_2) =$$

$$x^{1}y^{1}(\mathbf{e}_{1}, \mathbf{e}_{1}) + x^{1}y^{2}(\mathbf{e}_{1}, \mathbf{e}_{2}) + x^{2}y^{1}(\mathbf{e}_{2}, \mathbf{e}_{1}) + x^{2}y^{2}(\mathbf{e}_{2}, \mathbf{e}_{2}) = 3x^{1}y^{1} + 5x^{2}y^{2}.$$

One can see that all axioms are obeyed.

Remark Scalar product sometimes is called "inner" product or "dot" product. Later on we will use for scalar product $B(\mathbf{x}, \mathbf{y})$ just shorter notation (\mathbf{x}, \mathbf{y}) (or $\langle \mathbf{x}, \mathbf{y} \rangle$). Sometimes it is used for scalar product a notation $\mathbf{x} \cdot \mathbf{y}$. Usually this notation is reserved only for the canonical case (1.8).

Counterexample Consider again 2-dimensional vector space V with basis $\{e_1, e_2\}$.

Show that operation such that $(\mathbf{e}_1, \mathbf{e}_1) = (\mathbf{e}_2, \mathbf{e}_2) = 0$ and $(\mathbf{e}_1, \mathbf{e}_2) = 1$ does not define scalar product. *Solution*. For every two vectors $\mathbf{X} = x^1 \mathbf{e}_1 + x^2 \mathbf{e}_2$ and $\mathbf{Y} = y^1 \mathbf{e}_1 + y^2 \mathbf{e}_2$ we have that

$$(\mathbf{X}, \mathbf{Y}) = (x^1 \mathbf{e}_1 + x^2 \mathbf{e}_2, y^1 \mathbf{e}_1 + y^2 \mathbf{e}_2) = x^1 y^2 + x^2 y^1$$

hence for vector $\mathbf{X} = (1, -1)$ $(\mathbf{X}, \mathbf{X}) = -2 < 0$. Positive-definiteness is not fulfilled.

1.1.6 Orthonormal basis in Euclidean space

One can see that for scalar product (1.8) and for the basis $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ defined by the relation (1.6) the following relations hold:

$$(\mathbf{e}_i, \mathbf{e}_j) = \delta_{ij} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$$
 (1.9)

Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an ordered set of n vectors in n-dimensional Euclidean space which obeys the conditions (1.9). One can see that this ordered set is a basis ³.

³Indeed prove that conditions (1.9) imply that these n vectors are linear independent. Suppose that $\lambda_1 \mathbf{e}_1 + \lambda_2 \mathbf{e}_2 + \cdots + \lambda_n \mathbf{e}_n = 0$. For an arbitrary i multiply the left and right hand sides of this relation on a vector \mathbf{e}_i . We come to condition $\lambda_i = 0$. Hence vectors $(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n)$ are linearly dependent.

Definition-Proposition The ordered set of vectors $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ in n-dimensional Euclidean space which obey the conditions (1.9) is a basis. This basis is called an orthonormal basis.

One can prove that every (finite-dimensional) Euclidean space possesses orthonormal basis.

Later by default we consider only orthonormal bases in Euclidean spaces. Respectively scalar product will be defined by the formula (1.8). Indeed let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an orthonormal basis in Euclidean space. Then for an arbitrary two vectors \mathbf{x}, \mathbf{y} , such that $\mathbf{x} = \sum x^i \mathbf{e}_i$, $\mathbf{y} = \sum y^j \mathbf{e}_i$ we have:

$$(\mathbf{x}, \mathbf{y}) = \left(\sum x^i \mathbf{e}_i, \sum y^j \mathbf{e}_j\right) = \sum_{i,j=1}^n x^i y^j (\mathbf{e}_i, \mathbf{e}_j) = \sum_{i,j=1}^n x^i y^j \delta_{ij} = \sum_{i=1}^n x^i y^i$$

We come to the canonical scalar product (1.8). Later on we usually will consider scalar product defined by the formula (1.8) i.e. scalar product in orthonormal basis.

Remark We consider here general definition of scalar product then came to conclusion that in a special basis, (*orthonormal basis*), this is nothing but usual 'dot' product (1.8).

1.2 Affine spaces and vector spaces

AFFINE SPACE WITH ORIGIN IS A VECTOR SPACE

Let V be an arbitrary vector space.

Consider a set A whose elements will be called 'points' We say that A is an affine space associated with vector space V if the following rule is defined: to every point $P \in A$ and an arbitrary vector $\mathbf{x} \in V$ a point Q is assigned:

$$\forall P \in A, \quad \forall \mathbf{x} \in V, \quad (P, \mathbf{x}) \mapsto Q \in A$$
 (1.10)

We denote $Q = P + \mathbf{x}$.

The following properties must be satisfied:

- For arbitrary two vectors $\mathbf{x}, \mathbf{y} \in V$ and arbitrary point $P \in A$, $P + (\mathbf{x} + \mathbf{y}) = (P + \mathbf{x}) + \mathbf{y}$.
- For an arbitrary point P ∈ A, P + 0 = P.
 (Recall that 0 is the zero vector in the vector space V.)

• For arbitrary two points $P, Q \in A$ there exists unique vector $\mathbf{y} \in V$ such that $P + \mathbf{y} = Q$.

If $P + \mathbf{x} = Q$ we often denote the vector $\mathbf{x} = Q - P = \overrightarrow{PQ}$. We say that vector $\mathbf{x} = \overrightarrow{PQ}$ starts at the point P and it ends at the point Q.

One can see that if vector $\mathbf{x} = \vec{PQ}$, then $\vec{QP} = -\mathbf{x}$; if P, Q, R are three arbitrary points then $\vec{PQ} + \vec{QR} = \vec{PR}$.

One can reconstruct vector space V in terms of an affine space A, and vice versa. Namely, let A be an affine space associated with vector space V. Choose an arbitrary point $O \in A$ as an the origin, and consider the vectors starting at the origin: We come to the vector space V:

$$V = \text{set of vectors } \vec{OQ} \text{ where } Q \text{ is an arbitrary point in } A$$
,

which is associated with an affine space A.

Let V be an arbitrary vector space. We will define now an affine space associated with this vector space. Consider two copies of the vector space V. The elements of the first copy we will call "points", and the elements of the second copy we will call as usual "vectors":

first copy of
$$V$$
 second copy of V
$$\bigvee_{\text{elements of } V \text{ are points}} \text{elements of } V \text{ are vectors}$$
 (1.11)

Let $A = \mathbf{a}$ be an arbitrary point of the affine space, (i.e. an element of the *first* copy of vector space V) and let \mathbf{x} is an arbitrary vector of the vector space V (i.e. an element of the *second* copy of vector space V). We define the action (1.10) in the following way:

$$(A, \mathbf{x}) \mapsto B = A + \mathbf{x} = \mathbf{a} + \mathbf{x}, \quad \mathbf{x} = \vec{AB}.$$

The point B is the vector $\mathbf{a} + \mathbf{x} \in V$ belonging to the first copy of the vector space V.

We assign to two 'points' $A = \mathbf{a}, B = \mathbf{b}$ (elements of the *first* copy of vector space V) the vector $\mathbf{x} = \mathbf{b} - \mathbf{a}$ (elements of the *second* copy of vector space V).

For example vector space \mathbb{R}^n of *n*-tuples of real numbers can be considered as a set of points. If we choose arbitrary two points $A = (a^1, a^2, \dots, a^n)$ and

 $B = (b^1, b^2, \dots, b_n)$, then these two points define a vector \vec{AB} which is equal to $\vec{AB} = B - A = (b^1 - a^1, b^2 - a^2, \dots, b_n - a_n)$.

The associated with each other affine space and vector space \mathbb{R}^n we will usually denote by the same letter.

1.2.1 Eucldiean affine space.

Respectively one can consider Euclidean vector space as a set of points. Let \mathbf{E}^n be n-dimensional Euclidean vector space, i.e. vector space equipped with scalar product. Let $\{\mathbf{e}_i\}$ $(i=1,\ldots,n)$ be an arbitrary orthonormal basis in the vector space \mathbf{E}^n . Now consider this vector space as a set of points. Choose arbitrary two points (vectors of the *first* copy of the vector space \mathbf{E}^n), $A = a^1\mathbf{e}_1 + a^2\mathbf{e}_2 + \cdots + a^n\mathbf{e}_n$ and $B = b^1\mathbf{e}_1 + b^2\mathbf{e}_2 + \cdots + b^n\mathbf{e}_n$. These points define a vector \vec{AB} (in the *second* copy of the vector space \mathbf{E}^n) which is equal to

$$\vec{AB} = B - A = (b^1 - a^1)\mathbf{e}_1 + (b^2 - a^2)\mathbf{e}_2 + \dots + (b^n - a^n)\mathbf{e}_n$$

The distance between two points \vec{A} , \vec{B} is the length of corresponding vector \vec{AB} , and the length of the vector \vec{AB} is defined by the scalar product:

$$|\vec{AB}| = \sqrt{(\vec{AB}, \vec{AB})} = \sqrt{((b^1 - a^1) \mathbf{e}_1 + \dots + (b^n - a^n) \mathbf{e}_n, (b^1 - a^1) \mathbf{e}_1 + \dots + (b^n - a^n) \mathbf{e}_n)}$$

$$= \sqrt{(b^1 - a^1)^2 + \dots + (b^n - a^n)^2}.$$

We recall very important formula how scalar product is related with the angle between vectors: if φ is an angle between vectors \mathbf{x} and \mathbf{y} then

$$(\mathbf{x}, \mathbf{y}) = x^1 y^1 + x^2 y^2 + \dots + x^n y^n = |\mathbf{x}| |\mathbf{y}| \cos \varphi$$
 (1.12)

(We suppose that vectors \mathbf{x}, \mathbf{y} are defined in orthonormal basis.) In particularly it follows from this formula that

angle between vectors \mathbf{x}, \mathbf{y} is acute if scalar product (\mathbf{x}, \mathbf{y}) is positive angle between vectors \mathbf{x}, \mathbf{y} is obtuse if scalar product (\mathbf{x}, \mathbf{y}) is negative vectors \mathbf{x}, \mathbf{y} are perpendicular if scalar product (\mathbf{x}, \mathbf{y}) is equal to zero (1.13)

Remark The associated with each other affine space and Euclidean vector space \mathbf{E}^n we will denote by the same letter.

Remark Geometrical intuition tells us that cosinus of the angle between two vectors has to be less or equal to one and it is equal to one if and only if vectors \mathbf{x} , \mathbf{y} are collinear. Comparing with (1.12) we come to the inequality:

$$(\mathbf{x}, \mathbf{y})^2 = (x^1 y^1 + \dots + x^n y^n)^2 \le ((x^1)^2 + \dots + (x^n)^2) ((y^1)^2 + (\dots + (y^n)^2) = (\mathbf{x}, \mathbf{x})(\mathbf{y}, \mathbf{y})$$
and $(\mathbf{x}, \mathbf{y})^2 = (\mathbf{x}, \mathbf{x})(\mathbf{y}, \mathbf{y})$ if vectors are colinear, i.e. $x^i = \lambda y^i$

$$(1.14)$$

This is famous Cauchy–Buniakovsky–Schwarz inequality, one of most important inequalities in mathematics. (See for more details the last exercise in the Homework 0)

1.3 Transition matrices. Orthogonal bases and orthogonal matrices

1.3.1 Bases and transition matrices

One can consider different bases in vector space.

Let A be $n \times n$ matrix with real entries, $A = ||a_{ij}||, i, j = 1, 2, ..., n$:

$$A = \begin{pmatrix} a_{11} & a_{12} \dots & a_{1n} \\ a_{21} & a_{22} \dots & a_{2n} \\ a_{31} & a_{32} \dots & a_{3n} \\ \dots & \dots & \dots \\ a_{(n-1)1} & a_{(n-1)2} \dots & a_{(n-1)n} \\ a_{n1} & a_{n2} \dots & a_{nn} \end{pmatrix}$$

Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an arbitrary basis in *n*-dimensional vector space V.

The basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ can be considered as row of vectors, or $1 \times n$ matrix with entries—vectors.

Multiplying $1 \times n$ matrix $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ on matrix A we come to new row of vectors $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ such that

$$\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\} = \{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\} A =$$
 (1.15)

$$\{\mathbf{e}_{1}', \mathbf{e}_{2}', \dots, \mathbf{e}_{n}'\} = \{\mathbf{e}_{1}, \mathbf{e}_{2}, \dots, \mathbf{e}_{n}\} \begin{pmatrix} a_{11} & a_{12} \dots & a_{1n} \\ a_{21} & a_{22} \dots & a_{2n} \\ a_{31} & a_{32} \dots & a_{3n} \\ \dots & \dots & \dots \\ a_{(n-1)1} & a_{(n-1)2} \dots & a_{(n-1)n} \\ a_{n1} & a_{n2} \dots & a_{nn} \end{pmatrix}$$
(1.16)

$$\begin{cases} \mathbf{e}'_1 = a_{11}\mathbf{e}_1 + a_{21}\mathbf{e}_2 + a_{31}\mathbf{e}_3 + \dots + a_{(n-1)1}\mathbf{e}_{n-1} + a_{n1}\mathbf{e}_n \\ \mathbf{e}'_1 = a_{12}\mathbf{e}_1 + a_{22}\mathbf{e}_2 + a_{32}\mathbf{e}_3 + \dots + a_{(n-1)2}\mathbf{e}_{n-1} + a_{n2}\mathbf{e}_n \\ \mathbf{e}'_1 = a_{13}\mathbf{e}_1 + a_{23}\mathbf{e}_2 + a_{33}\mathbf{e}_3 + \dots + a_{(n-1)3}\mathbf{e}_{n-1} + a_{n1}\mathbf{e}_n \\ \dots = \dots + \dots + \dots + \dots + \dots \\ \mathbf{e}'_n = a_{1n}\mathbf{e}_1 + a_{2n}\mathbf{e}_2 + a_{3n}\mathbf{e}_3 + \dots + a_{(n-1)n}\mathbf{e}_{n-1} + a_{nn}\mathbf{e}_n \end{cases}$$

or shortly:

$$\mathbf{e}_i' = \sum_{k=1}^n \mathbf{e}_k a_{ki} \,. \tag{1.17}$$

Definition Matrix A which transforms a basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ to the row of vectors $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ (see equation (1.17)) is transition matrix from the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ to the row $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$.

What is the condition that the row $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ is a basis too? The row, ordered set of vectors, $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ is a basis if and only if vectors $(\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n)$ are linearly independent. Thus we come to

Proposition 1 Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be a basis in n-dimensional vector space V, and let A be an $n \times n$ matrix with real entries. Then

$$\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\} = \{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}A$$
 (1.18)

is a basis if and only if the transition matrix A has rank n, i.e. it is non-degenerate (invertible) matrix.

Recall that $n \times \text{matrix } A$ is nondegenerate (invertible) $\Leftrightarrow \det A \neq 0$.

Remark Recall that the condition that $n \times n$ matrix A is non-degenerate (has rank n) is equivalent to the condition that it is invertible matrix, or to the condition that $\det A \neq 0$.

Example let $\{\mathbf{e}_1, \mathbf{e}_2.\mathbf{e}_3\}$ be a basis in \mathbf{R}^3 . Consider set of vectors $\{\mathbf{e}_1, 3\mathbf{e}_1 + \lambda\mathbf{e}_2, 7\mathbf{e}_1 + 5\mathbf{e}_2 + 3\mathbf{e}_3\}$, where λ is an arboitrary parameter. The transition matrix from the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ to the row of vectors $\{\mathbf{e}_1, 3\mathbf{e}_1 + \lambda\mathbf{e}_2, 7\mathbf{e}_1 + 5\mathbf{e}_2 + 3\mathbf{e}_3\}$ is the following:

$$\{\mathbf{e}_1, 3\mathbf{e}_1 + \lambda \mathbf{e}_2, 7\mathbf{e}_1 + 5\mathbf{e}_2 + 3\mathbf{e}_3\} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}A = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\} \begin{pmatrix} 1 & 3 & 7 \\ 0 & \lambda & 5 \\ 0 & 0 & 3 \end{pmatrix}$$

We see that $\det A = 3\lambda$. In the case if $\mathbf{l} \neq 0$ then transition matrix is non-degenerate and the row $\{\mathbf{e}_1, 3\mathbf{e}_1 + \lambda\mathbf{e}_2, 7\mathbf{e}_1 + 5\mathbf{e}_2 + 3\mathbf{e}_3\}$ is a basis.

(See another examples in the Homework)

1.4 Orthonormal bases and orthogonal matirces

Now suppose that $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ is orthonoromal basis in *n*-dimensional Euclidean vector space. What is the condition that the new basis $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\} = \{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}A$ is an orthonormal basis too?

Definition We say that $n \times n$ matrix is orthogonal matrix if its product on transposed matrix is equal to unity matrix:

$$A^{T}A = I. (1.19)$$

Exercise. Prove that determinant of orthogonal matrix is equal to ± 1 :

$$A^{T} A = I \Rightarrow \det A = \pm 1. \tag{1.20}$$

Solution $A^TA = I$. Hence $\det(A^TA) = \det A^T \det A = (\det A)^2 = \det I = 1$. Hence $\det A = \pm 1$. We see that in particular orthogonal matrix is non-degenerate ($\det A \neq 0$). Hence it is a transition matrix from one basis to another. The following Proposition is valid:

Proposition 2 Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an orthonormal basis in n-dimensional Euclidean vector space. Then the new basis $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\} = \{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\} A$ is orthonormal basis if and only if the transition matrix A is orthogonal matrix.

Proof The basis $\{\mathbf{e}_1', \mathbf{e}_2', \dots, \mathbf{e}_n'\}$ is orthonormal means that $(\mathbf{e}_i', \mathbf{e}_j') = \delta_{ij}$. We have:

$$\delta_{ij} = (\mathbf{e}'_i, \mathbf{e}'_j) = \left(\sum_{m=1}^n \mathbf{e}_m A_{mi}, \mathbf{e}'_j = \sum_{n=1}^n \mathbf{e}_n A_{nj}\right) = \sum_{m,n=1}^n A_{mi} A_{nj}(\mathbf{e}_m, \mathbf{e}_n) =$$

$$\sum_{m,n=1}^{n} A_{mi} A_{nj} \delta_{mn} = \sum_{m=1}^{n} A_{mi} A_{mj} = \sum_{m=1}^{n} A_{im}^{T} A_{mj} = (A^{T} A)_{ij} \Rightarrow (A^{T} A)_{ij} = \delta_{ij}, \text{i.e. } A^{T} A = I.$$
(1.21)

Remark The set of orthogonal matrices form the group which is called O(n). This group is a subgroup of the group $GL(n, \mathbf{R})$ of linear invertible $n \times n$ matrices with real entries.

1.5 Linear operators.

1.5.1 Matrix of linear operator in a given basis

Recall here facts about linear operators in vector space

Let P be a linear operator in vector space V:

$$P: V \to V, \qquad P(\lambda \mathbf{x} + \mu \mathbf{y}) = \lambda P(\mathbf{x}) + \mu P(\mathbf{y}).$$

Let $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ be an arbitrary basis in *n*-dimensional vector space V. Consider the action of operator P on basis vectors: $\mathbf{e}'_i = P(\mathbf{e}_i)$. We denote by $p_{1k}, p_{2k}, \dots, p_{nk}$ coordinates of vector \mathbf{e}'_k in the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$:

$$\mathbf{e}'_{i} = P(\mathbf{e}_{i}) = \sum_{i} \mathbf{e}_{k} p_{ki},$$

$$\mathbf{e}'_{1} = P(\mathbf{e}_{1}) = \mathbf{e}_{1} p_{11} + \mathbf{e}_{2} p_{21} + \mathbf{e}_{3} p_{31} + \dots + \mathbf{e}_{n} p_{n1}$$

$$\mathbf{e}'_{2} = P(\mathbf{e}_{2}) = \mathbf{e}_{1} p_{12} + \mathbf{e}_{2} p_{22} + \mathbf{e}_{3} p_{32} + \dots + \mathbf{e}_{n} p_{n2}$$

$$\mathbf{e}'_{3} = P(\mathbf{e}_{3}) = \mathbf{e}_{1} p_{13} + \mathbf{e}_{2} p_{23} + \mathbf{e}_{3} p_{31} + \dots + \mathbf{e}_{n} p_{n3}$$

$$\dots$$

$$\mathbf{e}'_{n} = P(\mathbf{e}_{n}) = \mathbf{e}_{1} p_{1n} + \mathbf{e}_{2} p_{2n} + \mathbf{e}_{3} p_{3n} + \dots + \mathbf{e}_{n} p_{nn}$$

$$(1.22)$$

Definition-Proposition Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an arbitrary basis in n-dimensional vector space V, and let P be a linear operator in V. Then matrix $P = ||p_{ik}||$ in equation (1.22) is a matrix of linear transformation P in the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$. This matrix coincides with the transition matrix from the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ to the row of vectors $\{\mathbf{e}'_1, \mathbf{e}_2, \dots, \mathbf{e}'_n\}$.

In the case if linear operator P is non-degenerate (invertible) then vectors $\mathbf{e}'_1, \mathbf{e}'_2, \mathbf{e}'_3, \dots, \mathbf{e}'_n$, form also a basis.

Does matrix of linear operator change if we change the basis? See it:

Consider a new basis $\{\mathbf{f}_1, \dots, \mathbf{f}_n\}$ in the linear space V. Let A be transition matrix from the basis $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ to the new basis $\{\mathbf{f}_1, \dots, \mathbf{f}_n\}$:

$$\{\mathbf{f}_1, \dots, \mathbf{f}_n\} = \{\mathbf{e}_1, \dots, \mathbf{e}_n\} A$$
, i.e. $\mathbf{f}_i = \sum_{k=1}^n \mathbf{e}_k a_{ki}$ (1.23)

(see equation (1.17)). Find matrix for linear operator P considered above in (1.22), in the new basis $\{\mathbf{f}_i\}$. According to the formulae (1.23) and (1.22) we have

$$\mathbf{f}_{i}' = P(\mathbf{f}_{i}) = P\left(\sum_{q=1}^{n} \mathbf{e}_{q} a_{qi}\right) = \sum_{q=1}^{n} a_{qi} P\left(\mathbf{e}_{q}\right) = \sum_{q=1}^{n} a_{qi} \left(\sum_{r=1}^{n} \mathbf{e}_{r} p_{rq}\right) = \sum_{q,r=1}^{n} \mathbf{e}_{r} p_{rq} a_{qi} = \sum_{q=1}^{n} a_{qi} P\left(\mathbf{e}_{q}\right) = \sum$$

$$\sum_{r=1}^{n} \mathbf{e}_r (PA)_{ri} = \sum_{r,k=1}^{n} \mathbf{f}_k (A^{-1})_{kr=1} (PA)_{ri} = \sum_{k=1}^{n} \mathbf{f}_k (A^{-1}PA)_{ki}.$$

We see that in the new basis $\{\mathbf{f}_i\}$ a matrix of linear operator is equal to $A^{-1}PA$.

Proposition Let P be a linear operator acting in n-dimensional vector space V. Let $\{\mathbf{e}_i\}$ and $\{\mathbf{f}_j\}$ be two arbitrary bases in V. Let $P = ||p_{ik}||$ be a matrix of the operator P in the basis $\{\mathbf{e}_i\}$, and let $P' = ||p'_{ik}||$ be a matrix of the operator P in the basis $\{\mathbf{f}_i\}$:

basis
$$\{\mathbf{e}_i\}$$
 in $V - - - - ||p_{ik}||$ matrix of operator P in the basis $\{\mathbf{e}_i\}$

basis
$$\{\mathbf{f}_i\}$$
 in $V - - - - ||p'_{ik}||$ matrix of operator P in the basis $\{\mathbf{f}_i\}$

Then

$$p'_{ik} = (A^{-1} \circ P \circ A)_{ik} = \sum_{m,r=1}^{n} a_{im} p_{mr} a_{rk}.$$
 (1.24)

Remark Let a matrix $||p_{ij}||$ be a matrix of linear operator P in the basis $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$. Then for an arbitrary vector \mathbf{x}

$$\forall \mathbf{x} = \sum_{i=1}^{n} \mathbf{e}_{i} x^{i} = (\mathbf{e}_{1}, \mathbf{e}_{2}, \dots, \mathbf{e}_{n}) \cdot \begin{pmatrix} x^{1} \\ x^{2} \\ \dots \\ x^{n} \end{pmatrix}, \text{ then }$$

$$P(\mathbf{x}) = (\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n) \cdot P \cdot \begin{pmatrix} x^1 \\ x^2 \\ \dots \\ x^n \end{pmatrix} = \sum_{i=1}^n \mathbf{e}_i' y^i = \sum_{i,k=1}^n \mathbf{e}_k p_{ki} x^i.$$

If x^i are components of vector \mathbf{x} at the basis $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ and x'^i are components of the vector \mathbf{x} at the new basis $\{\mathbf{e}'_i\}$ then $x'^i = \sum_{k=1}^n p_{ik} x^k$.

1.5.2 Determinant and Trace of linear operator

We recall the definition of determinant and explain what is the trace of linear operator,

Definition-Proposition Let P be a linear operator in vector space V, let $\{\mathbf{e}_i\}$ be an arbitrary basis in V, and let $||p_{ik}||$ be a matrix of operator P

in this basis. Then we define determinant of linear operator as a determinant of its matrix:

$$\det P = \det (||p_{ik}||) ,$$

and in the same way we define we define trace of operator via trace of matrix:

$$\operatorname{Tr} P = \operatorname{Tr} (||p_{ik}||) = p_{11} + p_{22} + p_{33} + \dots + p_{nn}. \tag{1.25}$$

Determinant and trace of operator are well-defined. since due to the proposition above (see equation (1.24)), determinant and trace of transition matrice do not change if we change the basis in spite of the fact that transition matrix changes: $P \mapsto A^{-1}PA$, but

$$\det (A^{-1}PA) = \det A^{-1} \det P \det A = (\det A)^{-1} \det P \det A = \det P,$$

and

$$\operatorname{Tr}(A^{-1}PA) = \sum_{i} (A^{-1}PA)_{ii} = \sum_{i,k,p} (A^{-1})_{ik} p_{kp} A_{pi} = \sum_{i,k,p} A_{pi} (A^{-1})_{ik} p_{kp} = \sum_{i,k,p} (A \cdot A^{-1})_{pk} p_{kp} = \sum_{i,k} \delta_{kp} p_{kp} = \sum_{i} p_{kk} = \operatorname{Tr} P.$$

Trace of linear operator is an infinitesimal version of its determinant:

$$\det(1+tP) = 1 + t\operatorname{Tr} P + O(t^2).$$

This is infinitesimal version for the following famous formula which relates trace and det of linear operator:

$$\det e^{tA} = e^{t\operatorname{Tr} A}. \tag{1.26}$$

where
$$e^{tA} = \sum \frac{t^n A^n}{n!}$$
. E.g. if $A = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$, then $e^{tA} = \begin{pmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{pmatrix}$, $\det e^{tA} = 1$ and $e^{t\text{Tr } A} = e^0 = 1$.

1.5.3 Orthogonal linear operators

Now two words on orthogonal linear operators in Euclidean space.

Recall that linear operator P in Euclidean space \mathbf{E}^n is called orthogonal operator if it preserves scalar product:

$$(P\mathbf{x}, P\mathbf{y}) = (\mathbf{x}, \mathbf{y}), \text{ for arbitrary vectors } \mathbf{x}, \mathbf{y}$$
 (1.27)

In particular if $\{\mathbf{e}_i\}$ is orthonormal basis in Euclidean space then due to (1.27) the new basis $\{\mathbf{e}'_i = P(\mathbf{e}_i)\}$ is orthonormal too. Thus we see that

matrix of orthogonal operator P in a given orthogonal basis is orthogonal matrix:

$$P^T \cdot P = I \tag{1.28}$$

(see (1.19) in subsection 1.7). In particular we see that for orthogonal linear operator det $P = \pm 1$ (compare with (1.20)).

1.6 Orientation in vector space

You have heard a words 'orientation', you have heard expressions like:

A basis $\{a, b, c\}$ have the same orientation as the basis $\{a', b', c'\}$ if they both obey right hand rule or if they both obey left hand rule. In the other case we say that these bases have opposite orientation...

When you look in the mirror you know that 'left' is changing on the 'right' Try to give the exact meaning to these expressions.

1.6.1 Orientation in vector space. Oriented vector space

Consider the set of all bases in the given vector space V.

Let $(\mathbf{e}_1, \dots \mathbf{e}_n)$, $(\mathbf{e}'_1, \dots \mathbf{e}'_n)$ be two arbitrary bases in the vector space V and let T be transition matrix which transforms the basis $\{\mathbf{e}_i\}$ to the new basis $\{\mathbf{e}'_i\}$:

$$\{\mathbf{e}'_1, \dots \mathbf{e}'_n\} = \{\mathbf{e}_1, \dots \mathbf{e}_n\}T, \qquad (\mathbf{e}'_i = \sum_{k=1}^n \mathbf{e}_k t_{ki})$$
 (1.29)

(see also (1.16)).

Definition We say that two bases $\{\mathbf{e}_1, \dots \mathbf{e}_n\}$ and $\{\mathbf{e}'_1, \dots \mathbf{e}'_n\}$ in V have the same orientation if the determinant of transition matrix (1.29) from the first basis to the second one is positive: $\det T > 0$.

We say that the basis $\{\mathbf{e}_1, \dots \mathbf{e}_n\}$ has an orientation opposite to the orientation of the basis $\{\mathbf{e}'_1, \dots \mathbf{e}'_n\}$ (or in other words these two bases have opposite orientation) if the determinant of transition matrix from the first basis to the second one is negative: det T < 0.

Remark Transition matrix from basis to basis is non-degenerate, hence its determinant cannot be equal to zero. It can be or positive or negative.

Consider examples.

First the simplest example.

Example 0 Consider a line $\mathbf{R} = \mathbf{R}^1$ as 1-dimensional vector space with an origin at the point 0. Consider on \mathbf{R}^1 vectors

$$e = (2), e' = (-8), \tilde{e} = (10).$$

Vector \mathbf{e} is a basis of \mathbf{R} , as well as vector \mathbf{e}' is a basis, and vector $\tilde{\mathbf{e}}$ is a basis also. (Since space is 1-dimensional every non-zero vector is a basis!)

The basis $\{\mathbf{e}\}$ and the basis $\{\tilde{\mathbf{e}}\}$ have the same orientation since $\tilde{\mathbf{e}} = 5 \cdot \mathbf{e}$: transition matrix is 1×1 matrix, the determinant of transition matrix is equal to 5 and 5 > 0.

Respectively the basis $\{\mathbf{e}'\}$ and the basis $\{\mathbf{e}'\}$ have the opposite orientation since $\mathbf{e}' = -4 \cdot \mathbf{e}$: determinant of transition matrix is equal to -4 and -4 < 0.

Now example of 2-dimensional space:

Example 1 Consider two dimensional vector space \mathbb{R}^2 with a canonical basis

$$\{\mathbf{e}_1, \mathbf{e}_2\} = \left\{ \begin{pmatrix} 1\\0 \end{pmatrix}, \begin{pmatrix} 0\\1 \end{pmatrix} \right\}$$

Consider in \mathbb{R}^2 another basis

$$\{\mathbf{e}_1', \mathbf{e}_2'\} = \left\{ \begin{pmatrix} -2\\0 \end{pmatrix}, \begin{pmatrix} 0\\1 \end{pmatrix} \right\}$$

One can see that $\{\mathbf{e}_1', \mathbf{e}_2'\} = \{-2\mathbf{e}_1, \mathbf{e}_2\}$, transition matrix $T = \begin{pmatrix} -2 & 0 \\ 0 & 1 \end{pmatrix}$, and $\det T = -2 < 0$, i.e. bases $\{\mathbf{e}_1', \mathbf{e}_2'\}$ and $\{-2\mathbf{e}_1, \mathbf{e}_2\}$ have opposite orientation.

One can see that orientation establishes the equivalence relation in the set of all bases. Show it. We say that $\{\mathbf{e}_1, \dots \mathbf{e}_n\} \sim \{\mathbf{e}'_1, \dots \mathbf{e}'_n\}$, if two bases $\{\mathbf{e}_1, \dots \mathbf{e}_n\}$ and $\{\mathbf{e}'_1, \dots \mathbf{e}'_n\}$ have the same orientation, i.e. $\det T > 0$ for transition matrix.

Proposition Relation " \sim " is an equivalence relation, i.e. this relation is reflexive, symmetric and transitive.

Prove it:

• Proof of reflexivity

it is reflexive, i.e. for every basis $\{e_1, \dots e_n\}$

$$\{\mathbf{e}_1, \dots, \mathbf{e}_n\} \sim \{\mathbf{e}_1, \dots, \mathbf{e}_n\},$$
 (1.30)

because in this case transition matrix T = I and detI = 1 > 0.

• Proof of simmetricity

Prove, that relation " \sim " is symmetric, i.e. If $\{\mathbf{e}_1, \dots, \mathbf{e}_n\} \sim \{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}$ then $\{\mathbf{e}'_1, \dots, \mathbf{e}'_n\} \sim \{\mathbf{e}_1, \dots, \mathbf{e}_n\}$.

Let T be a transition matrix from the first basis $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}$ to the second basis $\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}$: $\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}=\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}T$, and $\det T>0$ since $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}\sim\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}$. Then the transition matrix from the second basis $\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}$ to the first basis $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}$ is the inverse matrix T^{-1} : $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}=\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}T^{-1}$. Hence $\det T^{-1}=\frac{1}{\det T}>0$ since $\det T>0$. Hence $\{\mathbf{e}_1',\ldots,\mathbf{e}_n'\}\sim\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}$. Symmetricity is proved.

• Proof of transitivity

We have to prove that if $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\} \sim \{\mathbf{e}'_1,\ldots,\mathbf{e}'_n\}$ and $\{\mathbf{e}'_1,\ldots,\mathbf{e}'_n\} \sim \{\tilde{\mathbf{e}}_1,\ldots,\tilde{\mathbf{e}}_n\}$, then $\{\mathbf{e}_1,\ldots,\mathbf{e}_n\} \sim \{\tilde{\mathbf{e}}_1,\ldots,\tilde{\mathbf{e}}_n\}$.

Do it in detail.

Formulate the following statement:

Proposition-Lemma Let $\{\mathbf{e}_i\}$, $\{\mathbf{e}'_i\}$ and $\{\tilde{\mathbf{e}}_i\}$ be arbitrary three bases in the vector space V. For convenience call a basis $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ the 'I-st' basis, call a basis $\{\mathbf{e}'_1, \ldots, \mathbf{e}'_n\}$ the 'II-nd' basis and call a basis $\{\tilde{\mathbf{e}}_1, \ldots, \tilde{\mathbf{e}}_n\}$ the 'III-rd' basis.

Let $T^{(12)}$ be a transition matrix from the I-st basis to the II-nd basis; let $T^{(13)}$ be a transition matrix from the I-st basis to the III-rd basis, and let $T^{(23)}$ be a transition matrix from the II-nd basis to the III-rd basis:

$$\begin{aligned}
\{\mathbf{e}'_{1}, \dots, \mathbf{e}'_{n}\} &= \{\mathbf{e}_{1}, \dots, \mathbf{e}_{n}\} T^{(12)} \\
\{\tilde{\mathbf{e}}_{1}, \dots, \tilde{\mathbf{e}}_{n}\} &= \{\mathbf{e}_{1}, \dots, \mathbf{e}_{n}\} T^{(13)} \\
\{\tilde{\mathbf{e}}_{1}, \dots, \tilde{\mathbf{e}}_{n}\} &= \{\mathbf{e}'_{1}, \dots, \mathbf{e}'_{n}\} T^{(23)}.
\end{aligned} (1.31)$$

Then

$$\underbrace{T^{(13)}}_{I-st \to III-rd} = \underbrace{T^{(12)}}_{I-st \to II-nd} \circ \underbrace{T^{(23)}}_{II-rd} \Rightarrow
\det T^{(13)} = \det(T^{(12)} \circ T^{(23)}) = \det T^{(12)} \cdot \det T^{(23)}.$$
(1.32)

Transitivity immediately follows from this statement: if I-st \sim II and II-nd \sim III-rd, then determinants of matrices $T^{(12)}$ and $T^{(23)}$ are positive. Hence according to relation (1.32) det $T^{(13)}$ is positive too, i.e. I-st \sim III-rd.

It remains to prove equation (1.32). This equation follows from equation (1.31): $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\} = \{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}T^{(23)} =$

$$(\{\mathbf{e}_1,\ldots,\mathbf{e}_n\}T^{(12)})T^{(23)} = \{\mathbf{e}_1,\ldots,\mathbf{e}_n\}T^{(12)} \circ T^{(23)} = \{\mathbf{e}_1,\ldots,\mathbf{e}_n\}T^{(13)}.$$

Thus we proved that relation \sim is equivalence realtion.

Since it is equivalence relation the set of all bases is a union if disjoint equivalence classes. Two bases are in the same equivalence class if and only if they have the same orientation.

How many equivalence classes exist? One, two or more?

Show first that there are at least two equivalence classes.

Example Let $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ be an arbitrary basis in n-dimensional vector space V. Swap the vectors $\mathbf{e}_1, \mathbf{e}_2$. We come to a new basis: $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$

$$\mathbf{e}_1' = \mathbf{e}_2, \mathbf{e}_2' = \mathbf{e}_1,$$
 all other vectors are the same: $\mathbf{e}_3 = \mathbf{e}_3', \dots, \mathbf{e}_n = \mathbf{e}_n'$ (1.33)

We have:

$$\{\mathbf{e}'_1, \mathbf{e}'_2, \mathbf{e}'_3, \dots, \mathbf{e}'_n\} = \{\mathbf{e}_2, \mathbf{e}_1, \mathbf{e}_3, \dots, \mathbf{e}_n\} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3, \dots, \mathbf{e}_n\} T_{\text{swap}}, \quad (1.34)$$

where one can easy see that the determinant for transition matrix T^{swap} is equal to -1, i.e. bases $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ and $\{\mathbf{e}_2, \mathbf{e}_1, \dots, \mathbf{e}_n\}$ have opposite orientation.

E.g. write down the transition matrix (1.34) in the case if dimension of vector space is equal to 5, n=5. Then we have $\{\mathbf{e}_1', \mathbf{e}_2', \mathbf{e}_3', \mathbf{e}_4', \mathbf{e}_5'\} = \{\mathbf{e}_2, \mathbf{e}_1, \mathbf{e}_3, \mathbf{e}_4, \mathbf{e}_5\} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3, \mathbf{e}_4, \mathbf{e}_5\}T$ where

$$T_{\text{swap}} = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \qquad (\det T_{\text{swap}} = -1). \tag{1.35}$$

We see that bases $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ and $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ have opposite orientation.

We see that there are at least two equivalence classes.

One can see that there are exactly two equivalence classes.

Proposition Let two bases $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ and $\{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}$ in vector space V have opposite orientation. Let $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\}$ be an arbitrary basis in V.

Then the basis $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\}$ and the first basis $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ have the same orientation or the basis $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\}$ and the second basis $\{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}$ have the same orientation.

In other words if bases $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$, $\{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}$ and $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\}$ are three bases in vector space V such that $\{\mathbf{e}_1, \dots, \mathbf{e}_n\} \not\sim \{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}$ then

$$\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\} \sim \{\mathbf{e}_1, \dots, \mathbf{e}_n\} \text{ or } \{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\} \sim \{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}.$$
 (1.36)

There are two equivalence classes of bases with respect to orientation.

In the case if bases $\{\tilde{\mathbf{e}}'_1, \dots, \tilde{\mathbf{e}}'_n\}$, $\{\tilde{\mathbf{e}}_1, \dots, \tilde{\mathbf{e}}_n\}$ have opposite orientation, then an arbitrary basis belongs to the equivalence class of the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$, or it belongs to the to the equivalence class of the basis $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$.

Proof of the statement immediately follows from statement (1.32).

In the same way like in statement (1.32) we call a basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$ the "I-st basis", a basis $\{\mathbf{e}'_1, \mathbf{e}'_2, \dots, \mathbf{e}'_n\}$ the "III-rd basis" and a basis $\{\tilde{\mathbf{e}}_1, \tilde{\mathbf{e}}_2, \dots, \tilde{\mathbf{e}}_n\}$ the "III-rd basis". We have to prove that the third basis has the same orientation as the first basis or it has the same orientation as the second basis.

Suppose the third basis has not the same orientation as the first basis, then for the transition matrix $T^{(13)}$ (see equation (1.31)) det $T^{(13)} < 0$. On the other hand det $T^{(12)} < 0$ also since the first and second bases have opposite orientation. Hence it follows from equation (1.32) that det $T^{(23)} < 0$, thus second and third bases have opposite orientation.

In the example considered above (see (1.33)) an arbitrary basis $\{\mathbf{e}'_1, \dots \mathbf{e}'_n\}$ have the same orientation as the basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$, i.e. belongs to the equivalence class of basis $\{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$, or it has the same orientation as the "swapped" basis $\{\mathbf{e}_2, \mathbf{e}_1, \dots, \mathbf{e}_n\}$, i.e. it belongs to the equivalence class of the "swappedd" basis $\{\mathbf{e}_2, \mathbf{e}_1, \dots, \mathbf{e}_n\}$.

The set of all bases is a union of two disjoint subsets.

Any two bases which belong to the same subset have the same orientation. Any two bases which belong to different subsets have opposite orientation.

Definition An orientation of a vector space is an equivalence class of bases in this vector space.

Note that fixing any basis we fix orientation, considering the subset of all bases which have the same orientation that the given basis.

There are two orientations. Every basis has the same orientation as a given basis or orientation opposite to the orientation of the given basis.

We choose an arbitrary basis, and call it 'left' basis. Then all bases which belong to the equivalence class of this basis may be called "left" bases and all the bases which do not belong to the equivalence class of this basis may be called "right" bases

Sure we could call this arbitrary basis "right" basis, (or any other term, this is just problem of consensus), then all the bases belonging to the equivalence class of this basis would be called by the same term.

Definition An oriented vector space is a vector space equipped with orientation.

Consider examples.

Example (Orientation in two-dimensional space). Let $\{\mathbf{e}_x, \mathbf{e}_y\}$ be arbitrary two bases in \mathbf{R}^2 and let \mathbf{a}, \mathbf{b} be arbitrary two vectors in \mathbf{R}^2 . Consider an ordered pair $\{\mathbf{a}, \mathbf{b}, \}$. The transition matrix from the basis $\{\mathbf{e}_x, \mathbf{e}_y\}$ to the ordered pair $\{\mathbf{a}, \mathbf{b}\}$ is $T = \begin{pmatrix} a_x & b_x \\ a_y & b_y \end{pmatrix}$:

$$\{\mathbf{a}, \mathbf{b}\} = \{\mathbf{e}_x, \mathbf{e}_y\}T = \{\mathbf{e}_x, \mathbf{e}_y\} \begin{pmatrix} a_x & b_x \\ a_y & b_y \end{pmatrix}, \quad \begin{cases} \mathbf{a} = a_x \mathbf{e}_x + a_y \mathbf{e}_y \\ \mathbf{b} = b_x \mathbf{e}_x + b_y \mathbf{e}_y \end{cases}$$

One can see that the ordered pair $\{\mathbf{a}, \mathbf{b}\}$ also is a basis, (i.e. these two vectors are linearly independent in \mathbf{R}^2) if and only if transition matrix is not degenerate, i.e. $\det T \neq 0$. The basis $\{\mathbf{a}, \mathbf{b}\}$ has the same orientation as the basis $\{\mathbf{e}_x, \mathbf{e}_y\}$ if $\det T > 0$ and the basis $\{\mathbf{a}, \mathbf{b}\}$ has the orientation opposite to the orientation of the basis $\{\mathbf{e}_x, \mathbf{e}_y\}$ if $\det T < 0$.

If we call the basis $\{\mathbf{e}_x, \mathbf{e}_y\}$ left basis then the basis $\{\mathbf{a}, \mathbf{b}\}$ will be called also left basis in the case if $\det T > 0$, and the basis $\{\mathbf{a}, \mathbf{b}\}$ will be called right basis in the case if $\det T < 0$; respectively if we call the basis $\{\mathbf{e}_x, \mathbf{e}_y\}$ right basis then the basis $\{\mathbf{a}, \mathbf{b}\}$ will be called also right basis in the case if $\det T > 0$, and the basis $\{\mathbf{a}, \mathbf{b}\}$ will be called left basis in the case if $\det T < 0$.

Example Let $\{e, f\}$ be a basis in 2-dimensional vector space. Consider bases $\{e, -f\}$, $\{f, -e\}$ and $\{f, e\}$.

1) We come to basis $\{\mathbf{e}, -\mathbf{f}\}$ reflecting the second basis vector. Transition matrix from initial basis $\{\mathbf{e}, \mathbf{f}\}$ to the basis $\{\mathbf{e}, -\mathbf{f}\}$ is $T_{\{\mathbf{e}, -\mathbf{f}\}} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. Its determinant is -1. Bases $\{\mathbf{e}, \mathbf{f}\}$ and $\{\mathbf{e}, -\mathbf{f}\}$ have opposite orientation. If $\{\mathbf{e}, \mathbf{f}\}$ is left basis then $\{\mathbf{e}, -\mathbf{f}\}$ is right basis and vice versa.

- 2) Transition matrix from initial basis $\{\mathbf{e}, \mathbf{f}\}$ to the basis $\{\mathbf{f}, -\mathbf{e}\}$ is $T_{\{\mathbf{f}, -\mathbf{e}\}} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$. Its determinant is 1. Bases $\{\mathbf{e}, \mathbf{f}\}$ and $\{\mathbf{f}, -\mathbf{e}\}$ have same orientation. They both are left bases or they both are right bases. Note that we come to basis $\{\mathbf{f}, -\mathbf{e}\}$ rotating the initial basis (on the angle $\pi/2$).
- 3) Transition matrix from initial basis $\{\mathbf{e}, \mathbf{f}\}$ to the basis $\{\mathbf{f}, \mathbf{e}\}$ is $T_{\{\mathbf{f}, \mathbf{e}\}} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. Its determinant is -1. Bases $\{\mathbf{e}, \mathbf{f}\}$ and $\{\mathbf{e}, -\mathbf{f}\}$ have opposite orientation. Basis $\{\mathbf{e}, -\mathbf{f}\}$ is right basis in the case if basis $\{\mathbf{e}, \mathbf{f}\}$ is left basis, and vice versa, Basis $\{\mathbf{e}, -\mathbf{f}\}$ is left basis in the case if basis $\{\mathbf{e}, \mathbf{f}\}$ is right basis.

Notice that we come to basis $\{\mathbf{f}, \mathbf{e}\}\ reflecting$ the initial basis.

(There are plenty exercises in the Homework 2.)

Example(Orientation in three-dimensional euclidean space.) Let $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ be any basis in \mathbf{E}^3 and $\mathbf{a}, \mathbf{b}, \mathbf{c}$ are arbitrary three vectors in \mathbf{E}^3 :

$$\mathbf{a} = a_x \mathbf{e}_x + a_y \mathbf{e}_y + a_z \mathbf{e}_z \ \mathbf{b} = b_x \mathbf{e}_x + b_y \mathbf{e}_y + b_z \mathbf{e}_z, \ \mathbf{c} = c_x \mathbf{e}_x + c_y \mathbf{e}_y + c_z \mathbf{e}_z.$$

Consider ordered triple $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$. The transition matrix from the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$

to the ordered triple $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ is $T = \begin{pmatrix} a_x & b_x & c_x \\ a_y & b_y & c_y \\ a_z & b_z & c_z \end{pmatrix}$:

$$\{\mathbf{a}, \mathbf{b}, \mathbf{c}\} = \{\mathbf{e}_{\mathbf{x}}, \mathbf{e}_{\mathbf{y}}, \mathbf{e}_{\mathbf{z}}\}\mathbf{T} = \{\mathbf{e}_{\mathbf{x}}, \mathbf{e}_{\mathbf{y}}, \mathbf{e}_{\mathbf{z}}\} \begin{pmatrix} a_x & b_x & c_x \\ a_y & b_y & c_y \\ a_z & b_z & c_z \end{pmatrix}$$

One can see that the ordered triple $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ also is a basis, (i.e. these three vectors are linearly independent) if and only if transition matrix is not degenerate det $T \neq 0$. The basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ has the same orientation as the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ if

$$\det T > 0. \tag{1.37}$$

The basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ has the orientation opposite to the orientation of the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ if

$$\det T < 0. \tag{1.38}$$

The usage of words "left" "right" is defined as always: if basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ is left basis, then basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ is also left if determinant of transition matrix is positive, and basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ is right if determinant of transition matrix is negative, and vice versa: if basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ is right basis, then basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ is also right if determinant of transition matrix is positive, and basis $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$ is left if determinant of transition matrix is negative.

Remark Note that in the example above we considered in \mathbf{E}^3 arbitrary bases not necessarily orthonormal bases.

I would like to emphasize again:

relations (1.37), (1.38) define equivalence relations in the set of bases. Orientation is equivalence class of bases. There are two orientations, every basis has the same orientation as a given basis or opposite orientation.

If two bases $\{\mathbf{e}_i\}$, $\{\mathbf{e}_{i'}\}$ have the same orientation then they can be transformed to each other by continuous transformation, i.e. there exists one-parametric family of bases $\{\mathbf{e}_i(t)\}$ such that $0 \le t \le 1$ and $\{\mathbf{e}_i(t)\}|_{t=0} = \{\mathbf{e}_i\}$, $\{\mathbf{e}_i(t)\}|_{t=1} = \{\mathbf{e}_{i'}\}$. (All functions $\mathbf{e}_i(t)$ are continuous) In the case of three-dimensional space the following statement is true: Let $\{\mathbf{e}_i\}$, $\{\mathbf{e}_{i'}\}$ (i=1,2,3) be two orthonormal bases in \mathbf{E}^3 which have the same orientation. Then there exists an axis \mathbf{n} such that basis $\{\mathbf{e}_i\}$ transforms to the basis $\{\mathbf{e}_{i'}\}$ under rotation around the axis. (This is Euler Theorem (see it later).

Exercise Show that bases $\{e, f, g\}$ and $\{f, e, g\}$ have opposite orientation but bases $\{e, f, g\}$ and $\{f, e, -g\}$ have the same orientation.

Solution. Transformation from basis $\{e, f, g\}$ to basis $\{f, e, g\}$ is "swapping" of vectors $((e, f) \mapsto (f, e)$. This is reflection and this transformation changes orientation. One can see it using transition matrix:

$$T: \{\mathbf{f}, \mathbf{e}, \mathbf{g}\} = \{\mathbf{e}, \mathbf{f}, \mathbf{g}\}T = \{\mathbf{e}, \mathbf{f}, \mathbf{g}\} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot \det T = -1$$

Transformation from basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ to basis $\{\mathbf{f}, \mathbf{e}, -\mathbf{g}\}$ is composition of two transformations: "swapping" of vectors $((\mathbf{e}, \mathbf{f}) \mapsto (\mathbf{f}, \mathbf{e}))$ and changing direction of vector $\mathbf{g}(\mathbf{g} \mapsto -\mathbf{g})$. We have two reflections:

$$\{e,f,g\} \quad \overset{\mathrm{reflection}}{\longrightarrow} \quad \{f,e,g\} \quad \overset{\mathrm{reflection}}{\longrightarrow} \quad \{f,e,-g\}$$

Any reflection changes orientation. Two reflections preserve orinetation. One may come to this result using transition matrix:

$$T: \{\mathbf{f}, \mathbf{e}, -\mathbf{g}\} = \{\mathbf{e}, \mathbf{f}, \mathbf{g}\} T = \{\mathbf{e}, \mathbf{f}, \mathbf{g}\} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} . \det T = 1. \text{ Orientation is not changed.}$$

$$(1.39)$$

(See also exercises in Homework 2)

1.6.2 Orientation of linear operator

. Let P be a linear operator acting in vector space V.

Let $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ be an arbitrary basis in V. Linear operator P transforms this basis to another basis $\{\mathbf{e}'_1, \ldots, \mathbf{e}'_n\}$ in the case if $\det P \neq 0$. Bearing in mind that determinant of transition matrix from basis $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ to the basis $\{\mathbf{e}'_1, \ldots, \mathbf{e}'_n\}$ is a matrix of operator P in the basis $\{\mathbf{e}_1, \ldots, \mathbf{e}_n\}$ we see that these both bases

$$\{\mathbf{e}_1, \dots, \mathbf{e}_n\}, \{\mathbf{e}'_1, \dots, \mathbf{e}'_n\}, \mathbf{e}'_i = P(\mathbf{e}_i)$$

have the same orientation if and only if $\det P > 0$ and they have opposite orientation if and only if $\det P < 0$.

In the case if $\det P = 0$, P is not invertible matrix, and it does not transform bases to bases.

If a linear operator P acting on the space V has positive determinant then under the action of this operator an arbitrary basis transforms to the basis with the same orientation. Respectively if a linear operator P acting on the space V has negative determinant then under the action of this operator an arbitrary basis transforms to the new basis which has opposite orientation.

Definition. Non-degenerate (invertible) linear operator P (det $P \neq 0$) acting in vector space V preserves an orientation of the vector space V if det P > 0. It changes the orientation if det P < 0.

1.7 Rotations and orthogonal operators preserving orientation of E^n (n=2,3)

Recall the notion of orthogonal operator (see 1.5.3). We study here orthogonal operators in \mathbf{E}^2 and \mathbf{E}^3 . In particular we will show that orthogonal operators preserving orientations define rotations.

1.7.1 Orthogonal operators in E²— Rotations and reflections

We show that an orthogonal operator in \mathbf{E}^2 'rotates the space' or makes a 'reflection'.

Let A be an orthogonal operator acting in Euclidean space \mathbf{E}^2 : $(A\mathbf{x}, A\mathbf{y}) = (\mathbf{x}, \mathbf{y})$. Let $\{\mathbf{e}, \mathbf{f}\}$ be an orthonormal basis in 2-dimensional Euclidean space \mathbf{E}^2 : $(\mathbf{e}, \mathbf{e}) = (\mathbf{f}, \mathbf{f}) = 1$ (i.e. $|\mathbf{e}| = |\mathbf{f}| = 1$) and $(\mathbf{e}, \mathbf{f}) = 0$ -vectors \mathbf{e}, \mathbf{f} have unit length and are orthogonal to each other.

Consider a new basis $\{\mathbf{e}', \mathbf{f}'\}$, an image of basis \mathbf{e}, \mathbf{f} under action of A: $\mathbf{e}' = A(\mathbf{e}), \mathbf{f}' = A(\mathbf{f})$. Let $\begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}$ be matrix of operator A in the basis \mathbf{e}, \mathbf{f} , (see equation (1.22) and definition after this equation):

$$\{\mathbf{e}', \mathbf{f}'\} = \{\mathbf{e}, \mathbf{f}\}A = \{\mathbf{e}, \mathbf{f}\}\begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}$$
, i.e. $\mathbf{e}' = \alpha \mathbf{e} + \gamma \mathbf{f}$, $\mathbf{f}' = \beta \mathbf{e} + \delta \mathbf{f}$

New basis is orthonormal basis also, $(\mathbf{e}', \mathbf{e}') = (\mathbf{f}', \mathbf{f}') = 1$, $(\mathbf{e}', \mathbf{f}') = 0$.

Operator A is orthogonal operator, and its matrix is orthogonal matrix:

$$A^{T}A = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}^{t} \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} = \begin{pmatrix} \alpha & \gamma \\ \beta & \delta \end{pmatrix} \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} = \begin{pmatrix} \alpha^{2} + \gamma^{2} & \alpha\beta + \gamma\delta \\ \alpha\beta + \gamma\delta & \beta^{2} + \delta^{2} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$
(1.40)

Remark With some abuse of notation, (if it is not a reason of confusion) we sometimes use the same letter for linear operator and the matrix of this operator in orthonormal basis.

Remark Note that condition (1.40) implies that det $A = \pm 1$.

We have $\alpha^2 + \gamma^2 = 1$, $\alpha\beta + \gamma\delta = 0$ and $\beta^2 + \delta^2 = 1$.

Hence one can choose angles $\varphi, \psi \colon 0 \leq 2\pi$ such that $\alpha = \cos \varphi, \gamma = \sin \varphi, \ \beta = \cos \psi, \ \delta = \sin \psi$. The condition $\alpha\beta + \gamma\delta = \text{means that}$

$$\cos\varphi\cos\psi + \sin\varphi\sin\psi = \cos(\varphi - \psi) = 0$$

We have

$$A = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} = \begin{pmatrix} \cos \varphi & \cos \psi \\ \sin \varphi & \sin \psi \end{pmatrix}, \text{ with } \cos(\varphi - \psi) = 0.$$

Condition $\cos(\varphi - \psi) = 0$ means $\psi - \varphi = \frac{\pi}{2} + \pi k \ (k = 0, \pm 1, \pm 2, ...)$

$$\begin{cases} \text{I-st case } \psi = \varphi + \frac{\pi}{2} + \pi m \ (m = 0, \pm 2, \pm 4 \dots), \text{ hence } \cos \psi = -\sin \varphi, \sin \psi = \cos \varphi \\ \text{II-nd case } \psi = \varphi + \frac{\pi}{2} + \pi k \ (m = \pm 1, \pm 3 \dots), \text{ hence } \cos \psi = \sin \varphi, \sin \psi = -\cos \varphi \end{cases}$$

In the I-st case $\cos \psi = -\sin \varphi$, $\sin \psi = \cos \psi$, and

$$A_{\varphi} = \begin{pmatrix} \cos \varphi & \cos \psi \\ \sin \varphi & \sin \psi \end{pmatrix} = \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix}, \quad \det A_{\varphi} = 1. \tag{1.41}$$

i.e. operator A preserves orientation.

In the II-nd case $\cos \psi = \sin \varphi$, $\sin \psi = -\cos \psi$, and

$$\tilde{A}_{\varphi} = \begin{pmatrix} \cos \varphi & \cos \psi \\ \sin \varphi & \sin \psi \end{pmatrix} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix}, \quad \det A_{\varphi} = -1. \tag{1.42}$$

i.e. operator A changes orientation.

In the first case matrix of operator A_{φ} is defined by the relation (1.41). In this case the new basis is:

$$(\mathbf{e}', \mathbf{f}') = (\mathbf{e}, \mathbf{f}) A_{\varphi} = (\mathbf{e}, \mathbf{f}) \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix}, \quad \mathbf{e}' = A_{\varphi}(\mathbf{e}) = \cos \varphi \, \mathbf{e} + \sin \varphi \, \mathbf{f}$$
$$\mathbf{f}' = A_{\varphi}(\mathbf{f}) = -\sin \varphi \, \mathbf{e} + \cos \varphi \, \mathbf{f}$$
$$(1.43)$$

For an arbitrary vector $\mathbf{x} = x\mathbf{e} + y\mathbf{f} \ \mathbf{x} \to A_{\varphi}(\mathbf{x}) = A_{\varphi}(x\mathbf{e} + y\mathbf{f}) = x'\mathbf{e} + y'\mathbf{f},$

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cos \varphi - y \sin \varphi \\ \sin \varphi + y \cos \varphi \end{pmatrix}. \tag{1.44}$$

Operator A_{φ} rotates basis vectors e,f and arbitrary vector x on an angle φ

In the second case a matrix of operator \tilde{A}_{φ} is defined by the relation (1.42). See how transforms the basis $\{\mathbf{e}, \mathbf{f}\}$ in this case. We have in analogy with (1.43) that in this case

$$(\tilde{\mathbf{e}}, \tilde{\mathbf{f}}') = (\mathbf{e}, \mathbf{f}) \tilde{A}_{\varphi} = (\mathbf{e}, \mathbf{f}) \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix}, \quad \tilde{\mathbf{e}} = \tilde{A}_{\varphi}(\mathbf{e}) = \cos \varphi \, \mathbf{e} + \sin \varphi \, \mathbf{f}$$

$$\tilde{\mathbf{f}} = \tilde{A}_{\varphi}(\mathbf{f}) = \sin \varphi \, \mathbf{e} - \cos \varphi \mathbf{f}$$

$$(1.45)$$

Comparing this equation with equation (1.43) we see that the difference between the basis $\{\tilde{e}, \tilde{f}\}$ in this equation with the basis $\{\mathbf{e}', \mathbf{f}'\}$ in equation rotation of basis on the angle is the following: the vectors \mathbf{e}' and $\tilde{\mathbf{e}}$ coincide, and vector $\tilde{\mathbf{f}} = -\mathbf{f}'$, i.e. these bases have opposite orientation.

One can see that

$$\tilde{A}_{\varphi} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix} = \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = A_{\varphi}R \tag{1.46}$$

where we denote by $R = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ a transition matrix from the basis $\{\mathbf{e}, \mathbf{f}\}$ to the basis $\{\mathbf{e}, -\mathbf{f}\}$ —"reflection"l.

We see that in the second case the orthogonal operator \tilde{A}_{φ} is composition of rotation and reflection: $\{\mathbf{e}, \mathbf{f}\} \stackrel{\tilde{A}_{\varphi} = A_{\varphi}R}{\longrightarrow} \{\tilde{\mathbf{e}}, \tilde{\mathbf{f}}\}$:

$$\{\mathbf{e}, \mathbf{f}\} \xrightarrow{A_{\varphi}} \{\mathbf{e}' = \cos \varphi \, \mathbf{e} + \sin \varphi \mathbf{f}, \mathbf{f}' = -\sin \varphi \, \mathbf{e} + \cos \varphi \, \mathbf{f}\} \xrightarrow{R} \{\tilde{\mathbf{e}} = \mathbf{e}', \tilde{\mathbf{f}} = -\mathbf{f}\}$$
 (1.47)

We come to proposition

Proposition. Let A be an arbitrary 2×2 orthogonal linear transformation, $A^TA = 1$, and in particularly det $A = \pm 1$. (As usual we consider matrix of orthogonal operator in the orthonormal basis.)

If det A=1 then there exists an angle $\varphi \in [0,2\pi)$ such that $A=A_{\varphi}$ is an operator which rotates basis vectors and any vector (1.41) on the angle φ .

If det A = -1 then there exists an angle $\varphi \in [0, 2\pi)$ such that $A = A_{\varphi}$ is a composition of rotation and reflection (see (1.47)).

Remark One can show that orthogonal operator \tilde{A}_{φ} is a reflection with respect to the axis which have the angle $\frac{\varphi}{2}$ with x-axis.

Consider just examples:

$$a)\varphi = 0, \ \tilde{A}_{\varphi} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \end{pmatrix} \mapsto \begin{pmatrix} \mathbf{e} \\ -\mathbf{f} \end{pmatrix}$$

(reflection with respect to x-axis)

$$b)\varphi = \pi, \ \tilde{A}_{\varphi} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \end{pmatrix} \mapsto \begin{pmatrix} -\mathbf{e} \\ \mathbf{f} \end{pmatrix}$$

(reflection with respect to y-axis)

$$b)\varphi = \frac{\pi}{2}, \ \tilde{A}_{\varphi} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \begin{pmatrix} \mathbf{e} \\ \mathbf{f} \end{pmatrix} \mapsto \begin{pmatrix} \mathbf{f} \\ \mathbf{e} \end{pmatrix}$$

(reflection with respect to axis y = x ("swapping" of basis vectors))

Try to do it in general case.

1.7.2 Orthogonal operators in E³ and rotations

We see in the previous paragraph that orthogonal operator preserving orientation of \mathbf{E}^2 is rotation operator. The same is true in \mathbf{E}^2 . The main result of this paragraph will be the Euler Theorem about rotation, that every orthogonal operator preserving orientation in \mathbf{E}^3 is rotation around some axis.

We will give an exact formulation of the Euler Theorem at the end of this paragraph. Now we will formulate just preliminary statement:

The Euler Theorem. (Preliminary statement) An orthogonal operator in \mathbf{E}^3 preserving orientation is rotation operator with respect to an axis l on the angle φ . The axis is directed along eigenvector \mathbf{N} of the operator P, $P(\mathbf{N}) = \mathbf{N}$, and angle of rotation is defined by equation

$$Tr P = 1 + 2\cos\varphi.$$

We will come to this statement gradually step by step, and then will formulate it completely.

Let \mathbf{E}^n be oriented vector space. Recall that oriented vector space means that it is chosen the equivalence class of bases: all bases in this class have the same orientation. We call all bases in the equivalence class defining orientation "left" bases. All "left" bases have the same orientation. To define an orientation in vector space V one may consider an arbitrary basis $\{\mathbf{e}_i^{(0)}\}$ in V and claim that this basis is "left" basis. The basis $\{\mathbf{e}_i^{(0)}\}$ defines equivalence class of "left" bases: all bases $\{\mathbf{e}_i\}$ such that $\{\mathbf{e}_i\} \sim \{\mathbf{e}_i^{(0)}\}$ will be called "left" bases. We can say that basis $\{\mathbf{e}_i^{(0)}\}$ defines the orientation.

Later on considering oriented vector space we often call all bases defining the orientation (i.e. belonging to the equivalence class of bases defining orientation) "left" bases.

Now we define rotation in \mathbf{E}^3 . Recall the definition of rotation in \mathbf{E}^2 (see 1.7.1):

Definition Let \mathbf{E}^2 be an oriented Euclidean space. We say that linear operator P rotates this space on an angle " φ " if for a given "left" orthonormal basis $\{\mathbf{e}, \mathbf{f}\}$

$$\begin{cases} \mathbf{e}' = P(\mathbf{e}) = \mathbf{e}\cos\varphi + \mathbf{f}\sin\varphi \\ \mathbf{f}' = P(\mathbf{f}) = -\mathbf{e}\sin\varphi + \mathbf{f}\cos\varphi \end{cases}$$
 i.e. $\{\mathbf{e}', \mathbf{f}'\} = \{\mathbf{e}, \mathbf{f}\} \begin{pmatrix} \cos\varphi & -\sin\varphi \\ \sin\varphi & \cos\varphi \end{pmatrix}$ (1.48)

i.e. transition matrix from basis $\{\mathbf{e}, \mathbf{f}\}$ to new basis $\{\mathbf{e}' = P(\mathbf{e}), \mathbf{f}' = P(\mathbf{f})\}$ is the rotation matrix (1.41) (see also (1.43)).

Remark One can show that the angle of rotation does not depend on the choice of "left" basis. If we will choose another left basis $\tilde{\mathbf{e}}, \tilde{\mathbf{f}}$ then the angle remains the same

Operator P rotates every vector rotates on the angle φ .

If we choose a basis with opposite orientation ("right" basis) then the angle will change: $\varphi \mapsto -\varphi$.

We already did it in 1.7.1 and we also see from formula (1.48) that the matrix of operator P is orthogonal matrix such that its determinant equals 1. In 2-dimensional case we came to simple Proposition (see Proposition in 1.7.1) which we will repeat again⁴:

Proposition Let P be an orthogonal operator in oriented 2-dimensional Euclidean space. If operator P preserves orientation (det P = 1) then it is a rotation operator (1.48) on some angle φ .

The situation is little bit more tricky in 3-dimensional case.

Let \mathbf{E}^3 be an Euclidean vector space. (Problem of orientation will be discussed latter.) Let $\mathbf{N} \neq 0$ be an arbitrary non-zero vector in \mathbf{E}^3 . Consider the line $l_{\mathbf{N}}$, spanned by vector \mathbf{N} . This is *axis* directed along the vector \mathbf{N} . Choose a unit vector

$$\mathbf{n} = \pm \frac{\mathbf{N}}{|\mathbf{N}|} \tag{1.49}$$

Vector **n** fixes an orientation on $l_{\mathbf{N}}$. Changing $\mathbf{n} \mapsto -\mathbf{n}$ changes an orientation on opposite).

Choose an arbitrary orthonormal basis such that first vector of this basis is directed along the axis: a basis $\{n, f, g\}$.

Definition We say that a linear operator P rotates the Euclidean space \mathbf{E}^3 on the angle φ with respect to an axis $l_{\mathbf{N}}$ directed along a vector \mathbf{N} if the following conditions are satisfied:

 $P(\mathbf{N}) = \mathbf{N}$

vector \mathbf{N} (and all vectors proportional to this vector) are eigenvectors of operator P with eigenvalue 1, i.e. axis remain intact

• for an orthonormal basis $\{n, f, g\}$ such that the first vector of this basis is equal to n, (n is a unit vector, proportional to N)

$$\begin{cases} \mathbf{f}' = P(\mathbf{f}) = \mathbf{f}\cos\varphi + \mathbf{g}\sin\varphi \\ \mathbf{g}' = P(\mathbf{g}) = -\mathbf{f}\sin\varphi + \mathbf{g}\cos\varphi \end{cases}$$
 i.e.
$$\{\mathbf{f}', \mathbf{g}'\} = \{\mathbf{f}, \mathbf{g}\} \begin{pmatrix} \cos\varphi & -\sin\varphi \\ \sin\varphi & \cos\varphi \end{pmatrix}.$$
 (1.50)

 $^{^4}$ Just here we denote the operator by letter 'P' instead letter 'A'

In other words plane (subspace) orthogonal to axis rotates on the angle φ : linear operator P rotates every vector orthogonal to axis on the angle φ in the plane (subspace) orthogonal to the axis.

Linear operator P transforms the basis $\{\mathbf{n}, \mathbf{f}, \mathbf{g}\}$ to the new basis $\{\mathbf{n}, \mathbf{f}', \mathbf{g}'\}$ = $\{\mathbf{n}, \mathbf{f} \cos \varphi + \mathbf{g} \sin \varphi, -\mathbf{f} \sin \varphi + \mathbf{g} \cos \varphi\}$. The matrix of operator P, i.e. the transition matrix from the basis $\{\mathbf{n}, \mathbf{f}', \mathbf{g}'\}$ to the basis $\{\mathbf{n}, \mathbf{f}', \mathbf{g}'\}$ is defined by the relation:

$$\{\mathbf{n}, \mathbf{f}', \mathbf{g}'\} = \{\mathbf{n}, \mathbf{f}\cos\varphi + \mathbf{g}\sin\varphi, -\mathbf{f}\sin\varphi + \mathbf{g}\cos\varphi\} = \{\mathbf{n}, \mathbf{f}, \mathbf{g}\} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\varphi & -\sin\varphi \\ 0 & \sin\varphi & \cos\varphi \end{pmatrix}$$
(1.51)

Recalling definition (1.25) of trace of linear operator we come to the following relation

$$TrP = 1 + 2\cos\varphi \tag{1.52}$$

where φ is angle of rotation. Note that Trace of the operator does not depend on the choice of the basis. This formula express cosine of the angle of rotation in terms of operator, irrelevant of the choice of the basis.

Remark This formula defines angle of rotation up to a sign.

If we change orientation then $\varphi \mapsto -\varphi$. For non-oriented Euclidean space rotation is defined up to a sign⁵

Careful reader maybe already noted that even fixing the orientation of \mathbf{E}^3 does not fix the "sign" of the angle: If we change the orientation of the axis (changing $\mathbf{n} \mapsto -\mathbf{n}$) then changing the corresponding "left" basis will imply that $\varphi \mapsto -\varphi$. In fact angle φ is the angle of rotation of oriented plane which is orthogonal to the axis of rotation. Orientation on the plane is defined by orientation in \mathbf{E}^3 and orientation of the axis which is orthogonal to this plane. In the case of 3-dimensional space sign of the angle depends not only on orientation of \mathbf{E}^3 but on orientation of axis. In what follows we will ignore this. This means that we define rotation on the angle $\pm \varphi$ up to a sign.... Rotation is defined for operators preserving orientation. The difference between angles of rotations φ and $-\varphi$ is depending not only on orientation of \mathbf{E}^3 but on orientation of axis too. But we ignore this difference. Note that $\cos \varphi$ in the formula is defined up to a sign

Rotation operator (1.51) evidently is orthogonal operator preserving orientation. Is it true converse implication? At the beginning of this paragraph we formulated the Euler Theorem. It gives the positive answer on this question. We will formulate this Theorem again in more detail:

⁵Does it recall you expressions such as "clockwise", "anticlock-wise" rotation?

Theorem (the Euler Theorem) Let P be an orthogonal operator, preserving an orientation of Euclidean space \mathbf{E}^3 , and which is not identical operator, i.e. operator P preserves the scalar product and orientation, and $P \neq \mathbf{id}$. Then it is a rotation operator with respect to an axis l on the angle φ , $(\varphi \neq 0)$.

Every vector \mathbf{N} directed along the axis does not change, i.e. the axis is 1-dimensional space of eigenvectors with eigenvalue 1, $P(\mathbf{N}) = \mathbf{N}$. Every vector orthogonal to axis rotates on the angle φ in the plane orthogonal to the axis, and

$$\operatorname{Tr} P = 1 + 2\cos\varphi$$
.

The angle φ is defined up to a sign. Changing orientation of the Euclidean space and of the axis change sign of φ .

This Theorem can be restated in the following way: every orthogonal operator P preserving orientation, (det $P \neq 0$) has an eigenvector $\mathbf{N} \neq 0$ with eigenvalue 1. This eigenvector defines the axis of rotation. In an orthonormal basis $\{\mathbf{n}, \mathbf{f}, \mathbf{g}\}$ where \mathbf{n} is a unit vector along the axis, the transition matrix of operator has an appearance (1.51). Angle of rotation can be defined via Trace of operator by formula $\operatorname{Tr} P = 1 + 2\cos\varphi$.

Remark If P is an identity operator, P = I then "there is no rotation", more precisely: any line can be considered as an axis of rotation (every vector is eigenvector of identity matrix with eigenvalue 1) and angle of rotation is equal to zero. If $P \neq I$ then axis of rotation is defined uniquely.

Proof of the Euler Theorem.

The proof of the Euler Theorem has two parts. First and central part is to prove the existence of the axis. The rest is easy: we take an arbitrary orthonormal basis \mathbf{n} , \mathbf{f} , \mathbf{g} such that \mathbf{n} is eigenvector, and we come to relations (1.50), (1.51).

There are many different proofs of existence of axis of rotation. We expose here sketchs of two proofs. The first which maybe is most beautiful proof which belongs to Coxeter. The second proof–using standard methods of linear algebra.

Coxeter's proof.

Let P be linear orthogonal operator preserving orientation. Note that for any two unit vectors \mathbf{e} , \mathbf{f} one can consider orthogonal operator $R_{\mathbf{e},\mathbf{f}}$ which swaps the vectors \mathbf{e} , \mathbf{f} , (it is reflection with respect to the plane spanned by the vectors $\mathbf{e} + \mathbf{f}$ and a vector $\mathbf{e} \times \mathbf{f}$).

Let $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ be an arbitrary orthonormal basis in \mathbf{E}^3 and let $\mathbf{e}', \mathbf{f}', \mathbf{g}'$ be image of this basis under operator P

$$P(\mathbf{e}) = \mathbf{e}', \ P(\mathbf{f}) = \mathbf{f}' \ P(\mathbf{g}) = \mathbf{g}'.$$

If $\mathbf{e} = \mathbf{e}'$ nothing to prove (\mathbf{e} is eigenvector with eigenvalue 1). If this is not the case, apply reflection operator $R_{\mathbf{e},\mathbf{e}'}$ to the initial basis $\{\mathbf{e},\mathbf{f},\mathbf{g}\}$ we come to the orthonormal basis $\{\mathbf{e}',\tilde{\mathbf{f}},\tilde{\mathbf{g}}\}$, Then applying reflection operator $R_{\tilde{\mathbf{f}},\mathbf{f}'}$ to this basis we come to the basis

 $\mathbf{e'}, \mathbf{f'}, \tilde{\mathbf{g}}$. The third vector has no choice: it has to be equal to \mathbf{g} . Indeed it may be equal to $\pm \mathbf{g}$, since all operator are orthogonal, but it cannot be equal to $-\mathbf{g'}$ since orientation will be opposite. orientation is opposite. Hence we see that operator P is the product of two reflections operators:

$$P = R_1 \circ R_2$$
.

Reflection operator is identical operator, on the plane.

Let α_1 be a plane such that R_1 is invariant on α_1 , and let α_2 be a plane such that R_2 is invariant on α_2 . Consider the line l, intersection of these planes, we come to eigenvectors with eigenvalue 1.

linear algebra proof

Any linear operator L in 3-dimensional vector space has at least one eigenvector \mathbf{x} : \mathbf{x} is non-zero solution of homogeneous equation $L\mathbf{x} = \lambda \mathbf{l}x$, where eigenvalue \mathbf{l} is a solution of cubic equation $\det(L - \mathbf{l}) = 0$, and this cubic equation has at least one root.

Hence orthogonal operator P has at least one eigenvector \mathbf{x} : $P\mathbf{x} = \lambda \mathbf{x}$, Since P is orthogonal operator, then $\lambda = \pm 1$. If $\lambda = 1$, then \mathbf{x} defines the axis since P preserves orientation. If $\lambda = -1$, $P\mathbf{x} = -\mathbf{x}$, then eigenvector with eigenvalue 1 belongs to the plane orthogonal to \mathbf{x} .

Example Consider linear operator P such that for orthonormal basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$

$$P(\mathbf{e}_x) = \mathbf{e}_y, P(\mathbf{e}_y) = \mathbf{e}_x, P(\mathbf{e}_z) = -\mathbf{e}_z \tag{1.53}$$

This is obviously orthogonal operator since it transforms orthogonal basis to orthogonal one. This operator swaps first two vectors and reflects the third one. It preserves orientation: matrix of operator in the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$, i.e. the transition matrix from the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ to the basis $\{P(\mathbf{e}_x), P(\mathbf{e}_y), P(\mathbf{e}_y)\}$ is defined by the relation:

$$\{P(\mathbf{e}_x), P(\mathbf{e}_y), P(\mathbf{e}_z)\} = \{\mathbf{e}_y, \mathbf{e}_x, -\mathbf{e}_z\} = \{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

det P = 1. This operator preserves orientation. Hence by Euler Theorem it is a rotation. Find first axis of rotation. It is easy to see from (1.53) that $\mathbf{N} = \lambda(\mathbf{e}_x + \mathbf{e}_y)$ is eigenvector with eigenvalue 1:

$$P(\mathbf{N}) = P(\mathbf{e}_x + \mathbf{e}_y) = \mathbf{e}_y + \mathbf{e}_x = \mathbf{N}$$
.

Hence axis of rotation is directed along the vector $\mathbf{e}_x + \mathbf{e}_y$. Tr $P = 1 + 2\cos\varphi = -1$. The angle of rotation $\varphi = \pi$.

One can calculate explicitly angle of rotation: Consider orthonormal basis $\{\mathbf{n}, f, \mathbf{g}\}$ adjusted to the axis $(\mathbf{n}||\mathbf{N})$. We have that $\mathbf{n} = \frac{\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}$ since \mathbf{n} is proportional to \mathbf{N} and it is unit vector. Choose $\mathbf{f} = \frac{-\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}$ and $\mathbf{g} = \mathbf{e}_z$. Then it is easy to see that

$$\{\mathbf{n}, \mathbf{f}, \mathbf{g}\} = \left\{ \frac{\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}, \frac{-\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}, \mathbf{g} \right\}$$

is orthonormal basis. Using (1.53) one can see that

$$P(\mathbf{n}) = P\left(\frac{\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}\right) = \frac{\mathbf{e}_y + \mathbf{e}_x}{\sqrt{2}} = \mathbf{n},$$

$$P(\mathbf{f}) = P\left(\frac{-\mathbf{e}_x + \mathbf{e}_y}{\sqrt{2}}\right) = \frac{-\mathbf{e}_y + \mathbf{e}_x}{\sqrt{2}} = -\mathbf{f}, \quad P(\mathbf{g}) = -\mathbf{g}$$

We see that

$$\{\mathbf{n}, \mathbf{f}, \mathbf{g}\} \xrightarrow{P} \{\mathbf{n}, -\mathbf{f}, -\mathbf{g}\}$$
.

Comparing with (1.50) and (1.51) we see that the operator P is rotation of \mathbf{E}^3 on the angle π with respect to the axis directed along the vector $\mathbf{e}_x + \mathbf{e}_y$.

1.8 Area of parallelogram, volume of parallelepiped, and determinant of linear operator

You know that area of parallelogram and volume of parallelepiped can be calculated in terms of vector (cross) product. These formulae explain geometrical meaning of determinant of linear operator.

1.8.1 Vector product in oriented E³

Now we give a definition of vector product of vectors in 3-dimensional Euclidean space equipped with orientation.

Let \mathbf{E}^3 be three-dimensional oriented Euclidean space, i.e. Euclidean space equipped with an equivalence class of bases with the same orientation. To define the orientation it suffices to consider just one orthonormal basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ which is claimed to be *left* basis. Then the equivalence class of the left bases is a set of all bases which have the same orientation as the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$.

Definition Vector product $L(\mathbf{x}, \mathbf{y}) = \mathbf{x} \times \mathbf{y}$ is a function of two vectors which takes vector values such that the following axioms (conditions) hold

• The vector $L(\mathbf{x}, \mathbf{y}) = \mathbf{x} \times \mathbf{y}$ is orthogonal to vector \mathbf{x} and vector \mathbf{y} :

$$(\mathbf{x} \times \mathbf{y}) \perp \mathbf{x}, \quad (\mathbf{x} \times \mathbf{y}) \perp \mathbf{y}$$
 (1.54)

In particular it is orthogonal to the plane spanned by the vectors \mathbf{x} , \mathbf{y} (in the case if vectors \mathbf{x} , \mathbf{y} are linearly independent)

$$\mathbf{x} \times \mathbf{y} = -\mathbf{y} \times \mathbf{x},$$
 (anticommutativity condition) (1.55)

•

$$(\lambda \mathbf{x} + \mu \mathbf{y}) \times \mathbf{z} = \lambda(\mathbf{x} \times \mathbf{z}) + \mu(\mathbf{y} \times \mathbf{z}),$$
 (linearity condition) (1.56)

• If vectors \mathbf{x} , \mathbf{y} are perpendicular each other then the magnitude of the vector $\mathbf{x} \times \mathbf{y}$ is equal to the area of the rectangle formed by the vectors \mathbf{x} and \mathbf{y} :

$$|\mathbf{x} \times \mathbf{y}| = |\mathbf{x}| \cdot |\mathbf{y}|, \quad \text{if } \mathbf{x} \perp \mathbf{y}, i.e.(\mathbf{x}, \mathbf{y}) = 0.$$
 (1.57)

• If the ordered triple of the vectors $\{\mathbf{x}, \mathbf{y}, \mathbf{z}\}$, where $\mathbf{z} = \mathbf{x} \times \mathbf{y}$ is a basis, then this basis and an orthonormal basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ defining orientation of \mathbf{E}^3 have the same orientation:

$$\{\mathbf{x}, \mathbf{y}, \mathbf{z}\} = \{\mathbf{e}, \mathbf{f}, \mathbf{g}\}T$$
, where for transition matrix T , $\det T > 0$. (1.58)

Vector product depends on orientation in Euclidean space.

Comments on conditions (axioms) (1.54)—(1.58):

1. The condition (1.56) of linearity of vector product with respect to the first argument and the condition (1.55) of anticommutativity imply that vector product is an operation which is linear with respect to the second argument too. Show it:

$$\mathbf{z} \times (\lambda \mathbf{x} + \mu \mathbf{y}) = -(\lambda \mathbf{x} + \mu \mathbf{y}) \times \mathbf{z} = -\lambda (\mathbf{x} \times \mathbf{z}) - \mu (\mathbf{y} \times \mathbf{z}) = \lambda (\mathbf{z} \times \mathbf{x}) + \mu (\mathbf{z} \times \mathbf{y}).$$

Hence vector product is bilinear operation. Comparing with scalar product we see that vector product is bilinear anticommutative (antisymmetric) operation which takes vector values, while scalar product is bilinear symmetric operation which takes real values.

2. The condition of anticommutativity immediately implies that vector product of two colinear (proportional) vectors \mathbf{x} , \mathbf{y} ($\mathbf{y} = \lambda \mathbf{x}$) is equal to zero. It follows from linearity and anticommutativity conditions. Show it: Indeed

$$\mathbf{x} \times \mathbf{y} = \mathbf{x} \times (\lambda \mathbf{x}) = \lambda(\mathbf{x} \times \mathbf{x}) = -\lambda(\mathbf{x} \times \mathbf{x}) = -\mathbf{x} \times (\lambda \mathbf{x}) = -\mathbf{x} \times \mathbf{y}.$$
 (1.59)

Hence $\mathbf{x} \times \mathbf{y} = 0$, if $\mathbf{y} = \lambda \mathbf{x}$

3. It is very important to emphasize again that vector product depends on orientation. According the condition (1.58) if $\mathbf{z} = \mathbf{x} \times \mathbf{y}$ and we change the orientation of Euclidean space, then $\mathbf{z} \to -\mathbf{z}$ since the basis $\{\mathbf{x}, \mathbf{y}, -\mathbf{z}\}$ as an orientation opposite to the orientation of the basis $\{\mathbf{x}, \mathbf{y}, \mathbf{z}\}$.

You may ask a question: Does this operation (taking the vector product) which obeys all the conditions (axioms) (1.54)—(1.58) exist? And if it exists is it unique? We will show that the vector product is well-defined by the axioms (1.54)—(1.58), i.e. there exists an operation $\mathbf{x} \times y$ which obeys the axioms (1.54)—(1.58) and these axioms define the operation uniquely.

We will assume first that there exists an operation $L(\mathbf{x}, \mathbf{y}) = \mathbf{x} \times \mathbf{y}$ which obeys all the axioms (1.54)—(1.58). Under this assumption we will construct explicitly this operation (if it exists!). We will see that the operation that we constructed indeed obeys all the axioms (1.54)—(1.58).

Let $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ be an arbitrary left orthonormal basis of oriented Euclidean space \mathbf{E}^3 , i.e. a basis which belongs to the equivalence class of the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{g}\}$ defining orientation of \mathbf{E}^3 . Then it follows from the considerations above for vector product that

$$\mathbf{e}_{x} \times \mathbf{e}_{x} = 0, \quad \mathbf{e}_{x} \times \mathbf{e}_{y} = \mathbf{e}_{z}, \quad \mathbf{e}_{x} \times \mathbf{e}_{z} = -\mathbf{e}_{y}$$

$$\mathbf{e}_{y} \times \mathbf{e}_{x} = -\mathbf{e}_{z}, \quad \mathbf{e}_{y} \times \mathbf{e}_{y} = 0, \quad \mathbf{e}_{y} \times \mathbf{e}_{z} = \mathbf{e}_{x}$$

$$\mathbf{e}_{z} \times \mathbf{e}_{x} = \mathbf{e}_{y}, \quad \mathbf{e}_{z} \times \mathbf{e}_{y} = -\mathbf{e}_{x}, \quad \mathbf{e}_{z} \times \mathbf{e}_{z} = 0$$

$$(1.60)$$

E.g. $\mathbf{e}_x \times \mathbf{e}_x = 0$, because of (1.55), $\mathbf{e}_x \times \mathbf{e}_y$ is equal to \mathbf{e}_z or to $-\mathbf{e}_z$ according to (1.57), and according to orientation arguments (1.58) $\mathbf{e}_x \times \mathbf{e}_y = \mathbf{e}_z$.

Now it follows from linearity and (1.60) that for two arbitrary vectors $\mathbf{a} = a_x \mathbf{e}_x + a_y \mathbf{e}_y + a_z \mathbf{e}_z$, $\mathbf{b} = b_x \mathbf{e}_x + b_y \mathbf{e}_y + b_z \mathbf{e}_z$

$$\mathbf{a} \times \mathbf{b} = (a_x \mathbf{e}_x + a_y \mathbf{e}_y + a_z \mathbf{e}_z) \times (b_x \mathbf{e}_x + b_y \mathbf{e}_y + b_z \mathbf{e}_z) = a_x b_y \mathbf{e}_x \times \mathbf{e}_y + a_x b_z \mathbf{e}_x \times \mathbf{e}_z + a_z \mathbf{e}_z$$

$$a_y b_x \mathbf{e}_y \times \mathbf{e}_x + a_y b_z \mathbf{e}_y \times \mathbf{e}_z + a_z b_x \mathbf{e}_z \times \mathbf{e}_x + a_z b_y \mathbf{e}_z \times \mathbf{e}_y =$$

$$(a_y b_z - a_z b_y) \mathbf{e}_x + (a_z b_x - a_x b_z) \mathbf{e}_y + (a_x b_y - a_y b_x) \mathbf{e}_z .$$

$$(1.61)$$

It is convenient to represent this formula in the following very familiar way:

$$L(\mathbf{a}, \mathbf{b}) = \mathbf{a} \times \mathbf{b} = \det \begin{pmatrix} \mathbf{e}_x & \mathbf{e}_y & \mathbf{e}_z \\ a_x & a_y & a_z \\ b_x & b_y & b_z \end{pmatrix}$$
(1.62)

We see that the operation $L(\mathbf{x}, \mathbf{y}) = \mathbf{x} \times \mathbf{y}$ which obeys all the axioms (1.54)—(1.58), if it exists, has an appearance (1.62), where $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ is an arbitrary orthonormal basis (with rightly chosen orientation). On the other hand using the properties of determinant and the fact that vectors are orthogonal if and only if their scalar product equals to zero one can easy see that the vector product defined by this formula indeed obeys all the conditions (1.54)—(1.58).

Thus we proved that the vector product is well-defined by the axioms (1.54)—(1.58) and it is given by the formula (1.62) in an arbitrary orthonormal basis (with rightly chosen orientation).

Remark In the formula above we have chosen an arbitrary orthonormal basis which belongs to the equivalence class of bases defining the orientation. What will happen if we choose instead the basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ an arbitrary orthonormal basis $\{\mathbf{f}_1, \mathbf{f}_2, \mathbf{f}_3\}$. We see that such that answer does not change if both bases $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ and $\{\mathbf{f}_1, \mathbf{f}_2, \mathbf{f}_3\}$ have the same orientation, Formulae (1.60) are valid for an arbitrary orthonormal basis which have the same orientation as the orthonormal basis $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$.— In oriented Euclidean space \mathbf{E}^3 we may take an arbitrary basis from the equivalence class of bases defining orientation. On the other hand if we will consider the basis with opposite orientation then according to the axiom (1.58) vector product will change the sign. (See also the question 6 in Homework 4)

1.8.2 Vector product—area of parallelogram

The following Proposition states that vector product can be considered as area of parallelogram:

Proposition 2 The modulus of the vector $\mathbf{z} = \mathbf{x} \times \mathbf{y}$ is equal to the area of parallelogram formed by the vectors \mathbf{x} and \mathbf{y} :

$$S(\mathbf{x}, \mathbf{y}) = S(\Pi(\mathbf{x}, \mathbf{y})) = |\mathbf{x} \times \mathbf{y}|, \qquad (1.63)$$

where we denote by $S(\mathbf{x}, \mathbf{y})$ the area of parallelogram $\Pi(\mathbf{x}, \mathbf{y})$ formed by the vectors \mathbf{x}, \mathbf{y} .

Proof: Consider the expansion $\mathbf{y} = \mathbf{y}_{||} + \mathbf{y}_{\perp}$, where the vector \mathbf{y}_{\perp} is orthogonal to the vector \mathbf{x} and the vector $\mathbf{y}_{||}$ is parallel to to vector \mathbf{x} . The area of the parallelogram formed by vectors \mathbf{x} and \mathbf{y} is equal to the product of the length of the vector \mathbf{x} on the height. The height is equal to the length of the vector \mathbf{y}_{\perp} . We have $S(\mathbf{x}, \mathbf{y}) = |\mathbf{x}||\mathbf{y}_{\perp}|$. On the other $\mathbf{z} = \mathbf{x} \times \mathbf{y} = \mathbf{x} \times (\mathbf{y}_{||} + \mathbf{y}_{\perp}) = \mathbf{x} \times \mathbf{y}_{||} + \mathbf{x} \times \mathbf{y}_{\perp}$. But $\mathbf{x} \times \mathbf{y}_{||} = 0$, because these vectors are colinear. Hence $\mathbf{z} = \mathbf{x} \times \mathbf{y}_{\perp}$ and $|\mathbf{z}| = |\mathbf{x}||\mathbf{y}_{\perp}| = S(\mathbf{x}, \mathbf{y})$ because vectors $\mathbf{x}, \mathbf{y}_{\perp}$ are orthogonal to each other.

This Proposition is very important to understand the meaning of vector product. Shortly speaking vector product of two vectors is a vector which is orthogonal to the plane spanned by these vectors, such that its magnitude is equal to the area of the parallelogram formed by these vectors. The direction is defined by orientation.

Remark It is useful sometimes to consider area of parallelogram not as a positive number but as an real number positive or negative (see the next subsubsection.)

It is not worthless to recall the formula which we know from the school that area of parallelogram formed by vectors \mathbf{x} , \mathbf{y} equals to the product of the base on the height. Hence

$$|\mathbf{x} \times \mathbf{y}| = |\mathbf{x}| \cdot |\mathbf{y}| |\sin \theta|, \qquad (1.64)$$

where θ is an angle between vectors \mathbf{x}, \mathbf{y} .

Finally I would like again to stress:

Vector product of two vectors is equal to zero if these vectors are colinear (parallel). Scalar product of two vectors is equal to zero if these vector are orthogonal.

Exercise †Show that the vector product obeys to the following identity:

$$((\mathbf{a} \times \mathbf{b}) \times \mathbf{c}) + ((\mathbf{b} \times \mathbf{c}) \times \mathbf{a}) + ((\mathbf{c} \times \mathbf{a}) \times \mathbf{b}) = 0. \quad (Jacoby identity)$$
 (1.65)

This identity is related with the fact that heights of the triangle intersect in the one point.

Exercise \dagger Show that $\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = \mathbf{b}(\mathbf{a}, \mathbf{c}) - \mathbf{c}(\mathbf{a}, \mathbf{b})$.

1.8.3 Area of parallelogram in E^2 and determinant of 2×2 matrices

Let \mathbf{a}, \mathbf{b} be two vectors in 2-dimensional vector space \mathbf{E}^2 .

One can consider \mathbf{E}^2 as a plane in 3-dimensional Euclidean space \mathbf{E}^3 . Our aim is to calculate the area of the parallelogram $\Pi(\mathbf{a}, \mathbf{b})$ formed by vectors

 \mathbf{a}, \mathbf{b} . Let \mathbf{n} be a unit vector in \mathbf{E}^3 which is orthogonal to \mathbf{E}^2 . Then it is obvious that the vector product $\mathbf{a} \times \mathbf{b}$ is proportional to the normal vector \mathbf{n} to the plane \mathbf{E}^2 :

$$\mathbf{a} \times \mathbf{b} = A(\mathbf{a}, \mathbf{b})\mathbf{n}, \tag{1.66}$$

and the area of the parallelogram $\Pi(\mathbf{a}, \mathbf{b})$ equals to the modulus of the coefficient $A(\mathbf{c}, \mathbf{b})$:

$$S(\Pi(\mathbf{a}, \mathbf{b})) = |\mathbf{a} \times \mathbf{b}| = |A(\mathbf{a}, \mathbf{b})|. \tag{1.67}$$

The normal unit vector \mathbf{n} and coefficient $A(\mathbf{a}, \mathbf{b})$ are defined up to a sign: $\mathbf{n} \to -\mathbf{n}$, $A \to -A$. On the other hand the vector product $\mathbf{a} \times \mathbf{b}$ is defined up to a sign too: vector product depends on orientation. The answer for $\mathbf{a} \times \mathbf{b}$ is not changed if we perform calculations for vector product in an arbitrary basis $\{\mathbf{e}'_x, \mathbf{e}'_y, \mathbf{e}'_z\}$ which have the same orientation as the the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ and $\mathbf{a} \times \mathbf{b} \mapsto -\mathbf{a} \times \mathbf{b}$. If we consider an arbitrary basis $\{\mathbf{e}'_x, \mathbf{e}'_y, \mathbf{e}'_z\}$ which have the orientation opposite to the orientation of the basis $\{\mathbf{e}, \mathbf{f}, \mathbf{n}\}$ (e.g. the basis $\{\mathbf{e}, \mathbf{f}, -\mathbf{n}\}$) then $A(\mathbf{a}, \mathbf{b}) \to -A(\mathbf{a}, \mathbf{b})$. The magnitude $A(\mathbf{a}, \mathbf{b})$ is so called algebraic area of parallelogram. It can positive and negative.

If (a_1, a_2) , (b_1, b_2) are coordinates of the vectors \mathbf{a} , \mathbf{b} in the orthonormal basis $\{\mathbf{e}, \mathbf{f}\}$: $\mathbf{a} = a_1\mathbf{e} + a_2\mathbf{f}$, $\mathbf{b} = b_1\mathbf{e} + b_2\mathbf{f}$ and according to (1.62)

$$\mathbf{a} \times \mathbf{b} = \det \begin{pmatrix} \mathbf{e} & \mathbf{f} & \mathbf{n} \\ a_1 & a_2 & 0 \\ b_1 & b_2 & 0 \end{pmatrix} = \mathbf{n} \det \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix}$$
(1.68)

Thus $A(\mathbf{a}, \mathbf{b})$ in equation (1.67) is equal to $\det \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix}$, and we come to the following formula for area of parallelogram

$$S(\Pi(\mathbf{a}, \mathbf{b})) = |\mathbf{a} \times \mathbf{b}| = \left| \det \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix} \right|. \tag{1.69}$$

This is an important formula for relation between determinant of 2×2 matrix, area of parallelogram and vector product.

One can deduce this relation in other way:

Let \mathbf{E}^2 be a 2-dimensional Euclidean space. The function $A(\mathbf{a}, \mathbf{b})$ defined by the relation (1.69) obeys the following conditions:

• It is anticommutative:

$$A(\mathbf{a}, \mathbf{b}) = -A(\mathbf{a}, \mathbf{b}) \tag{1.70}$$

• It is bilinear

$$A(\lambda \mathbf{a} + \mu \mathbf{b}, \mathbf{c}) = \lambda A(\mathbf{a}, \mathbf{c}) + \mu A(\mathbf{b}, \mathbf{c}); \ A(\mathbf{c}, \lambda \mathbf{a} + \mu \mathbf{b}) = \lambda A(\mathbf{c}, \mathbf{a}) + \mu A(\mathbf{c}, \mathbf{b}). \ (1.71)$$

• and it obeys normalisation condition:

$$A(\mathbf{e}, \mathbf{f}) = \pm 1 \tag{1.72}$$

for an arbitrary orthonormal basis.

(Compare with conditions (1.54)—(1.58).)

One can see that these conditions define uniquely $A(\mathbf{a}, \mathbf{b})$ and these are the conditions which define the determinant of the 2×2 matrix.

1.8.4 Areas of parallelograms and determinants of linear operators in E^2

Let A be an arbitrary linear operator in \mathbf{E}^2 . One can see that the following formula holds.

Let \mathbf{a}, \mathbf{b} be two arbitrary vectors in \mathbf{E}^2 . Let \mathbf{a}', \mathbf{b}' be two vectors such that

$$\mathbf{a}' = A(\mathbf{a}), \quad \mathbf{b} = A(\mathbf{b}').$$

Consider two parallelograms: Parallelogram $\Pi(\mathbf{a}, \mathbf{b})$ formed by vectors \mathbf{a}, \mathbf{b} , and the second parallelogram $\Pi(\mathbf{a}', \mathbf{b}')$ formed by vectors $\alpha'.b'$. Then one can deduce from equation (1.69) that

Area of
$$\Pi(\mathbf{a}', \mathbf{b}') = |\det A| \cdot \text{Area of } \Pi(\mathbf{a}, \mathbf{b})$$
 . (1.73)

This formula relates volumes of parallelograms $\Pi(\mathbf{a}, \mathbf{b})$, $\Pi(\mathbf{a}', \mathbf{b}')$ with determinant of linear operator which transforms the first parallelogram to the second one. (See also exercise 9 in Homework 4).

Prove straightforwardly equation (1.73). Let vectors **a**, **b** be linearly independent (if they are dependent, then area of both parallelograms in (??) evidently vanish). We have:

$$\mathbf{a}' = A(\mathbf{a}) \,, \quad \text{i.e. } \begin{pmatrix} a_1' \\ a_2' \end{pmatrix} = A \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \begin{pmatrix} A_{11}a_1 + A_{12}a_2 \\ A_{21}a_1 + A_{22}a_2 \end{pmatrix} \,,$$

$$\mathbf{b}' = A(\mathbf{b}), \text{ i.e. } \begin{pmatrix} b_1' \\ b_2' \end{pmatrix} = A \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} A_{11}b_1 + A_{12}b_2 \\ A_{21}b_1 + A_{22}b_2 \end{pmatrix}.$$

Hence

$$\begin{pmatrix} \mathbf{a}' \\ \mathbf{b}' \end{pmatrix} = \begin{pmatrix} a_1' & a_2' \\ b_1' & b_2' \end{pmatrix} = \begin{pmatrix} A_{11}a_1 + A_{12}a_2 & A_{21}a_1 + A_{22}a_2 \\ A_{11}b_1 + A_{12}b_2 & A_{21}b_1 + A_{22}b_2 \end{pmatrix} = \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix} \circ \begin{pmatrix} A_{11} & A_{21} \\ A_{12} & A_{22} \end{pmatrix} = \begin{pmatrix} \mathbf{a} \\ \mathbf{b} \end{pmatrix} A^T \,.$$

Taking determinants we come to: Area of $\Pi a', b' =$

$$\left| \det \begin{pmatrix} \mathbf{a}' \\ \mathbf{b}' \end{pmatrix} \right| = \left| \det \begin{pmatrix} \begin{pmatrix} \mathbf{a} \\ \mathbf{b} \end{pmatrix} A^T \end{pmatrix} \right| = \left| \det \begin{pmatrix} \begin{pmatrix} \mathbf{a} \\ \mathbf{b} \end{pmatrix} A^T \end{pmatrix} \right| = \left| \det \begin{pmatrix} \begin{pmatrix} \mathbf{a} \\ \mathbf{b} \end{pmatrix} \det A \right| = \text{Area of } \Pi(\mathbf{a}, \mathbf{b}) | \det A |.$$

1.8.5 Volume of parallelepiped

The vector product of two vectors is related with area of parallelogram. What about a volume of parallelepiped formed by three vectors $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$?

Consider parallelepiped $\Pi(\mathbf{a}, \mathbf{b}, \mathbf{c})$ formed by vectors $\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}$. The parallelegram $\Pi(\mathbf{a}, \mathbf{b})$ formed by vectors \mathbf{b}, \mathbf{c} can be considered as a base of this parallelepiped.

Let θ be an angle between height and vector \mathbf{a} . It is just the angle between the vector $\mathbf{b} \times \mathbf{c}$ and the vector \mathbf{a} . Then the volume is equal to the length of the height multiplied on the area of the parallelogram, $V = Sh = S|\mathbf{a}|\cos\theta$, i.e. volume is equal to scalar product of the vectors \mathbf{a} on the vector product of vectors \mathbf{b} and \mathbf{c} :

$$V(\{\mathbf{a}, \mathbf{b}, \mathbf{c}\}) = |(\mathbf{a}, \mathbf{b} \times \mathbf{c})| = \left| \left(a_x \mathbf{e}_x + a_y \mathbf{e}_y + a_z \mathbf{e}_z, \det \begin{pmatrix} \mathbf{e}_x & \mathbf{e}_y & \mathbf{e}_z \\ b_x & b_y & b_z \\ c_x & c_y & c_z \end{pmatrix} \right) \right|$$

$$= |(a_x \mathbf{e}_x + a_y \mathbf{e}_y + a_z \mathbf{e}_z, (b_y c_z - b_z c_y) \mathbf{e}_x + (b_z c_x - b_x c_z) \mathbf{e}_y + (b_x c_y - b_y c_x) \mathbf{e}_z)| =$$

$$|a_x (b_y c_z - b_z c_y) + a_y (b_z c_x - b_x c_z) + a_z (b_x c_y - b_y c_x)| = \left| \det \begin{pmatrix} a_x & a_y & a_z \\ b_x & b_y & b_z \\ c_x & c_y & c_z \end{pmatrix} \right|.$$

We come to beautiful and useful formula:

volume of
$$\Pi(\mathbf{a}, \mathbf{b}, \mathbf{c}) = |(\mathbf{a}, [\mathbf{b} \times \mathbf{c}])| = \left| \det \begin{pmatrix} a_x & a_y & a_z \\ b_x & b_y & b_z \\ c_x & c_y & c_z \end{pmatrix} \right|.$$
 (1.74)

Compare this formula with the formula (1.69) for the area of parallelogram. **Remark** In these formulae we consider the volume of the parallelepiped as a positive number. It is why we put the sign of 'modulus' in all the formulae above. On the other hand often it is very useful to consider the volume as a real number (it could be positive and negative).

1.8.6 Volumes of parallelepipeds and determinants of linear operators in E³

Write down an equation for the volumes of parallelepipeds analogous to equation (1.73) for the the areas of parallelograms. Now instead parallelogram we consider parallelepiped, and instead linear operator A in \mathbf{E}^2 we consider linear operator A in \mathbf{E}^3 .

Let A be an arbitrary linear operator in \mathbf{E}^3 . In the same way as in formula (1.73) the following formula holds:

Let $\mathbf{a}, \mathbf{b}, \mathbf{c}$ be three arbitrary vectors in \mathbf{E}^3 . Linear operator A transforms these three vectors to three vectors $\mathbf{a}', \mathbf{b}', \mathbf{c}'$ where

$$\mathbf{a}' = A(\mathbf{a}), \quad \mathbf{b} = A(\mathbf{b}'), \mathbf{c}' = P(\mathbf{c}').$$

Consider two parallelepipeds: Parallelepiped $\Pi(\mathbf{a}, \mathbf{b} c)$ formed by vectors $\mathbf{a}, \mathbf{b}, \mathbf{c}$ and the second parallelepiped $\Pi(\mathbf{a}', \mathbf{b}' c')$ formed by vectors $\alpha'.b', \mathbf{c}'$. Then it follows from (1.74) the following formula and determinant of operator A:

Volume of
$$\Pi(\mathbf{a}', \mathbf{b}', \mathbf{c}') = |\det A| \cdot \text{Volume of } \Pi(\mathbf{a}, \mathbf{b}, \mathbf{c})$$
. (1.75)

This formula relates volumes of parallelepipeds $\Pi(\mathbf{a}, \mathbf{b}, \mathbf{c})$, $\Pi(\mathbf{a}', \mathbf{b}', \mathbf{c}')$ with determinant of linear operator which transforms the first parallelepiped to the second one. (See also exercise 9 in Homework 4).

2 Differential forms

2.1 Tangent vectors, curves, velocity vectors on the curve

Tangent vector is a vector \mathbf{v} applied at the given point $\mathbf{p} \in \mathbf{E}^n$.

The set of all tangent vectors at the given point \mathbf{p} is a vector space. It is called tangent space of \mathbf{E}^n at the point \mathbf{p} and it is denoted $T_{\mathbf{p}}(\mathbf{E}^n)$.

One can consider vector field on \mathbf{E}^n , i.e.a function which assigns to every point \mathbf{p} vector $\mathbf{v}(\mathbf{p}) \in T_{\mathbf{p}}(\mathbf{E}^n)$.

Here we consider on an equal footing vectors of vector space \mathbf{E}^n and points of associated affine space \mathbf{E}^n (as usual we denote them by the same letter (see for details subsection 1.2)

It is instructive to study the conception of tangent vectors and vector fields on the curves and surfaces embedded in \mathbf{E}^n . In this course we mainly consider tangent vectors to curves.

A curve in \mathbf{E}^n with parameter $t \in (a, b)$ is a continuous map

C:
$$(a,b) \to \mathbf{E}^n$$
 $\mathbf{r}(t) = (x^1(t), \dots, x^n(t)), \quad a < t < b$ (2.1)

For example consider in \mathbf{E}^2 the curve

C:
$$(0, 2\pi) \to \mathbf{E}^2$$
 $\mathbf{r}(t) = (R \cos t, R \sin t), \ 0 \le t < 2\pi$.

The image of this curve is the circle of the radius R. It can be defined by the equation:

$$x^2 + y^2 = R^2.$$

To distinguish between curve and its image we say that curve C in (2.1) is parameterised curve or path. We will call the image of the curve unparameterised curve (see for details the next subsection). It is very useful to think about parameter t as a "time" and consider parameterised curve like point moving along a curve. Unparameterised curve is the trajectory of the moving point. It is locus of the points. The using of word "curve" without adjective "parameterised" or "nonparameterised" sometimes is ambiguous.

Vectors tangent to curve—velocity vector

Let $\mathbf{r}(t)$ $\mathbf{r} = \mathbf{r}(t)$ be a curve in \mathbf{E}^n . Velocity $\mathbf{v}(t)$ it is the vector

$$\mathbf{v}(t) = \frac{d\mathbf{r}}{dt} = (\dot{x}^1(t), \dots, \dot{x}^n(t)) = (v^1(t), \dots, v^n(t))$$

in \mathbf{E}^n . Velocity vector is tangent vector to the curve.

Let $C: \mathbf{r} = \mathbf{r}(t)$ be a curve and $\mathbf{r}_0 = \mathbf{r}(t_0)$ any given point on it. Then the set of all vectors tangent to the curve at the point $\mathbf{r}_0 = \mathbf{r}(t_0)$ is one-dimensional vector space $T_{\mathbf{r}_0}C$. It is linear subspace in vector space $T_{\mathbf{r}_0}C$: $T_{\mathbf{r}_0}C < T_{\mathbf{r}_0}\mathbf{E}^n$. The points of the tangent space $T_{\mathbf{r}_0}C$ are the points of tangent line.

Remark We consider by default only *smooth*, regular curves. Curve $\mathbf{r}(t) = (x^1(t), \dots, x^n(t))$ is called smooth if all functions $x^i(t)$, $(i = 1, 2, \dots, n)$ are smooth functions (Function is called smooth if it has derivatives of arbitrary order.) Curve $\mathbf{r}(t)$ is called regular if velocity vector $\mathbf{v}(t) = \frac{d\mathbf{r}(t)}{dt}$ is not equal to zero at all t.

2.2 Reparameterisation

One can move along trajectory with different velocities, i.e. one can consider different parameterisation. E.g. consider

$$C_1$$
: $\begin{cases} x(t) = t \\ y(t) = t^2 \end{cases}$ $0 < t < 1$, C_2 : $\begin{cases} x(t) = \sin t \\ y(t) = \sin^2 t \end{cases}$ $0 < t < \frac{\pi}{2}$

Images of these two parameterised curves are the same. (These curves have the same loci.) In both cases point moves along a piece of the same parabola but with different velocities.

Definition

Two smooth curves C_1 : $\mathbf{r}_1(t)$: $(a_1, b_1) \to \mathbf{E}^n$ and C_2 : $\mathbf{r}_2(\tau)$: $(a_2, b_2) \to \mathbf{E}^n$ are called equivalent if there exists reparameterisation map:

$$t(\tau): (a_2, b_2) \to (a_1, b_1),$$

such that

$$r_2(\tau) = r_1(t(\tau)) \tag{2.2}$$

Reparameterisation $t(\tau)$ is diffeomorphism, i.e. function $t(\tau)$ has derivatives of all orders and first derivative $t'(\tau)$ is not equal to zero.

E.g. curves in (2.2) are equivalent because a map $\varphi(t) = \sin t$ transforms first curve to the second.

 $\label{lem:equivalence} Equivalence\ class\ of\ equivalent\ parameterised\ curves\ is\ called\ non-parameterised\ curve.$

Equivalent curves have the same image. (They have the same loci.)

It is useful sometimes to distinguish curves in the same equivalence class which differ by orientation.

Definition Let C_1, C_2 be two equivalent curves. We say that they have same orientation (parameterisations $\mathbf{r}_1(t)$ and $\mathbf{r}_1(\tau)$ have the same orientation) if reparameterisation $t = t(\tau)$ has positive derivative, $t'(\tau) > 0$. We say that they have opposite orientation (parameterisations $\mathbf{r}_1(t)$ and $\mathbf{r}_1(\tau)$ have the opposite orientation) if reparameterisation $t = t(\tau)$ has negative derivative, $t'(\tau) < 0$.

Changing orientation means changing the direction of "walking" around the curve.

Equivalence class of equivalent curves splits on two subclasses with respect to orientation.

Non-formally: Two curves are equivalent curves (belong to the same equivalence class) if these parameterised curves (paths) have the same images. Two equivalent curves have the same image. They define the same set of points in \mathbf{E}^n . Different parameters correspond to moving along curve with different velocity. Two equivalent curves have opposite orientation If two parameterisations correspond to moving along the curve in different directions then these parameterisations define opposite orientation.

What happens with velocity vector if we change parameterisation? It changes its value, but it can change its direction only on opposite (If these parameterisations have opposite orientation of the curve):

$$\mathbf{v}(\tau) = \frac{d\mathbf{r}_2(\tau)}{d\tau} = \frac{d\mathbf{r}(t(\tau))}{d\tau} = \frac{dt(\tau)}{d\tau} \cdot \frac{d\mathbf{r}(t)}{dt} \Big|_{t=t(\tau)}$$
(2.3)

Or shortly: $\mathbf{v}(\tau)|_{\tau} = t_{\tau}(\tau)\mathbf{v}(t)|_{t=t(\tau)}$

We see that velocity vector is multiplied on the coefficient (depending on the point of the curve), i.e. velocity vectors for different parameterisations are collinear vectors.

(We call two vectors \mathbf{a}, \mathbf{b} collinear, if they are proportional each other, i,e, if $\mathbf{a} = \lambda \mathbf{b}$.)

Example Consider following three curves in \mathbf{E}^2 :

$$C_{1}: \begin{cases} x = R\cos\theta \\ y = R\sin\theta \end{cases}, 0 < \theta < \pi, \quad \mathbf{v} = \begin{pmatrix} v_{x} \\ v_{y} \end{pmatrix} = \begin{pmatrix} x_{\theta} \\ y_{\theta} \end{pmatrix} = \begin{pmatrix} -R\sin\theta \\ R\cos\theta \end{pmatrix}, |\mathbf{v}| = R, \end{cases}$$

$$C_{2}: \begin{cases} x = R\cos 4\varphi \\ y = R\sin 4\varphi \end{cases}, 0 < \varphi < \frac{\pi}{4}, \quad \mathbf{v} = \begin{pmatrix} v_{x} \\ v_{y} \end{pmatrix} = \begin{pmatrix} x_{\varphi} \\ y_{\varphi} \end{pmatrix} = \begin{pmatrix} -4R\sin\theta \\ 4R\cos\theta \end{pmatrix}, |\mathbf{v}| = 4R, \end{cases}$$

$$C_{3}: \begin{cases} x = Ru \\ y = R\sqrt{1 - u^{2}} \end{cases}, -1 < u < 1, \quad \mathbf{v} = \begin{pmatrix} v_{u} \\ v_{u} \end{pmatrix} = \begin{pmatrix} x_{u} \\ y_{u} \end{pmatrix} = \begin{pmatrix} R \\ \frac{-Ru}{\sqrt{1 - u^{2}}} \end{pmatrix}, \tag{2.4}$$

These three parameterised curves, (paths) define the same non-parameterised curve: the upper piece of the circle: $x^2 + y^2 = 1, y > 0$. The reparameterisation $\theta = 4\varphi$ transforms the first curve, to the second curve. The reparameterisation $u(\theta) = \cos \theta$ transforms the third curve to the first one.

Curves C_1, C_2 have the same orientation, because $\theta'(\varphi) = 4 > 0$. Curves C_1, C_3 have opposite orientation because $u'(\theta) < 0$.

Curves C_2 and C_2 have opposite orientations too since the curves C_2 and C_1 have the same orientation, and the curves C_3 and C_1 have the opposite orientation.

In the first case point moves with constant pace $|\mathbf{v}(\theta)| = R$ anti clock-wise "from right to left" from the point A = (R, 0) to the point B = (-R, 0).

In the second case point moves with constant pace $|\mathbf{v}(\theta)| = 4R$ anti clockwise "from right to left" from the point A = (R, 0) to the point B = (-R, 0).

In the third case pace is not constant, but $v_x = 1$ is constant. Point moves clock-wise "from left to right", from the point B = (-R, 0) to the point A = (R, 0). In the third case point also moves clock-wise "from the left to right".

There are other examples in the Homeworks.

2.3 Differential 0-forms and 1-forms

2.3.1 Definition and examples of 0-forms and 1-forms

Most of considerations of this and next subsections can be considered only for \mathbf{E}^2 . All examples for differential forms is only for \mathbf{E}^2 .

0-form on \mathbf{E}^n it is just function on \mathbf{E}^n (all functions under consideration are differentiable)

Now we define 1-forms.

Definition Differential 1-form ω on \mathbf{E}^n is a function on tangent vectors of \mathbf{E}^n , such that it is linear at each point:

$$\omega(\mathbf{r}, \lambda \mathbf{v}_1 + \mu \mathbf{v}_2) = \lambda \omega(\mathbf{r}, \mathbf{v}_1) + \mu \omega(\mathbf{r}, \mathbf{v}_2). \tag{2.5}$$

Here $\mathbf{v}_1, \mathbf{v}_2$ are vectors tangent to \mathbf{E}^n at the point \mathbf{r} , $(\mathbf{v}_1, \mathbf{v}_2 \in T_x \mathbf{E}^n)$ (We recall that vector tangent at the point \mathbf{r} means vector attached at the point \mathbf{r}). We suppose that ω is smooth function on points \mathbf{r} .

If $\mathbf{X}(\mathbf{r})$ is vector field and ω -1-form then evaluating ω on $\mathbf{X}(\mathbf{r})$ we come to the function $\omega(\mathbf{r}, \mathbf{X}(\mathbf{r}))$ on \mathbf{E}^n .

Let $\mathbf{e}_1, \dots, \mathbf{e}_n$ be a basis in \mathbf{E}^n and (x^1, \dots, x^n) corresponding coordinates: an arbitrary point with coordinates (x^1, \dots, x^n) is assigned to the vector $\mathbf{r} = x^1 \mathbf{e}_1 + x^2 \mathbf{e}_2 + \dots x^n \mathbf{e}_n$ starting at the origin.

Translating basis vectors \mathbf{e}_i (i = 1, ..., n) from the origin to other points of \mathbf{E}^n we come to vector field which we also denote \mathbf{e}_i (i = 1, ..., n). The value of vector field \mathbf{e}_i at the point $(x^1, ..., x^n)$ is the vector \mathbf{e}_i attached at this point (tangent to this point).

Let ω be an 1-form on \mathbf{E}^n . Consider an arbitrary vector field $\mathbf{A}(\mathbf{r}) = \mathbf{A}(x^1, \dots, x^n)$:

$$\mathbf{A}(\mathbf{r}) = A^{1}(\mathbf{r})\mathbf{e}_{1} + \dots + A^{n}(\mathbf{r})\mathbf{e}_{n} = \sum_{i=1}^{n} A^{i}(\mathbf{r})\mathbf{e}_{i}$$

Then by linearity

$$\omega(\mathbf{r}, \mathbf{A}(\mathbf{r})) = \omega\left(\mathbf{r}, A^{1}(\mathbf{r})\mathbf{e}_{1} + \dots + A^{n}(\mathbf{r})\mathbf{e}_{n}\right) = A^{1}\omega(\mathbf{r}, \mathbf{e}_{1}) + \dots + A^{n}\omega(\mathbf{r}, \mathbf{e}_{n}).$$

Consider basic differential forms dx^1, dx^2, \dots, dx^n such that

$$dx^{i}(\mathbf{e}_{j}) = \delta_{j}^{i} = \begin{cases} 1 \text{ if } i = j \\ 0 \text{ if } i \neq j \end{cases}$$
 (2.6)

Then it is easy to see that

$$dx^{1}(\mathbf{A}) = \mathbf{A}^{1}, dx^{2}(\mathbf{A}) = \mathbf{A}^{2}, ..., i.e. dx^{i}(\mathbf{A}) = A^{i}$$

Hence

$$\omega(\mathbf{r}, \mathbf{A}(\mathbf{r})) = (\omega_1(\mathbf{r})dx^1 + \omega_2(\mathbf{r})dx^2 + \dots + \omega_n(\mathbf{r})dx^n) (\mathbf{A}(\mathbf{r}))$$

where components $\omega_i(\mathbf{r}) = \omega(\mathbf{r}, \mathbf{e}_i)$.

In the same way as an arbitrary vector field on \mathbf{E}^n can be expanded over the basis $\{\mathbf{e}_i\}$ (see (2.3.1)), an arbitrary differential 1-form ω can be expanded over the basis forms (2.3.1)

$$\omega = \omega_1(x^1, \dots, x^n)dx^1 + \omega_2(x^1, \dots, x^n)dx^2 + \dots + \omega_n(x^1, \dots, x^n)dx^n.$$

Example Consider in \mathbf{E}^2 a basis $\mathbf{e}_x, \mathbf{e}_y$ and corresponding coordinates (x, y).

Then

$$dx(\mathbf{e}_x) = 1, dx(\mathbf{e}_y) = 0$$

$$dy(\mathbf{e}_x) = 0, dy(\mathbf{e}_y) = 1$$
(2.7)

The value of a differential 1-form $\omega = a(x,y)dx + b(x,y)dy$ on vector field $\mathbf{X} = A(x,y)\mathbf{e}_x + B(x,y)\mathbf{e}_y$ is equal to

$$\omega(\mathbf{r}, \mathbf{X}) = a(x, y)dx(\mathbf{X}) + b(x, y)dx(\mathbf{X}) =$$

$$a(x, y)A(x, y) + b(x, y)B(x, y).$$

It is very useful (see below) introduce for basic vectors new notations:

$$\mathbf{e}_i \mapsto \frac{\partial}{\partial x^i}$$
 for basic vectors $\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z$ in $\mathbf{E}^3 \mathbf{e}_x \mapsto \frac{\partial}{\partial x} \mathbf{e}_y \mapsto \frac{\partial}{\partial y} \mathbf{e}_z \mapsto \frac{\partial}{\partial z}$. (2.8)

In these new notations the formula (2.3.1) looks like

$$dx^{i} \left(\frac{\partial}{\partial x^{j}} \right) = \delta^{i}_{j} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$$

and the formula (2.7) looks like

$$dx\left(\frac{\partial}{\partial x}\right) = 1, dx\left(\frac{\partial}{\partial y}\right) = 0$$
$$dy\left(\frac{\partial}{\partial x}\right) = 0, dy\left(\frac{\partial}{\partial y}\right) = 1$$

The notations which we introduce look 'odd'. However they are powerful, since uisng these notations we can work in *arbitrary coordinates*.

We will try to demonstrate it later. Now in the next subsection we will consider the directional derivative of function along vector fields. The formula which will be introduced can be written in arbitrary coordinates; it will be another justification of notations (2.8).

2.3.2 Vectors—directional derivatives of functions

Let **R** be a vector in \mathbf{E}^n tangent to the point $\mathbf{r} = \mathbf{r}_0$ (attached at a point $\mathbf{r} = \mathbf{r}_0$). Define the operation of derivative of an arbitrary (differentiable) function at the point \mathbf{r}_0 along the vector \mathbf{R} —directional derivative of function f along the vector \mathbf{R}

Definition

Let $\mathbf{r}(t)$ be a curve such that

$$\bullet |\mathbf{r}(t)|_{t=0} = \mathbf{r}_0$$

• Velocity vector of the curve at the point \mathbf{r}_0 is equal to \mathbf{R} : $\frac{d\mathbf{r}(t)}{dt}\Big|_{t=0} = \mathbf{R}$

Then directional derivative of function f with respect to the vector \mathbf{R} at the point $\mathbf{r}_0 \partial_{\mathbf{R}} f|_{\mathbf{r}_0}$ is defined by the relation

$$\partial_{\mathbf{R}} f \big|_{\mathbf{r}_0} = \frac{d}{dt} \left(f \left(\mathbf{r}(t) \right) \right) \big|_{t=0}.$$
 (2.9)

Using chain rule one come from this definition to the following important formula for the directional derivative:

If
$$\mathbf{R} = \sum_{i=1}^{n} R^{i} \mathbf{e}_{i}$$
 then $\partial_{\mathbf{R}} f \big|_{\mathbf{r}_{0}} = \sum_{i=1}^{n} R^{i} \frac{\partial}{\partial x^{i}} f(x^{1}, \dots, x^{n}) \big|_{\mathbf{r} = \mathbf{r}_{0}}$ (2.10)

It follows form this formula that

One can assign to every vector $\mathbf{R} = \sum_{i=1}^{n} R^{i} \mathbf{e}_{i}$ the operation $\partial_{\mathbf{R}} = R^{1} \frac{\partial}{\partial x^{1}} + R^{2} \frac{\partial}{\partial x^{2}} + \cdots + R^{n} \frac{\partial}{\partial x^{n}}$ of taking directional derivative:

$$\mathbf{R} = \sum_{i=1}^{n} R^{i} \mathbf{e}_{i} \mapsto \partial_{\mathbf{R}} = \sum_{i=1}^{n} R^{i} \frac{\partial}{\partial x^{i}}$$
 (2.11)

Thus we come to notations (2.8). The symbols $\partial_x, \partial_y, \partial_z$ correspond to partial derivative with respect to coordinate x or y or z. Later we see that these new notations are very illuminating when we deal with arbitrary coordinates, such as polar coordinates or spherical coordinates, The conception of orthonormal basis is ill-defined in arbitrary coordinates, but one can still consider the corresponding partial derivatives. Vector fields $\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z$ (or in new notation $\partial_x, \partial_y, \partial_z$) can be considered as a basis⁶ in the space of all vector fields on \mathbf{E}^3 .

An arbitrary vector field (2.3.1) can be rewritten in the following way:

$$\mathbf{A}(\mathbf{r}) = A^{1}(\mathbf{r})\mathbf{e}_{1} + \dots + A^{n}(\mathbf{r})\mathbf{e}_{n} = A^{1}(\mathbf{r})\frac{\partial}{\partial x^{1}} + A^{2}(\mathbf{r})\frac{\partial}{\partial x^{2}} + \dots + A^{n}(\mathbf{r})\frac{\partial}{\partial x^{n}}$$
(2.12)

2.3.3 Differential acting 0-forms \rightarrow 1-forms

Now we introduce very important operation: Differential d which acts on 0-forms and transforms them to 1 forms.

Later we will learn how differential acts on 1-forms transforming them to 2-forms.

Definition Let f = f(x)-be 0-form, i.e. function on \mathbf{E}^n . Then

$$df = \sum_{i=1}^{n} \frac{\partial f(x^{1}, \dots, x^{n})}{\partial x^{i}} dx^{i}.$$
 (2.13)

 $^{^6}$ Coefficients of expansion are functions, elements of algebra of functions, not numbers ,elements of field. To be more careful, these vector fields are basis of the *module* of vector fields on \mathbf{E}^3

The value of 1-form df on an arbitrary vector field (2.12) is equal to

$$df(\mathbf{A}) = \sum_{i=1}^{n} \frac{\partial f(x^{1}, \dots, x^{n})}{\partial x^{i}} dx^{i}(\mathbf{A}) = \sum_{i=1}^{n} \frac{\partial f(x^{1}, \dots, x^{n})}{\partial x^{i}} A^{i} = \partial_{\mathbf{A}} f \quad (2.14)$$

We see that value of differential of 0-form f on an arbitrary vector field \mathbf{A} is equal to directional derivative of function f with respect to the vector \mathbf{A} .

The formula (2.14) defines df in invariant way without using coordinate expansions. Later we check straightforwardly the coordinate-invariance of the definition (2.13).

Exercise Check that

$$dx^{i}(\mathbf{A}) = \partial_{\mathbf{A}}x^{i} \tag{2.15}$$

Example If f = f(x, y) is a function (0 - form) on \mathbf{E}^2 then

$$df = \frac{\partial f(x,y)}{\partial x} dx + \frac{\partial f(x,y)}{\partial y} dy$$

and for an arbitrary vector field $\mathbf{A} = \mathbf{A} = A_x \mathbf{e}_x + \mathbf{A}_y \mathbf{e}_y = A_x(x,y)\partial_x + A_y(x,y)\partial_y$

$$df(\mathbf{A}) = \frac{\partial f(x,y)}{\partial x} dx(\mathbf{A}) + A_y(x,y) \frac{\partial f(x,y)}{\partial y} dy(\mathbf{A}) = \mathbf{A}_x(x,y) \frac{\partial f(x,y)}{\partial x} + A_y(x,y) \frac{\partial f(x,y)}{\partial y} = \partial_{\mathbf{A}} f.$$

Example Find the value of 1-form $\omega = df$ on the vector field $\mathbf{A} = x\partial_x + y\partial_y$ if $f = \sin(x^2 + y^2)$.

 $\omega(\mathbf{A}) = df(\mathbf{A})$. One can calculate it using formula (2.13) or using formula (2.14).

Solution (using (2.13)):

$$\omega = df = \frac{\partial f}{\partial x}dx + \frac{\partial f}{\partial y}dy = 2x\cos(x^2 + y^2)dx + 2y\cos(x^2 + y^2)dy.$$

$$\omega(\mathbf{A}) = 2x\cos(x^2 + y^2)dx(\mathbf{A}) + 2y\cos(x^2 + y^2)dy(\mathbf{A}) =$$

$$2x\cos(x^2+y^2)A_x + 2y\cos(x^2+y^2)dA_y = 2(x^2+y^2)\cos(x^2+y^2),$$

Another solution (using (2.14))

$$df(\mathbf{A}) = \partial_{\mathbf{A}} f = A_x \frac{\partial f}{\partial x} + A_x \frac{\partial f}{\partial y} = 2(x^2 + y^2)\cos(x^2 + y^2).$$

See other examples in Homeworks.

2.3.4 Exact forms

1-form ω is called exact if there exists a function f such that $\omega = df$.

E.g. a form $\omega = xdy + ydx = d(xy)$ is an exact form, a form $\omega = xdx + ydy$ is also exact form: $\omega = xdx + ydy = d\left(\frac{x^2}{2} + \frac{y^2}{2}\right)$.

Of course not any form is an exact form (see exercises in Homeworks.) E.g. 1-form $\omega = xdy$ is not an exact form. Indeed suppose that this is an exact form, i.e. $xdy = dF = F_x dx + F_y dy$, then $F_y = x$ and $F_x = 0$. We see that on one hand $F_{xy} = (f_x)_y = 0$ and on the other hand $f_{yx} = (f_y)_x = 1$. Contradiction.

Another example:

Example Consider 1-form $\omega = 2ydx + xdy$ and another 1-form $\sigma = x\omega = 2xydx + x^2dy$. One can easy to see that 1-form ω is not exact 1-form, and 1-form σ is an exact 1-form.

Later we will see that exact 1-forms are easy to integrate over curves.

2.3.5 Differential forms in arbitrary coordinates

We learnt how to calculate directional derivative of functions along vector fields, we learnt how to calculate values of differential 1-forms on vector fields, We did the calculations in $Cartesian\ coordinates$ in \mathbf{E}^n (In examples above we considered Cartesian coordinates (x,y) in \mathbf{E}^2 .) One of the reasons why differential forms are so important is that in fact our calculations may be performed in arbitrary coordinates. The power of applications of differential forms is that the constructions are invariant, they do not depend on choice of coordinates we are working with. Here we consider just few examples (see for more details Appendices.).

Example Calculate the value of differential forms $\omega = xdy - ydx$, $\sigma = xdx + ydy$ on vector fields $\mathbf{A} = x\partial_x + y\partial_y$ and $\mathbf{B} = x\partial_y - y\partial_x$

We solved this execise (see 2 in Homework 5). Now we will do the same exerices but in *polar coordinates*

$$\begin{cases} x = r \cos \varphi \\ y = r \sin \varphi \end{cases} \Leftrightarrow \begin{cases} r = \sqrt{x^2 + y^2} \\ \varphi = \arctan \frac{y}{x} \end{cases} . \tag{2.16}$$

We have for differential forms

$$\omega = xdy - ydx = r\cos\varphi d(r\sin\varphi) - r\sin\varphi d(r\cos\varphi) =$$

 $r\cos\varphi(\sin\varphi dr + r\cos\varphi d\varphi) - r\sin\varphi(\cos\varphi dr - r\sin\varphi d\varphi) = r^2(\cos^2\varphi + \sin^2\varphi)d\varphi = r^2d\varphi.$ (2.17)

and

$$\sigma = xdx + ydy = r\cos\varphi d(r\cos\varphi) + r\sin\varphi d(r\sin\varphi) =$$

 $r\cos\varphi(\cos\varphi dr - r\sin\varphi d\varphi) + r\sin\varphi(\sin\varphi dr + r\cos\varphi d\varphi) = r(\cos^2\varphi + \sin^2\varphi) dr = rdr.$ (2.18)

and for vector fields we have:

$$\mathbf{A} = x\partial_x + y\partial_y = x\frac{\partial}{\partial x} + y\frac{\partial}{\partial y} = x\left(\frac{\partial r}{\partial x}\frac{\partial}{\partial r} + \frac{\partial \varphi}{\partial x}\frac{\partial}{\partial \varphi}\right) + y\left(\frac{\partial r}{\partial y}\frac{\partial}{\partial r} + \frac{\partial \varphi}{\partial y}\frac{\partial}{\partial \varphi}\right) = x\left(\frac{x}{r}\frac{\partial}{\partial r} - \frac{y}{x^2 + y^2}\frac{\partial}{\partial \varphi}\right) + y\left(\frac{y}{r}\frac{\partial}{\partial r} + \frac{y}{x^2 + y^2}\frac{\partial}{\partial \varphi}\right) = r\frac{\partial}{\partial r}, \quad (2.19)$$

and

$$\mathbf{B} = x\partial_y - y\partial_x = x\frac{\partial}{\partial y} - y\frac{\partial}{\partial x} = x\left(\frac{\partial r}{\partial y}\frac{\partial}{\partial r} + \frac{\partial \varphi}{\partial y}\frac{\partial}{\partial \varphi}\right) - y\left(\frac{\partial r}{\partial x}\frac{\partial}{\partial r} + \frac{\partial \varphi}{\partial x}\frac{\partial}{\partial \varphi}\right) = x\left(\frac{y}{r}\frac{\partial}{\partial r} + \frac{x}{x^2 + y^2}\frac{\partial}{\partial \varphi}\right) - y\left(\frac{x}{r}\frac{\partial}{\partial r} - \frac{x}{x^2 + y^2}\frac{\partial}{\partial \varphi}\right) = \frac{\partial}{\partial \varphi}.$$
 (2.20)

We have

$$\omega(\mathbf{A}) = r^2 d\varphi \left(r \frac{\partial}{\partial r} \right) = 0,,$$

$$\omega(\mathbf{B}) = r^2 d\varphi \left(\frac{\partial}{\partial \varphi} \right) = r^2 = x^2 + y^2,$$

$$\sigma(\mathbf{A}) = r dr \left(r \frac{\partial}{\partial r} \right) = r^2,,$$

$$\sigma(\mathbf{B}) = r dr \left(\frac{\partial}{\partial \varphi} \right) = 0,$$

Now we see we can calculate the values of differential forms on vector fields in Cartesian and in polar coordinates, and answers are the same:

We used in the calculations for vector fields the chain rule. Pay attention how useful for these calculations are notations of vector fields as derivatives. We also used the matrices for partial derivatives of changing of coordinates (polarcoordinates) Calculations for vector fields were not very easy, but answers are beautiful⁷!

2.4 Integration of differential 1-forms over curves

Differential forms are natural objects of integration over surfaces. We consider integration of differential 1-forms over curves.

Let $\omega = \omega_1(x^1, \dots, x^n)dx^1 + \dots + \omega_1(x^1, \dots, x^n)dx^n = \sum_{i=1}^n \omega_i dx^i$ be an arbitrary 1-form in \mathbf{E}^n

and $C: \mathbf{r} = \mathbf{r}(t), t_1 \le t \le t_2$ be an arbitrary smooth curve in \mathbf{E}^n .

One can consider the value of one form ω on the velocity vector field $\mathbf{v}(t) = \frac{d\mathbf{r}(t)}{dt}$ of the curve:

$$\omega(\mathbf{v}(t)) = \sum_{i=1}^{n} \omega_i \left(x^1(t), \dots, x^n(t) \right) dx^i(\mathbf{v}(t)) = \sum_{i=1}^{n} \omega_i \left(x^1(t), \dots, x^n(t) \right) \frac{dx^i(t)}{dt}$$

We define now integral of 1-form ω over the curve C.

Definition The integral of the form $\omega = \omega_1(x^1, \dots, x^n)dx^1 + \dots + \omega_n(x^1, \dots, x^n)dx^n$ over the curve C: $\mathbf{r} = \mathbf{r}(t)$ $t_1 \le t \le t_2$ is equal to the integral of the function $\omega(\mathbf{v}(t))$ over the interval $t_1 \le t \le t_2$:

$$\int_{C} \omega = \int_{t_{1}}^{t_{2}} \omega(\mathbf{v}(t)) dt = \int_{t_{1}}^{t_{2}} \left(\sum_{i=1}^{n} \omega_{i} \left(x^{1}(t), \dots, x^{n}(t) \right) \frac{dx^{i}(t)}{dt} \right) dt . \quad (2.22)$$

Proposition The integral $\int_C \omega$ does not depend on the choice of coordinates on \mathbf{E}^n . It does not depend (up to a sign) on parameterisation of the curve: if C: $\mathbf{r} = \mathbf{r}(t)$ $t_1 \leq t \leq t_2$ is a curve and $t = t(\tau)$ is an arbitrary reparameterisation, i.e. new curve C': $\mathbf{r}'(\tau) = \mathbf{r}(t(\tau))$ $\tau_1 \leq \tau \leq \tau_2$, then $\int_C \omega = \pm \int_C' \omega$:

$$\int_{C} \omega = \int_{C'} \omega, \quad \text{if orientation is not changed, i.e. if } t'(\tau) > 0$$

⁷Paris, vaut bien une messe!

and

$$\int_C \omega = -\int_{C'} \omega, \quad \text{if orientaion is changed, i.e. if } t'(\tau) < 0$$

If reparameterisation changes the orientation then starting point of the curve becomes the ending point and vice versa.

Proof of the Proposition Show that integral does not depend (up to a sign) on the parameterisation of the curve. Let $t(\tau)$ ($\tau_1 \leq t \leq \tau_2$) be reparameterisation. We come to the new curve C': $\mathbf{r}'(\tau) = \mathbf{r}(t(\tau))$. Note that the new velocity vector $\mathbf{v}'(\tau) = \frac{dr(t(\tau))}{d\tau} = t'(\tau)\mathbf{v}(t(\tau))$. Hence $\omega(\mathbf{v}'(\tau)) = w(\mathbf{v}(t(\tau)))t'(\tau)$. For the new curve C'

$$\int_{C'} \omega = \int_{\tau_1}^{\tau_2} \omega(\mathbf{v}'(\tau)) d\tau = \int_{\tau_1}^{\tau_2} \omega(\mathbf{v}(t(\tau))) \frac{dt(\tau)}{d\tau} d\tau = \int_{t(\tau_1)}^{t(\tau_2)} \omega(\mathbf{v}(t)) dt$$

 $t(\tau_1) = t_1$, $t(\tau_2) = t_2$ if reparameterisation does not change orientation and $t(\tau_1) = t_2$, $t(\tau_2) = t_1$ if reparameterisation changes orientation.

Hence $\int_{C'} w = \int_{t_1}^{t_2} \omega(\mathbf{v}(t)) dt = \int_{C} \omega$ if orientation is not changed and $\int_{C'} w = \int_{t_2}^{t_1} \omega(\mathbf{v}(t)) dt = -\int_{t_1}^{t_2} \omega(\mathbf{v}(t)) dt = -\int_{C} \omega$ is orientation is changed.

Example

Let

$$\omega = a(x, y)dx + b(x, y)dy$$

be 1-form in \mathbf{E}^2 (x, y)—are usual Cartesian coordinates). Let C: $\mathbf{r} = \mathbf{r}(t)$ $\begin{cases} x = x(t) \\ y = y(t) \end{cases}$, $t_1 \le t \le t_2$ be a curve in \mathbf{E}^2 .

Consider velocity vector field of this curve

$$\mathbf{v}(t) = \frac{d\mathbf{r}(t)}{dt} = \begin{pmatrix} v_x(t) \\ v_y(t) \end{pmatrix} = \begin{pmatrix} x_t(t) \\ y_t(t) \end{pmatrix} = x_t \partial_x + y_t \partial_y$$
 (2.23)

$$(x_t = \frac{dx(t)}{dt}, y_t = \frac{dy(t)}{dt}).$$

One can consider the value of one form ω on the velocity vector field $\mathbf{v}(t)$ of the curve: $\omega(\mathbf{v}) = a(x(t), y(t))dx(\mathbf{v}) + b(x(t), y(t))dy(\mathbf{v}) =$

$$a(x(t), y(t))x_t(t) + b(x(t), y(t))y_t(t)$$
.

The integral of the form $\omega = a(x, y)dx + b(x, y)dy$ over the curve C: $\mathbf{r} = \mathbf{r}(t)$ $t_1 \le t \le t_2$ is equal to the integral of the function $\omega(\mathbf{v}(t))$ over the

interval $t_1 \leq t \leq t_2$:

$$\int_{C} \omega = \int_{t_{1}}^{t_{2}} \omega(\mathbf{v}(t))dt = \int_{t_{1}}^{t_{2}} \left(a(x(t), y(t)) \frac{dx(t)}{dt} + b(x(t), y(t)) \frac{dy(t)}{dt} \right) dt.$$
(2.24)

Example Consider an integral of the form $\omega = 3dy + 3y^2 dx$ over the curve $C : \mathbf{r}(t) \begin{cases} x = \cos t \\ y = \sin t \end{cases}$, $0 \le t \le \pi/2$. (C is the arc of the circle $x^2 + y^2 = 1$ defined by conditions $x, y \ge 0$).

Velocity vector $\mathbf{v}(t) = \frac{d\mathbf{r}(t)}{dt} = \begin{pmatrix} v_x(t) \\ v_y(t) \end{pmatrix} = \begin{pmatrix} x_t(t) \\ y_t(t) \end{pmatrix} = \begin{pmatrix} -\sin t \\ \cos t \end{pmatrix}$. The value of the form on velocity vector is equal to

 $\omega(\mathbf{v}(t)) = 3y^2(t)v_x(t) + 3v_y(t) = 3\sin^2 t(-\sin t) + 3\cos t = 3\cos t - 3\sin^3 t$ and

$$\int_C \omega = \int_0^{\frac{\pi}{2}} \omega(\mathbf{v}(t)) dt = \int_0^{\frac{\pi}{2}} (3\cos t - 3\sin^3 t) dt = 3\left(\sin t + \cos t - \frac{\cos^3 t}{3}\right) \Big|_0^{\frac{\pi}{2}}$$

Change parameterisation: $t=2\tau$. Reparameterisation does not change orientation $(t_{\tau}=2>0)$, hence the answer has to be the same. Check it: We come to C': $\mathbf{r}(t(\tau))$ $\begin{cases} x=\cos 2\tau \\ y=\sin 2\tau \end{cases}$, $0 \le \tau \le \pi/4$. (C' has the same locus.)

Velocity vector $\mathbf{v}(\tau) == \begin{pmatrix} -2\sin 2\tau \\ 2\cos 2\tau \end{pmatrix}$. The value of the form on velocity vector is equal to

$$\omega(\mathbf{v}(\tau)) = 3y^2(t)v_x(\tau) + 3v_y(\tau) = 3\sin^2 2\tau(-2\sin 2\tau) + 6\cos 2\tau = 6\cos 2\tau - 6\sin^3 2\tau$$
 and

$$\int_{C}' \omega = \int_{0}^{\frac{\pi}{4}} \omega(\mathbf{v}(\tau)) d\tau = \int_{0}^{\frac{\pi}{4}} (6\cos 2\tau - 6\sin^{3} 2\tau) d\tau = 3\left(\sin 2\tau + \cos 2\tau - \frac{\cos^{3} 2\tau}{3}\right)\Big|_{0}^{\frac{\pi}{4}}$$

One can see that the answer is the same,

Example Now consider the integral of 1-form over the curve C which is the upper half of the circle $x^2+y^2=R^2$: C: $\begin{cases} x^2+y^2=1\\ y\geq 0 \end{cases}$. Curve is

given as an image. We have the image of the curve not the parameterised curve. We have to define a parameterisation ourself.

We consider three different parameterisations of this curve. Sure to calculate the integral it suffices to calculate $\int_C \omega$ in an arbitrary given parameterisation $\mathbf{r} = \mathbf{r}(t)$ of the curve C, then note that for an arbitrary reparameterisation $t = t(\tau)$, the integral will remain the same or it will change a sign depending on the reparameterisation $t = t(\tau)$ preserves orientation or not.

$$C_1 - \mathbf{r}_1(t) : \begin{cases} x = R \cos t \\ y = R \sin t \end{cases}, 0 \le t \le \pi, \quad C_2 - \mathbf{r}_2(\tau) : \begin{cases} x = R \cos \Omega \tau \\ y = R \sin \Omega \tau \end{cases}, 0 \le \tau \le \frac{\pi}{\Omega}, (\Omega > 0)$$

and

$$C_3 - \mathbf{r}_3(u):$$

$$\begin{cases} x = u \\ y = \sqrt{R^2 - u^2} \end{cases}, -R \le u \le R,,$$
 (2.25)

All these curves are the same image. If $\Omega=1$ the second curve coincides with the first one. First and second curve have the same orientation (reparameterisation $t=\Omega\tau$) The third curve has orientation opposite to first and second (reparameterisation $u=\cos t$, the derivative $\frac{d\cos t}{dt}<0$).

second (reparameterisation $u = \cos t$, the derivative $\frac{d \cos t}{dt} < 0$). Calculate integrals $\int_{C_1} \omega$, $\int_{C_2} \omega$, $\int_{C_3} \omega$ in the case if $\omega = xdy - ydx$ and check straightforwardly that these integrals coincide if orientation is the same or they have different signs if orientation is opposite.

I Calculation for the first curve C_1 .

We have $\mathbf{v} = x_t \partial_x + y_t \partial_y$. For the form $\omega = x dy - y dx \ \omega(\mathbf{v}) = x y_t - y x_t = R \cos t (R \sin t) - R \sin t (-R \cos t) = R^2$. We have

$$\int_{C_1} \omega = \int_0^{\pi} (xy_t - yx_t)dt = \int_0^{\pi} R^2 dt = \pi R^2.$$

II Calculation for the second curve C_2 .

We have $\mathbf{v} = x_{\tau} \partial_x + y_{\tau} \partial_y = -R\Omega \sin \Omega \tau + R\Omega \cos \Omega \tau$. Thus $\omega(\mathbf{v}) = (xdy - ydx)(\mathbf{v}) = xy_{\tau} - yx_{\tau} = R \cos \Omega \tau R\Omega \cos \Omega \tau - R \sin \Omega \tau (-R\Omega \sin \Omega \tau = R^2\Omega$. We have

$$\int_{C_2} \omega = \int_0^{\frac{\pi}{\Omega}} (xy_\tau - yx_\tau) d\tau = \int_0^{\frac{\pi}{\Omega}} R^2 \Omega = \pi R^2.$$

These answers coincide: both parameterisation have the same orientation.

III Calculation for the third curve C_3 .

We have $\mathbf{v} = x_u \partial_x + y_u \partial_y = \partial_x - \frac{u \partial_y}{\sqrt{R^2 - u^2}}$. Thus $\omega(\mathbf{v}) = (x dy - y dx)(\mathbf{v}) = x y_u - y x_u = -\frac{R^2}{\sqrt{R^2 - u^2}}$. We have that for the third parameterisation:

$$\int_{C_3} \omega = \int_{-R}^{R} (xy_u - yx_u) du = \int_{-R}^{R} \left(-\frac{R^2}{\sqrt{R^2 - u^2}} \right) du =$$

$$-2R^2 \int_{0}^{R} \frac{du}{\sqrt{R^2 - u^2}} = -2R^2 \int_{0}^{1} \frac{dz}{\sqrt{1 - z^2}} = -\pi R^2.$$

We see that the sign is changed.

Note that one can consider the integral of the form $\omega = xdy - ydx$ over the semicircle in polar coordinates instead Cartesian coordinates. We have that in polar coordinates semicircle is $\begin{cases} r(t) = R \\ \varphi(t) = t \end{cases}, \ 0 \leq t \leq \pi. \text{ The form } \omega = xdy - ydx = r\cos\varphi d(r\sin\varphi) - r\sin\varphi d(r\sin\varphi) = r^2d\varphi \text{ and } \mathbf{v}(t) = (r_t,\varphi_t) = (0,1), \text{ i.e. } \mathbf{v}(t) = \partial_\varphi. \text{ We have that } \omega(\mathbf{v}(t)) = r(t)^2d\varphi(\partial_\varphi) = R^2. \text{ Hence } \int_C \omega = \int_0^\pi R^2dt = \pi R^2. \text{ Answer is the same: The value of integral does not change if we change coordinates in the plane.}$

For other examples see Homeworks.

2.5 Integral over curve of exact form

Recall that 1-form ω is called exact if there exists a function f such that $\omega = df$. Of course not any form is an exact form (see exercises in Homeworks.) see subsection 2.3.4 above and exercises in Homeworks.).

Theorem

Let ω be an exact 1-form in \mathbf{E}^n , $\omega = df$.

Then the integral of this form over an arbitrary curve C: $\mathbf{r} = \mathbf{r}(t)$ $t_1 \leq t \leq t_2$ is equal to the difference of the values of the function f at starting and ending points of the curve C:

$$\int_{C} \omega = f \big|_{\partial C} = f(\mathbf{r}_2) - f(\mathbf{r}_1), \quad \mathbf{r}_1 = \mathbf{r}(t_1), \mathbf{r}_2 = \mathbf{r}(t_2). \tag{2.26}$$

Proof: According definition of the integral of 1-form over curve we have that $\int_C df = \int_{t_1}^{t_2} df(\mathbf{v}(t))dt$. On the other hand according definition of directional derivative (2.9) we have that

$$df(\mathbf{v}(t)) = \partial_{\mathbf{v}}(t)f(\mathbf{r})\big|_{\mathbf{r}(t)} = \frac{d}{dt}(f(\mathbf{r}(t))),$$

hence we come to

$$\int_{C} df = \int_{t_{1}}^{t_{2}} df(\mathbf{v}(t))dt = \int_{t_{1}}^{t_{2}} \frac{d}{dt} f(\mathbf{r}(t))dt = f(\mathbf{r}(t))|_{t_{1}}^{t_{2}}.$$
 The proof is finished.

Example Calculate an integral of the form $\omega = 3x^2(1+y)dx + x^3dy$ over the arc of the semicircle $x^2 + y^2 = 1, y > 0$.

One can calculate the integral naively using just the formula (2.24): Choose a parameterisation of C,e.g., $x = \cos t$, $y = \sin t$, then $\mathbf{v}(t) = -\sin t\partial_x + \cos t\partial_x$ and $\omega(\mathbf{v}(t)) = (3x^2(1+y)dx + x^3dy)(-\sin t\partial_x + \cos t\partial_y) = -3\cos^2 t(1+\sin t)\sin t + \cos^3 t \cdot \cos t$ and

$$\int_{C} \omega = \int_{0}^{\pi} (-3\cos^{2}t \sin t - 3\cos^{2}t \sin^{2}t + \cos^{4}t) dt = \dots$$

Calculations are boring and they are not short.

On the other hand for the form $\omega = 3x^2(1+y)dx + x^3dy$ one can calculate the integral in a much more efficient way noting that it is an exact form:

$$\omega = 3x^{2}(1+y)dx + x^{3}dy = d\left(x^{3}(1+y)\right)$$
(2.27)

Hence it follows from the Theorem that

$$\int_{C} \omega = f(\mathbf{r}(\pi)) - f(\mathbf{r}(0)) = x^{3}(1+y) \Big|_{x=1,y=0}^{x=-1,y=0} = -2$$
 (2.28)

Remark If we change the orientation of curve then the starting point becomes the ending point and the ending point becomes the starting point.—The integral changes the sign in accordance with general statement, that integral of 1-form over parameterised curve is defined up to reparameterisation.

Corollary The integral of an exact form over an arbitrary closed curve is equal to zero.

Proof. According to the Theorem $\int_C \omega = \int_C df = f|_{\partial C} = 0$, because the starting and ending points of closed curve coincide.

Example. Calculate the integral of 1-form $\omega = x^5 dy + 5x^4 y dx$ over the ellipse $x^2 + \frac{y^2}{9} = 1$.

The form $\omega = x^5 dy + 5x^4 y dx$ is exact form because $\omega = x^5 dy + 5x^4 y dx = d(x^5 y)$. Hence the integral over ellipse is equal to zero, because it is a closed curve.

2.6 Calculation of integral of 1-form over curve in arbitrary coordinates

In subsection 2.3.5 we considered examples of calculations with differential forms in arbitrary coordinates. In the subsection 2.4 we defined the integral of 1-form over curve. In fact the definition is valid for an arbitrary coordinates,⁸. We have not time to go in details in this question, and in this subsection we just consider example of calculations of integral for differential form $\omega = xdy - ydx$ considered in subsection?? in polar coordinates.

Let
$$C$$
:
$$\begin{cases} x = x(t) \\ y = y(t) \end{cases}$$
, $t_1 < t < t_2$ be a curve in \mathbf{E}^2 . We will demostrate

explicitly that the result of claculation of $\int_C \omega$ in polar coordinates coicide with the result of calculation of this integral in Cartesian coordinates. In Cartesian coordinates the calculations for $\int_C \omega$ are the following (see

subsection 2.4 or Homeworks):
$$\mathbf{v}(t) = \begin{pmatrix} x_t \\ y_t \end{pmatrix} = x_t \partial_x + y_t \partial_y, \ \omega(\mathbf{v}(t)) = (xdy - ydx)(x_t \partial_x + y_t \partial_y) = xy_t - yx_t \text{ and}$$

$$\int_{C} \omega = \int_{t_1}^{t_2} \omega(\mathbf{v}(t)) dt = \int_{t_1}^{t_2} \left(x(t) \frac{dy(t)}{dt} - y(t) \frac{dx(t)}{dt} \right) dt.$$
 (2.29)

Calculate the same integral in polar coordinates: The differential form $\omega = xdy - ydx$ in polar coordinates has an appearance $\omega = r^2d\varphi$ (see subsection 2.3.5.)

$$C \colon \begin{cases} x = x(t) \\ y = y(t) \end{cases} \Rightarrow \begin{cases} r = r(t) \\ \varphi = \varphi(t) \end{cases}, t_1 < t < t_2.$$

For velocity vector

$$\mathbf{v}(t) = \begin{pmatrix} \frac{dr(t)}{dt} \\ \frac{d\varphi(t)}{dt} \end{pmatrix} = \frac{dr(t)}{dt} \frac{\partial}{\partial r} + \frac{d\varphi(t)}{dt} \frac{\partial}{\partial \varphi},$$

and

$$\omega(\mathbf{v}) = r^2 d\varphi \left(r_t \partial_r + \varphi_t \partial_\varphi \right) = r^2 \varphi_t \,.$$

(as usual we use on an equal footing notations $\frac{\partial}{\partial r} \leftrightarrow \partial_r$, $\frac{\partial}{\partial \varphi} \leftrightarrow \partial_{\varphi}$, $\frac{dr}{dt} \leftrightarrow r_t$, and $\frac{d\varphi}{dt} \leftrightarrow \varphi_t$). We come to

$$\int_C \omega = \int_{t_1}^{t_2} \omega(\mathbf{v}(t))dt = \int_{t_1}^{t_2} \left(r^2(t)\frac{d\varphi(t)}{dt}\right)dt. \tag{2.30}$$

⁸this is why the differential forms are so powerful in geometry

Show straightforwardly that integrals (2.30) and (2.29) coincide. We have that $\begin{cases} x = r\cos\varphi \\ y = r\sin\varphi \end{cases}$, thus we have from (2.29) that

$$\int_{C} \omega = \int_{t_{1}}^{t_{2}} \omega(\mathbf{v}(t))dt = \int_{t_{1}}^{t_{2}} \left(x(t) \frac{dy(t)}{dt} - y(t) \frac{dx(t)}{dt} \right) dt =$$

$$\int_{t_{1}}^{t_{2}} \left(r(t) \cos \varphi(t) \frac{d}{dt} \left(r(t) \sin \varphi(t) \right) - r(t) \sin \varphi(t) \frac{d}{dt} \left(r(t) \cos \varphi(t) \right) \right) dt =$$

$$\int_{t_{1}}^{t_{2}} \left(r \cos \varphi \left(r_{t} \sin \varphi + r \cos \varphi \varphi_{t} \right) - r \sin \varphi \left(r_{t} \cos \varphi - r \sin \varphi \varphi_{t} \right) \right) dt =$$

$$\int_{t_{1}}^{t_{2}} r^{2} \left(\cos^{2} \varphi + \sin^{2} \varphi \right) \varphi_{t} = \int_{t_{1}}^{t_{2}} \left(r^{2}(t) \frac{d\varphi(t)}{dt} \right) dt.$$

See another examples of caclulatios of integrals in polar coordinates in Homework 6.

3 Conic sections and Projective Geometry

In this section we consider very famous and important curves, ellipses, hyperbolas and parabolas.

We first consider geometrical definitions of these curves without using the analytical methods, then we will show that in Cartesian coordinates these curves can be expressed by well-known standard formulae:

an ellipse,
$$\frac{x}{a^2} + \frac{y^2}{b^2} = 1$$
, $a \ge b > 0$,

a hyperbola, $\frac{x}{a^2} - \frac{y^2}{b^2} = 1$,

• and parabola, $y^2 = 2px$.

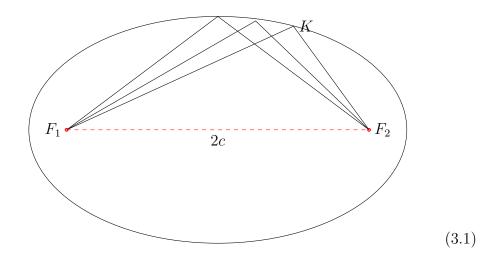
We study properties of these curves, and will show that these curves are sections of conic surfaces (it is why they are called conic sections.)

Finally we will consider some elements of Projective geometry and we will look at these curves from the point of view of the projective geometry.

3.1 Geometrical definitions of conic sections

3.1.1 Ellipse on the Euclidean plane

Ellipse on the plane is the locus of all the points such that the sum of distances from these points to two fixed points F_1, F_2 is equal to given constant.



$$|F_1F_2| = 2c$$
, $a > c > 0$.

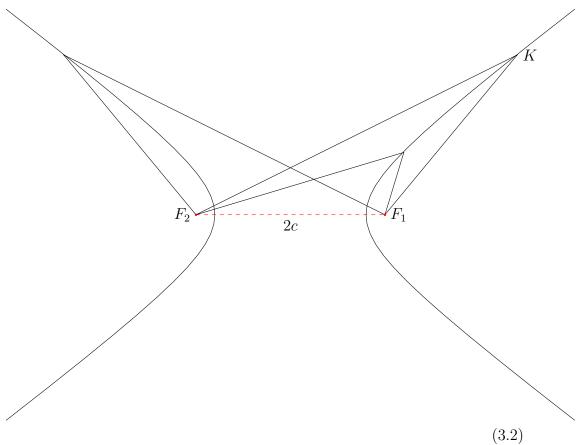
Ellipse =
$$\{K: |KF_1| + |KF_2| = 2a\}$$

 F_1 , F_2 are foci of ellipse.

If c = 0, ellipse becomes circle.

3.1.2 Hyperbola on the Euclidean plane

Hyperbola on the plane is the locus of all the points such that the difference of distances from these points to two fixed points F_1, F_2 is equal to given constant.



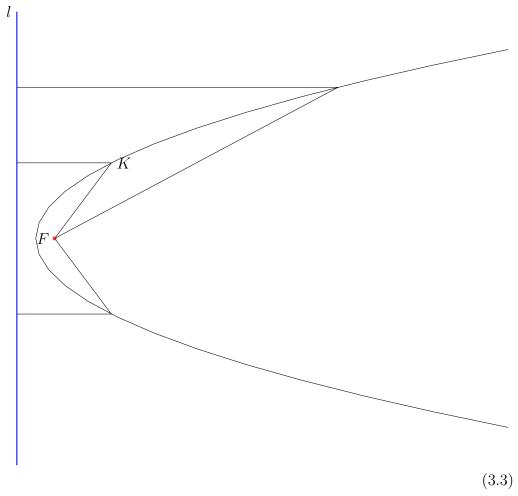
Hyperbola = $\{K: ||KF_1| - |KF_2|| = 2a\}$

The points F_1, F_2 are foci of hyperbola.

Remark Notice that for ellipse we denote "left" focus F_1 and 'right focus" F_2 , and for hyperbola vice versa: "right" focus F_1 and 'left" focus F_2 . The difference between these two notations will be important only when we will consider analytical definitions of hyperbola, and we will note it later.

3.1.3 Parabola on the Euclidean plane

Parabola on the plane is the locus of all the points such that they are at the same distance from the given point F and the given line l.



 ${\it Parabola} = \{K\colon \quad d(K,l) = |KF|\}$

The point F is called the focus of the parabola, and the line l is called the directrix of the parabola.

One can consider directrices for hyperbola and ellipse also. See later subsection??

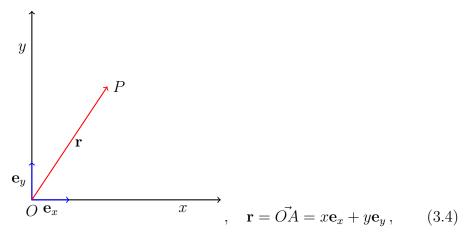
3.2 Cartesian and affine coordinates in the plane and in the space

3.2.1 Cartesian coordinates in plane E^2

Consider affine space \mathbf{E}^2 .

Recall that fixing the point in affine space we come to the vector space \mathbf{E}^2 (see 1.2).

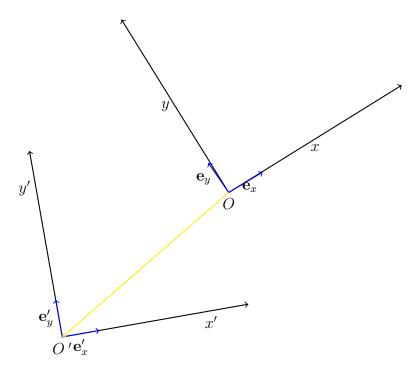
Let $\{\mathbf{e}_x, \mathbf{e}_y\}$ be an orthonormal basis attached at a point O. The point O is an origin of this frame of reference. For every point $A \in \mathbf{E}^2$ we have



$$|\mathbf{e}_x| = |\mathbf{e}_y| = 1$$
 and $\angle (\mathbf{e}_x, \mathbf{e}_y) = \frac{\pi}{2}$, i.e. $(\mathbf{e}_x, \mathbf{e}_x) = (\mathbf{e}_y, \mathbf{e}_y) = 1$, $(\mathbf{e}_x, \mathbf{e}_y) = 0$.

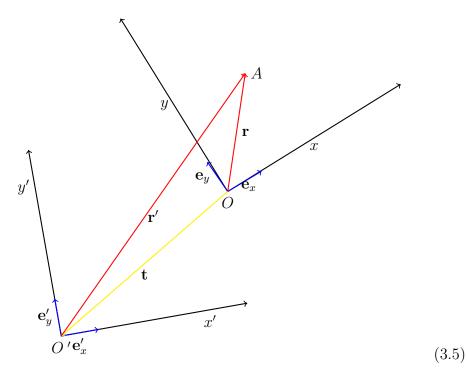
Coordinates x, y are Cartesian coordinates of the point P with respect to frame XOY

Consider now two frames of references. Let $\{\mathbf{e}_x, \mathbf{e}_y\}$ be an orthonormal basis attached at the point O, and let $\{\mathbf{e}'_x, \mathbf{e}'_y\}$ be an orthonormal basis attached at the point O'. O.



Let A be an arbitrary point in the plane \mathbf{E}^2 . We denote:

$$\mathbf{r} = \vec{OP}, \mathbf{r}' = \vec{O'P}, \text{ and } \mathbf{t} = \vec{OO'}.$$



We have

$$\mathbf{r} = \mathbf{t} + \mathbf{r}'$$

i.e.

$$\underbrace{(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} x \\ y \end{pmatrix}}_{\mathbf{r}} = \underbrace{(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} a \\ b \end{pmatrix}}_{\mathbf{t}} + \underbrace{(\mathbf{e}_x', \mathbf{e}_y') \begin{pmatrix} x' \\ y' \end{pmatrix}}_{\mathbf{r}'}$$

, i.e.

$$x\mathbf{e}_x + y\mathbf{e}_y = \mathbf{r} = \mathbf{t} + \mathbf{r}' = a\mathbf{e}_x + b\mathbf{e}_y + x'\mathbf{e}_x' + y'\mathbf{e}_y'$$

The bases $\{\mathbf{e}_x, \mathbf{e}_y\}$, $\{\mathbf{e}_x', \mathbf{e}_y'\}$ are related by orthogonal matrix:

$$(\mathbf{e}_x', \mathbf{e}_y') = (\mathbf{e}_x, \mathbf{e}_y) \underbrace{\begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix}}_{\text{orthog. matr.}}.$$
(3.6)

$$(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} x \\ y \end{pmatrix} = (\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} a \\ b \end{pmatrix} + (\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix}$$

i.e.

$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \\ b \end{pmatrix} + \underbrace{\begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix}}_{\text{orthog.matrix}} \begin{pmatrix} x' \\ y' \end{pmatrix} .$$

$$\begin{cases} x = a + p_{11}x' + p_{12}y' \\ y = b + p_{21}x' + p_{22}y' \end{cases}$$

One can say that changing of coordinates is translation + rotation (or rotation+ reflection.)

Example Consider in (3.6) an orthogonal matrix

$$P = \begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix} = \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix}, \quad (P^+ \circ P = \mathbf{id}, \quad \text{and } \det P = 1)$$

(P preserves orientation. This is rotation on the angle φ)

Another example:

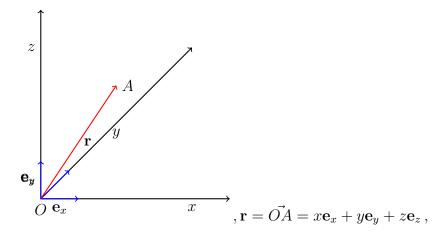
Example Consider in (3.9), an orthogonal matrix

$$P = \begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{pmatrix} , \quad (P^+ \circ P = \mathbf{id} \,, \quad \text{and } \det P = -1)$$

(P changes orientation. This is rotation and reflection)

3.2.2 Cartesian coordinates in E^3

The analogous considerations in E^3 :



$$|\mathbf{e}_x| = |\mathbf{e}_y| = |\mathbf{e}_z| = 1 \text{ and } \angle(\mathbf{e}_x, \mathbf{e}_y) = \angle(\mathbf{e}_x, \mathbf{e}_z) = \angle(\mathbf{e}_y, \mathbf{e}_z)) = \frac{\pi}{2},$$

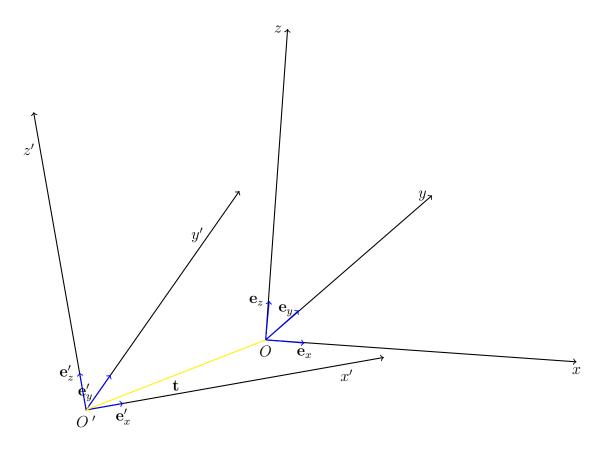
i.e.

$$(\mathbf{e}_x, \mathbf{e}_x) = (\mathbf{e}_y, \mathbf{e}_y) = 1, (\mathbf{e}_z, \mathbf{e}_z) = 1, (\mathbf{e}_x, \mathbf{e}_y) = (\mathbf{e}_x, \mathbf{e}_z) = (\mathbf{e}_y, \mathbf{e}_z) = 0.$$

We have that $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$ is an orthonormal basis and x, y, z are Cartesian coordinates of point A with respect to frame OXYZ.

Now coonsider two different Cartesian coordinates

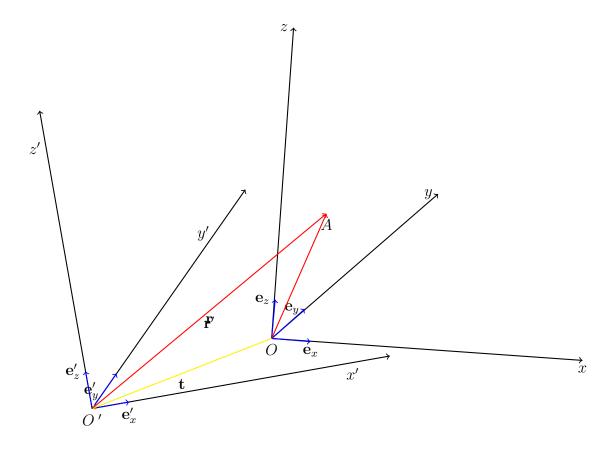
Let in \mathbf{E}^3 , \mathbf{e}_x , \mathbf{e}_y , \mathbf{e}_z be an orthonormal frame with the origin at the point O and let \mathbf{e}'_x , \mathbf{e}'_y . \mathbf{e}'_z be an another orthonormal frame with the origin at the the point O':



Let A be an arbitrary point in \mathbf{E}^3 .

Denote by

$$\mathbf{r} = \vec{OA}, \mathbf{r}' = \vec{O'A}, \text{ and } \mathbf{t} = \vec{OO'}.$$



We have in the same way as in E^2 :

$$\mathbf{r} = \mathbf{t} + \mathbf{r}'$$

i.e.

$$\underbrace{(\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z) \begin{pmatrix} x \\ y \\ z \end{pmatrix}}_{\mathbf{r}} = \underbrace{(\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z) \begin{pmatrix} a \\ b \\ c \end{pmatrix}}_{\mathbf{t}} + \underbrace{(\mathbf{e}'_x, \mathbf{e}'_y, \mathbf{e}'_z) \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}}_{\mathbf{r}'}$$

, i.e.

$$x\mathbf{e}_x + y\mathbf{e}_y + z\mathbf{e}_z = \mathbf{r} = \mathbf{t} + \mathbf{r}' = a\mathbf{e}_x + b\mathbf{e}_y + c\mathbf{e}_z = x'\mathbf{e}_x' + y'\mathbf{e}_y' + z'\mathbf{e}_z'$$
.

The orthonormal bases $\{\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z\}$, $\{\mathbf{e}_x', \mathbf{e}_y', \mathbf{e}_z\}$ are related by orthogonal

matrix:

$$(\mathbf{e}'_x, \mathbf{e}'_y, \mathbf{e}'_z) = (\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z) \underbrace{\begin{pmatrix} p_{11} & p_{12} & p_{13} \\ p_{21} & p_{22} & p_{23} \\ p_{31} & p_{32} & p_{33} \end{pmatrix}}_{\text{orthog. matr.}}.$$

Thus

$$(\mathbf{e}_{x}, \mathbf{e}_{y}, \mathbf{e}_{z}) \begin{pmatrix} x \\ y \\ z \end{pmatrix} = (\mathbf{e}_{x}, \mathbf{e}_{y}, \mathbf{e}_{z}) \begin{pmatrix} a \\ b \\ c \end{pmatrix} + (\mathbf{e}_{x}, \mathbf{e}_{y}, \mathbf{e}_{z}) \begin{pmatrix} p_{11} & p_{12} & p_{13} \\ p_{21} & p_{22} & p_{23} \\ p_{31} & p_{32} & p_{33} \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} ,$$

i.e.

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} a \\ b \\ c \end{pmatrix} + \underbrace{\begin{pmatrix} p_{11} & p_{12} & p_{13} \\ p_{21} & p_{22} & p_{23} \\ p_{31} & p_{32} & p_{33} \end{pmatrix}}_{\text{orthog.matrix}} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} ,$$

$$\begin{cases} x = a + p_{11}x' + p_{12}y' + p_{13}z' \\ y = b + p_{21}x' + p_{22}y' + p_{23}z' \\ z = c + p_{31}x' + p_{32}y' + p_{33}z' \end{cases}$$
(3.7)

Recalling Euler Theorem one One can say that changing of coordinates is translation + rotation (or rotation+ reflection.)

Consider examples of transformation (3.7)

Rotation around axis Oz on the angle φ :

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos \varphi & -\sin \varphi & 0 \\ \sin \varphi & \cos \varphi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}.$$

Another example:

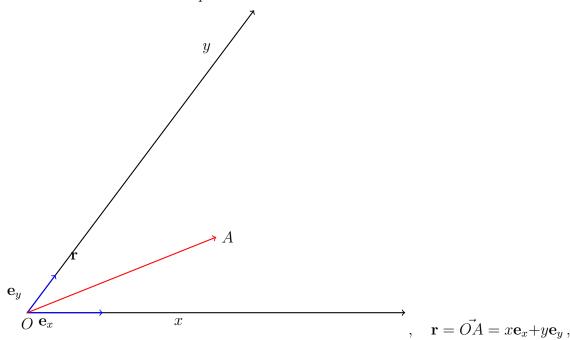
Rotation around axis Oz on the angle $\frac{\pi}{6}$ and translation:

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} a \\ b \\ c \end{pmatrix} + \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{\sqrt{3}}{2} & -\frac{1}{2} \\ 0 & \frac{1}{2} & \frac{\sqrt{3}}{2} \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} ,$$

3.2.3 Arbitrary affine coordinates

We consider in (3.9) (3.7) changing of Cartesian coordinates from one orthonormal basis to another. In the general case if bases are not orthonormal we have coordinates hich are not Cartesian, they are just so called *affine* coordinates. We consider affine coordinates in the plane \mathbf{E}^2 .

Affine coordinates of the point:



In this case $\{\mathbf{e}_x, \mathbf{e}_y\}$ is not in general an orthonormal basis: the vectors $\mathbf{e}_x, \mathbf{e}_y$ may have an arbitrary length, and the angle θ between them may be an abritrary. Sure the vectors $\mathbf{e}_x, \mathbf{e}_y$ have to be linearly independent, since $\{\mathbf{e}_x, \mathbf{e}_y\}$ is a basis, this is why $|\mathbf{e}_x| \neq 0$, $|\mathbf{e}_y| \neq 0$ and $\theta \neq 0$. In the special case if these vectors are unit vectors and angle between them is equal to $\frac{\pi}{2}$ we come to Cartesian coordinates (3.4)

Write down the formulae of changing coordinates if we change the coordinate systems.

Let OXY be coordinate frame with the origin at the point O and with basic vectors $\{\mathbf{e}_x, \mathbf{e}_y\}$, and let O'X'Y' be coordinate frame with origin at the point O' and with basic vectors $\{\mathbf{e}_x', \mathbf{e}_y'\}$, (Compare with (3.5)).

In the same way as in (3.5) we have

$$r = t + r'$$

i.e.

$$\underbrace{(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} x \\ y \end{pmatrix}}_{\mathbf{r}} = \underbrace{(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} a \\ b \end{pmatrix}}_{\mathbf{r}} + \underbrace{(\mathbf{e}_x', \mathbf{e}_y') \begin{pmatrix} x' \\ y' \end{pmatrix}}_{\mathbf{r}'}$$

, i.e.

$$x\mathbf{e}_x + y\mathbf{e}_y = \mathbf{r} = \mathbf{t} + \mathbf{r}' = a\mathbf{e}_x + b\mathbf{e}_y + x'\mathbf{e}_x' + y'\mathbf{e}_y'$$

In this case the bases $\{\mathbf{e}_x, \mathbf{e}_y\}$, $\{\mathbf{e}_x', \mathbf{e}_y'\}$ are related by non-degenerate matrix:

$$(\mathbf{e}'_x, \mathbf{e}'_y) = (\mathbf{e}_x, \mathbf{e}_y) \underbrace{\begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix}}_{\det P \neq 0},$$
 (3.8)

and

$$(\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} x \\ y \end{pmatrix} = (\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} a \\ b \end{pmatrix} + (\mathbf{e}_x, \mathbf{e}_y) \begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix},$$

i.e.

$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \\ b \end{pmatrix} + \underbrace{\begin{pmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{pmatrix}}_{} \begin{pmatrix} x' \\ y' \end{pmatrix}.$$

non-degenerate matrix

$$\begin{cases}
 x = a + p_{11}x' + p_{12}y' \\
 y = b + p_{21}x' + p_{22}y'
\end{cases}$$
(3.9)

The difference with the case of changing of Cartesian coordinates considered in subsection 3.2.1 (see (3.6)) is the following: in the case of changing of Cartesian coordinates, the matrix P in equation (3.6) is orthogonal matrix, since it is transition matrix between two orthonormal bases, and the matrix P in equation (3.8) is just non-degenerate transition matrix between two bases.

3.2.4 Affine coordinates and area of ellipse

We consider in this subsection one very important example.

Consider an ellipse

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1. (3.10)$$

in Cartesian coordinates. (We will consider this equation of ellipse later in detail. Now we just use very simple properties of this formula to come to the

formula for area of ellipse (more precisely the area of the domain restircted by this ellipse).) Under affine transformations

$$\begin{cases} x = au \\ y = bv \end{cases} \tag{3.11}$$

, (a,b;0) the ellipse becomes the circle $u^2+v^2=1$, and the area of the interior of the unit circle circle is equal to π .

Area of the interior of the ellipse (3.10) =
$$\int_{\frac{x^2}{a^2} + \frac{y^2}{b^2} \le 1} dx dy = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial(u, v)} \right| du dv = \int_{u^2 + v^2 \le 1} \left| \frac{\partial(x, u)}{\partial$$

$$\int_{u^2+v^2 \le 1} \left| \det \begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} \\ \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} \end{pmatrix} \right| du dv = \int_{u^2+v^2 \le 1} \left| \det \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \right| du dv = ab \int_{u^2+v^2 \le 1} du dv = ab \cdot \text{ {Area of the interior of the unit circle}} = \pi ab.$$

Indeed one can see it without using integration formulae: the affine transformation (3.11) changes size in OX direction a times and it changes the size in OY direction in b times, hence area of changing ab times, i.e. area of the ellipse is equal to πab .