

Homework 5. Solutions.

Christoffel symbols and Lagrangians

1 Consider the Lagrangian of "free" particle $L = \frac{1}{2}g_{ik}\dot{x}^i\dot{x}^k$ for Riemannian manifold with a metric $G = g_{ik}dx^i dx^k$.

Write down Euler-Lagrange equations of motion for this Lagrangian and compare them with differential equations for geodesics on this Riemannian manifold.

In fact show that

$$\underbrace{\frac{\partial L}{\partial x^i} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}^i}}_{\text{Euler-Lagrange equations}} \Leftrightarrow \underbrace{\frac{d^2 x^i}{dt^2} = \Gamma_{km}^i \dot{x}^k \dot{x}^m}_{\text{Equations for geodesics}}, \quad (1)$$

where

$$\Gamma_{km}^i = \frac{1}{2}g^{ij} \left(\frac{\partial g_{jk}}{\partial x^m} + \frac{\partial g_{jm}}{\partial x^k} - \frac{\partial g_{km}}{\partial x^j} \right). \quad (2)$$

Solution: see the lecture notes.

2 a) Write down the Lagrangian of "free" particle $L = \frac{1}{2}g_{ik}\dot{x}^i\dot{x}^k$ for Euclidean plane in polar coordinates. Calculate Christoffel symbols for canonical flat connection in polar coordinates using Euler-Lagrange equations for this Lagrangian. Compare with answers which you obtained by the direct use of the formula (2). b) Do the same for cylindrical coordinates in \mathbf{E}^3 .

Solution. Canonical flat connection is Levi-Civita connection of Euclidean metric $G = dx^2 + dy^2$. Hence we can calculate Christoffel symbols using Lagrangian method.

Euclidean metric in polar coordinates is $dr^2 + r^2 d\varphi^2$. Hence the Lagrangian of the free particle is

$$L = \frac{\dot{r}^2 + r^2 \dot{\varphi}^2}{2}$$

Euler-Lagrange equations:

1) for r :

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) = \ddot{r} = \frac{\partial L}{\partial r} = r\dot{\varphi}^2$$

i.e.

$$\ddot{r} - r\dot{\varphi}^2 = 0 \Rightarrow \Gamma_{rr}^r = \Gamma_{\varphi r}^r = \Gamma_{r\varphi}^r = 0, \Gamma_{\varphi\varphi}^r = -r.$$

2) for φ :

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\varphi}} \right) = \frac{d}{dt} (r^2 \dot{\varphi}) = r^2 \ddot{\varphi} + 2r\dot{r}\dot{\varphi} = \frac{\partial L}{\partial \varphi} = 0,$$

i.e.

$$\ddot{\varphi} + \frac{2}{r}\dot{r}\dot{\varphi} = 0 \Rightarrow \Gamma_{rr}^\varphi = \Gamma_{\varphi\varphi}^\varphi = 0, \Gamma_{r\varphi}^\varphi = \Gamma_{\varphi r}^\varphi = \frac{1}{r}.$$

b) cylindrical coordinates in \mathbf{E}^3 . Calculations almost the same as for polar coordinates in \mathbf{E}^2 . $G = dr^2 + r^2 d\varphi^2 + dh^2$,

$$L = \frac{\dot{r}^2 + r^2 \dot{\varphi}^2 + \dot{h}^2}{2}$$

for r : $\ddot{r} - r\dot{\varphi}^2 = 0 \Rightarrow$

$$\Gamma_{rr}^r = \Gamma_{\varphi r}^r = \Gamma_{r\varphi}^r = \Gamma_{rh}^r = \Gamma_{hr}^r = \Gamma_{h\varphi}^r = \Gamma_{\varphi h}^r = \Gamma_{hh}^r = 0, \Gamma_{\varphi\varphi}^r = -r.$$

for φ , $r^2\ddot{\varphi} + 2r\dot{r}\dot{\varphi} = \frac{\partial L}{\partial \varphi} = 0$, i.e. $\ddot{\varphi} + \frac{2}{r}\dot{r}\dot{\varphi} = 0 \Rightarrow$

$$\Gamma_{rr}^\varphi = \Gamma_{rh}^\varphi = \Gamma_{hr}^\varphi = \Gamma_{\varphi\varphi}^\varphi = \Gamma_{\varphi h}^\varphi = \Gamma_{h\varphi}^\varphi = \Gamma_{hh}^\varphi = 0, \Gamma_{r\varphi}^\varphi = \Gamma_{\varphi r}^\varphi = \frac{1}{r}$$

3) for h , $\ddot{h} = 0$,

$$\Gamma_{rr}^h = \Gamma_{r\varphi}^h = \Gamma_{\varphi r}^h = \Gamma_{rh}^h = \Gamma_{hr}^h = \Gamma_{\varphi\varphi}^h = \Gamma_{\varphi h}^h = \Gamma_{h\varphi}^h = \Gamma_{hh}^h = 0,$$

3 Write down the Lagrangian of "free" particle $L = \frac{1}{2}g_{ik}\dot{x}^i\dot{x}^k$ for the sphere of radius R in \mathbf{E}^3 in spherical coordinates. Calculate Christoffel symbols of Levi-Civita connection in spherical coordinates using Euler-Lagrange equations for this Lagrangian. (The induced Riemannian metric on the sphere equals $G = R^2d\theta^2 + R^2\sin^2\theta d\varphi^2$.)

Solution.

Riemannian metric on sphere in spherical coordinates is $R^2d\theta^2 + R^2\sin^2\theta d\varphi^2$. Hence the Lagrangian of the free particle is

$$L = \frac{R^2\dot{\theta}^2 + R^2\sin^2\theta\dot{\varphi}^2}{2}$$

Euler-Lagrange equations:

for θ : $\ddot{\theta} = \sin\theta\cos\theta\dot{\varphi}^2$, hence

$$\ddot{\theta} = \sin\theta\cos\theta\dot{\varphi}^2 \Rightarrow \Gamma_{\theta\theta}^\theta = \Gamma_{\theta\varphi}^\theta = \Gamma_{\varphi\theta}^\theta = 0, \Gamma_{\varphi\varphi}^\theta = -\sin\theta\cos\theta.$$

for φ , $\frac{d}{dt}(R^2\sin^2\theta\dot{\varphi}) = 0$, i.e. $\sin^2\theta\ddot{\varphi} + 2\sin\theta\cos\theta\dot{\theta}\dot{\varphi} = 0$, hence

$$\ddot{\varphi} + 2\cotan\theta\dot{\theta}\dot{\varphi} = 0 \Rightarrow \Gamma_{\theta\theta}^\varphi = \Gamma_{\varphi\varphi}^\varphi = 0, \Gamma_{\varphi\theta}^\varphi = \Gamma_{\theta\varphi}^\varphi = \cotan\theta.$$

4 Calculate Christoffel symbols of Levi-Civita connection for Riemannian metric $G = adu^2 + b dv^2$. Lagrangian

$$L = \frac{a\dot{u}^2 + b\dot{v}^2}{2}$$

Euler-lagrange equations

for u :

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{u}}\right) = \frac{d}{dt}(a\dot{u}) = a_u\dot{u}^2 + a_v\dot{v}\dot{u} + a\ddot{u} = \frac{\partial L}{\partial u} = \frac{a_u\dot{u}^2 + b_u\dot{v}^2}{2}$$

hence

$$\ddot{u} + \frac{1}{2}\frac{a_u}{a}\dot{u}^2 + \frac{a_v}{a}\dot{v}\dot{u} - \frac{1}{2}\frac{b_u}{a}\dot{v}^2 \Rightarrow \Gamma_{uu}^u = \frac{1}{2}\frac{a_u}{a}, \Gamma_{uv}^u = \Gamma_{vu}^u = \frac{1}{2}\frac{a_v}{a}, \Gamma_{vv}^u = -\frac{1}{2}\frac{b_u}{a},$$

for v :

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{v}}\right) = \frac{d}{dt}(b\dot{v}) = b_v\dot{v}^2 + b_u\dot{u}\dot{v} + b\ddot{v} = \frac{\partial L}{\partial v} = \frac{a_v\dot{u}^2 + b_v\dot{v}^2}{2}$$

hence

$$\ddot{v} + \frac{1}{2}\frac{b_v}{b}\dot{v}^2 + \frac{b_u}{b}\dot{u}\dot{v} - \frac{1}{2}\frac{a_v}{b}\dot{u}^2 \Rightarrow \Gamma_{vv}^v = \frac{1}{2}\frac{b_v}{b}, \Gamma_{vu}^v = \Gamma_{uv}^v = \frac{1}{2}\frac{b_u}{b}, \Gamma_{uu}^v = -\frac{1}{2}\frac{a_v}{b}.$$