

Recall that generating function $S(x, q) = S(x^i, q_a)$ defines a map

$$\Phi_S^*: C(N) \rightarrow C(M)$$

such that

$$\forall g(y) \in CM(N)C(M) \ni f(x) = \Phi_S^*(g) = g(y) + S(x, q) - y^a q_a ,$$

where functions y^a, q_a are such that

$$y^a = \frac{\partial S(x, q)}{\partial q_a} \Big|_{q_a = \frac{\partial g}{\partial y^a}} . \quad \Rightarrow \quad \frac{\partial f}{\partial x^i} = p_i = \frac{\partial S(x, q)}{\partial x^i}$$

Now consider the case when $M = N$.

Let $H = H(x, p)$ be Hamiltonian on T^*M .

s Hamiltonian generates infinitesimal thick morphism

$$\Phi_{\varepsilon S}^*: g(x) \mapsto g(x) + \varepsilon H(x, p) \Big|_{p = \frac{\partial g}{\partial y}} ,$$

This infinitesimal map is nothing but the Hamilton-Jacobi vector field on $C(M)$

$$\mathbf{X}_S = \int dx H \left(x, \frac{\partial g}{\partial x} \right) \frac{\delta}{\delta x} .$$

This we see that for finite interval of the time

$$\Phi_S^{(t)*}(g) = A(g|x, t): \quad \begin{cases} \frac{\partial A}{\partial t} = H \left(x, \frac{\partial A}{\partial x} \right) \\ A|_{t=0} = g \end{cases}$$

This is up to a sign Hamilton Jacobi equation.

Thus we see that the action of inverse image of thick diffeomorphism generated by the Hamiltonian H on the arbitrary function $g(y)$ is the action $A(g|x, t)$, which is equal to $g(x)$ at the moment $t = 0$.

It is difficult to avoid a temptation to claim that the quantum version of this equation has to be Shrodinger equation.