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Here I will try to calculate the continual integral for free particle using some linear algebra stuff

To calculate the continual integral for free function we deal with exponent of the polynomial

$$F = (x_0 - x_1)^2 + (x_1 - x_2)^2 + (x_2 - x_3)^2 + (x_3 - x_2)^2 + \dots + (x_{N-1} - x_N)^2,$$

where initial and final points are fixed

$$x_0 = a, x_N = b$$

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Instead doing standard calculations (see the book of Feynman) we will present here two calculations based on the linear algebra.

§1 **Straightforward calculations** The function  $F$  is the polynomial of order 2 on the  $N - 1$  variables  $x_1, x_2, \dots, x_N$  it is equal to

$$F = M_{ik}x^i x^k + L_i x^i + N, (i, k = 1, \dots, N - 1)$$

where the  $(N - 1) \times (N - 1)$  matrix  $M_{ik}$  is the symmetrical matrix which is equal to

$$M = \begin{pmatrix} 2 & -1 & 0 & 0 & 0 \dots & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \dots & 0 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \dots & 0 & 0 & 0 \\ 0 & 0 & -1 & 2 & -1 \dots & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 2 \dots & 0 & 0 & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & 0 & 1 \dots & 2 & -1 & 0 \\ 0 & 0 & 0 & 0 & 1 \dots & -1 & 2 & -1 \\ 0 & 0 & 0 & 0 & 1 \dots & 1 & -1 & 2 \end{pmatrix},$$

the vector  $L$  is equal to

$$L = (-2a, 0, 0, 0, 0, \dots, 0, 0, -2b)$$

and the scalar  $N$  is equal to

$$N = a^2 - b^2$$

The orthogonal transformations make the matrix  $M$  the diagonal matrix,

We see that

$$\int e^{-cF} dx^1 dx^2 \dots dx^{N-1} = \int e^{-c(M_{ik}x^i x^k + L_i x^i + N)} dx^1 dx^2 \dots dx^{N-1}.$$

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\* One have to calculate the gaussian integral  $e^{iCF}$

Now choose the vector  $\mathbf{a}$  such that  $2M\mathbf{a} = -\mathbf{L}$ , i.e.

$$\mathbf{a} = -\frac{1}{2}M^{-1}\mathbf{L}.$$

and perform the translation  $x^i = a^i + y^i$ : we come to

$$\begin{aligned} \int e^{-cF} dx^1 dx^2 \dots dx^{N-1} &= \int e^{-c(M_{ik}x^i x^k + L_i x^i + N)} dx^1 dx^2 \dots dx^{N-1} = \\ &= \int e^{-c(M_{ik}(y^i + a^i)(y^k + a^k) + L_i(y^i + a^i) + N)} dx^1 dx^2 \dots dx^{N-1} = \\ &= \int e^{-c(M_{ik}y^i y^k + M_{ik}a^i a^k + L_i a^i + N)} dx^1 dx^2 \dots dx^{N-1} = \\ &= \left(\frac{\pi}{c}\right)^{N-1} \frac{1}{\sqrt{\det M}} e^{-c(M_{ik}a^i a^k + L_i a^i + N)} \quad \text{where } \mathbf{a} = -\frac{1}{2}M^{-1}\mathbf{L} \text{ and } N = b^2 - a^2. \end{aligned}$$

## §2 More sophisticated linear algebra calculations

Consider linear transformation

$$\begin{cases} \xi_1 = x_1 - x_0 = x_1 - a \\ \xi_2 = x_2 - x_1 \\ \xi_3 = x_3 - x_2 \\ \dots \\ \xi_{N-1} = x_{N-1} - x_N \end{cases} \Leftrightarrow \begin{cases} x_1 = \xi_1 + a \\ x_2 = \xi_2 + \xi_1 + a \\ x_3 = \xi_3 + \xi_2 + \xi_1 + a \\ \dots \\ x_{N-1} = \xi_{N-1} + \dots + \xi_1 + a \end{cases}$$

This is not linear transformation, and its “linear” part is not orthogonal transformation, but it is unimodular transformation. In the new coordinates

$$F = \sum_{i,k=1}^{N-1} M_{ik} x^i x^k + \sum_{i=1}^{N-1} L_i x^i + N = \xi_1^2 + \xi_2^2 + \xi_3^2 + \dots + \xi_{N-1}^2 + (b - a - \xi_1 - \xi_2 - \dots - \xi_{N-1})^2.$$

Find translation which kills linear terms:  $\xi_i = \eta_i + z_i$  such that only constants and quadratic terms will survive. One can see that this is

$$\xi_i = \eta_i + \frac{b-a}{N}.$$

Check it:

$$\begin{aligned} F &= \sum_{i,k=1}^{N-1} M_{ik} x^i x^k + \sum_{i=1}^{N-1} L_i x^i + N = \xi_1^2 + \xi_2^2 + \xi_3^2 + \dots + \xi_{N-1}^2 + (a - b - \xi_1 - \xi_2 - \dots - \xi_{N-1})^2 = \\ &= (\eta_1^2 + \eta_2^2 + \dots + \eta_{N-1}^2) + 2(\eta_1 + \dots + \eta_{N-1}) \frac{b-a}{N} + \frac{(N-1)(b-a)^2}{N^2} + \end{aligned}$$

$$\begin{aligned}
& + \left( \frac{a-b}{N} - \eta_1 - \eta_2 - \dots - \eta_{N-1} \right)^2 = \\
& = 2(\eta_1^2 + \eta_2^2 + \dots + \eta_{N-1}^2) + 2 \sum_{i < j} \eta_i \eta_j + \frac{(b-a)^2}{N}.
\end{aligned}$$

The linear terms are cancelled, and we have that

$$F = \tilde{M}_{ik} u^i u^k + \frac{(b-a)^2}{N},$$

where

$$\tilde{M}_{ik} = \delta_{ik} + t_i t_k, \quad (t_i = (1, \dots, 1)).$$

Now we can calculate the integral:

$$\begin{aligned}
I &= \int e^{-cF} dx^1 dx^2 \dots dx^{N-1} = \int e^{-c(M_{ik} x^i x^k + L_i x^i + N)} dx^1 dx^2 \dots dx^{N-1} = \\
&= \int e^{-c\left(\sum_{k=1}^{N-1} \xi_k^2 + (b-a - \sum_{m=1}^{N-1} \xi_m^2)^2\right)} \underbrace{\left(\frac{\partial(x_1, \dots, x_{N-1})}{\partial(\xi_1, \dots, \xi_{N-1})}\right)}_{\text{equals to 1}} d\xi_1 d\xi_2 \dots d\xi_{N-1} = \\
&= \int e^{-c\left(2\sum_{k=1}^{N-1} \xi_k^2 + (b-a - \sum_{m=1}^{N-1} \xi_m^2)^2\right)} d\xi_1 d\xi_2 \dots d\xi_{N-1} = \\
&= \int e^{-c\left(2\sum_{i,k=1}^{N-1} \eta_i \eta_k + \frac{(b-a)^2}{N}\right)} \underbrace{\left(\frac{\partial(\xi_1, \dots, \xi_{N-1})}{\partial(\eta_1, \dots, \eta_{N-1})}\right)}_{\text{equals to 1}} d\eta_1 d\eta_2 \dots d\eta_{N-1} = \\
&= \int e^{-c\left(N_{ik} \eta^i \eta^k + \frac{(b-a)^2}{N}\right)} d\eta_1 d\eta_2 \dots d\eta_{N-1} = \\
&= \left(\frac{\pi}{c}\right)^{\frac{N-1}{2}} \sqrt{\frac{1}{\det(\tilde{M})}} e^{-c\left(\frac{(b-a)^2}{N}\right)}.
\end{aligned}$$

Notice that the matrix  $\tilde{M}$  has eigenvector  $\mathbf{t}$  with eigenvalue  $N$ , and all other  $N-2$  eigenvectors which are orthogonal to the vector  $\mathbf{t}$  with the eigenvalue 1. Hence

$$\det \tilde{M} = N,$$

and we have for integral:

$$I = \int e^{-c\left(N_{ik} \eta^i \eta^k + \frac{(b-a)^2}{N}\right)} d\eta_1 d\eta_2 \dots d\eta_{N-1} =$$

$$\sqrt{\frac{\pi}{\det(c\tilde{M})}} e^{-c\left(\frac{(b-a)^2}{N}\right)} = \left(\frac{\pi}{c}\right)^{\frac{N-1}{2}} \frac{1}{\sqrt{N}} e^{-c\left(\frac{(b-a)^2}{N}\right)}.$$

### §3 Returning to Quantum mechanics

Now return to Quantum Mechanics:

We calculate now the wave function of free particle using continual integral and the classial action.

Let particle starts at the point  $x_0, t_0$  and ends at the point  $x_1, t_1$  Divide  $[t_0, t_1]$  on

$$[t_0, t_1, t_2, \dots, t_{N-1}, t_N], \quad (t_N \mapsto t_1)$$

and consider the integral

$$\begin{aligned} & \int \int \int \dots \int e^{\frac{i}{\hbar} S_{\text{free}}(x_0, t_0; x_1, t_1)} e^{\frac{i}{\hbar} S_{\text{free}}(x_1, t_1; x_2, t_2)} e^{\frac{i}{\hbar} S_{\text{free}}(x_2, t_2; x_3, t_3)} \dots \\ & e^{\frac{i}{\hbar} S_{\text{free}}(x_{N-2}, t_{N-2}; x_{N-1}, t_{N-1})} e^{\frac{i}{\hbar} S_{\text{free}}(x_{N-1}, t_{N-1}; x_N, t_N)} dx^1 dx^2 dx^3 \dots dx^{N-2} dx^{N-1} = \\ & \int \prod_{i=1}^N \exp \left[ \frac{i}{\hbar} S_{\text{free}}(x_{i-1}, t_{i-1}; x_i, t_i) \right] \prod_{j=1}^{N-1} dx^j, \end{aligned}$$

where

$$S_{\text{free}}(x_0, t_0; x_1, t_1) = \frac{m(x_1 - x_0)^2}{2t}$$

is the classical action.

This we have for the integral

$$I = \int \prod_{i=1}^N \exp \left[ \frac{i}{\hbar} S_{\text{free}}(x_{i-1}, t_{i-1}; x_i, t_i) \right] \prod_{j=1}^{N-1} dx^j = \int \prod_{i=1}^N \exp \left[ \frac{im}{2\varepsilon\hbar} (x_i - x_{i-1})^2 \right] \prod_{j=1}^{N-1} dx^j,$$

where

$$\varepsilon = \frac{t_N - t_0}{N + 1},$$

We have that according the results of calculations in the previous paragraph

$$\begin{aligned} G(x_s, t_s, x_f, t_f) &= \int \exp \left[ \frac{im}{2\varepsilon\hbar} \sum (x_i - x_{i-1})^2 \right] dx^1 dx^2 \dots dx^n = \\ & \left( \frac{\pi}{2c} \right)^{\frac{N}{2}} e^{-c\left(\frac{(b-a)^2}{N}\right)}, \text{ with } c = \frac{m}{2\varepsilon\hbar}, \end{aligned}$$

i.e.

$$G(x_s, t_s, x_f, t_f) = \left( \frac{\pi}{2c} \right)^{\frac{N}{2}} e^{-c\left(\frac{(b-a)^2}{N}\right)} =$$

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#### §4 One linear problem question

Making this calculation we in fact calculated the matrix

$$M = \begin{pmatrix} 2 & -1 & 0 & 0 & 0 \dots & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \dots & 0 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \dots & 0 & 0 & 0 \\ 0 & 0 & -1 & 2 & -1 \dots & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 2 \dots & 0 & 0 & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & 0 & 1 \dots & 2 & -1 & 0 \\ 0 & 0 & 0 & 0 & 1 \dots & -1 & 2 & -1 \\ 0 & 0 & 0 & 0 & 1 \dots & 1 & -1 & 2 \end{pmatrix},$$

Indeed the quadratic form of this matrix is

$$Q_M = \sum_{i,k=1}^{N-1} M_k x^i x^k = (x_1 - a)^2 + \sum_{i=2}^{N-1} (x_k - x_{k-1})^2 + (b - x_{N-1})^2 + 2ax_1 - a^2 + 2bx_{N-1} - b^2.$$

In the new coordinates

$$\begin{cases} \xi_1 = x_1 - x_0 = x_1 \\ \xi_2 = x_2 - x_1 \\ \xi_3 = x_3 - x_2 \\ \dots \\ \xi_{N-1} = x_{N-1} - x_N \end{cases} \Leftrightarrow \begin{cases} x_1 = \xi_1 \\ x_2 = \xi_2 + \xi_1 \\ x_3 = \xi_3 + \xi_2 + \xi_1 \\ \dots \\ x_{N-1} = \xi_{N-1} + \dots + \xi_1 \end{cases}$$

(this is linear unimodular transformation, and it is different from the trasforation considered in §2.) we have

$$Q_M = \sum_{i,k=1}^{N-1} M_k x^i x^k = (x_1 - a)^2 + \sum_{i=2}^{N-1} (x_k - x_{k-1})^2 + (b - x_{N-1})^2 + 2ax_1 - a^2 + 2bx_{N-1} - b^2 =$$

$$(\xi_1 - a)^2 + \sum_{i=2}^{N-1} \xi_k^2 + (b - \xi_1 - \dots - \xi_{N-1})^2 + 2a\xi_1 - a^2 + 2b(\xi_1 + \dots + \xi_{N-1}) - b^2 =$$

$$\sum_{i=1}^{N-1} \xi_i^2 + \sum_{i,k=1}^{N-1} \xi_i \xi_k = \xi^i (\delta_{ik} + t_i t_k) \xi^k,$$

with

$$\mathbf{t} = (1, 1, \dots, 1)$$

We have

$$M = P^* M_\xi P$$

In particular since the transformation is unimodular we can calculate the the characteristic polynomial of  $M$ . It is equal to

$$P_M(\lambda) = \det(M - \lambda) = \det((1 - \lambda) \delta_{ik} + t_i t_k) .$$

To calculate the determinant of this matrix notice the vector  $\mathbf{t}$  is the eigenvector with the eigenvalue  $(1 - \lambda + (N - 1)) = \lambda + N$ , and all the vectors which are orthogonal to the vector  $\mathbf{t}$  has eigenvalue  $1 - \lambda$ . Hence

$$P_M(\lambda) = \det(M - \lambda) = \det((1 - \lambda) \delta_{ik} + t_i t_k) = (\lambda + N)(1 - \lambda)^{N-2}$$