

Very important example

In thick morphisms very important role play the case when thick morphism is just usual morphism. In this case

$$S(x, q) = \varphi^i(x) q_i$$

and the Hamiltonian is linear function.

In some sence it is degenerate case.

There is also the following (may be not less important?) example: **free particle**.

Consider the free particle. Its action is

$$S_t(x, y) = \frac{m(y - x)^2}{2t}$$

and

$$S_t(x, q) = x^i q_i + \frac{q^2 t}{2m}$$

The LHS is the action, but it is in (x, q) -representation, it is the Legendre transform of the classical action (see the texts for July):

$$\mathcal{S}_t(x, q) = y^i q_i - S_t(x, y) \text{ with } q_i = \frac{\partial S_t(x, y)}{\partial y^i}.$$

It is easy to see that in this case one can easy to calculate the quantum case properly (ot just in stationary limit!)

The classical case is

$$f(x) = (\Phi^* g)(x) = g(x) + x^i q_i + \frac{q^2 t}{2m} - y^i q_i \text{ with } y^i = \frac{\partial S_t(x, q)}{\partial q_i} = x^i + \frac{q^i t}{m}, q_i = \frac{\partial g(y)}{\partial y^i},$$

i.e.

$$\begin{aligned} f(x) &= g\left(x^i + \frac{t}{m} q^i\right) = g\left(x^i + \frac{t}{m} \left(x^i + \frac{t}{m} q^i\right)\right) = \\ &= g\left(x^i + \frac{t}{m} \left(x^i + \frac{t}{m} \left(x^i + \frac{t}{m} q^i\right)\right)\right) = \dots = \\ &= \lim_{N \rightarrow \infty} g\left(x^i + \frac{t}{m} \left(x^i + \frac{t}{m} \left(x^i + \dots\right) \dots\right)\right), \text{ where } N \text{ is 'number' of brackets.} \end{aligned}$$

This iteration procedure has simple integral representation since it can be described properly using quantum case:

$$\begin{aligned} \exp\left(\frac{i}{\hbar} f(x)\right) &= \int \exp\left(\frac{i}{\hbar} \left(g(y) + S_t(x, q) - y^i q_i\right)\right) \mathcal{D}q \mathcal{D}y = \\ &= \int \exp\left(\frac{i}{\hbar} \left(g(y) + x^i q_i + \frac{q^2 t}{2m} - y^i q_i\right)\right) \mathcal{D}q \mathcal{D}y \end{aligned}$$

For free particle one can calculate this itnegral for an arbitrary \hbar (If $\hbar \rightarrow 0$ then this integral is á la Legendre transform, andn this case it describes the pull back in classical case).

We have

$$\exp\left(\frac{i}{\hbar}f(x)\right) = \int \exp\left(\frac{i}{\hbar}\left(g(y) + x^i q_i + \frac{q^2 t}{2m} - y^i q_i\right)\right) \mathcal{D}q \mathcal{D}y =$$

$$\int \exp\left(\frac{i}{\hbar}(g(y))\right) K_t(x, y) \mathcal{D}y,$$

whew $K_t(x, y)$ is Green function of heat equation:

$$K_t(x, y) = \int \exp\left(\frac{i}{\hbar}\left(x^i q_i + \frac{q^2 t}{2m} - y^i q_i\right)\right) \mathcal{D}q =$$

$$\int \exp\left(\frac{i}{\hbar}\left(x^i q_i + \frac{t}{2m}\left(q^i - \frac{m}{t}(y^i - x^i)\right)^2 - \frac{m(y-x)^2}{2t}\right)\right) \mathcal{D}q =$$

$$\frac{C}{\sqrt{t}} \exp\left(\frac{i}{\hbar}\left(-\frac{m(y-x)^2}{2t}\right)\right)$$