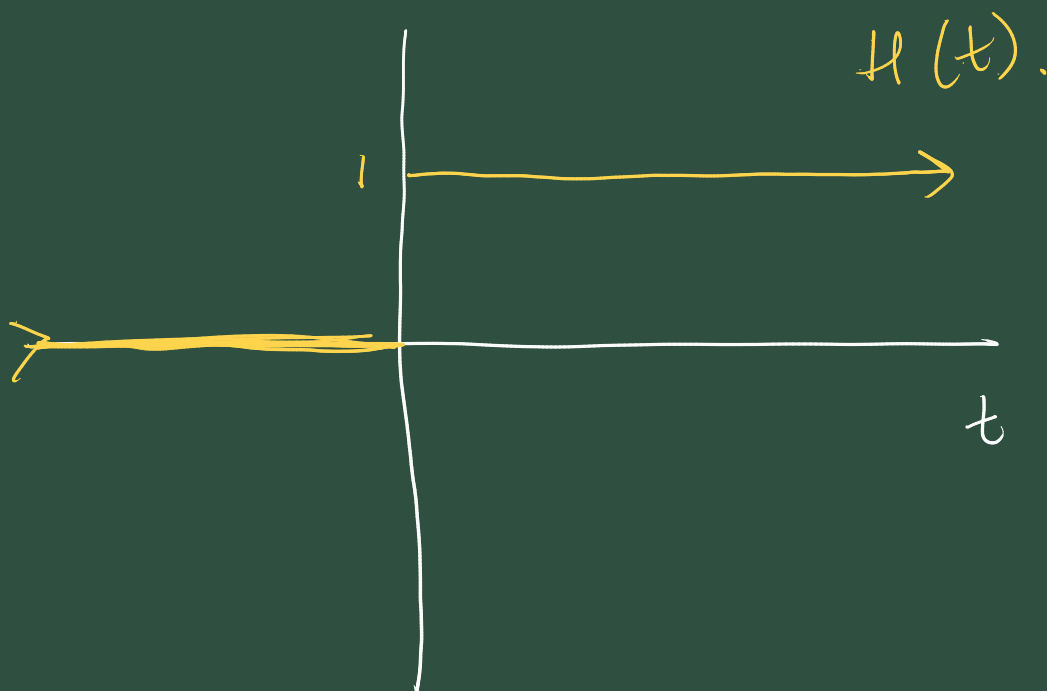


Dirac delta function δ
Heaviside step function H .



$$H(t) = \begin{cases} 0, & t < 0 \\ 1, & t \geq 0 \end{cases}$$

$$H(t-a) = \begin{cases} 0, & t-a < 0 \Leftrightarrow t < a \\ 1, & t-a \geq 0 \Leftrightarrow t \geq a \end{cases}$$

Examples exploiting H

1. $g(t) = H(t-a) - H(t-b)$

where $a < b$.

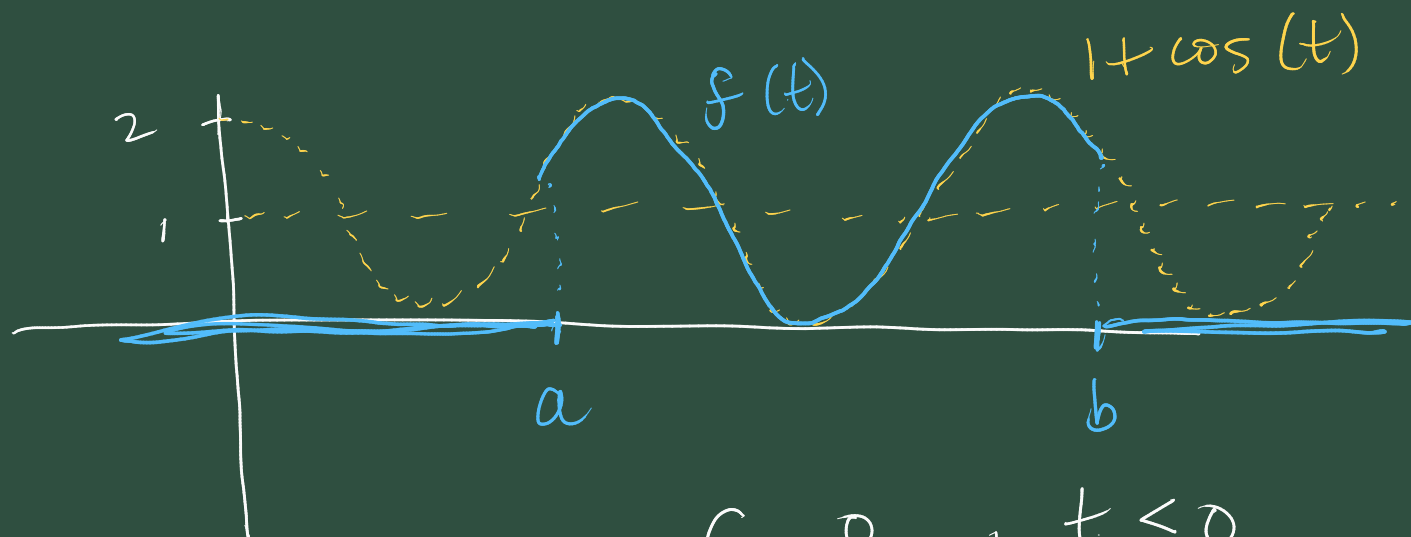


$$g(t) = \begin{cases} 0, & t < a \\ 1, & a \leq t < b \\ 0, & t \geq b \end{cases}$$

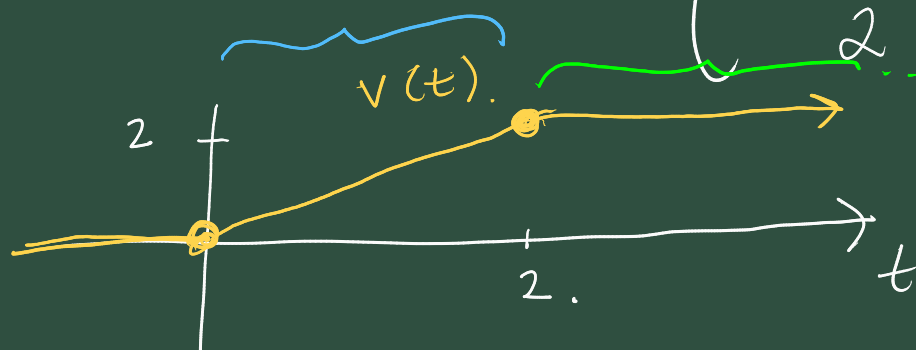
2. let $a < b$

$$f(t) = \underbrace{(1 + \cos(t))}_{\text{blue}} \underbrace{(H(t-a) - H(t-b))}_{\text{red}}$$

$$f(t) = \begin{cases} 0, & t < a \\ 1 + \cos(t), & a \leq t < b \\ 0, & t \geq b \end{cases}$$



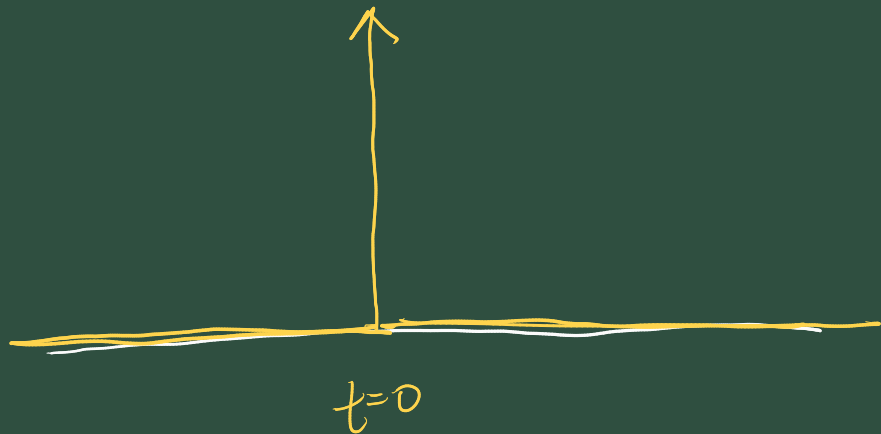
$$3. \quad v(t) = \begin{cases} 0, & t < 0 \\ t, & 0 \leq t < 2 \\ 2, & t \geq 2 \end{cases}$$



$$v(t) = t \left(H(t) - H(t-2) \right) + \underline{2 H(t-2)}$$

$$= \underbrace{t H(t)}_{\text{same shifted}} - \underbrace{(t-2) H(t-2)}_{\text{same shifted argument}}$$

$$\delta(t) = \begin{cases} \infty, & t=0 \\ 0, & t \neq 0 \end{cases}$$



To solve ODEs featuring H , δ we need to understand their transforms.

Theorem 4.7.9

$$\mathcal{L} \{ H(t-a) \} = \frac{e^{-as}}{s}$$

Pf: Using the definition

$$\mathcal{L}\{H(t-a)\} = \int_0^{\infty} e^{-st} H(t-a) dt.$$

Assuming $a > 0$

$$= \int_0^a e^{-st} \cdot 0 dt + \int_a^{\infty} e^{-st} \cdot 1 dt$$

$$= \int_a^{\infty} e^{-st} dt.$$

$$= \left[\frac{e^{-st}}{-s} \right]_{t=a}^{t \rightarrow \infty}$$

$$= 0 - \frac{e^{-sa}}{-s}$$

, for $s > 0$

$$= \frac{e^{-as}}{s}, \text{ as claimed.}$$

□ $\left\{ \int_0^{\infty} e^{-st} f(t) dt \right\}$
↓
~
~

Theorem 4.7.10

If $\mathcal{L}\{f(t)\} = F(s)$ then.

$$\mathcal{L}\{f(t-a)H(t-a)\} = e^{-as} \underbrace{F(s)}_{\sim}$$

Proof: Again, using def of \mathcal{L} .

$$\mathcal{L} \{ f(t-a) \mathcal{H}(t-a) \} \quad , \text{ assuming } a > 0.$$

$$= \int_0^{\infty} e^{-st} f(t-a) \underbrace{\mathcal{H}(t-a)}_{0 \text{ or } 1} dt.$$

$$= \underbrace{\int_0^a e^{-st} f(t-a) \cdot 0 dt}_{=0} + \int_a^{\infty} e^{-st} f(\underline{t-a}) \underline{dt}$$

$$= \int_a^{\infty} e^{-st} f(t-a) dt.$$

Perform the substitution
 $u = t - a.$
as $t \rightarrow \infty$, $u \rightarrow \infty$
when $t = a$, $u = 0$

$$= \int_0^{\infty} e^{-s(u+a)} f(\underline{u}) \underline{du}$$

$$= \int_0^{\infty} e^{-su} \underbrace{e^{-sa}}_{\substack{\text{indep.} \\ \text{of } u}} f(u) \underline{du}$$

$$= e^{-sa} \int_0^{\infty} e^{-su} f(u) du$$

$$= e^{-sa} F(s).$$

$$F(s) = \int_0^{\infty} e^{-st} f(t) dt$$

Ex 4.7.12

rl. Sometimes some prep. is needed before using the theorem.

$$R(t) = t^2 H(t-2)$$

$$\mathcal{L}\{R(t)\} = ?$$

We know

$$\mathcal{L}\{t^2\} = \frac{2}{s^3}$$

Can't immediately apply the previous theorem. First we must express

t^2 in terms of the argument $t-2$.

$$t^2 = (t-2)^2 + 4t - 4.$$

$$= (t-2)^2 + 4(t-2) + 4$$

$$= g(t-2)$$

where g is the polynomial
defined by $g(t) = t^2 + 4t + 4$

This allows to write → ready for Th 4.7.10

$$R(t) = \boxed{g(t-2) H(t-2)}$$

$$\text{So } \mathcal{L}\{R(t)\} = \mathcal{L}\{g(t-2) H(t-2)\}$$

$$= e^{-2s} G(s)$$

$$\text{where } G(s) = \mathcal{L}\{g(t)\}$$

$$= \mathcal{L}\{t^2 + 4t + 4\}$$

$$= \mathcal{L}\{t^2\} + 4\mathcal{L}\{t\} + 4\mathcal{L}\{1\}$$

$$= \frac{2}{s^3} + \frac{4}{s^2} + \frac{4}{s}$$

$$\text{So } \mathcal{L}\{R(t)\}$$

$$= e^{-2s} \left(\frac{2}{s^3} + \frac{4}{s^2} + \frac{4}{s} \right)$$

Ex 4.7.13 $\dot{y} = \frac{dy}{dt}$, $\ddot{y} = \frac{d^2y}{dt^2}$

Solve $\ddot{y} - 4\dot{y} + 3y = \underline{6H(t-5)}$

subject to initial conditions $y(0) = \dot{y}(0) = 0$

Follow the 4-steps.

$$\mathcal{L}\{\ddot{y} - 4\dot{y} + 3y\} = \mathcal{L}\{6H(t-5)\}$$

By linearity of \mathcal{L}

$$\mathcal{L}\{\ddot{y}\} - 4\mathcal{L}\{\dot{y}\} + 3\mathcal{L}\{y\} = 6\mathcal{L}\{H(t-5)\}$$

We'll write $\bar{y} = \mathcal{L}\{y\}$ and use transform of derivatives properties.

$$s^2 \bar{y} - s y(0) - \dot{y}(0) - 4(s \bar{y} - y(0))$$

$$+ 3\bar{y} = \underline{\underline{6 \frac{e^{-5s}}{s}}}$$

Inserting the initial conditions on y, \dot{y} gives.

$$s^2 \bar{y} - 4s \bar{y} + 3\bar{y} = \frac{6e^{-5s}}{s}$$

$$\Rightarrow (s^2 - 4s + 3) \bar{y} = \underline{\underline{6e^{-5s}}}$$

$$\Rightarrow \bar{y} = \frac{6e^{-5s}}{s(s^2 - 4s + 3)}$$

The solution to the ODE will be

$$y(t) = \mathcal{L}^{-1} \{ \bar{y} \}$$

$$= \mathcal{L}^{-1} \left\{ e^{-5s} \right.$$

$$\left. \frac{6}{s(s-3)(s-1)} \right\}$$

$\bar{f}(s)$

So by Theorem 4.7.10

$$y(t) = \underbrace{f(t-5)}_? H(t-5)$$

So we'll find f by finding

$$f(t) = \mathcal{L}^{-1} (\bar{f}(s))$$

$$= \mathcal{L}^{-1} \left\{ \frac{6}{s(s-3)(s-1)} \right\}$$

First we need the partial fraction expansion of $\bar{f}(s)$

$$\bar{f}(s) = 6 \frac{1}{s(s-3)(s-1)}$$

(*)

$$= 6 \left(\frac{\alpha}{s} + \frac{\beta}{s-3} + \frac{\gamma}{s-1} \right)$$

$$= 6 \left(\frac{\alpha(s^2 - 4s + 3) + \beta(s^2 - s) + \gamma(s^2 - 3s)}{s(s-3)(s-1)} \right)$$

$$= 6 \left(\frac{(\alpha + \beta + \gamma)s^2 + (-4\alpha - \beta - 3\gamma)s + 3\alpha}{s(s-3)(s-1)} \right)$$

**) (circled)

So numerator poly of (**) = 1.

$$\Rightarrow \begin{cases} \alpha + \beta + \gamma = 0 \\ -4\alpha - \beta - 3\gamma = 0 \\ 3\alpha = 1 \end{cases}$$

$$\Rightarrow \begin{cases} \alpha = 1/3 \\ \beta + \gamma = -1/3 \\ -\beta - 3\gamma = 4/3 \end{cases}$$

$$\Rightarrow -2\gamma = 1$$

$$\Rightarrow \gamma = -1/2$$

$$\Rightarrow \beta - 1/2 = -1/3$$

$$\Rightarrow \beta = -\frac{1}{3} + \frac{1}{2} = \frac{-2+3}{6} = \frac{1}{6}$$

So

$$\bar{f}(s) = 6 \left(\frac{1/3}{s} + \frac{1/6}{s-3} - \frac{1/2}{s-1} \right)$$

$$\text{So } \boxed{f(t) = \mathcal{L}^{-1}\{\bar{f}(s)\}}$$

$$= 2 \mathcal{L}^{-1}\left\{\frac{1}{s}\right\} + \mathcal{L}^{-1}\left\{\frac{1}{s-3}\right\}$$

$$- 3 \mathcal{L}^{-1}\left\{\frac{1}{s-1}\right\}$$

$$= \boxed{2 + e^{3t} - 3e^t}$$

by linearity
reading transforms
from table.

So finally we can express
our ODE solution as

$$\begin{aligned} y(t) &= f(t-5) u(t-5) \\ &= \left(2 + e^{3t-15} - 3e^{t-5} \right) u(t-5) \end{aligned}$$