

# **MONTE CARLO RENORMALIZATION GROUP**

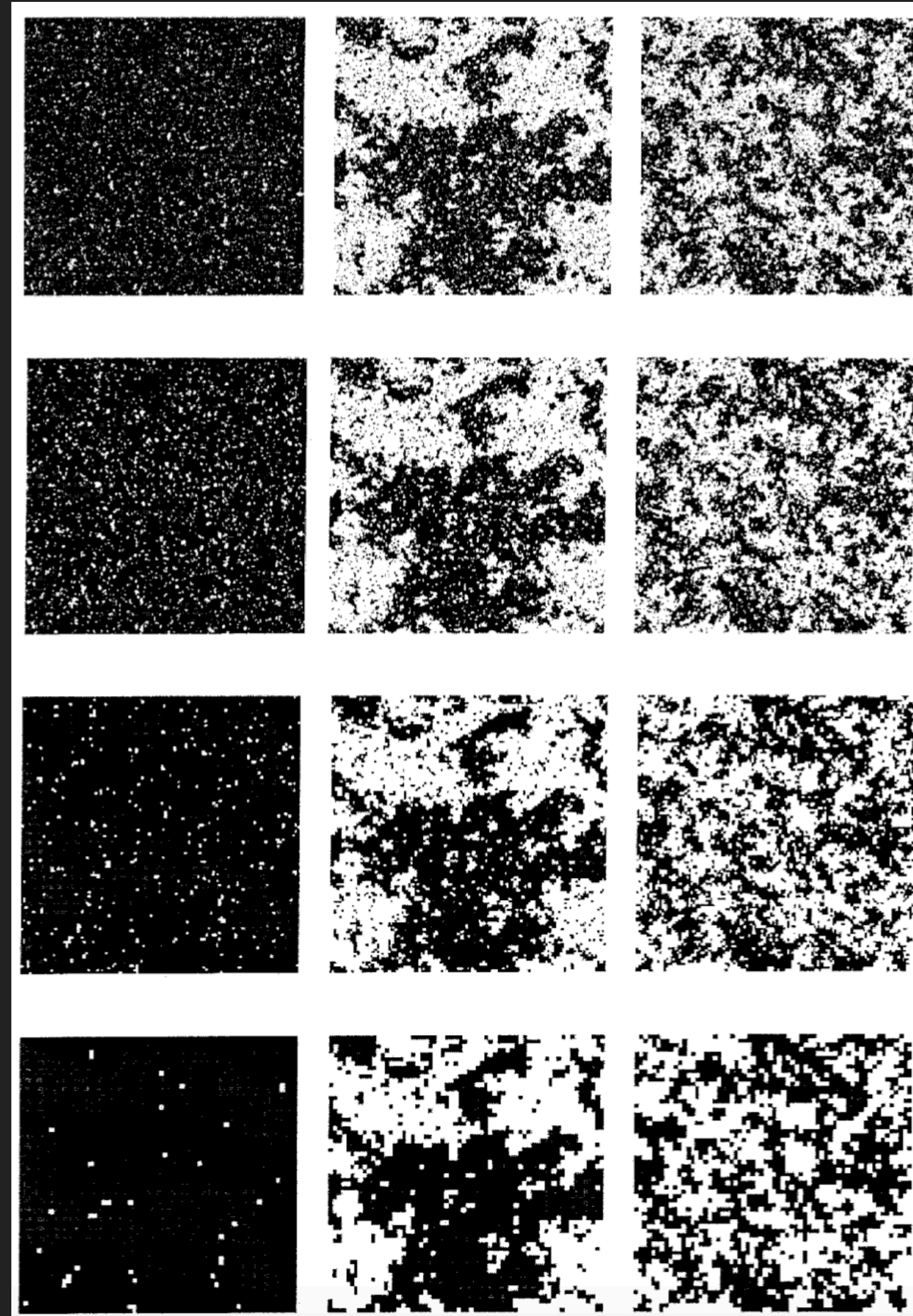
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# OUTLINE

- ▶ Review of key concepts
- ▶ Swendsen's MCRG method
  - ▶ Calculate critical coupling
  - ▶ Extract critical exponents
  - ▶ Monte Carlo simulation
- ▶ Example: 2D Ising model



Landau and Binder, A Guide to Monte Carlo Simulations in Statistical Physics (2005).

## REVIEW OF KEY CONCEPTS

- ▶ Ising-type Hamiltonian:

$$\mathcal{H} = \beta H = \sum_{\alpha} K_{\alpha} S_{\alpha}$$

- ▶ Recall:

$$H = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j, \quad J > 0, \quad \implies \quad K_1 = -\beta J, \quad S_1 = \sum_{\langle i,j \rangle} \sigma_i \sigma_j$$

- ▶ Starting Hamiltonian:

$$\mathcal{H}^{(0)} = K_1^{(0)} S_1^{(0)}, \quad K_{\alpha} = 0 \quad \forall \quad \alpha > 1$$



# REVIEW OF KEY CONCEPTS

- ▶ Starting Hamiltonian:

$$\mathcal{H}^{(0)} = K_1^{(0)} S_1^{(0)}, \quad K_\alpha = 0 \quad \forall \quad \alpha > 1$$

- ▶ After  $n$  RG transformations:

$$\mathcal{H}^{(n)} = \sum_{\alpha} K_{\alpha}^{(n)} S_{\alpha}^{(n)}$$

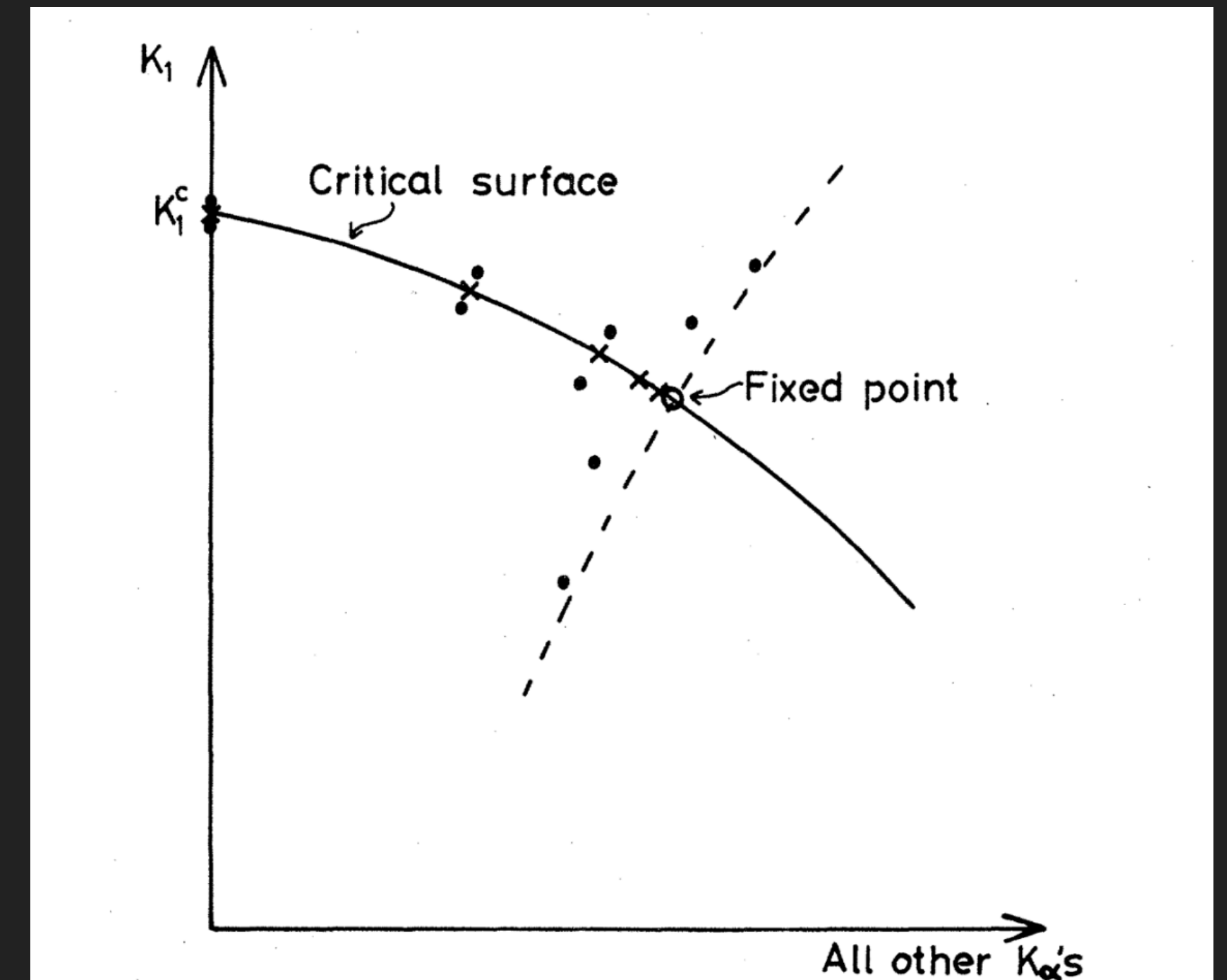


FIG. 1. Schematic diagram for RG flows in the many-dimensional space of coupling constants, indicating the transient into the fixed point starting from the critical value  $K_1^c$ , and the instability of the fixed point to perturbations out of the critical surface.

## REVIEW OF KEY CONCEPTS

- ▶ Choose a renormalization transformation (with scale factor  $b$ ):

$$\mathcal{H}^{(n+1)} = R_b \mathcal{H}^{(n)}$$

- ▶ Linearize transformation about fixed-point Hamiltonian  $\mathcal{H}^*$ :

$$\delta K_{\alpha}^{(n+1)} = T_{\alpha\beta}^* \delta K_{\beta}^{(n)}, \text{ where } T_{\alpha\beta}^* = \left. \frac{\partial K_{\alpha}^{(n+1)}}{\partial K_{\beta}^{(n)}} \right|_{\mathcal{H}=\mathcal{H}^*}$$

- ▶ Goal: Calculate the eigenvalues of  $T_{\alpha\beta}^*$ .

## REVIEW OF KEY CONCEPTS

- ▶ External magnetic field is zero at the critical point.
- ▶  $T_{\alpha\beta}^*$  decouples into a block diagonal form.

$$\begin{aligned}\lambda_1^{(e)} &= \text{largest eigenvalue of } T_{\alpha\beta}^{*(e)} &\implies \lambda_1^{(e)} &= b^{y_t} \\ \lambda_1^{(o)} &= \text{largest eigenvalue of } T_{\alpha\beta}^{*(o)} &\implies \lambda_1^{(o)} &= b^{y_h}\end{aligned}$$

- ▶ Critical exponents written in terms of RG exponents  $y_t, y_h$ .

## REVIEW OF KEY CONCEPTS

- ▶ For  $d$  dimensions:

$$\alpha = 2 - \frac{d}{y_t}, \quad \beta = \frac{d - y_h}{y_t}, \quad \gamma = \frac{2y_h - d}{y_t}, \quad \delta = \frac{y_h}{d - y_h}, \quad \eta = d + 2 - 2y_h, \quad \nu = \frac{1}{y_t}$$

- ▶ For 2D Ising model:

$$y_t = 1, \quad y_h = \frac{15}{8} \quad \Longrightarrow \quad \alpha = 0, \quad \beta = \frac{1}{8}, \quad \gamma = \frac{7}{4}, \quad \delta = 15, \quad \eta = \frac{1}{4}, \quad \nu = 1$$

## MA'S METHOD

- ▶ First MCRG approach (1976).
- ▶ Basic idea:
  - ▶ Explicitly flow towards fixed point in truncated coupling space.
  - ▶ Construct  $T_{\alpha\beta}^*$  directly from renormalized Hamiltonian.
- ▶ Drawbacks:
  - ▶ Truncation errors accumulate quickly.
  - ▶ Monte Carlo simulation involves many couplings.



## MA'S METHOD

Approximate  $K^*$

Equilibrate spins at  $K^*$

Sample configurations at  $K^*$

Calculate  $T_{\alpha\beta}^* \rightarrow$  critical exponents

## SWENDSEN'S METHOD

Approximate  $K_1^c$

Equilibrate spins at  $K = (K_1^c, 0)$

Sample configurations at  $K$

Apply  $n$  RG transformations

Calculate  $T_{\alpha\beta}^* \rightarrow$  critical exponents

## SWENDSEN'S METHOD – CALCULATING CRITICAL COUPLING

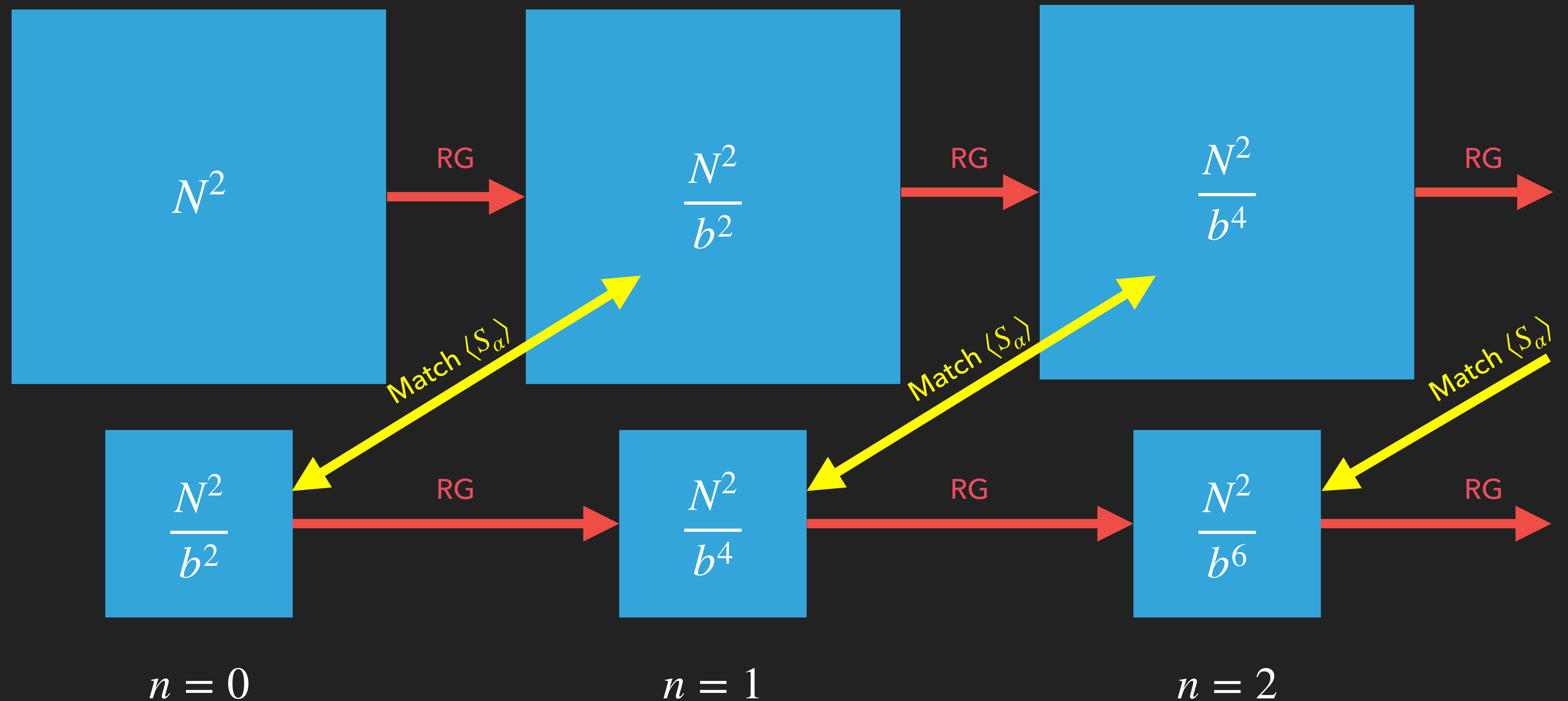
- ▶ Match correlation functions for transformed and untransformed systems.
- ▶ Cancel unknown finite size effects by comparing lattices of different sizes but the same number of lattice sites:

$$\langle S_{\alpha}^{(n+m)} \rangle_L = \langle S_{\alpha}^{(n)} \rangle_S,$$

where  $L = N^2$  and  $S = N^2/b^{2m}$ .

- ▶ Most effective for large lattices and large  $n$ .

## SWENDSEN'S METHOD – CALCULATING CRITICAL COUPLING



## SWENDSEN'S METHOD – CALCULATING CRITICAL COUPLING

- ▶ Linear approximation:

$$\langle S_{\alpha}^{(n+m)} \rangle_L - \langle S_{\alpha}^{(n)} \rangle_S = \left[ \frac{\partial \langle S_{\alpha}^{(n+m)} \rangle_L}{\partial K_1^{(0)}} - \frac{\partial \langle S_{\alpha}^{(n)} \rangle_S}{\partial K_1^{(0)}} \right] \delta K_1^{(0)},$$

where

$$\frac{\partial \langle S_{\alpha}^{(n)} \rangle}{\partial K_1^{(0)}} = \langle S_{\alpha}^{(n)} S_1^{(0)} \rangle - \langle S_{\alpha}^{(n)} \rangle \langle S_1^{(0)} \rangle.$$

# SWENDSEN'S METHOD – CALCULATING CRITICAL EXPONENTS

- ▶ Chain rule and invert:

$$\frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\beta^{(n)}} = \frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\alpha^{(n+1)}} \frac{\partial K_\alpha^{(n+1)}}{\partial K_\beta^{(n)}} \implies T_{\alpha\beta}^{(n)} = \left[ \frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\alpha^{(n+1)}} \right]^{-1} \frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\beta^{(n)}}$$

- ▶ Calculate derivatives using MC samples:

$$\frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\alpha^{(n+1)}} = \langle S_\gamma^{(n+1)} S_\alpha^{(n+1)} \rangle - \langle S_\gamma^{(n+1)} \rangle \langle S_\alpha^{(n+1)} \rangle$$

$$\frac{\partial \langle S_\gamma^{(n+1)} \rangle}{\partial K_\beta^{(n)}} = \langle S_\gamma^{(n+1)} S_\beta^{(n)} \rangle - \langle S_\gamma^{(n+1)} \rangle \langle S_\beta^{(n)} \rangle$$



## SWENDSEN'S METHOD – CALCULATING CRITICAL EXPONENTS

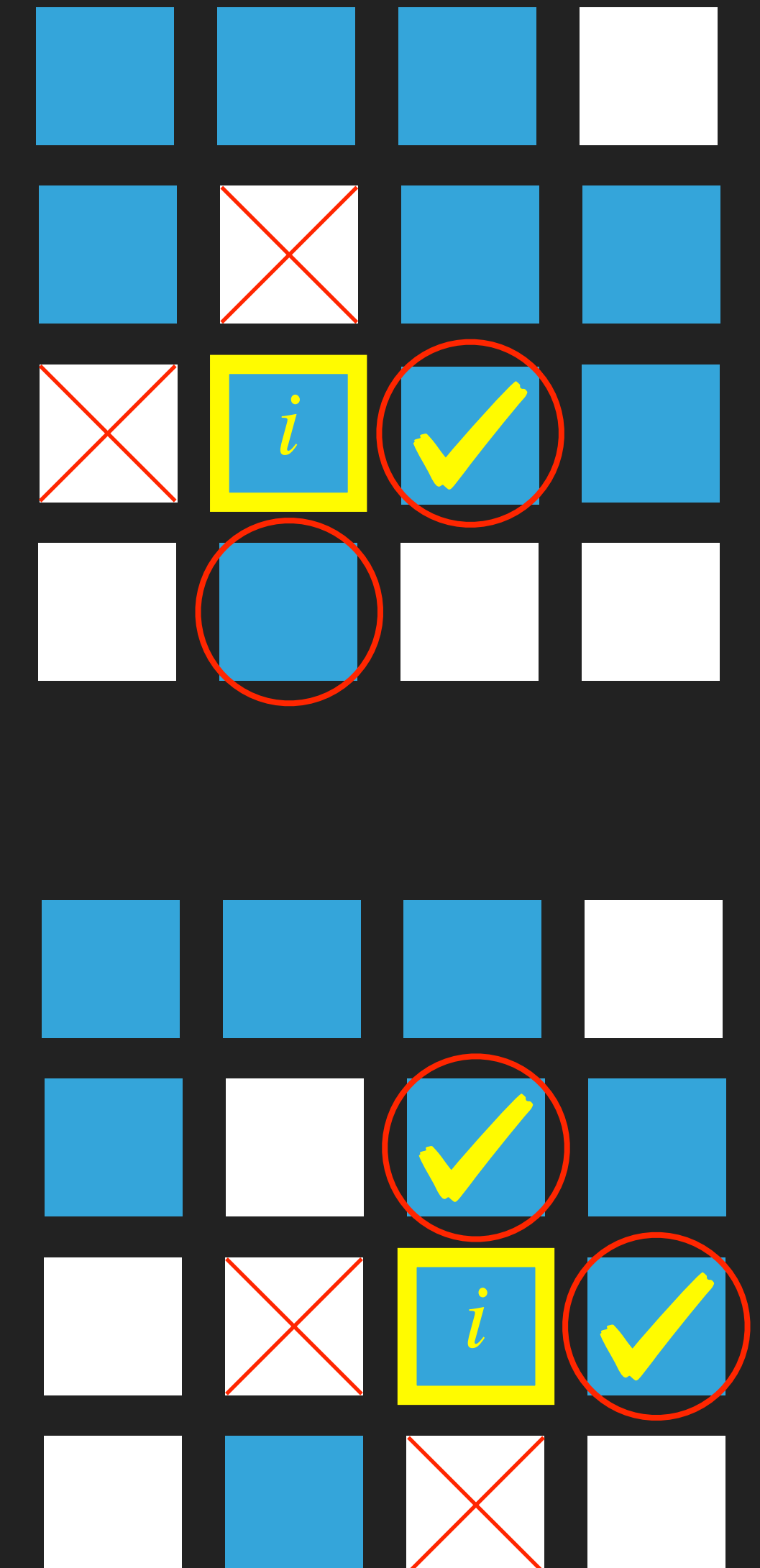
- ▶ Need to construct both even and odd blocks for  $\beta$  and  $\gamma$ .
- ▶ Only need even block for  $\alpha$  and  $\nu$ .
- ▶ Only need odd block  $\delta$  and  $\eta$ .
- ▶ For 2D Ising model example, consider nearest-neighbor and next-nearest-neighbor spin interactions only.
- ▶ Again, calculation of  $T_{\alpha\beta}^{(n)}$  most accurate for large lattices and large  $n$ .

## MONTE CARLO SIMULATION

- ▶ At criticality, correlation length diverges.
- ▶ Standard Metropolis updates are too local to efficiently simulate Ising system.
- ▶ Need cluster algorithms:
  - ▶ Swendsen-Wang algorithm
  - ▶ Wolff algorithm
- ▶ Wolff update has a larger probability of flipping a larger cluster of spins compared to a Swendsen-Wang update.

## WOLFF ALGORITHM

1. Choose random site  $i$  to add to the initially empty cluster.
2. Search through nearest neighbor sites  $j$ . If  $\sigma_i = \sigma_j$ , add site  $j$  to the cluster with probability  $P = 1 - e^{2K_1}$ .
3. Flip spin at site  $i$  and remove from the cluster.
4. If cluster is still empty, go back to step 1.
5. If cluster is not empty, repeat steps 2-5 until cluster stops growing.



## EXAMPLE: 2D ISING MODEL

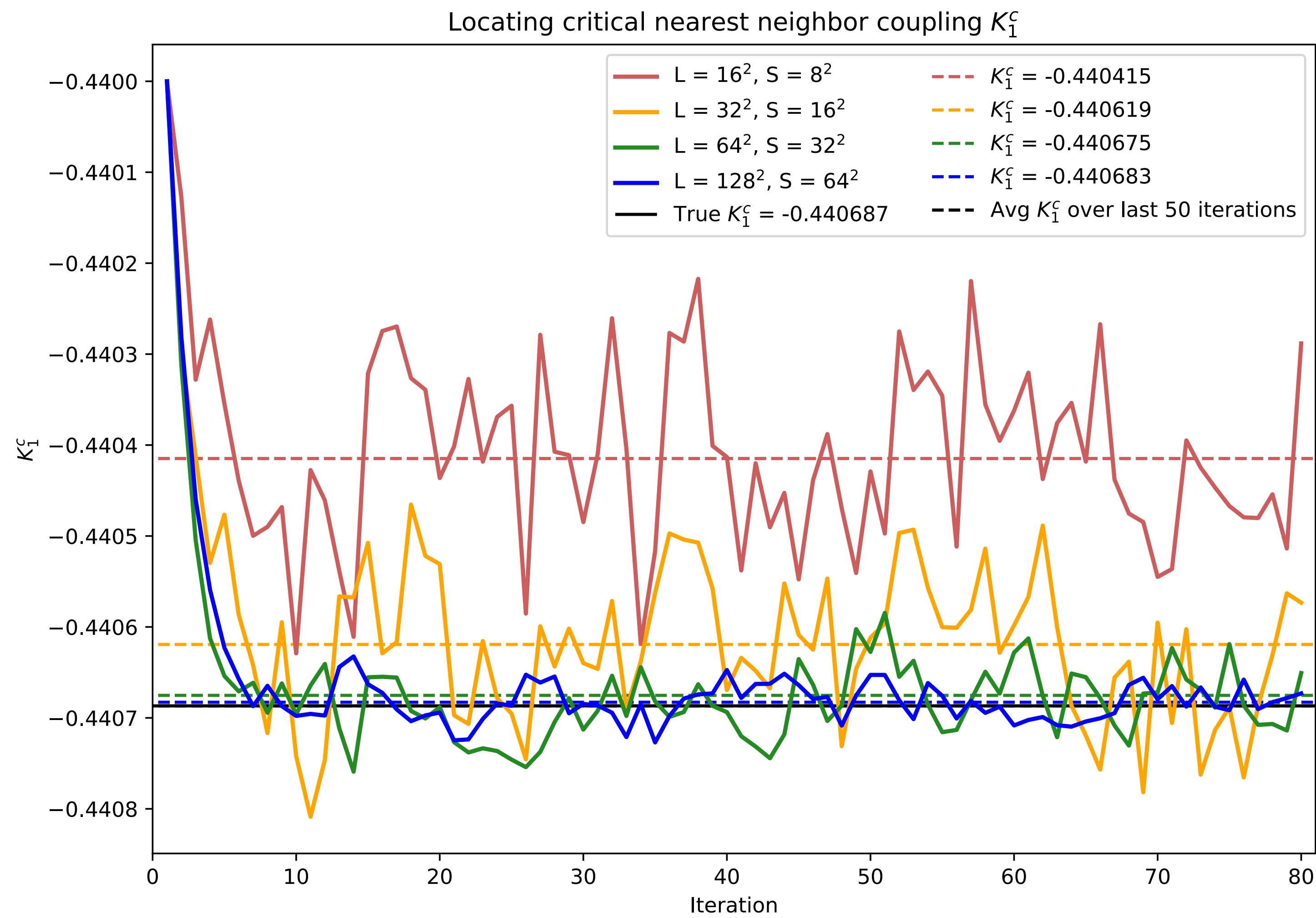
- ▶ Simple implementation:
  - ▶ Sample from pure nearest-neighbor Ising Hamiltonian using Wolff algorithm.
  - ▶ Majority-rule block spin transformation with scale factor  $b$ .
  - ▶ Use only nearest-neighbor spin correlation functions to tune  $K_1^c$ .
  - ▶ Use nearest-neighbor and next-nearest-neighbor spin correlations to construct truncated  $T_{\alpha\beta}^{(e)}$ .
  - ▶ Extract largest even eigenvalue  $\lambda_1^{(e)}$  and calculate  $\nu = \ln b / \ln \lambda_1^{(e)}$ .

## EXAMPLE: 2D ISING MODEL

- ▶ 1E4 samples during equilibration.
- ▶ 1E6 samples for locating critical coupling.
- ▶ 1E7 samples for calculating critical exponent.
- ▶ Calculated  $\nu$  at  $K_1^c(N)$  and  $K_1^c(N \rightarrow \infty)$ .
- ▶ Recall Onsager's solution:

$$T_c = \frac{2J}{\ln(1 + \sqrt{2})} \implies K_c \equiv K_1^c(N \rightarrow \infty) = -\frac{\ln(1 + \sqrt{2})}{2} \approx -0.440687$$





## RESULTS

- First iteration of locating critical coupling  $K_1^c$  from starting value  $K_1 = -0.44$  for various blocking levels and lattice sizes:

n	N = 16	N = 32	N = 64	N = 128
0	-0.4420811188	-0.4417429398	-0.4415445759	-0.4413292658
1	-0.4404791358	-0.4404577112	-0.4404935597	-0.4404266727
2	-0.4401272033	-0.4402638848	-0.4403468642	-0.4403131062
3	-	-0.4402975527	-0.4403307834	-0.4403016019
4	-	-	-0.4403116430	-0.4402985574
5	-	-	-	-0.4402807798

BENCHMARK

Table 9.1 *Variation of the thermal eigenvalue exponent for the Ising square lattice with the number of couplings  $N_c$ , the number of iterations  $N_r$ , and for different lattice sizes. From Swendsen (1982).*

$N_r$	$N_c$	$L = 64$	$L = 32$	$L = 16$
1	1	0.912(2)	0.904(1)	0.897(3)
1	2	0.967(3)	0.966(2)	0.964(3)
1	3	0.968(2)	0.968(2)	0.966(3)
1	4	0.969(4)	0.968(2)	0.966(3)
2	1	0.963(4)	0.953(2)	0.937(3)
2	2	0.999(4)	0.998(2)	0.993(3)
2	3	1.001(4)	1.000(2)	0.994(3)
2	4	1.002(5)	0.998(2)	0.984(4)
3	1	0.957(2)	0.936(3)	0.921(5)
3	2	0.998(2)	0.991(3)	1.013(4)
3	3	0.999(2)	0.993(3)	1.020(3)
3	4	0.997(2)	0.987(4)	...

## RESULTS

- Critical exponent  $\nu$  for  $N = 16, 32$  and various blocking levels at known critical coupling  $K_1^c$  and calculated critical coupling  $K_1^c(N)$ :

n	N = 16 @ $K_1^c(N)$	N = 16 @ $K_1^c$	N = 32 @ $K_1^c(N)$	N = 32 @ $K_1^c$
0	1.0366474702	1.0371936855	1.0359594271	1.0364461518
1	1.0055084088	1.0058615220	1.0030411491	1.0024276306
2	0.9850458988	0.9832937022	1.0042751683	1.0040962187
3	-	-	0.9833658883	0.9817702370



## RESULTS

- Critical exponent  $\nu$  for  $N = 64, 128$  and various blocking levels at known critical coupling  $K_1^c$  and calculated critical coupling  $K_1^c(N)$ :

n	N = 64 @ $K_1^c(N)$	N = 64 @ $K_1^c$	N = 128 @ $K_1^c(N)$	N = 128 @ $K_1^c$
0	1.0340840722	1.0356603405	1.0343762634	1.0349921920
1	1.0013550363	1.0025377191	1.0018397044	1.0006737634
2	1.0021687106	1.0029573471	1.0019303651	1.0022232598
3	1.0029619937	1.0039006140	1.0025651221	1.0023411428
4	0.9820802731	0.9819506190	1.0035789853	1.0040042374
5	-	-	0.9814148379	0.9819137809



## CONCLUSION

- ▶ Swendsen's MCRG method can efficiently extract critical exponents, even in an aggressively truncated coupling space.
- ▶ Truncation errors are small and systematically improvable by including higher order couplings and by simulating larger lattices.
- ▶ Computational costs are kept low by only requiring the simulations of pure nearest-neighbor Ising systems.

## BIBLIOGRAPHY

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**QUESTIONS?**