

# Calibrating Tight Binding for Biaxially Strained Graphene

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We study the effect of **uniform biaxial strain** on monolayer graphene using a minimal **tight-binding model** with nearest and next-nearest neighbor hopping, treating the carbon–carbon distance  $a$  as a strain parameter while keeping hopping amplitudes  $t$  and  $t_2$  fixed. This strain-independent parametrization fails to match experimental trends. We introduce a distance-dependent hopping  $t(a)$  of **exponential form at equilibrium**. This modification significantly enhances the model's accuracy in describing the strain dependence of low-energy observables in graphene over a realistic biaxial strain range.

## How can we strain graphene?

Graphene is a semimetal with a honeycomb lattice of carbon atoms. At equilibrium, the carbon–carbon bond length is ( $a_0 = 1.42 \text{ \AA}$ ) and the band structure is gapless. We varied from  $a = 1.36 \text{ \AA}$  (compressive strain) to  $a = 1.56 \text{ \AA}$  (tensile strain).

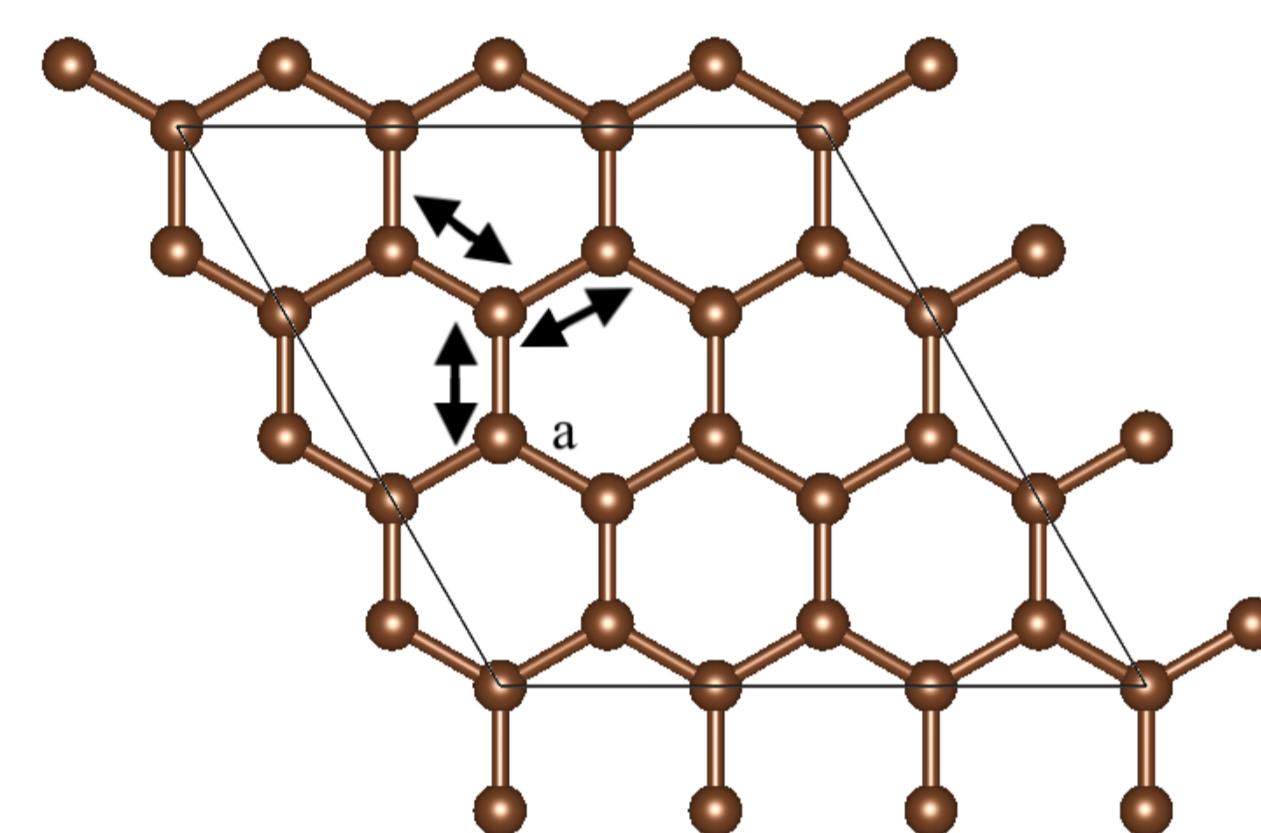


Figure 1. Graphene: a honeycomb lattice of carbon atoms.

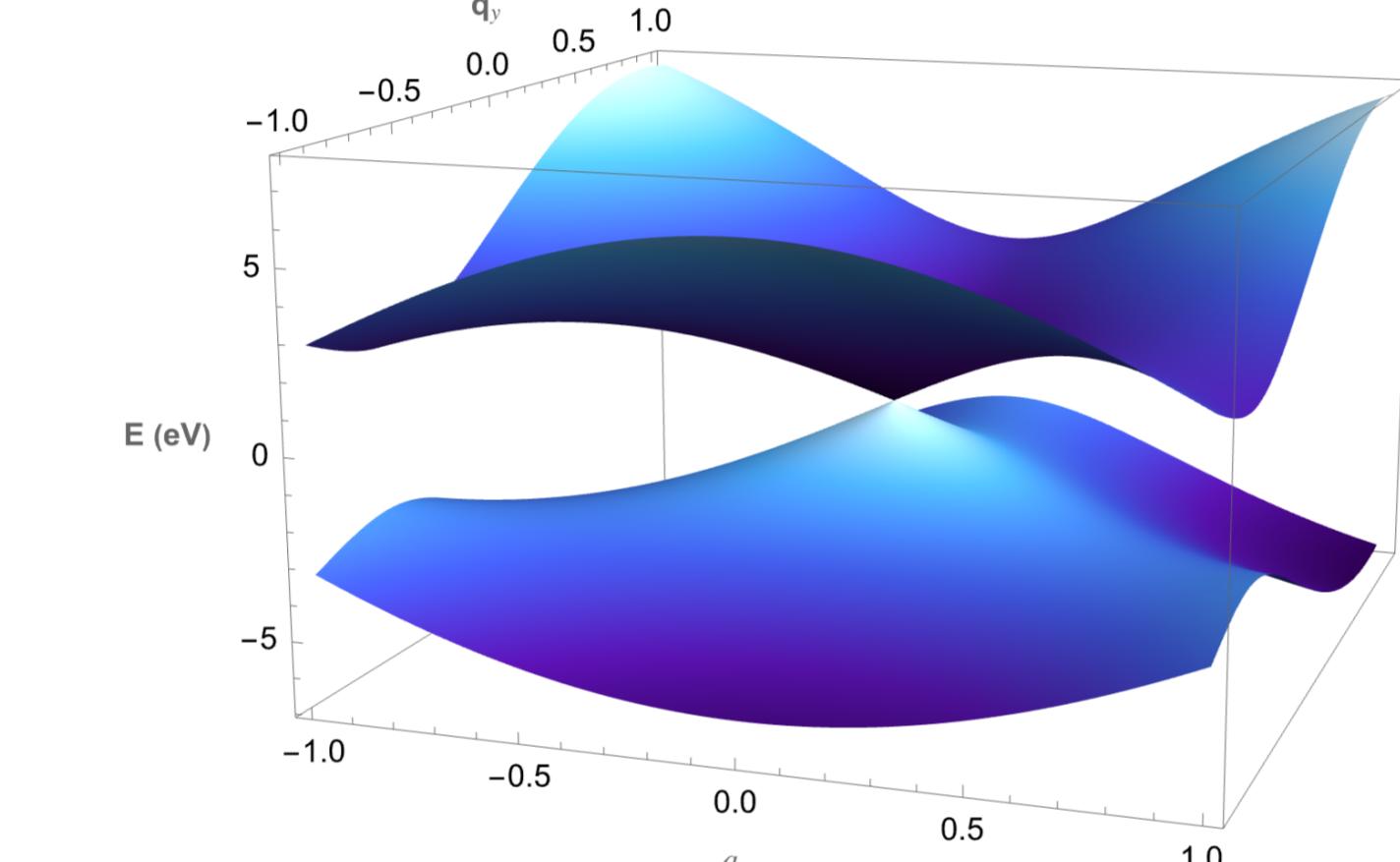


Figure 2. Energy  $E$  versus crystal momentum  $\mathbf{k}$  near the  $K$  point of graphene's Brillouin zone ( $a = 1.42 \text{ \AA}$ ).

## Two Tight Binding Models

The band structure of graphene is distinctive because it exhibits a **Dirac cone** at the  $K$  point. Under uniform biaxial strain, we study how this Dirac cone changes within the tight-binding approximation using two models:

- **Model A (geometry-only strain):** the bond length  $a$  changes with strain, while the hopping parameters  $t$  and  $t_2$  are kept fixed.
- **Model B (distance-dependent hopping):** the bond length  $a$  changes with strain, and the hopping parameters decay with distance,

$$t(a) = t_0 \exp\left[-\beta\left(\frac{a}{a_0} - 1\right)\right], \quad t_2(a) = t_{2,0} \exp\left[-\beta\left(\frac{a}{a_0} - 1\right)\right].$$

We employ a minimal tight-binding model that includes both nearest- and next-nearest-neighbor interactions.

## Why strain matters

### around the Dirac cone...

Under uniform biaxial strain, the carbon–carbon bond length  $a$  changes. In a next-nearest neighbor (NNN) tight binding (TB) model with distance-dependent hopping, stretching ( $a \uparrow$ ) reduces orbital overlap and typically lowers the slope of the cone (Fermi velocity), while compression ( $a \downarrow$ ) increases it. Graphene is gapless at  $K$ , so the conduction and valence bands touch and form a cone-like dispersion.

Comparison of graphene Dirac cones using the two parameterizations: constant  $t$  versus  $t(a)$ .

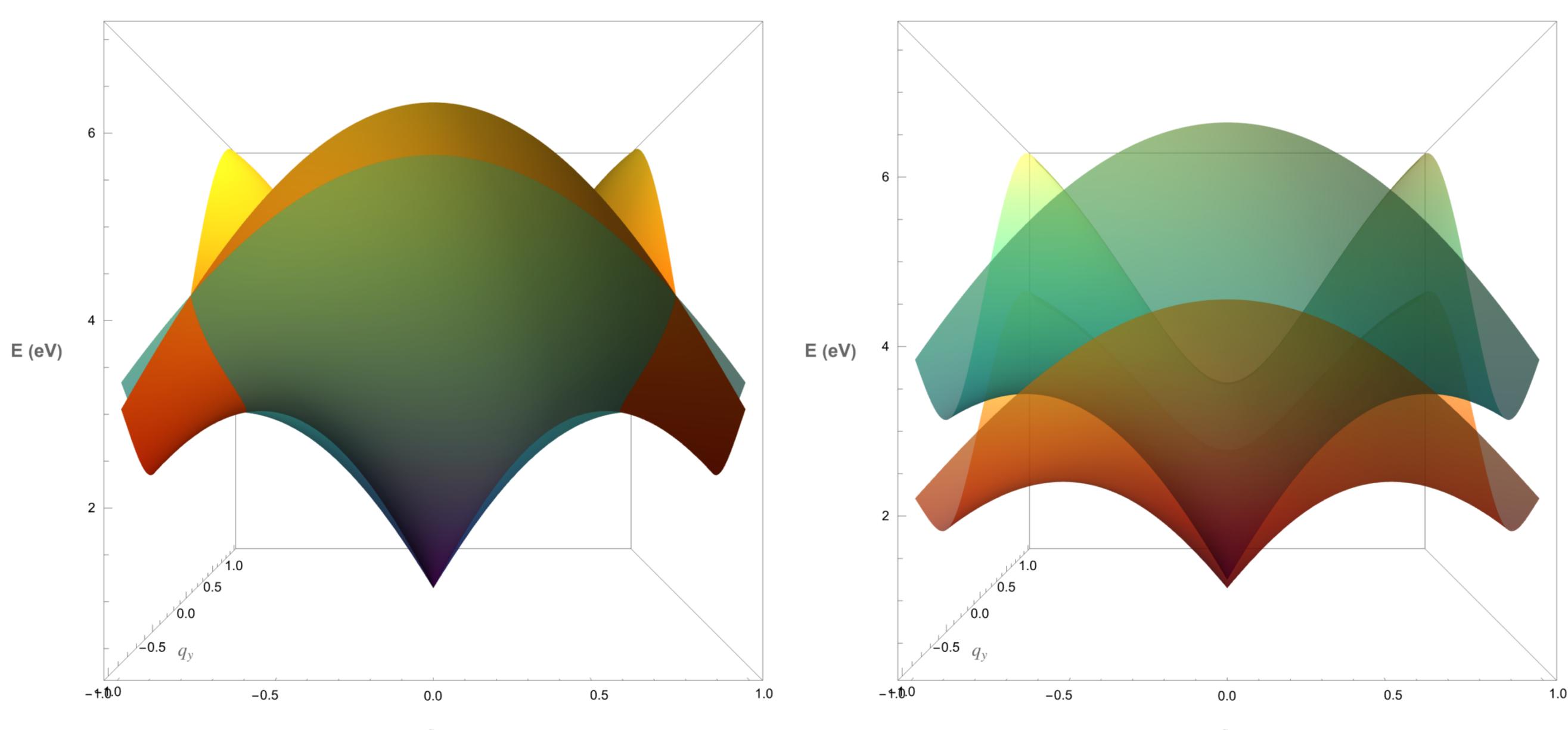


Figure 3. 3D Dirac cones for extreme biaxial strain ( $a = 1.36 \text{ \AA}$  compression and  $a = 1.56 \text{ \AA}$  tension) using the NNN tight-binding model with constant  $t$

Figure 4. 3D Dirac cones for the same strains using distance-dependent hopping  $t(a)$  (and  $t_2(a)$ )

Calculations were performed in Wolfram Mathematica, and the code is available at <https://kimreyesg.github.io/research.html>.

**Fermi velocity from the Dirac cone (tight binding)** Near a Dirac point  $K$ , the dispersion is approximately linear:

$$E_+(\mathbf{K} + \mathbf{q}) \approx \hbar v_F |\mathbf{q}|.$$

Therefore, the Fermi velocity is the slope of the conduction cone at  $K$ :

$$v_F = \frac{1}{\hbar} \frac{\partial E_+}{\partial |\mathbf{q}|} \Big|_{\mathbf{q}=0}$$

In the nearest-neighbor tight-binding model this gives the standard result:

$$v_F = \frac{3at}{2\hbar}$$

where  $a$  is the C–C bond length and  $t$  is the nearest-neighbor hopping.

Table 1. Fermi velocity for different biaxial strains in graphene.

$a (\text{\AA})$	$\varepsilon = (a - a_0)/a_0 (\%)$	$v_F$ (constant $t$ , $10^5 \text{ m s}^{-1}$ )	$v_F$ ( $t(a)$ , $10^5 \text{ m s}^{-1}$ )
1.36	-4.2	8.37	9.65
1.42	0.0	8.74	8.74
1.56	+9.9	9.60	6.89

## Final Results

- **Unstrained check:** We obtain  $v_F(a_0) \approx 8.7 \times 10^5 \text{ m s}^{-1}$ , within  $\sim 10\%-20\%$  of the commonly quoted  $v_F^{\text{exp}} \sim (1.0-1.1) \times 10^6 \text{ m s}^{-1}$ .
- **Why the offset is expected:** A minimal nearest and next-nearest neighbor tight binding model (with  $t_0 = 2.7 \text{ eV}$ ) omits many-body and environmental effects (substrates, screening).
- **Model B is better:** Distance-dependent hopping  $t(a)$  accounts for reduced orbital overlap under tension and corrects the qualitative strain trend (the constant- $t$  model can predict an unrealistic increase of  $v_F$  with tensile strain).

## References

- [1] A. H. Castro Neto, F. Guinea, N. M. R. Peres, K. S. Novoselov, and A. K. Geim. The electronic properties of graphene. *Reviews of Modern Physics*, 81(1):109–162, 2009. doi: 10.1103/RevModPhys.81.109.