



UNIVERSITY OF AMSTERDAM

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# On Statistical Inference and Modelling Data with the Ising Model

Applications to Neuronal Activity and US Supreme Court Voting Patterns

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**Student:**

Tycho Stam (13303147)

**Lecturer:**

dr. C.M.C. de de Mulatier

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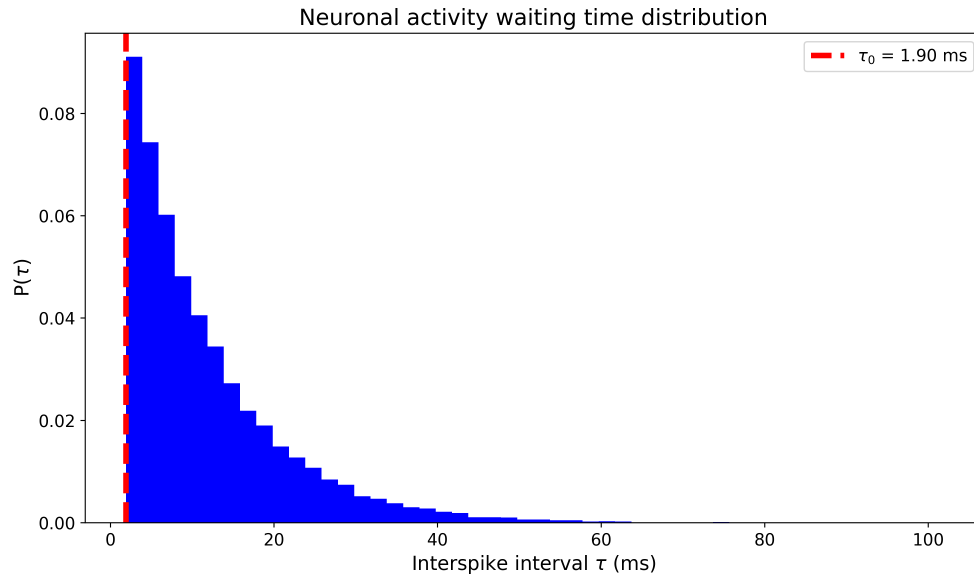
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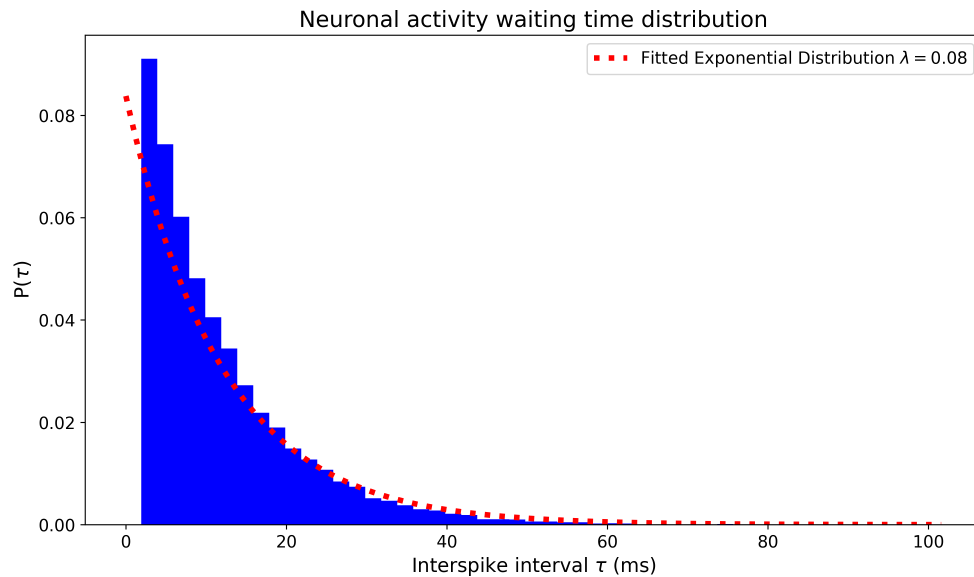
# 1 Modeling the activity of a single neuron

**Q1.** Fig. 1 displays the distribution  $P(\tau)$  of time intervals  $\tau$  between consecutive spikes, as derived from the initial dataset.  $\tau_0$  is hereby equal to 1.90ms.



**Figure 1:** The distribution  $P(\tau)$  of time intervals  $\tau$  between consecutive spikes. Hereby  $\tau_0$  shows the refractory period, in which time a neuron does not spike again. For this plot 50 bins were used to group the values.

**Q2.** When fitting the data on an exponential function, it gives back a pattern similar to the distribution  $P(\tau)$ , based on  $\lambda = 0.08$ , as seen in Fig. 2.

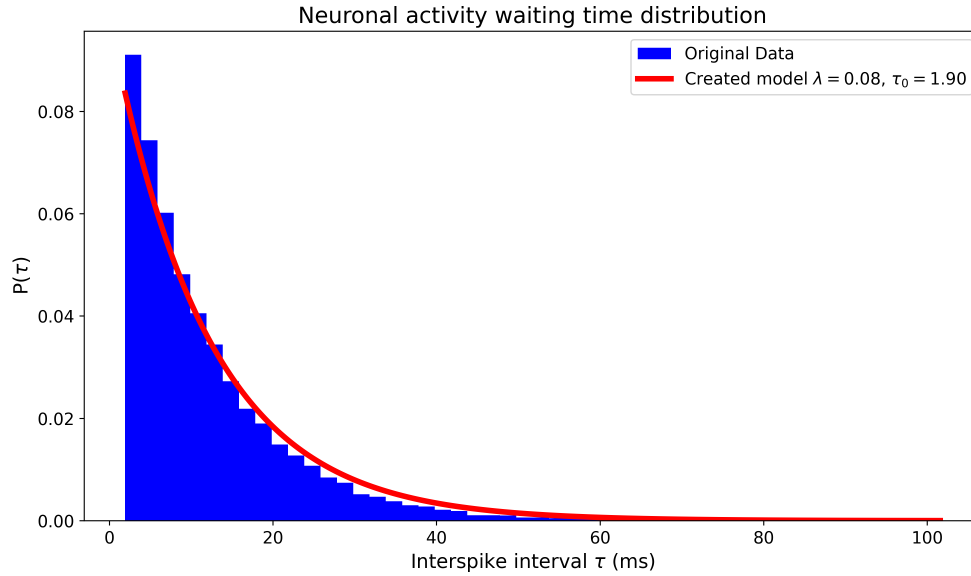


**Figure 2:** Fitted exponential function of distribution  $P(\tau)$ , resulting in  $\lambda = 0.08$ .

**Q3.** Let  $\tau$  be the time between two spikes, given what we know about the refractory period we get the following inter-spike interval distribution:

$$P(\tau) = \begin{cases} \lambda \exp[-\lambda(\tau - \tau_0)] & \text{if } \tau \geq \tau_0 \\ 0 & \text{if } \tau < \tau_0 \end{cases} \quad (1)$$

When comparing the created model with the original data we get the following, as seen in Fig. 3.



**Figure 3:** The model data (in red) is plotted over the original dataset, clearly illustrating that the model closely aligns with the observed data.

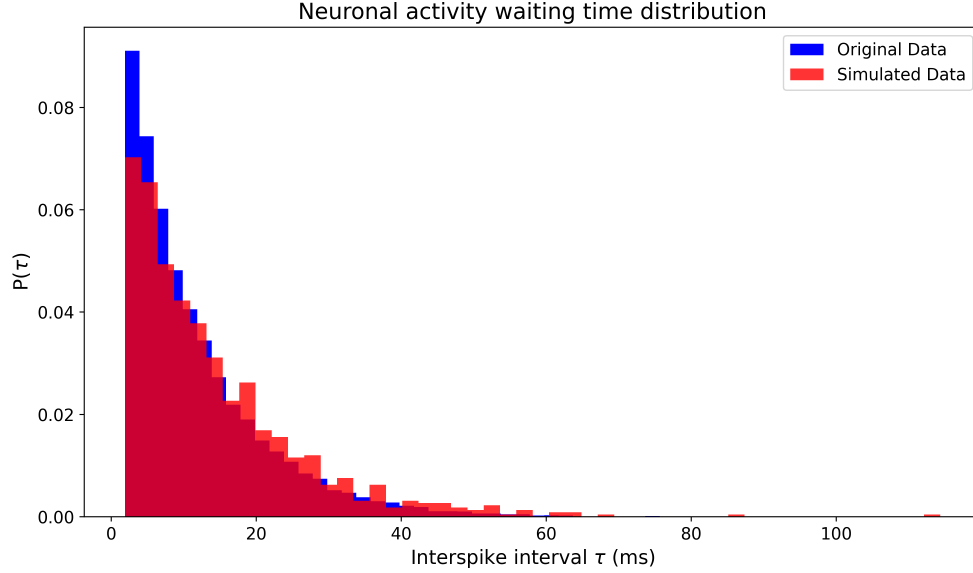
**Q4.** To sample from the PDF, we will use inverse transformation sampling. For this the CDF of  $P(\tau)$  was created:

$$F(\tau) = \begin{cases} 1 - \exp[-\lambda(\tau - \tau_0)] & \text{if } \tau \geq \tau_0 \\ 0 & \text{if } \tau < \tau_0 \end{cases} \quad (2)$$

Hereby we sample from  $u_k \sim \text{Uniform}(0, 1)$  and solve  $u_k = F(\tau)$ , giving us:

$$\tau_k = \tau_0 - \frac{\ln(1 - u_k)}{\lambda} \quad (3)$$

Here,  $k$  denotes the index of the  $k$ -th spike out of a total of  $N = 1000$  spikes. This results in Fig. 4.



**Figure 4:** Comparison of the simulated data and original data, sampled using a inverse transformation. Whereby the simulated dataset contains  $N = 1000$  samples.

**Q5.** For the exponential distribution we know that  $\mathbb{E}[X] = \frac{1}{\lambda}$ . Given Eq. 1, this will result in  $\mathbb{E}[\tau] = \frac{1}{\lambda} + \tau_0$ . The average firing rate is then defined as:

$$f = \frac{1}{\mathbb{E}[\tau]} = \frac{1}{\tau_0 + \frac{1}{\lambda}} = \frac{\lambda}{1 + \tau_0 \lambda} \quad (4)$$

When plugging in  $\lambda$  and  $\tau_0$  we get the average spiking rate of 72.3 Hz

## 2 Modeling binary data with the Ising model

### 2.1 Pairwise spin model

For the pairwise spin model we will take the probability distribution to have the general form of an Ising model:

$$p_{\mathbf{g}}(s) = \frac{1}{Z(\mathbf{g})} \exp \left( \sum_{i=1}^n h_i s_i + \sum_{\text{pair}(i,j)} J_{ij} s_i s_j \right) \quad (5)$$

where  $n$  is the number of spin variables, where  $\text{pair}(i,j)$  denotes a summation over all possible pairs of distinct spin variables, where  $\mathbf{g} = (h_1, \dots, h_n, J_{1,2}, \dots, J_{n-1,n})$  is a vector of (real) parameters, and where  $Z(\mathbf{g})$  is a normalization factor.

**Q1.1.** Given the sum over  $\text{pair}(i,j)$ , where  $J_{ij}$  is the strength of the coupling between. We will have  $n - 1$  terms, where  $n$  is the number of spins in the model. This is because each  $i$  will be coupled with all of the other spins, and this scales with the number of spins. Because of the same reasoning the vector  $\mathbf{g}$  will have  $n + \binom{n}{2}$  parameters. The first term corresponds to the number of  $h$  parameters in the model, while  $\binom{n}{2}$  represents the number of  $J_{i,j}$  interactions.

**Q1.2.** When  $n = 3$  in Eq. 5 we get the following expression:

$$\mathbf{g} = (h_1, h_2, h_3, J_{1,2}, J_{1,3}, J_{2,3})$$

$$p_{\mathbf{g}}(s) = \frac{1}{Z(\mathbf{g})} \exp (h_1 s_1 + h_2 s_2 + h_3 s_3 + J_{12} s_1 s_2 + J_{13} s_1 s_3 + J_{23} s_2 s_3) \quad (6)$$

**Q1.3.** The Boltzmann distribution is defined as:

$$P(s) = \frac{\exp[-\beta E(s)]}{Z} \quad (7)$$

where  $\beta = \frac{1}{k_b T}$ . Given that  $\beta = 1$ , the energy function ( $E(s)$ ) is defined as follows:

$$E(s) = - \left( \sum_{i=1}^n h_i s_i + \sum_{pair(i,j)} J_{ij} s_i s_j \right) \quad (8)$$

and the partition function  $Z$ :

$$Z = \sum_s \exp \left( \sum_{i=1}^n h_i s_i + \sum_{pair(i,j)} J_{ij} s_i s_j \right) \quad (9)$$

The goal of the partition function is to normalize the constants over all  $s$ .

**Q1.4**  $h_i$  represents the external influences on the individual judges, such as political bias. If  $h_i$  is positive, the direction of  $s_i$  will also want to be positive, to give the lowest associated energy  $-h_i s_i$ . Hereby  $J_{ij}$  represents the interaction between the connected judges, if it is positive the judges will most likely vote the same way. So if  $J_{ij}$  is positive the alignment of the spins should be in the same direction (i.e. both negative or positive), as this will make the product 1 and will minimize the coupling energy  $-J_{ij} s_i s_j$ .

## 2.2 Observables

**Q2.1.** Given the stationary probability distribution of the state  $p_g(s)$ , we can define  $\langle s_i \rangle$  and  $\langle s_i s_j \rangle$  as follows:

$$\langle s_i \rangle = \sum_s s_i p_g(s) \quad (10)$$

$$\langle s_i s_j \rangle = \sum_s s_i s_j p_g(s) \quad (11)$$

**Q2.2.** To calculate the empirical averages denoted by  $\langle s_i \rangle_D$  and  $\langle s_i s_j \rangle_D$  over the whole dataset, we can use the following equations:

$$\langle s_i \rangle_D = \frac{1}{N} \sum_{k=1}^N s_i^k$$

$$\langle s_i s_j \rangle_D = \frac{1}{N} \sum_{k=1}^N s_i^k s_j^k$$

Where  $N$  is the amount of samples in the dataset  $\hat{s}$ .

**Q2.3.** The samples  $s^k$  are defined as independent and identically distributed (i.i.d.) and are drawn from the distribution  $p(s)$ . Given the Law of Large Numbers we can state that:

$$\lim_{N \rightarrow \infty} \langle s_i \rangle_D = \mathbb{E}_{p(s)}[s_i] = \langle s_i \rangle \quad (12)$$

The same property holds true for  $\langle s_i s_j \rangle_D$ , giving us:

$$\lim_{N \rightarrow \infty} \langle s_i s_j \rangle_D = \mathbb{E}_{p(s)}[s_i s_j] = \langle s_i s_j \rangle \quad (13)$$

## 2.3 Maximum Entropy models

**Q3.1.** The equation for Shannon entropy is defined as follows:

$$S[p(s)] = - \sum_s p(s) \log p(s) \quad (14)$$

Whereby  $p(s) = p_{\mathbf{g}}(s)$  and  $p_{\mathbf{g}}(s)$  is defined in Eq. 5.

**Q3.2.** Based the following sets constraints:

$$\underbrace{\sum_s p(s) = 1}_{\text{normalization constraint}} \quad \text{and} \quad \underbrace{\sum_s p(s) s_i(s) = \langle s_i \rangle_D}_{\text{magnetization constraint}} \quad \text{and} \quad \underbrace{\sum_s p(s) s_i(s) s_j(s) = \langle s_i s_j \rangle_D}_{\text{correlation constraint}} \quad (15)$$

We will define the total number of constraints as  $m$ . The normalization constraint contributes 1 constraint. The magnetization constraints contribute  $n$  constraints, one for each  $i$ . Finally, the correlation constraints contribute  $\binom{n}{2}$  constraints, one for each unique pair. Therefore, the total number of constraints is:

$$m = 1 + n + \binom{n}{2}$$

For  $n = 3$ , this will result in:

$$\begin{aligned} m &= 4 + \frac{3!}{2!(3-2)!} \\ &= 4 + \frac{6}{2} \\ &= 7 \end{aligned}$$

**Q3.3.** Given the auxiliary function:

$$\begin{aligned} U[p(s)] &= S[p(s)] + \lambda_0 \left( \sum_s p(s) - 1 \right) + \sum_{i=1}^n \alpha_i \left( \sum_s p(s) s_i(s) - \langle s_i \rangle_D \right) \\ &\quad + \sum_{\text{pair}(i,j)}^n \eta_{ij} \left( \sum_s p(s) s_i(s) s_j(s) - \langle s_i s_j \rangle_D \right) \end{aligned} \quad (16)$$

We can substitute Eq. 14 in Eq. 16 to get:

$$\begin{aligned} U[p(s)] &= - \sum_s p(s) \log p(s) + \lambda_0 \left( \sum_s p(s) - 1 \right) + \sum_{i=1}^n \alpha_i \left( \sum_s p(s) s_i(s) - \langle s_i \rangle_D \right) \\ &\quad + \sum_{\text{pair}(i,j)}^n \eta_{ij} \left( \sum_s p(s) s_i(s) s_j(s) - \langle s_i s_j \rangle_D \right) \end{aligned} \quad (17)$$

We will set the fixed state of  $p(s)$  to  $p_s$  and give it one value of  $s$ , removing the summations and giving us:

$$\begin{aligned} U[p_s] &= -p_s \log p_s + \lambda_0 (p_s - 1) + \sum_{i=1}^n \alpha_i (p_s s_i(s) - \langle s_i \rangle_D) \\ &\quad + \sum_{\text{pair}(i,j)}^n \eta_{ij} (p_s s_i(s) s_j(s) - \langle s_i s_j \rangle_D) \end{aligned} \quad (18)$$

When taking the derivative  $\frac{\partial}{\partial p_s}$  over each term in Eq. 18:

$$\begin{aligned} \frac{\partial U[p_s]}{\partial p_s} &= \frac{\partial}{\partial p_s} - p_s \log p_s + \frac{\partial}{\partial p_s} \lambda_0 (p_s - 1) + \frac{\partial}{\partial p_s} \sum_{i=1}^n \alpha_i (p_s s_i(s) - \langle s_i \rangle_D) \\ &\quad + \frac{\partial}{\partial p_s} \sum_{pair(i,j)}^n \eta_{ij} (p_s s_i(s) s_j(s) - \langle s_i s_j \rangle_D) \\ &= -\log(p_s) - 1 + \lambda_0 + \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \end{aligned} \quad (19)$$

Proven that the equation stated in the assignment holds.

**Q3.4.** To maximize the equation we set Eq. 19 to 0:

$$\begin{aligned} 0 &= \frac{\partial U[p_s]}{\partial p_s} \\ 0 &= -\log(p_s) - 1 + \lambda_0 + \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \end{aligned}$$

Solving this equation for  $p_s$  gives us:

$$\begin{aligned} \log(p_s) &= -1 + \lambda_0 + \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \\ p_s &= \exp \left( -1 + \lambda_0 + \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right) \\ p_s &= \exp(-1 + \lambda_0) \exp \left( \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right) \end{aligned} \quad (20)$$

When substituting the normalization constraint

$$\sum_s p_s = 1$$

in Eq. 20, we get:

$$\begin{aligned} \sum_s p_s &= 1 \\ \sum_s \exp(-1 + \lambda_0) \exp \left( \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right) &= 1 \\ \exp(-1 + \lambda_0) \sum_s \exp \left( \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right) &= 1 \\ \exp(-1 + \lambda_0) &= \frac{1}{\sum_s \exp \left( \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right)} \end{aligned}$$

Where

$$Z = \sum_s \exp \left( \sum_{i=1}^n \alpha_i s_i + \sum_{pair(i,j)}^n \eta_{ij} s_i(s) s_j(s) \right)$$



Resulting in:

$$\exp(-1 + \lambda_0) = \frac{1}{Z} \quad (21)$$

and

$$\alpha_i = h_i, \quad \eta_{ij} = J_{ij}$$

## 2.4 Statistical inference: model with no couplings

**Q4.1.** We know that the product of powers with the same base is equal to the base raised to the sum of the exponents, based on this relation we can rewrite

$$p_{\mathbf{g}}(s) = \frac{1}{Z(\mathbf{g})} \exp\left(\sum_{i=1}^n h_i s_i\right)$$

as:

$$p_{\mathbf{g}}(s) = \frac{1}{Z(\mathbf{g})} \prod_{i=1}^n \exp(h_i s_i)$$

We can rewrite that as:

$$p_{\mathbf{g}}(s) = \prod_{i=1}^n p_{h_i}(s_i) \quad (22)$$

Given that the distribution for each spin is normalized, a corresponding partition function  $Z_i$  exists for each spin:

$$Z_i = \exp(h_i) + \exp(-h_i) = 2 \cosh(h_i)$$

Results in:

$$p_{h_i} \equiv \frac{\exp(h_i s_i)}{2 \cosh(h_i)} \quad (23)$$

**Q4.2.** As stated the probability distribution  $p_{h_i}$  is the distribution for a spin variable  $s_i$ , resulting in:

$$\langle s_i \rangle = \sum_s p_{h_i} s_i$$

Where  $s \in \{-1, 1\}$ . Using Eq. 23 and the properties of the spin, we will get:

$$\begin{aligned} \langle s_i \rangle &= (1)p_{h_i} + (-1)p_{h_i} \\ &= \frac{\exp(h_i)}{2 \cosh(h_i)} - \frac{\exp(-h_i)}{2 \cosh(h_i)} \\ &= \frac{\exp(h_i)}{\exp(h_i) + \exp(-h_i)} - \frac{\exp(-h_i)}{\exp(h_i) + \exp(-h_i)} \\ &= \frac{\exp(h_i) - \exp(-h_i)}{\exp(h_i) + \exp(-h_i)} \\ &= \tanh(h_i) \end{aligned} \quad (24)$$

Applying the inverse of the function in Eq. 24 yields:

$$h_i = \tanh^{-1}(\langle s_i \rangle)$$

Given the constraint  $\langle s_i \rangle = \langle s_i \rangle_D$

$$h_i = \tanh^{-1}(\langle s_i \rangle_D) \quad (25)$$

**Q4.3.** Given that when  $\langle s_i \rangle_D > 0$ ,  $h_i$  will be positive, indicating that the  $i$ -th judge tends to vote positively overall, a similar pattern as described in question 1.4 will emerge. However, now each judge's vote will be influenced by the average votes of all the other judges.

## 2.5 Statistical inference: maximizing the log-likelihood function

**Q5.1.** The dataset  $\hat{s} = (s^1, s^2, \dots, s^N)$  consists of  $N$  independent samples drawn from the distribution  $p_{\mathbf{g}}(s)$ . Because of this property, we can rewrite the probability that the model  $p_{\mathbf{g}}(s)$  produces the dataset

$$\mathcal{L} = \log P_s(\hat{s})$$

as:

$$\begin{aligned} \mathcal{L} &= \log \left[ \prod_{k=1}^N p_{\mathbf{g}}(s^k) \right] \\ &= \sum_{k=1}^N \log [p_{\mathbf{g}}(s^k)] \end{aligned}$$

Instead of summing over all  $N$ , we can sum over the unique states  $s$ , resulting in a similar sum:

$$\mathcal{L} = \sum_s K(s) \log [p_{\mathbf{g}}(s^k)] \quad (26)$$

By definition we know that:

$$\begin{aligned} p_D &= \frac{K(s)}{N} \\ K(s) &= N p_D \end{aligned}$$

When substituting this in Eq. 26 we get:

$$\begin{aligned} \mathcal{L} &= \sum_s N p_D \log [p_{\mathbf{g}}(s^k)] \\ &= N \sum_s p_D \log [p_{\mathbf{g}}(s^k)] \end{aligned} \quad (27)$$

**Q5.2.** First we will rewrite the log-likelihood function in terms of the Ising model:

$$\begin{aligned} \mathcal{L} &= N \sum_s p_D \log p_{\mathbf{g}}(s) \\ &= N \sum_s p_D \left( \sum_{i=1}^n h_i s_i + \sum_{\text{pair}(i,j)} J_{ij} s_i s_j - \log(Z) \right) \end{aligned} \quad (28)$$

First we will take the derivative of Eq. 28 with a respect to  $h_i$ :

$$\frac{\partial \mathcal{L}}{\partial h_i} = N \sum_s p_D \frac{\partial}{\partial h_i} \left( \sum_{i=1}^n h_i s_i + \sum_{\text{pair}(i,j)} J_{ij} s_i s_j - \log(Z) \right) \quad (29)$$

Hereby the first term:

$$\frac{\partial}{\partial h_i} \sum_{i=1}^n h_i s_i = s_i$$

The second term:

$$\frac{\partial}{\partial h_i} \sum_{\text{pair}(i,j)} J_{ij} s_i s_j = 0$$

And the last term, given Eq. 9:

$$\begin{aligned}
\frac{\partial}{\partial h_i} - \log(Z) &= \frac{1}{Z} \frac{\partial Z}{\partial h_i} \\
&= \frac{1}{Z} \frac{\partial}{\partial h_i} \left[ \sum_s \exp \left( \sum_{i=1}^n h_i s_i + \sum_{\text{pair}(i,j)}^n J_{ij} s_i s_j \right) \right] \\
&= \frac{1}{Z} \left[ \sum_s s_i \exp \left( \sum_{i=1}^n h_i s_i + \sum_{\text{pair}(i,j)}^n J_{ij} s_i s_j \right) \right] \\
&= \sum_s s_i p_{\mathbf{g}}(s) \\
&= \langle s_i \rangle
\end{aligned} \tag{30}$$

Making that:

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial h_i} &= N \sum_s p_D(s_i - \langle s_i \rangle) \\
&= N (\langle s_i \rangle_D - \langle s_i \rangle)
\end{aligned} \tag{31}$$

When setting the maximum of  $\mathcal{L}$ :

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial h_i} &= 0 \\
N (\langle s_i \rangle_D - \langle s_i \rangle) &= 0 \\
(\langle s_i \rangle_D - \langle s_i \rangle) &= 0 \\
\langle s_i \rangle_D &= \langle s_i \rangle
\end{aligned} \tag{32}$$

The same can be done with taking the derivative of Eq. 28 with a respect to  $J_{ij}$ , giving us the first term:

$$\frac{\partial}{\partial J_{ij}} \sum_{i=1}^n h_i s_i = 0$$

The second term:

$$\frac{\partial}{\partial J_{ij}} \sum_{\text{pair}(i,j)}^n J_{ij} s_i s_j = s_i s_j$$

and

$$\frac{\partial}{\partial h_i} - \log(Z) = \langle s_i s_j \rangle \tag{33}$$

for the third term with the same steps as Eq. 30, but than with a respect to  $J_{ij}$ . Giving us the maximum:

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial h_i} &= 0 \\
N (\langle s_i s_j \rangle_D - \langle s_i s_j \rangle) &= 0 \\
(\langle s_i s_j \rangle_D - \langle s_i s_j \rangle) &= 0 \\
\langle s_i s_j \rangle_D &= \langle s_i s_j \rangle
\end{aligned} \tag{34}$$

### 3 Application to the analysis of the US supreme Court

#### Q6.1.

**Q6.2.**

**Q6.3.**

**Q6.4.**

**Q6.5.**

**Q6.6.**

**Q6.7.**

**Q6.8.**

**Q6.9.**